

Magnetism in materials

Solutions - Week 04

1. Consider the case of two interacting spin- $\frac{1}{2}$ electrons. The good quantum numbers are $S = 0$ and 1 so that there is a triplet state and a singlet state which will be separated by an energy gap Δ . We define the sign of Δ so that when $\Delta > 0$ the singlet state ($S = 0$) is the lower state and when $\Delta < 0$ the triplet state is the lower state. These situations are shown in Fig.1. For small B field and small χ , show that the susceptibility in this model is given by

$$\chi = \frac{2ng^2\mu_B^2\mu_0}{k_B T \left(3 + e^{\frac{\Delta}{k_B T}} \right)}$$

which is known as the Bleaney-Bowers equation. Use the script Bleaney_Bowers.ipynb to plot this function. Estimate the minimal temperature above which the susceptibility can be fitted with the Curie law.

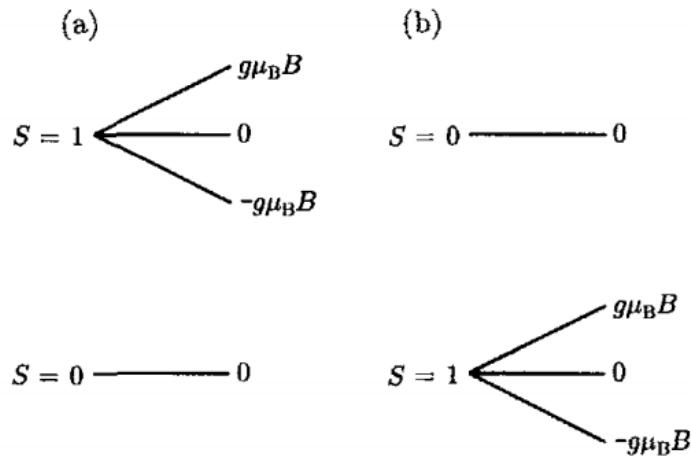


Figure 1: Energy levels if : (a) $\Delta > 0$, (b) $\Delta < 0$

Solution The energy of each state is the following :

$$\begin{aligned} E_s &= -\Delta \\ E_t(m_s = 1) &= g\mu_B B \\ E_t(m_s = 0) &= 0 \\ E_t(m_s = -1) &= -g\mu_B B \end{aligned}$$

Using the partition function :

$$Z = \sum_i e^{-E_i \beta}$$

with E_i the energy of the state i and $\beta = \frac{1}{k_B T}$, the mean magnetic moment can be found :

$$\begin{aligned} \langle M \rangle &= \frac{\sum_i n g \mu_B m_i e^{-E_i \beta}}{Z} \\ &= n g \mu_B \frac{e^{g \mu_B B \beta} - e^{-g \mu_B B \beta}}{e^{\Delta \beta} + e^{-g \mu_B B \beta} + 1 + e^{g \mu_B B \beta}} \\ &\approx \frac{2 n g^2 \mu_B^2 \beta B}{3 + e^{\Delta \beta}} \end{aligned}$$

Where the last step was done using that $B \ll 1$. In the small susceptibility limit : $B = \mu_0(1 + \chi)H \approx \mu_0 H$. So we can found the magnetic susceptibility :

$$\begin{aligned} \chi &= \frac{M}{H} \approx \frac{\mu_0 M}{B} \\ &\approx \frac{2 n g^2 \mu_B^2 \mu_0 \beta}{3 + e^{\Delta \beta}} \end{aligned}$$

In the limit $T \gg \Delta/k_B$ the exponential part becomes 1 and we found the Curie law.

2. Calculate the magnitude of the magnetic field 1 Å and 10 Å from a proton in a direction (a) parallel and (b) perpendicular to the proton spin direction.

Solutions The proton magnetic moment is $\mu_p = 2.793 \mu_n$ with $\mu_n = \frac{e\hbar}{2m_p}$. The dipole magnetic field is :

$$\vec{B} = \frac{\mu_0}{4\pi|\vec{r}|^3} (3(\vec{\mu}_p \hat{r})\hat{r} - \vec{\mu}_p)$$

(a)

$$\begin{aligned} |\vec{B}(\vec{r})| &= \frac{\mu_0 \mu_p}{2\pi|\vec{r}|^3} \\ |\vec{B}(1\text{\AA})| &= 2.82 \text{mT} \\ |\vec{B}(10\text{\AA})| &= 2.82 \cdot 10^{-3} \text{mT} \end{aligned}$$

(b)

$$\begin{aligned} |\vec{B}(\vec{r})| &= \frac{\mu_0 \mu_p}{4\pi|\vec{r}|^3} \\ |\vec{B}(1\text{\AA})| &= 1.41 \text{mT} \\ |\vec{B}(10\text{\AA})| &= 1.41 \cdot 10^{-3} \text{mT} \end{aligned}$$

3. Estimate the ratio of the exchange and dipolar coupling of two adjacent Fe atoms in metallic Fe.

Hint Iron has a BCC structure with a lattice constant of 286.65 pm. The magnetic dipole for each atom is $\mu = 2.2 \mu_B$. The exchange constant in Fe can be crudely estimated by setting it equal to $k_B T_c$ where T_c is the Curie temperature. The curie temperature for iron is $T_c = 1043$ K.

Solution The dipole interaction energy is given by

$$E = \frac{\mu_0}{4\pi} \left[\frac{\vec{\mu}_1 \vec{\mu}_2}{r^3} - \frac{3(\vec{\mu}_1 \vec{r})(\vec{\mu}_2 \vec{r})}{r^5} \right] \approx \frac{\mu_0 \mu_{Fe}^2}{4\pi r^3}$$

Using the hints and the fact that the closest neighbour in a bcc structure is given by $r = \frac{3a}{\sqrt{2}}$ with a the lattice constant, we found that $E \approx 1.15 \mu\text{eV}$. The exchange constant being $J \approx 0.09 \text{eV}$, the exchange dominate in pure iron.

4. Provide a rough estimate of the size of the exchange constant in a magnetic oxide which is coupled by superexchange using the measured value of the electronic bandwidth (determined by inelastic neutron scattering) of 0.05 eV. Take the Coulomb energy to be 1 eV. Hence estimate the antiferromagnetic ordering temperature.

Solution $J \approx \frac{-t^2}{U}$ with $t = 0.05 \text{eV}$ and $U = 1 \text{eV}$. Hence the Néel temperature can be estimated using $J \approx k_B T_N$:

$$T_N \approx \frac{t^2}{U k_B} \approx 29 \text{K}$$

5. Dzyaloshinskii–Moriya interaction is a magnetic exchange interaction between two neighboring magnetic spins, \vec{S}_i and \vec{S}_j , described by the Hamiltonian $\vec{D} \cdot (\vec{S}_i \times \vec{S}_j)$. In the two spin system depicted here, \vec{D} is along \hat{x} , and $\vec{S}_1 = \frac{1}{\sqrt{2}}(-\hat{x} + \hat{z})$ and $\vec{S}_2 = -\hat{y}$. Find and sketch the direction of $\vec{S}_1 \times \vec{S}_2$. Now, assume that $\vec{S}_1 = \frac{1}{\sqrt{2}}(\hat{x} - \hat{z})$ and repeat the same exercise. Which of these two configurations have lower energy?

Use the script DMI.ipynb to visualize how the vector \vec{D} changes when the central ligand is moved around. What symmetries are broken when you move it?

Solution

$$\vec{S}_1 \times \vec{S}_2 = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{1}{\sqrt{2}}(-1) & 0 & \frac{1}{\sqrt{2}}(1) \\ 0 & -1 & 0 \end{vmatrix} = \frac{1}{\sqrt{2}}(\hat{x} + \hat{z})$$

Now, if we assume that $\vec{S}_1 = \frac{1}{\sqrt{2}}(\hat{x} - \hat{z})$, we can repeat the same steps to find $\vec{S}_1 \times \vec{S}_2$:

$$\vec{S}_1 \times \vec{S}_2 = \frac{1}{\sqrt{2}}(-\hat{x} - \hat{z})$$

Comparing the two configurations, we see that the direction of $\vec{S}_1 \times \vec{S}_2$ differs in sign for the \hat{x} component. The second configuration has lower energy as the \hat{x} component is negative.

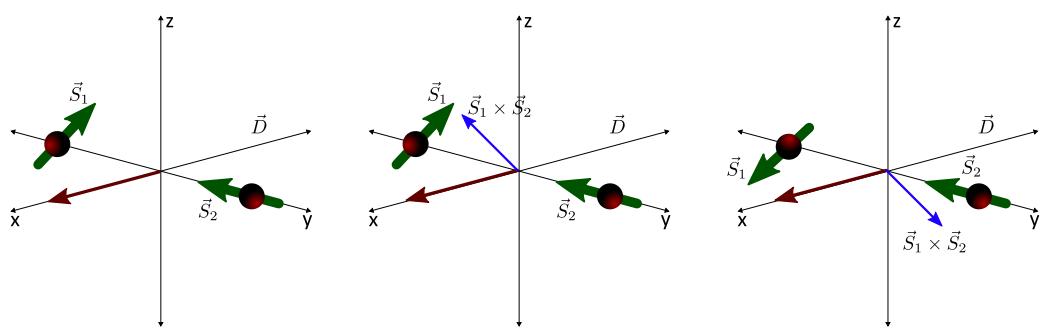


Figure 2: