

MICRO-523: Optical Detectors

Week Two: Detector Formalism and Noise

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Outline

- 2.1 Quantum efficiency
- 2.2 Responsivity and depth
- 2.3 Optimal internal gain
- 2.4 Light source statistics

Exercise 2.1: Quantum Efficiency and Detectivity

Consider a semiconductor photodiode with a band gap E_g and an ideal quantum efficiency.

Sketch:

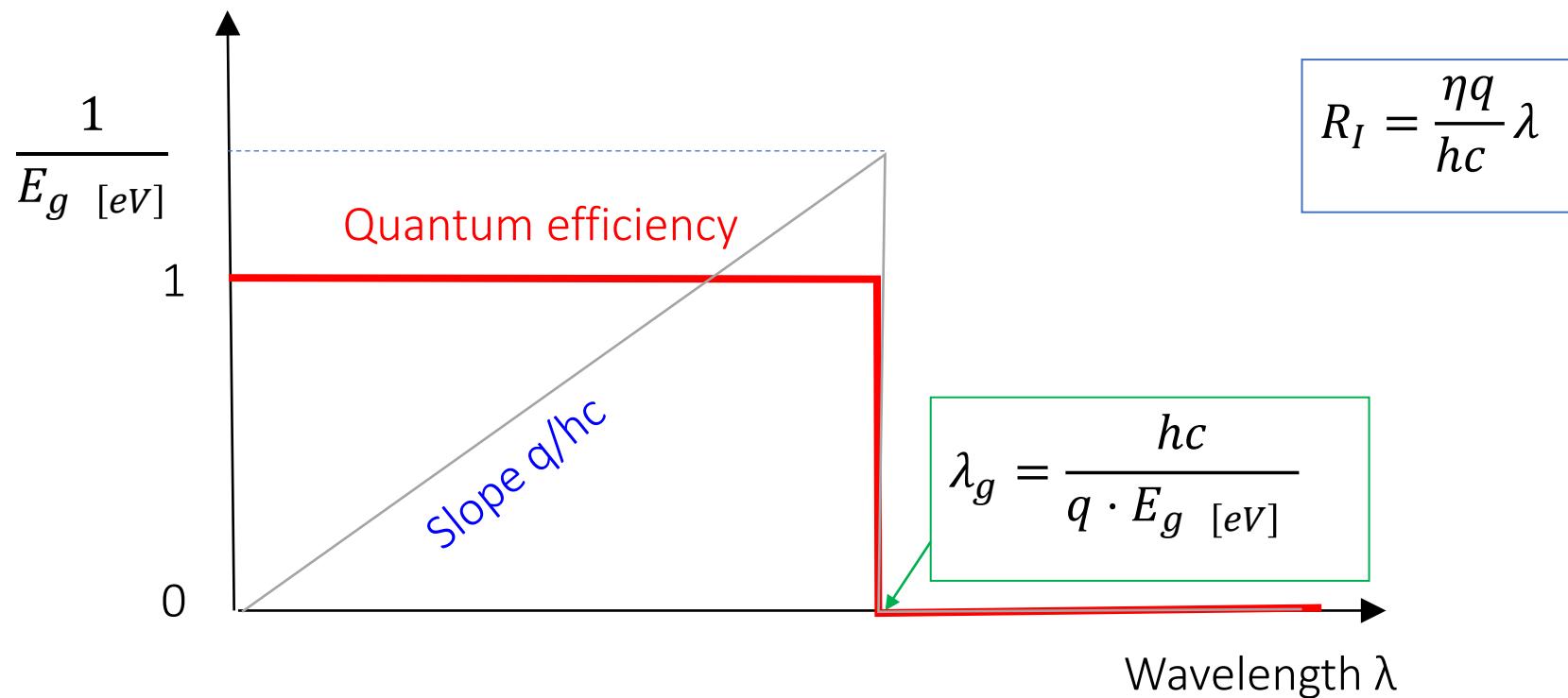
- its quantum efficiency η and
- its responsivity R_I

as a function of the wavelength of the incident photons.

Consider the noise N to be independent of wavelength.

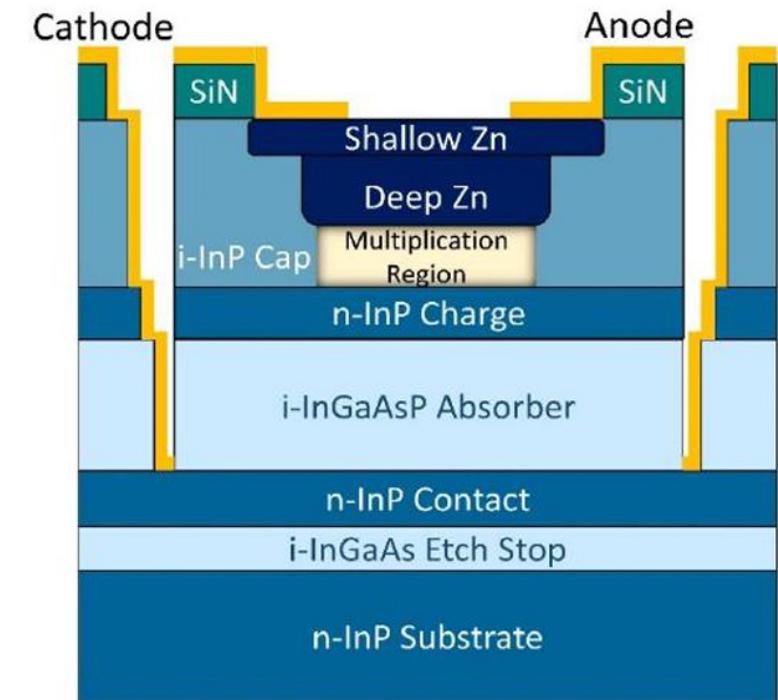
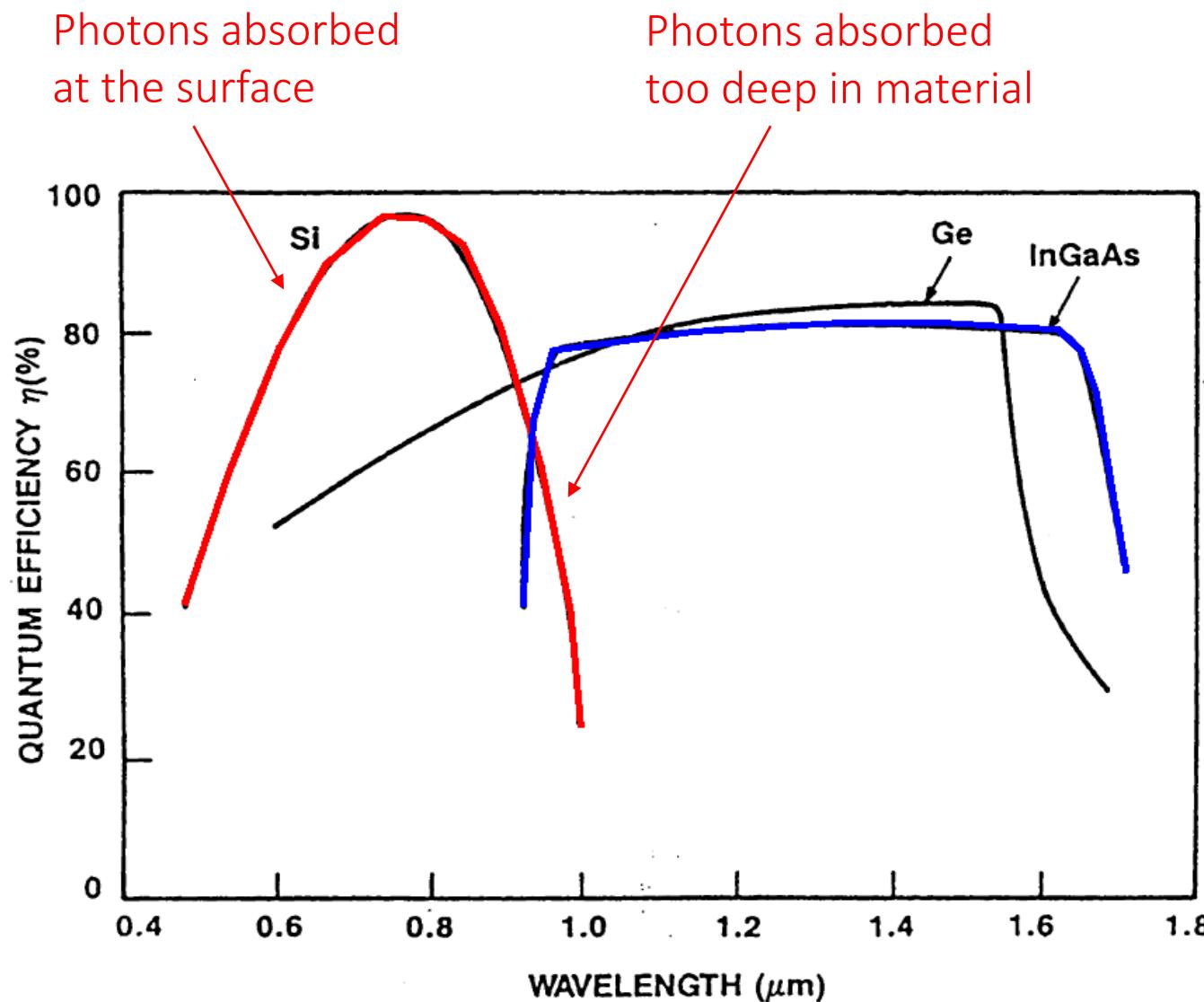
Sketch its detectivity as a function of its wavelength.

Exercise 2.1: Quantum Efficiency

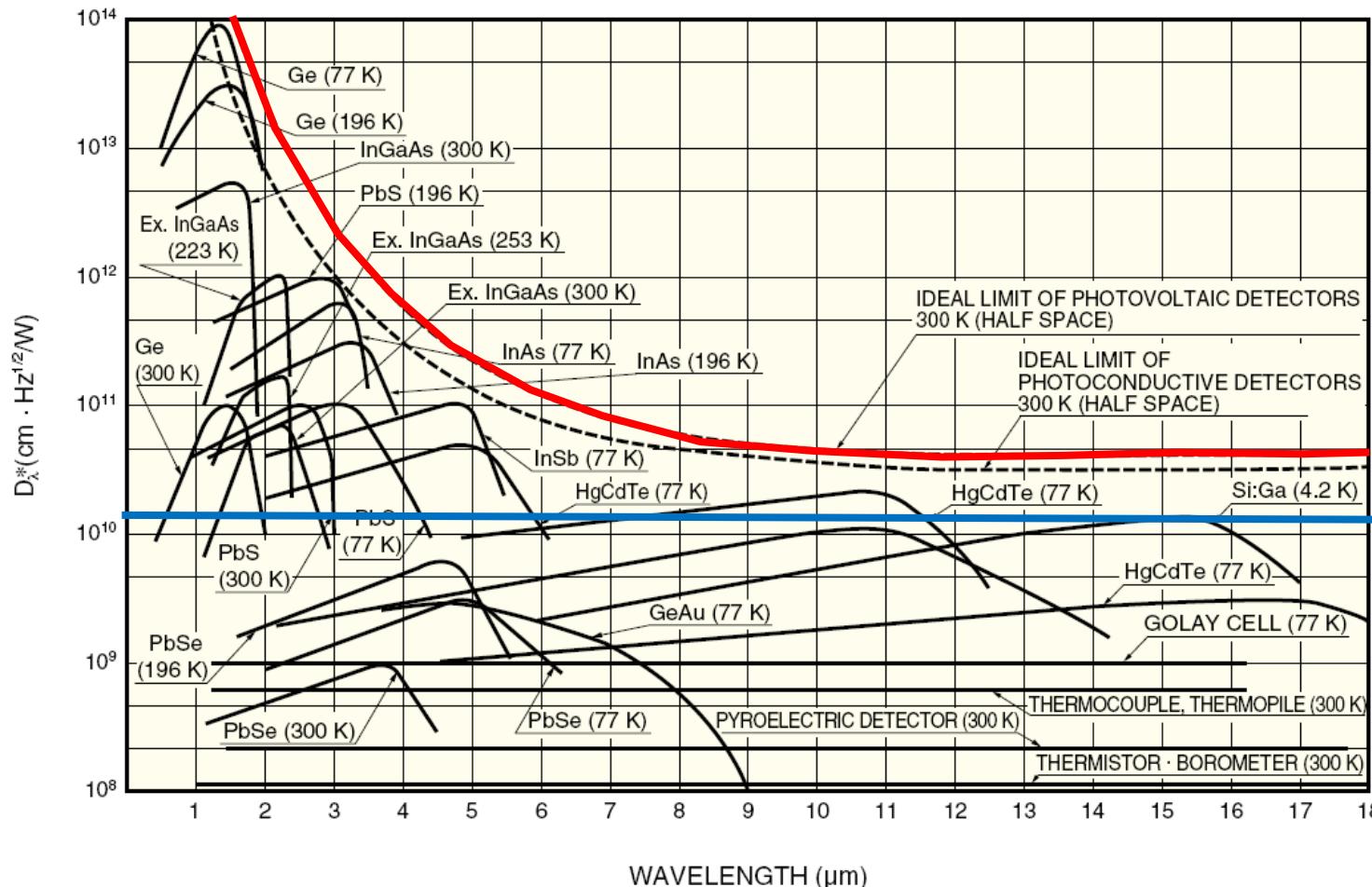


$$\frac{S}{N} = 1 \Rightarrow \frac{R_I \cdot NEP}{N} = 1 \iff D \equiv \frac{1}{NEP} = \frac{R_I}{N} \iff \text{The spectrum of } D \text{ is similar to that of the responsivity.}$$

Exercise 2.1: Quantum Efficiency: Example



Exercise 2.1: Spectral Dependence of the Responsivity



Thermocouple (300 K)

Hamamatsu Catalog

Exercise 2.2: Responsivity and Depth

The absorption coefficient of silicon can be approximated as $\alpha_{cm^{-1}}(\lambda_{\mu m}) \cong 10^{7.2-5.5\lambda}$

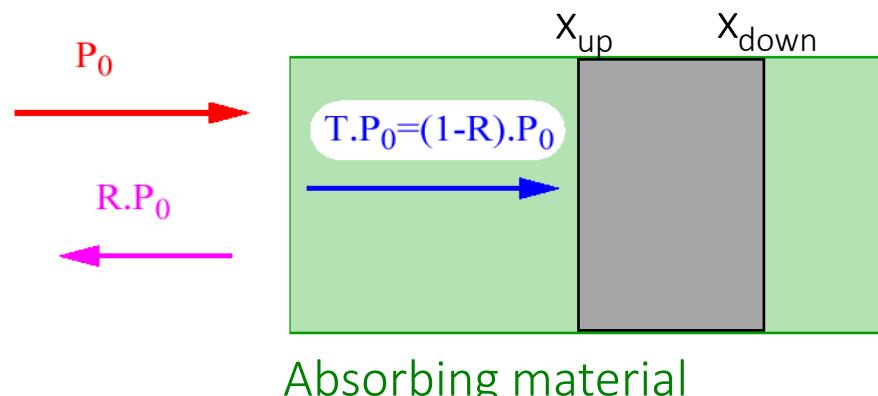
Consider two detectors and two wavelengths ($\lambda=450\text{nm}$ and $\lambda = 600\text{nm}$).

The first detector is sensitive between $x_{1\text{up}} = 0.05\mu\text{m}$ and $x_{1\text{down}} = 0.3\mu\text{m}$.

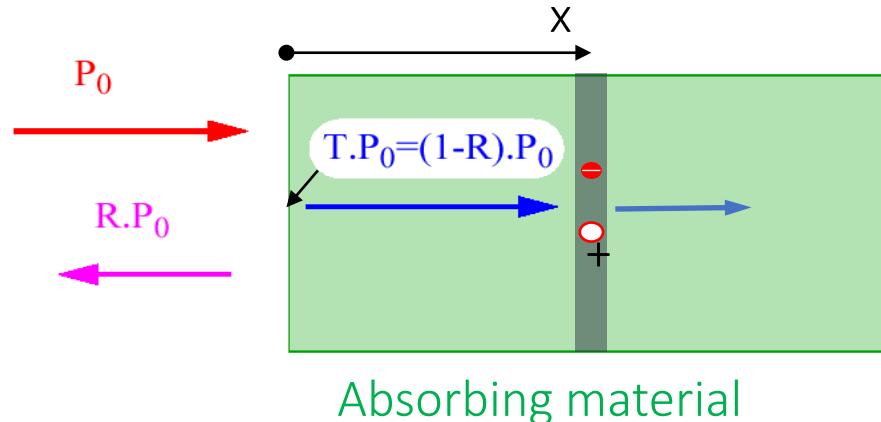
The second detector is sensitive between $x_{2\text{up}} = 0.9\mu\text{m}$ and $x_{2\text{down}} = 4\mu\text{m}$.

The reflection coefficient is 10%.

Calculate the responsivity R_i and the quantum efficiency of both detectors at the abovementioned wavelengths.



Exercise 2.2: Absorption and Generation Rate



R = reflection coefficient

T = transmission coefficient

$g(x)$ = generation rate of carriers

$$g(x) = \frac{P(x)}{h\nu} \cdot \alpha = \frac{P_0}{h\nu} \cdot (1 - R) \cdot e^{-\alpha x} \cdot \alpha \quad \left[\frac{1}{cm \cdot s} \right]$$

Exercise 2.2: Quantum Efficiency

For photodetectors the quantum efficiency is defined as follows:

$$\eta = \frac{\text{number of optically generated electrons}}{\text{number of incident photons}}$$

It takes into account:

- reflection
- absorption
- recombination
- and electron scattering

It does not consider:

- internal gain
- avalanche phenomena, ...

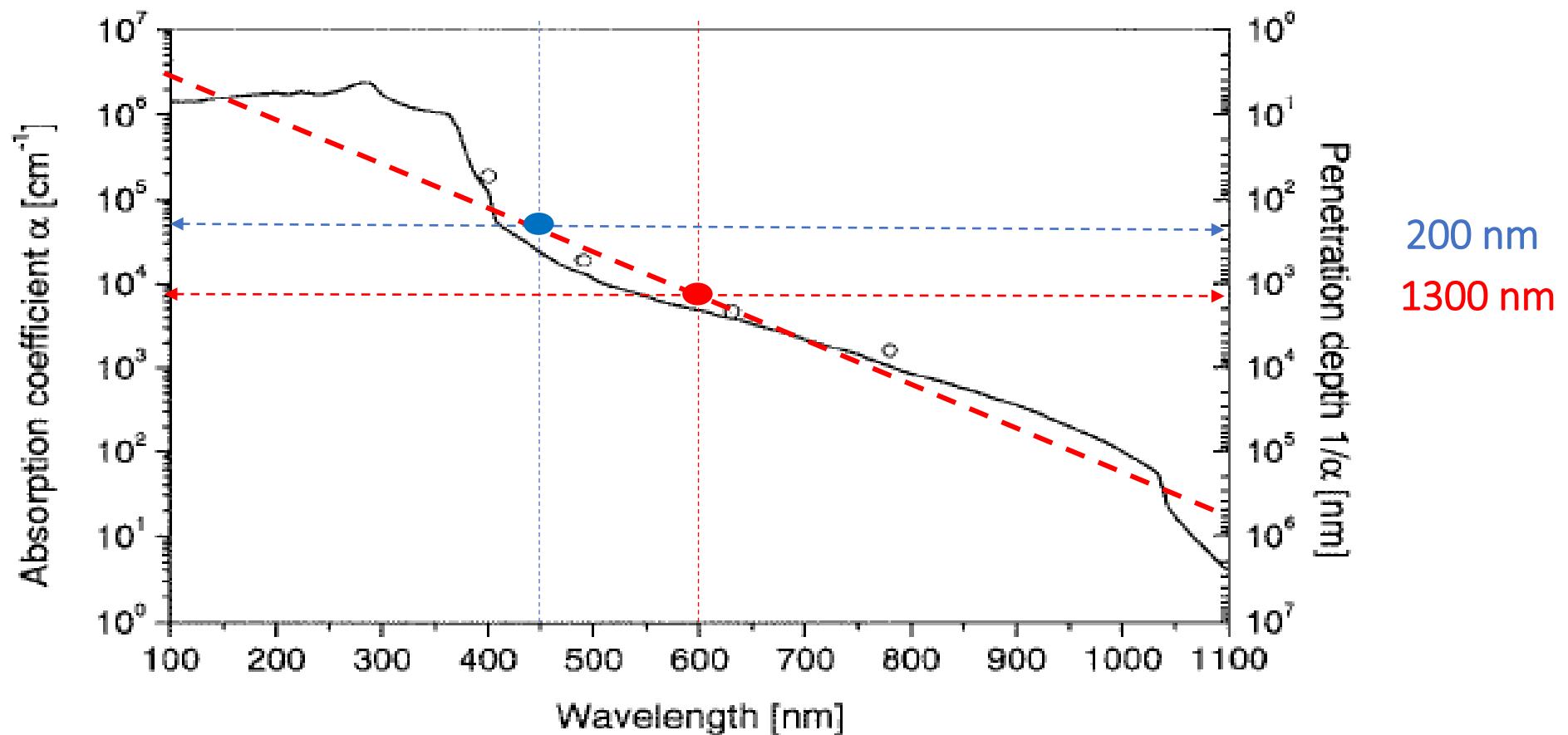
Relationship with responsivity

$$\eta = \frac{I_s/q}{P_s/h\nu}$$



$$R_I = \frac{I_s}{P_s} = \frac{\eta q}{h \nu}$$

Exercise 2.2: Absorption by silicon



Approximation: $\alpha_{cm^{-1}}(\lambda_{\mu m}) \cong 10^{7.2 - 5.5\lambda}$

Exercise 2.2: Absorption, Responsivity and Quantum Efficiency

Absorption:

$$\alpha_{cm^{-1}}(\lambda_{\mu m}) \cong 10^{7.2-5.5\lambda}$$

Generation rate:

$$g(x) = \frac{P(x)}{h \nu} \cdot \alpha = \frac{P_0}{h \nu} \cdot (1 - R) \cdot e^{-\alpha x} \cdot \alpha \quad \left[\frac{1}{cm \cdot s} \right]$$

Photocurrent:

$$I = q \cdot \int_{x_{up}}^{x_{down}} g(x) \cdot dx = \frac{q}{h\nu} \cdot (1 - R) \cdot (e^{-\alpha x_{up}} - e^{-\alpha x_{down}}) \cdot P_0$$

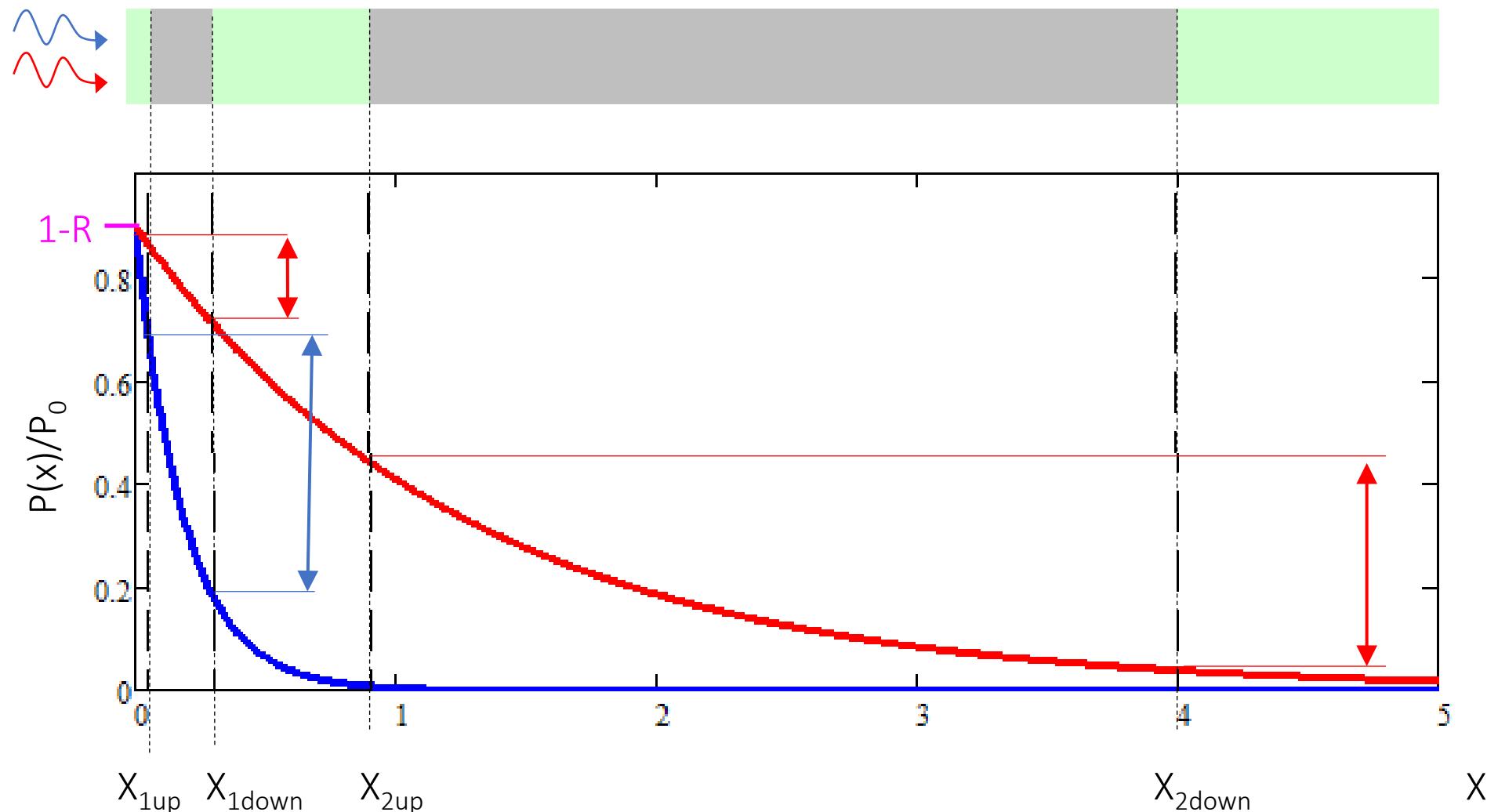
Responsivity:

$$R_I = \frac{I}{P_0} = \frac{q}{h\nu} \cdot (1 - R) \cdot (e^{-\alpha x_{up}} - e^{-\alpha x_{down}})$$

Quantum efficiency

$$\eta = R_I \cdot \frac{h\nu}{q} = (1 - R) \cdot (e^{-\alpha x_{up}} - e^{-\alpha x_{down}})$$

Exercise 2.2: Interpretation: Quantum Efficiency



Exercise 2.2: Numerical Values

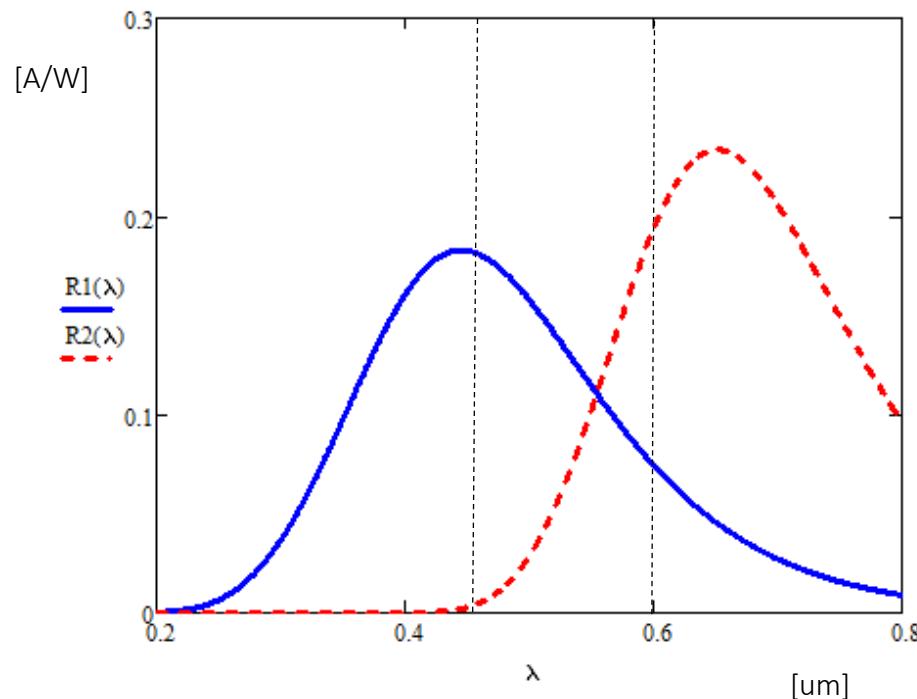
$$\alpha(450) = 5.3 \cdot 10^4 \text{ cm}^{-1}$$

$$\alpha(600) = 7.9 \cdot 10^3 \text{ cm}^{-1}$$

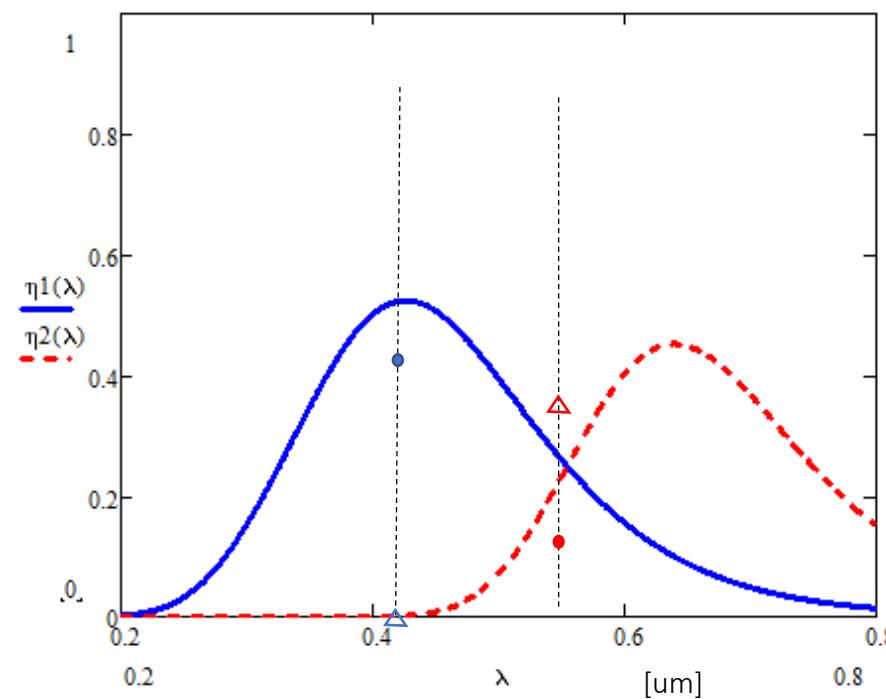
	“Superficial” detector		“Deep” detector	
	R_I [A/W]	η %	R_I [A/W]	η %
$\lambda=450\text{nm}$	0.183	51 %	0.003	0.8 %
$\lambda=600\text{nm}$	0.074	15.5 %	0.194	40 %
R_{450}/R_{600} η_{450}/η_{600}	2.5	3.3	0.014	0.02

Exercise 2.2: Responsivity and Quantum Efficiency

Responsivity



Quantum efficiency

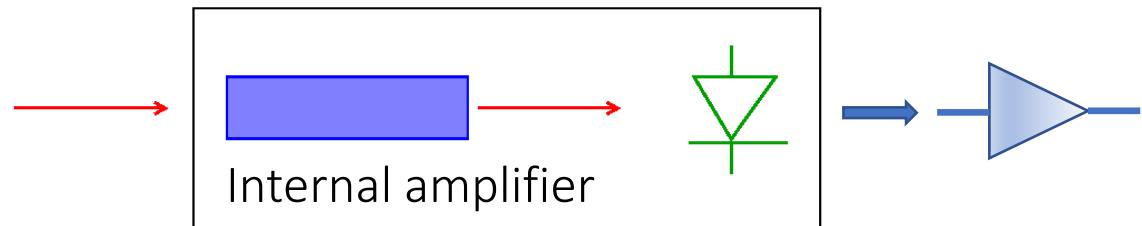


— Superficial detector
- - - Deep detector

Exercise 2.3: Optimal Internal Gain

We would like to detect an optical signal with $P_0=25$ pW using a detector with a variable internal gain G (for example an avalanche photodiode).

Its responsivity is $R_i=0.4$ A/W,
its bandwidth is $\Delta f=1$ MHz,
and its excess noise factor is $F=G^{0.3}$.



Determine as a function of its gain G :

- The amplified photocurrent $I_{\text{sig}}(G)$.
- The shot noise of this current.

The electronics generate an rms noise of $\Delta I_{\text{el}}=100$ pA.

- Calculate the total noise ΔI_{tot} as a function of the gain.
- Determine the signal to noise ratio as a function of gain.
- What is the optimal gain G_{opt} ?

Exercise 2.3: Internal Gain (1)

Primary photocurrent:

$$I_0 = R_I \cdot P_0 = 10 \text{ pA}$$

Amplified photocurrent:

$$I_{sig}(G) = G \cdot I_0$$

Shot noise of the
amplified signal:

$$\Delta I_{sig}(G) = \sqrt{2qI_0\Delta f \cdot G^2 \cdot F(G)}$$

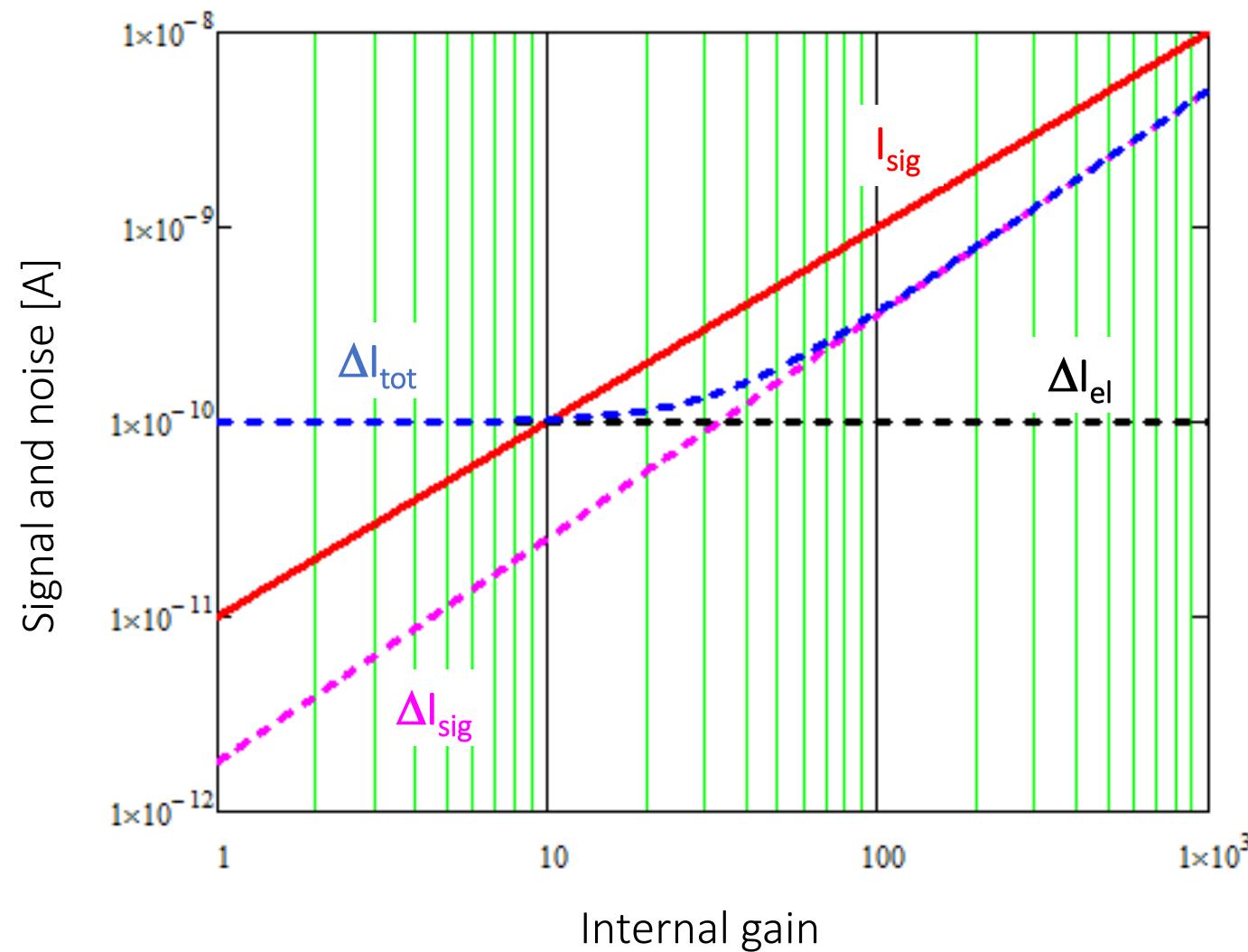
Total noise:

$$\Delta I_{tot} = \sqrt{\Delta I_{el}^2 + \Delta I_{sig}^2(G)}$$

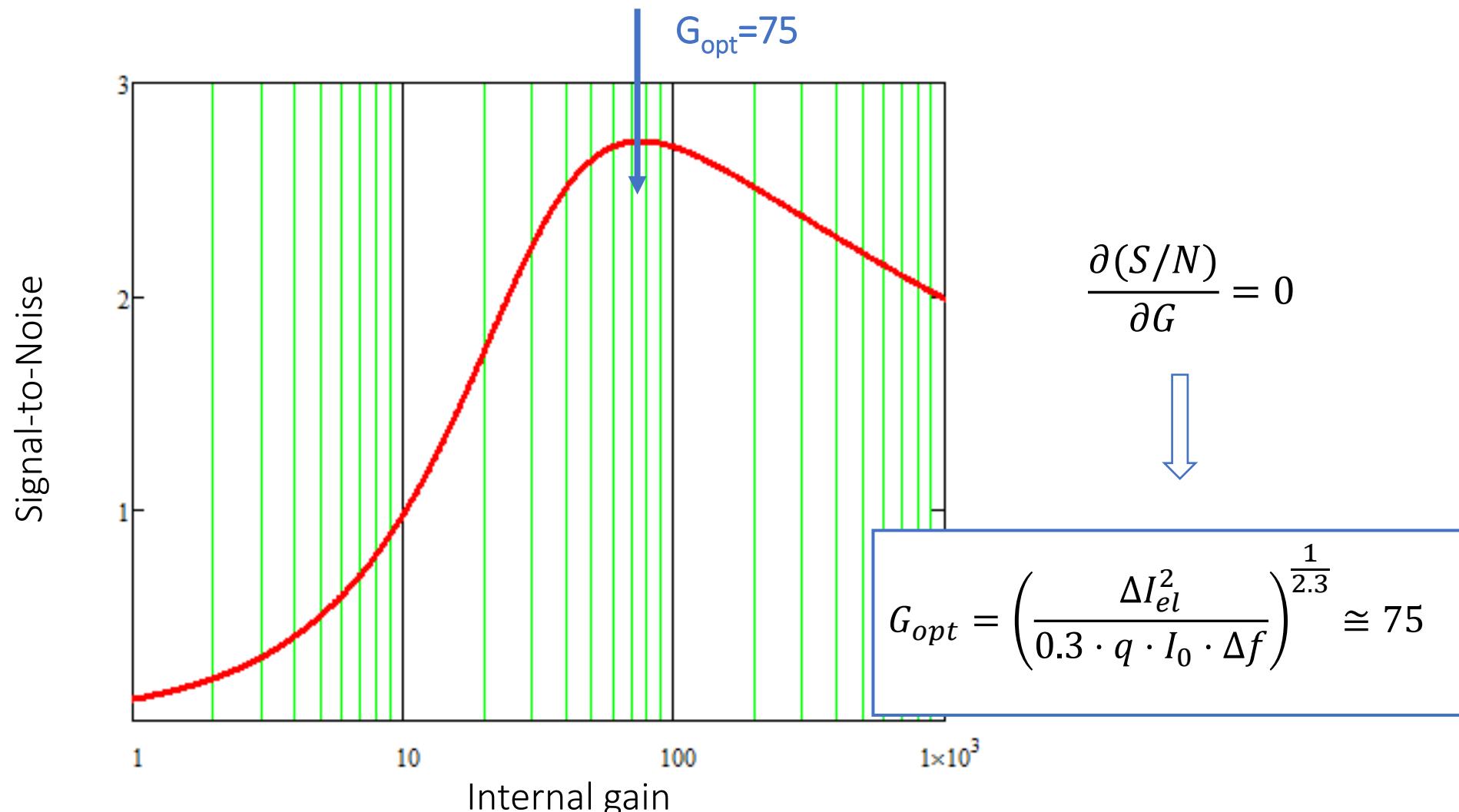
Signal-to-Noise ratio:

$$S/N = \frac{G \cdot I_0}{\sqrt{\Delta I_{el}^2 + 2qI_0\Delta f \cdot G^2 \cdot G^{0.3}}}$$

Exercise 2.3: Internal Gain (2)



Exercise 2.3: Internal Gain (3)



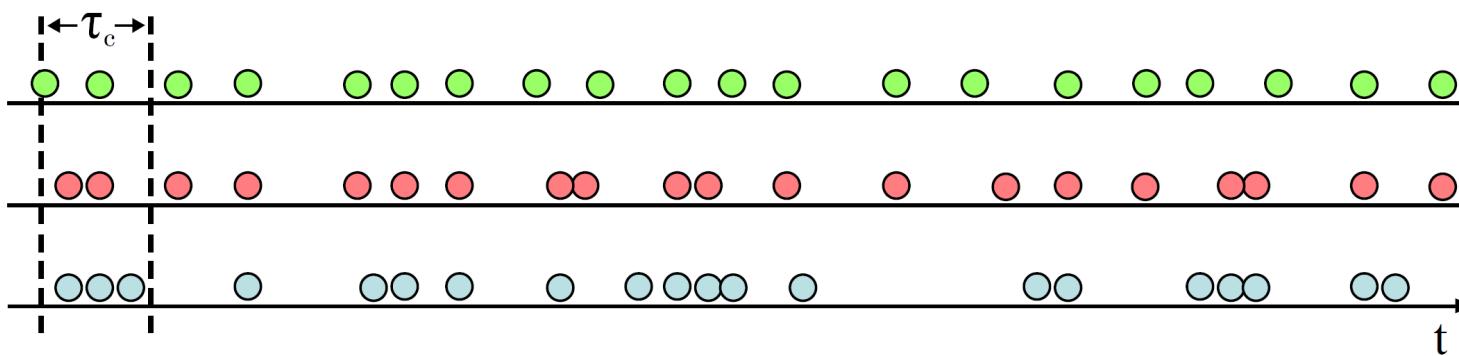
Exercise 2.4: Light sources statistics

Questions

- Which kind of light sources exist?
- How are their statistical emission properties?
- *Recap: shot noise*

$$N^2(f) = \frac{\langle \Delta I_{shot}^2 \rangle}{\Delta f} = 2q \cdot |\langle I \rangle|$$

$$N^2(f) = \frac{\langle \Delta P_{shot}^2 \rangle}{\Delta f} = 2h\nu \cdot \langle P \rangle$$



Photon detections as function of time for a) antibunched, b) random, and c) bunched light

By Ajbura - Vectorised version of File:Photon bunching.png, CC BY-SA 4.0,
<https://commons.wikimedia.org/w/index.php?curid=73299604>

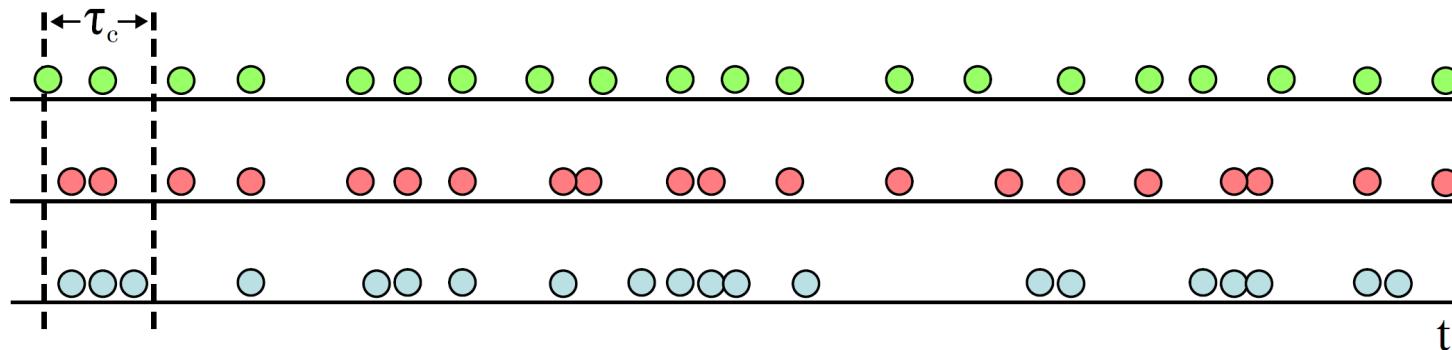
Exercise 2.4: Poisson Distribution vs. Light Sources

- Non-classical light: Sub-Poissonian \rightarrow antibunched (anticorrelated)
- Coherent light source (Laser): Poissonian, random spacing (uncorrelated)
- Thermal Light: Super-Poissonian, Bose-Einstein distribution with zero counts as most probable count (bunched, positively correlated)

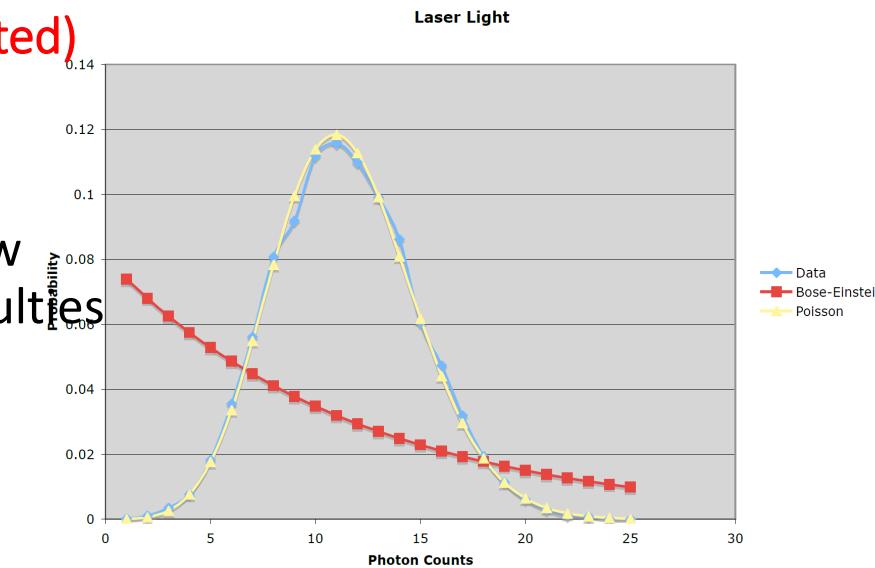
However, in practice it defaults to Gaussian due to the very low coherence time, $O(ps)$, and the corresponding experimental difficulties

Experimentally one can use pseudothermal light*.

<https://demonstrations.wolfram.com/PhotonNumberDistributions/>



Photon detections as function of time for a) antibunched, b) random, and c) bunched light



By Ajbura - Vectorised version of File:Photon bunching.png, CC BY-SA 4.0,
<https://commons.wikimedia.org/w/index.php?curid=73299604>

*E.g. scattering of a laser beam on a rotating ground glass disc

http://physics.gu.se/~tfkhj/lecture_X_differential_transmission-2.pdf

https://www.stmarys-ca.edu/sites/default/files/attachments/files/GriderJordanFinalReport_0.pdf

Exercise 2.4: Light sources statistics

\bar{n} = average photon number

Non-classical light: Sub-Poissonian $< \sqrt{\bar{n}}$

Coherent light source (Laser): Poissonian

$$P(n) = \frac{\bar{n}^n}{n!} e^{-\bar{n}}, \sigma = \sqrt{\bar{n}}$$

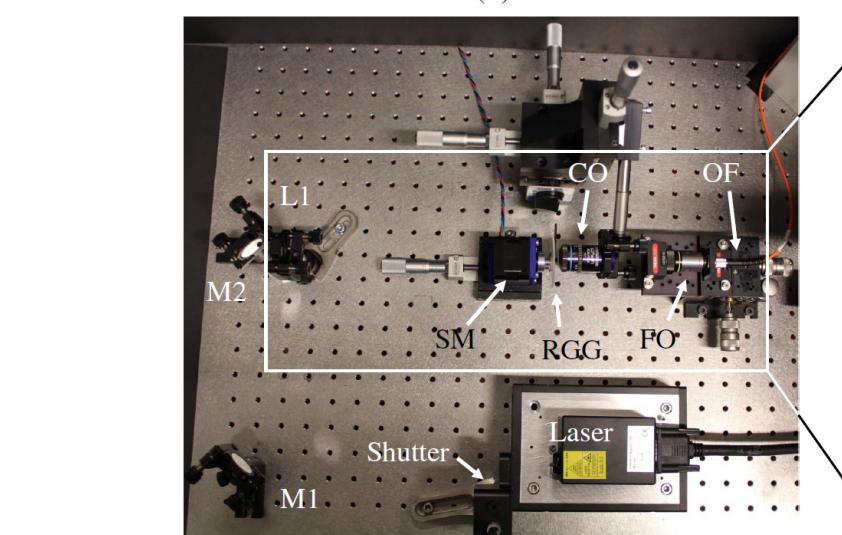
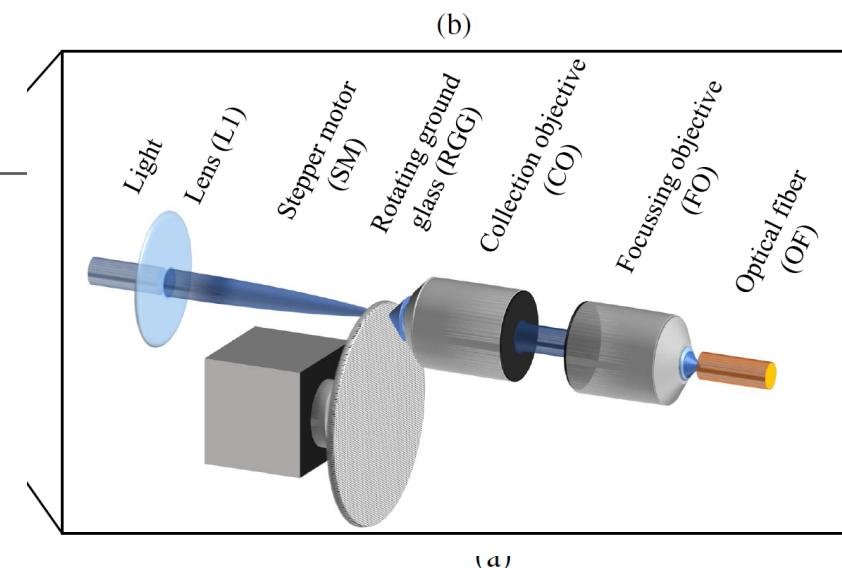
For large photon numbers, the relative fluctuations σ/\bar{n} tend to 0

Thermal Light: Super-Poissonian, Bose-Einstein distribution

$$P(n) = (1 - e^{-\hbar\omega/k_B T}) e^{-n\hbar\omega/k_B T} = \frac{\bar{n}^n}{(\bar{n} + 1)^{n+1}}, \bar{n} = (e^{\hbar\omega/k_B T} - 1)^{-1},$$

$$\sigma = \sqrt{\bar{n}^2 + \bar{n}} \quad (\text{for } T \ll \tau_c) > \sqrt{\bar{n}}$$

For large photon numbers, the relative fluctuations σ/\bar{n} tend to 1



Pseudothermal light source

T. Stagner et al., *Step-by-step guide to reduce spatial Coherence of laser light using a rotating ground glass diffuser*, OSA Applied Optics 56 (2017).

Advanced Lab Course (F-Praktikum), Exp. 45, *Photon Statistics*, v. Aug. 21 2017

http://physics.gu.se/~tfkhj/lecture_X_differential_transmission-2.pdf

https://www.stmarys-ca.edu/sites/default/files/attachments/files/GriderJordanFinalReport_0.pdf