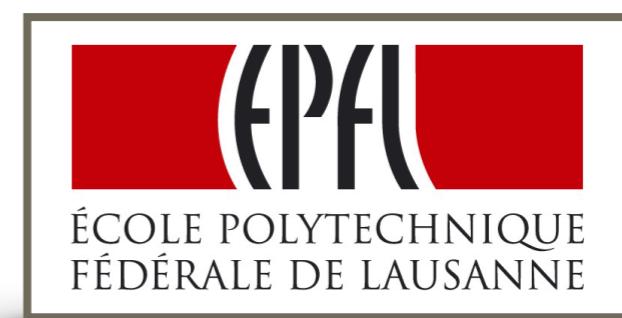


Transmission Matrices



Konstantinos Makris

University of Crete, Greece

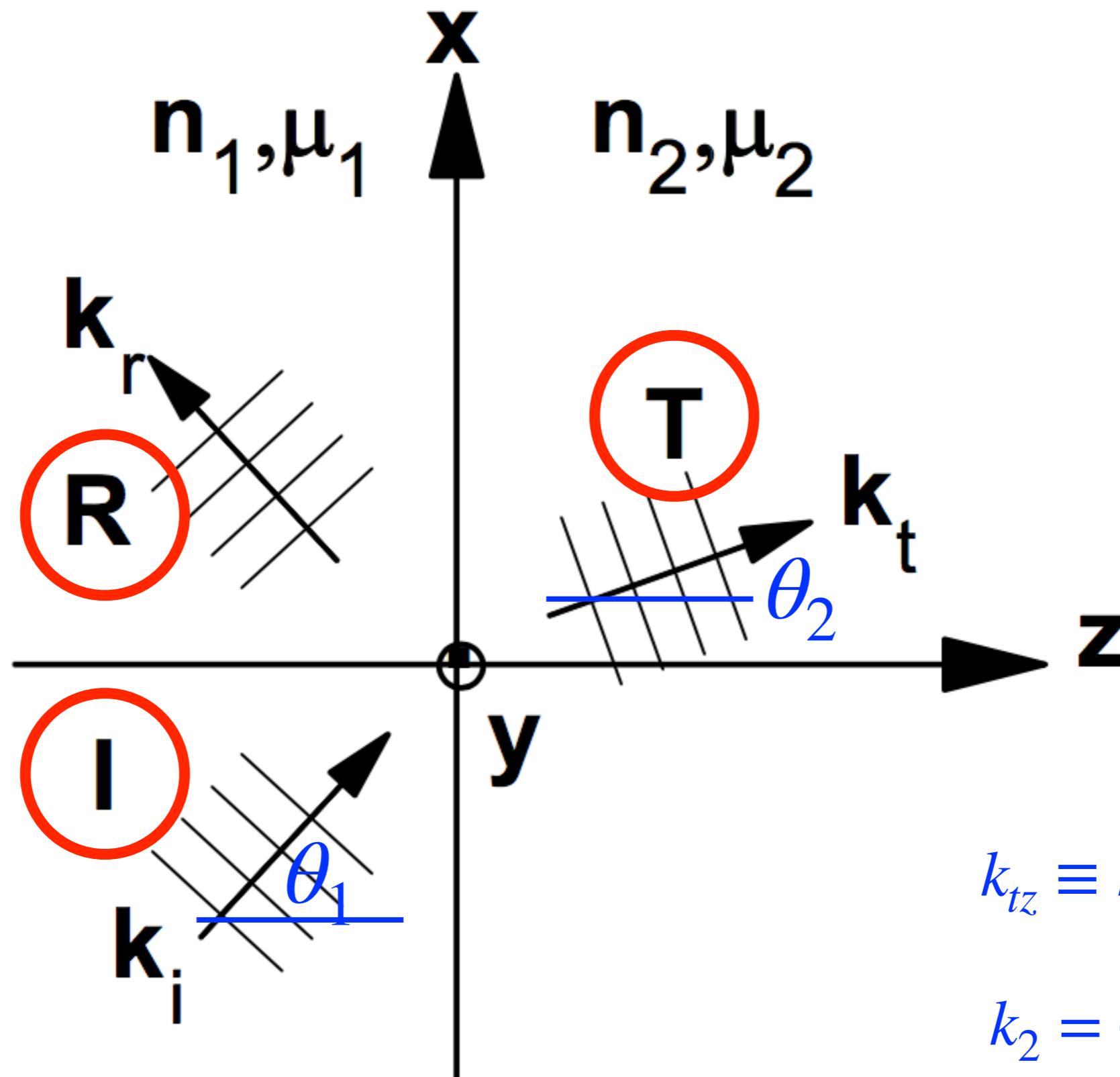


Demetri Psaltis

Optics Lab-EPFL, Switzerland

SCATTERING IN COMPLEX MEDIA

Single Interface



$$k_{iz} \equiv k_1 \cos \theta_1$$

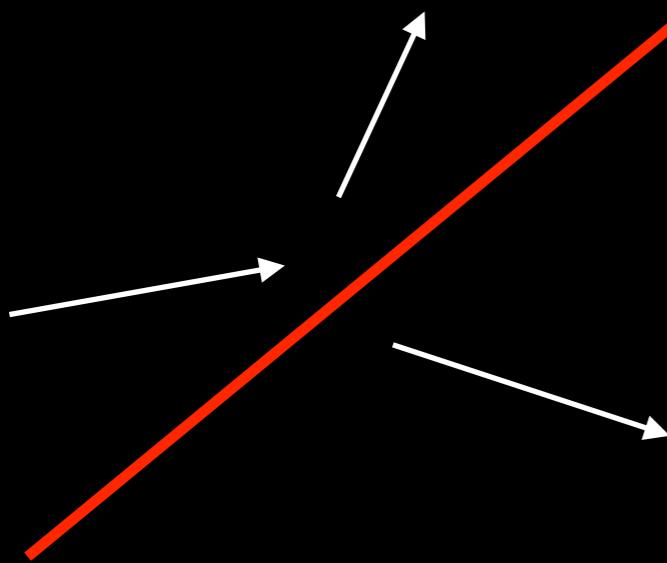
$$k_1 = \frac{2\pi n_1}{\lambda}$$

$$k_{tz} \equiv k_2 \cos \theta_2$$

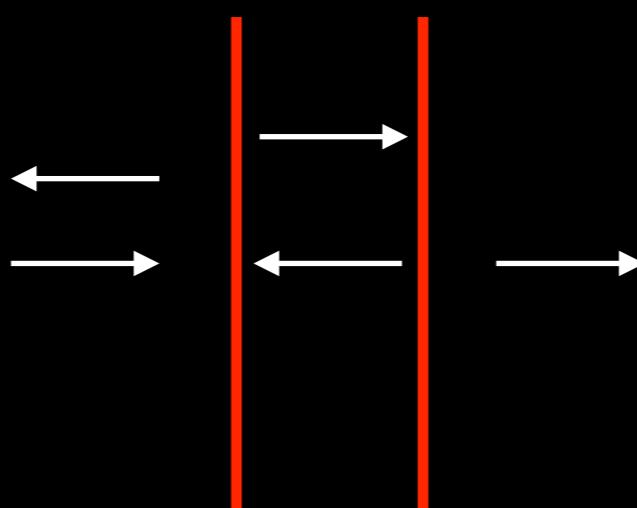
$$k_2 = \frac{2\pi n_2}{\lambda}$$

Scattering elements

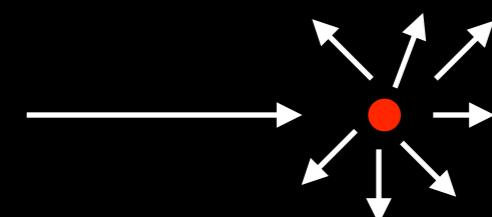
SINGLE
INTERFACE



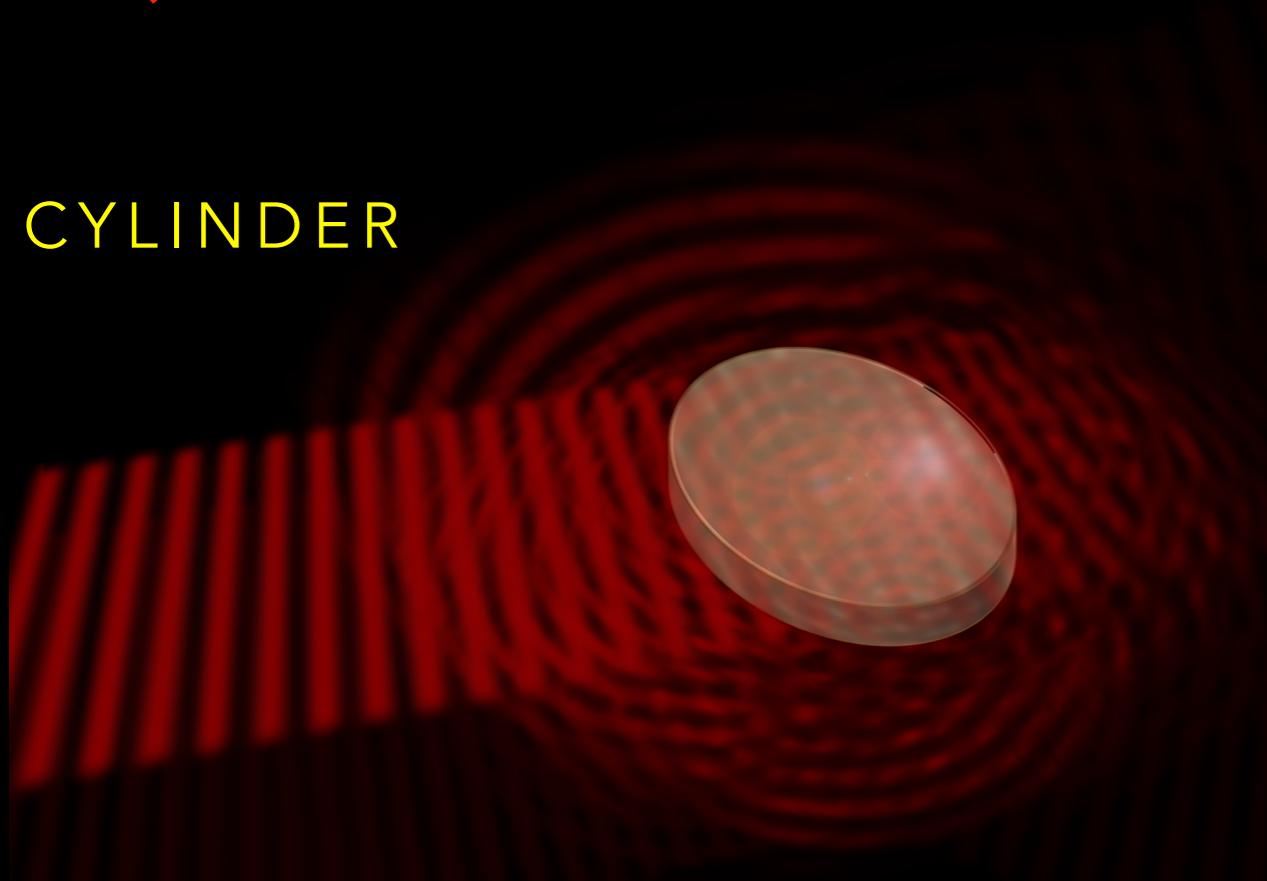
DOUBLE
INTERFACE



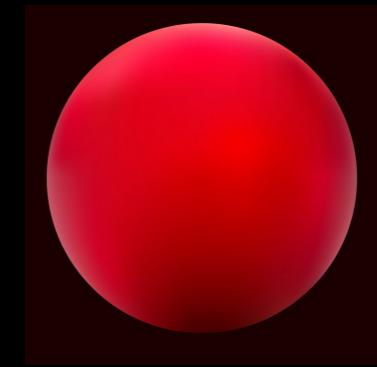
POINT SCATTERER



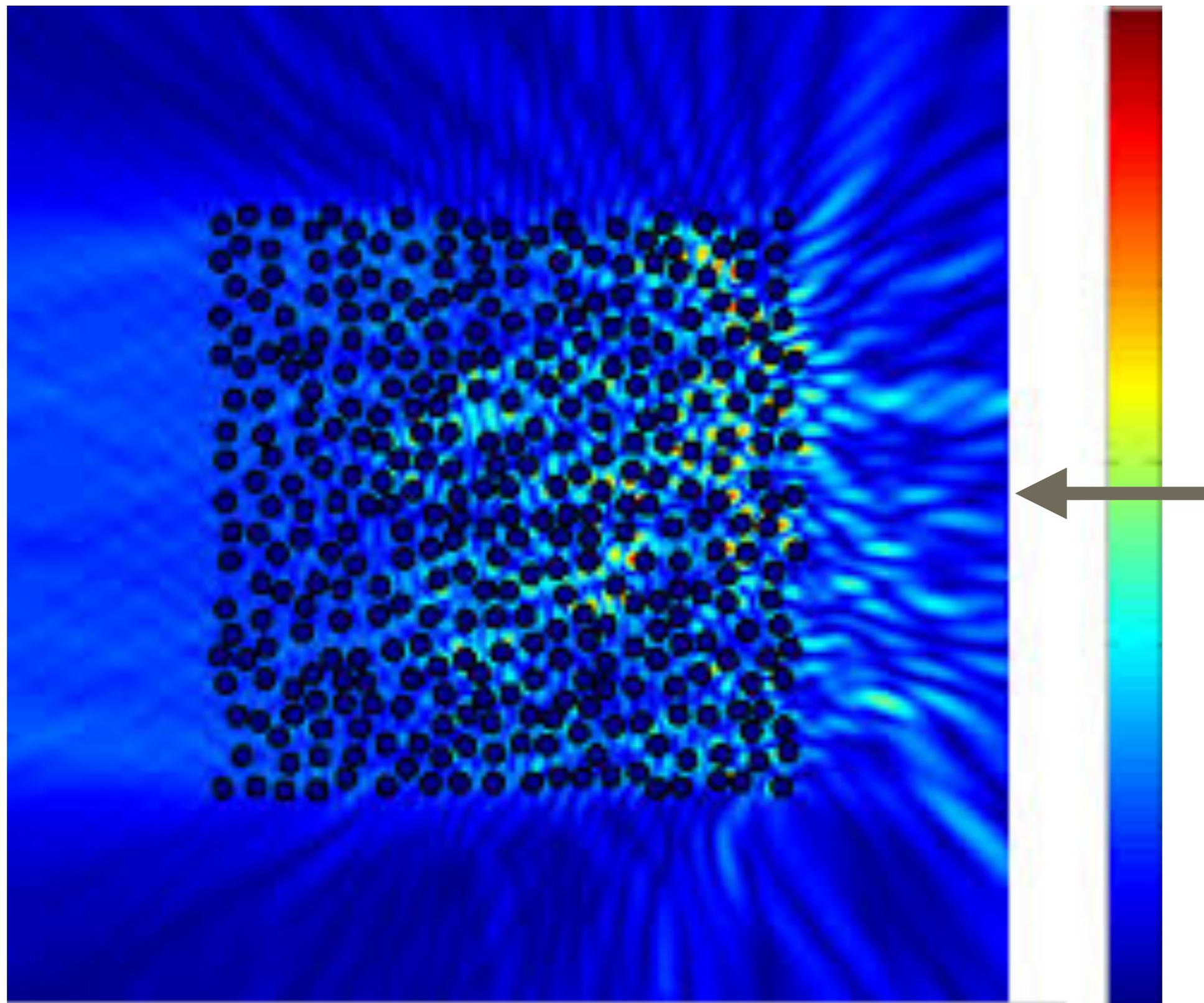
CYLINDER



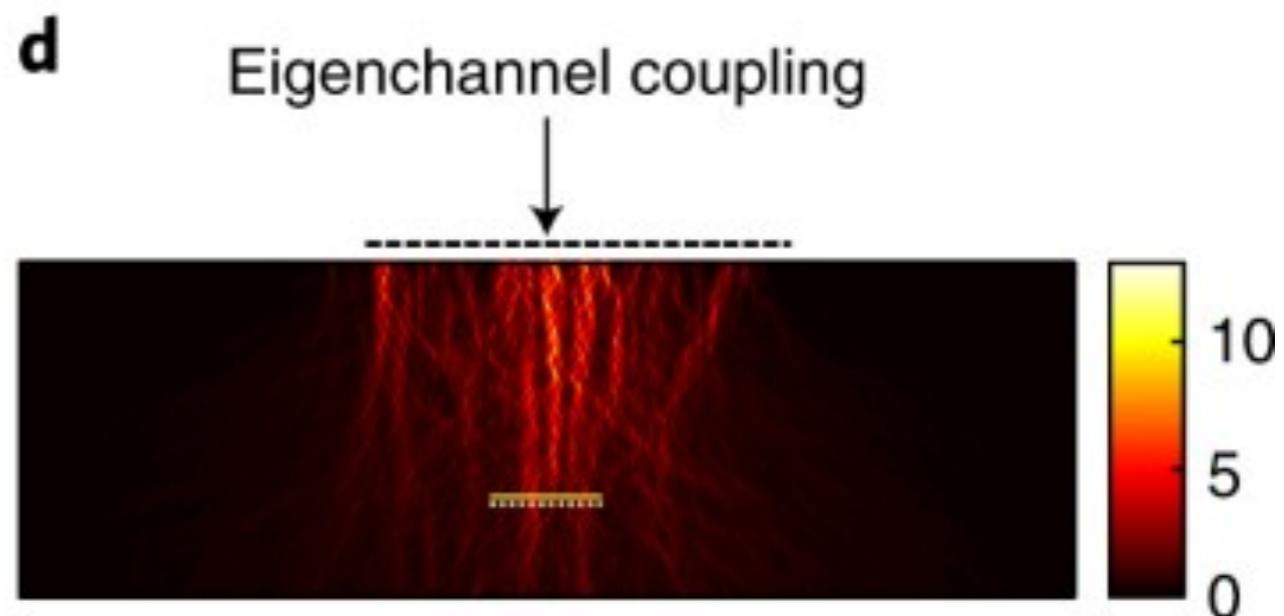
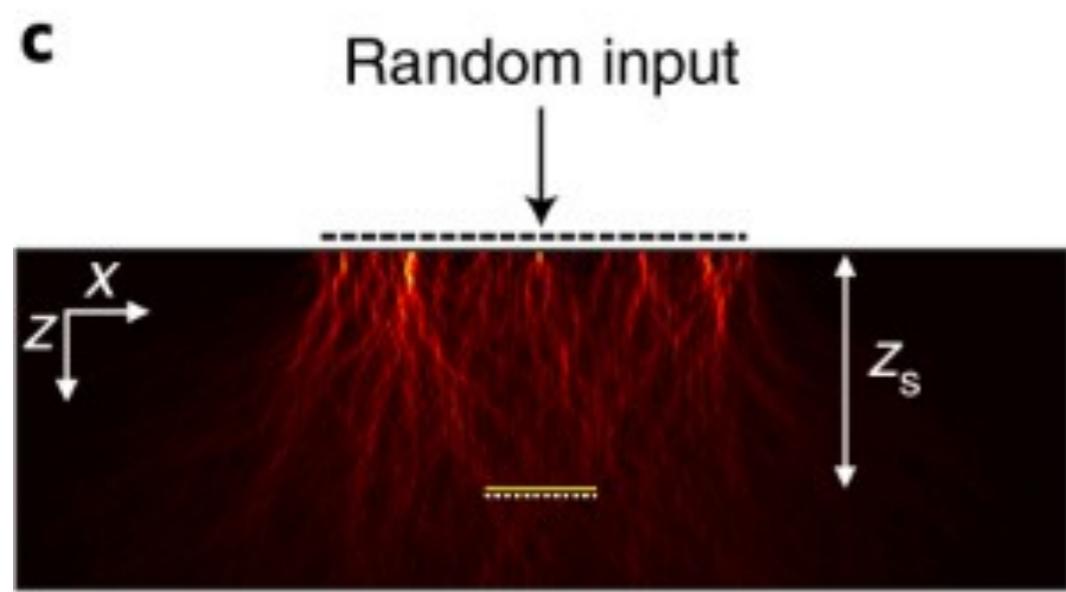
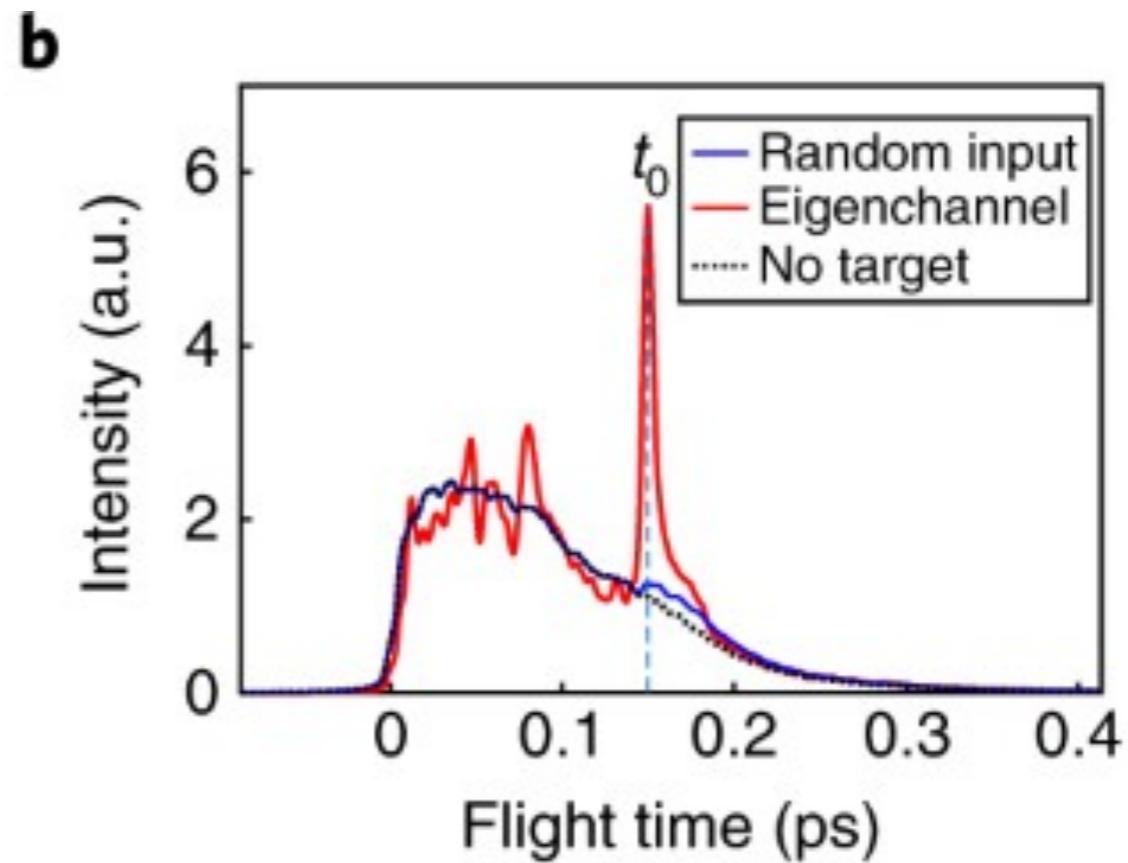
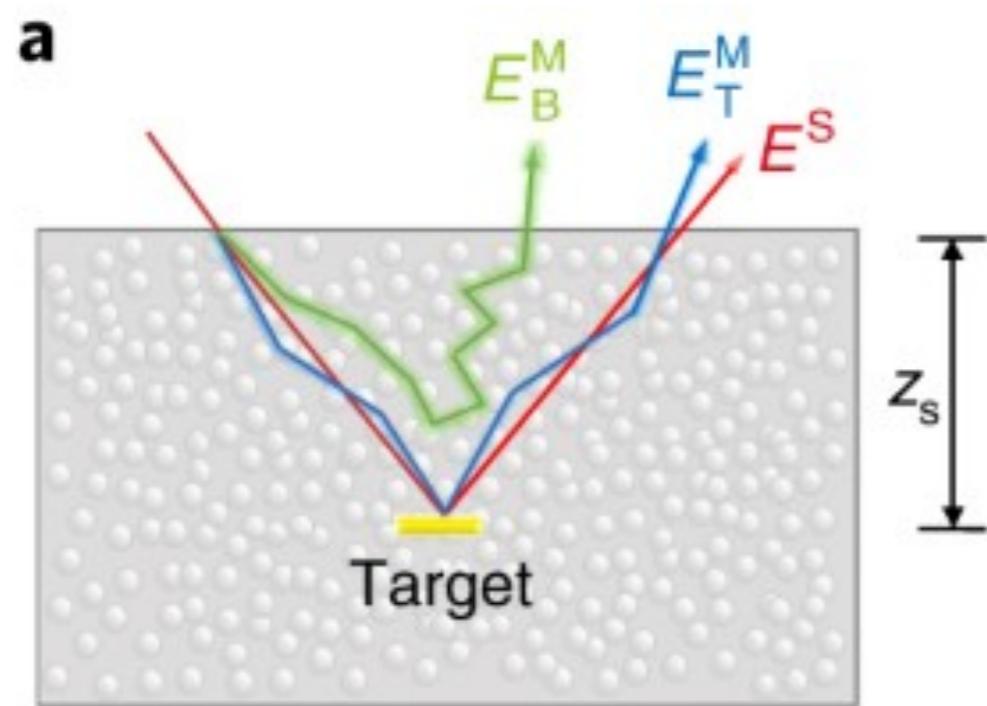
SPHERE



Multiple scattering



Penetrating Disorder



The fog clears

A technique has been developed to image a fluorescent object hiding behind a light-scattering screen without the need for a detector behind the screen. The approach could find applications in imaging biological tissue. [SEE LETTER P.232](#)

DEMETRI PSALTIS
& IOANNIS N. PAPADOPOULOS

A golfer faced with the problem of hitting a ball out of the woods after an errant shot sometimes makes a brave choice: she aims straight for the trees, swinging the club as hard as possible in the hope that the ball will bounce off the trees and miraculously emerge from the woods. On page 232 of this issue, Bertolotti *et al.*¹ describe a technique for imaging objects through light-scattering media, such as fog and human tissue, that overcomes a challenge that is in some ways similar to this one.

Consider light from a torch passing through a human hand. Information about the shapes of the bones, or even the cells, that make up the hand is thought to be encrypted in this transmitted light (Fig. 1), but a simple device such as a lens cannot be used to image the hand's interior. Numerous attempts have been made to retrieve the shapes of objects that hide behind or are within media that transmit and scatter light. Some of the photons that travel through a light-scattering medium do so without interacting with any of the medium's constituent matter. Such 'ballistic' photons exit the medium a little earlier than their

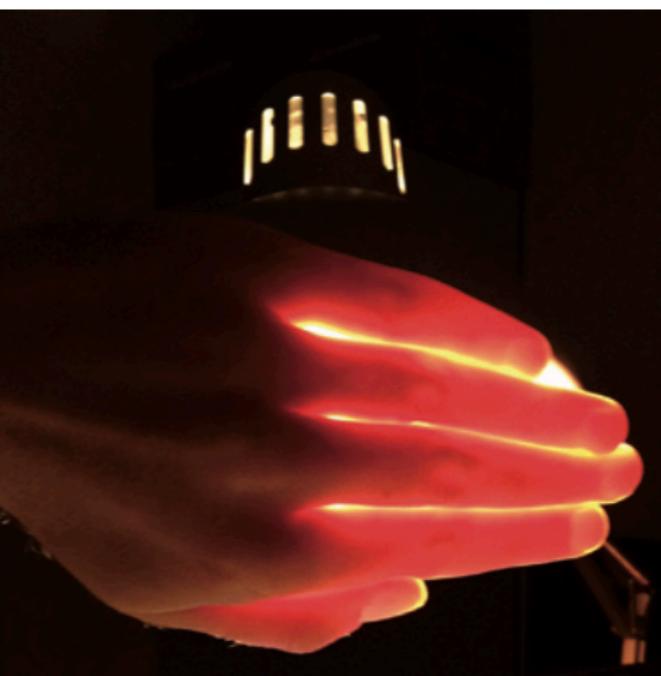
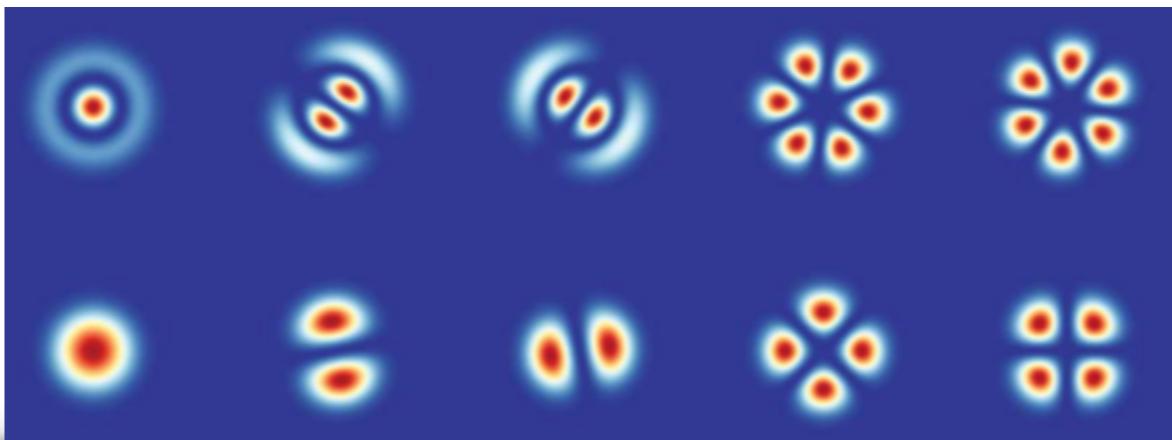
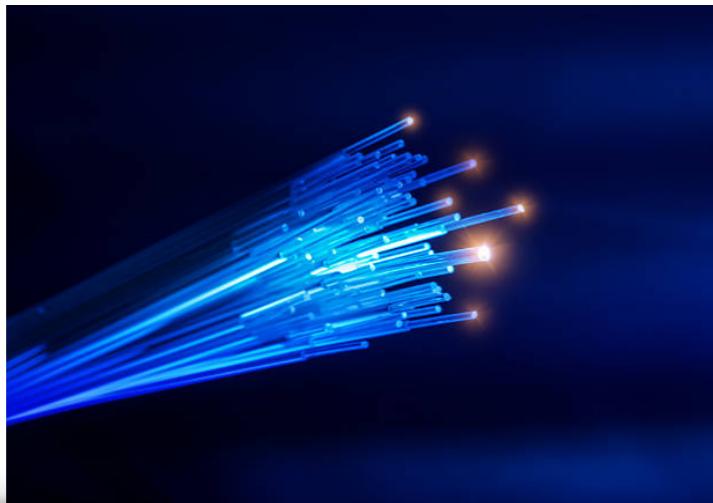


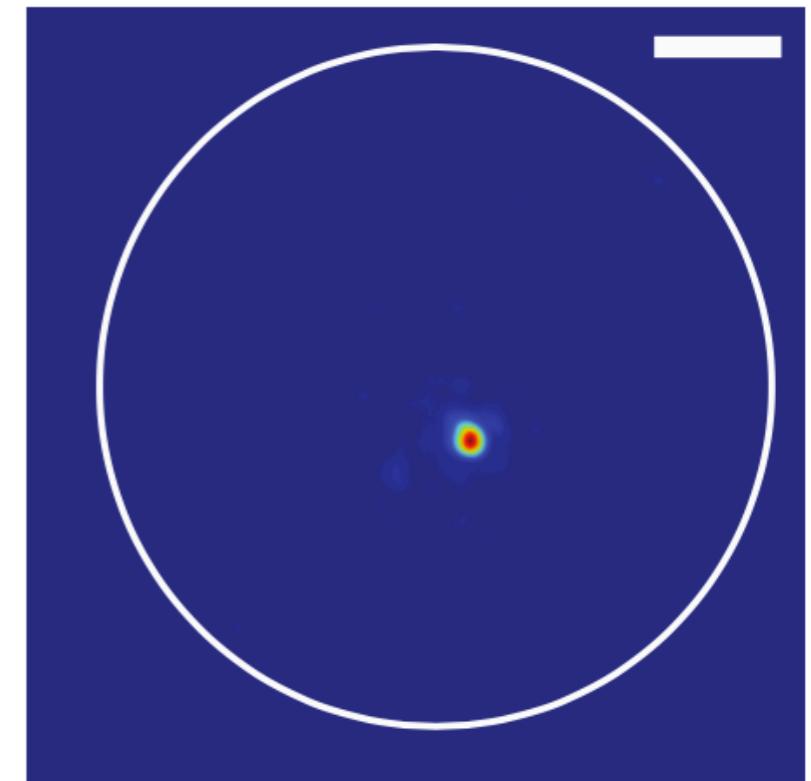
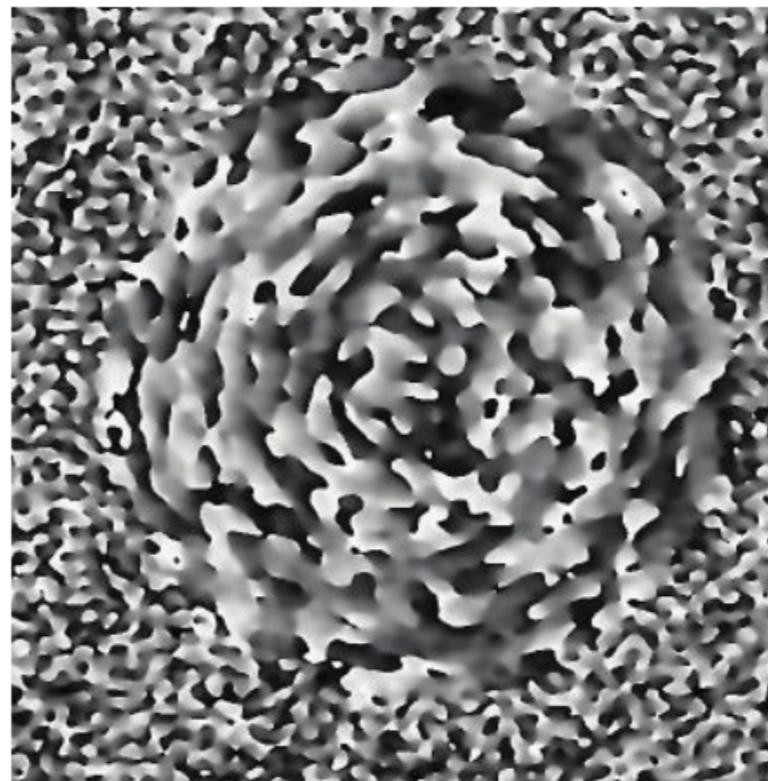
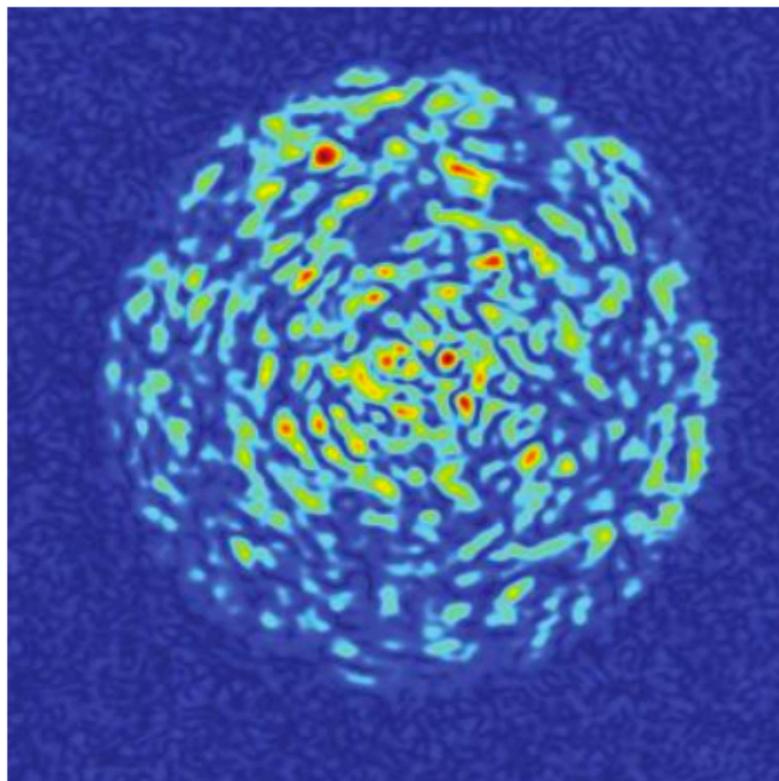
Figure 1 | Seeing through. A light source placed behind a human hand emits photons that travel through the hand. Information about the hand's interior, such as its bones or cells, is encoded in the light's field but cannot be directly retrieved because it is scrambled by the scattering properties of human tissue. Bertolotti *et al.*¹ propose a method for retrieving such information.

non-ballistic counterparts, which bounce off the matter as they pass through the scattering medium. If the ballistic photons alone are captured in a detector, the blurring effects of scattering can be avoided². However, for

Focusing by phase conjugation



In many applications we have multimode operation



Input Intensity

Input Phase

Output Intensity

Into the Labyrinth



TRANSMISSION MATRIX - GENERAL

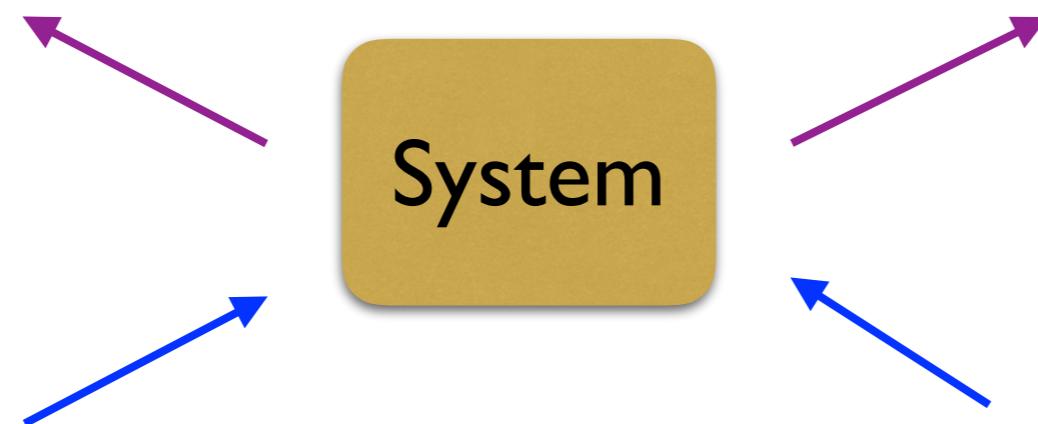
Scattering Matrix

$$\begin{pmatrix} u_1 \\ u_2 \end{pmatrix} = S \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} \equiv \begin{pmatrix} r_L & t \\ t & r_R \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix}$$

Outgoing waves

Scattering matrix

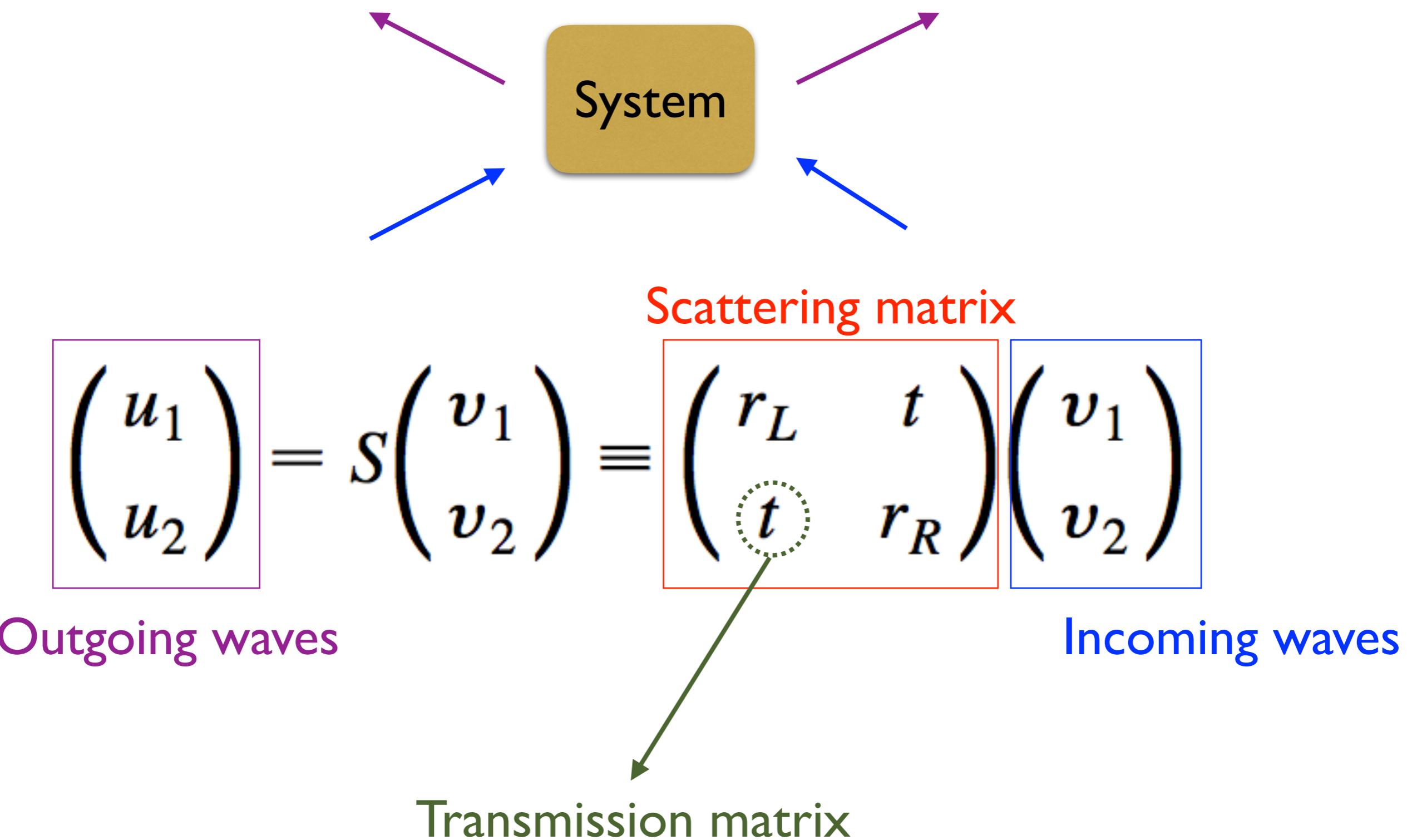
Incoming waves



Energy conservation - Unitarity: $S^\dagger S = I$

Reciprocity - Symmetry: $S^T = S$

Transmission Matrix



Transmission Matrix - Multimode Fibers

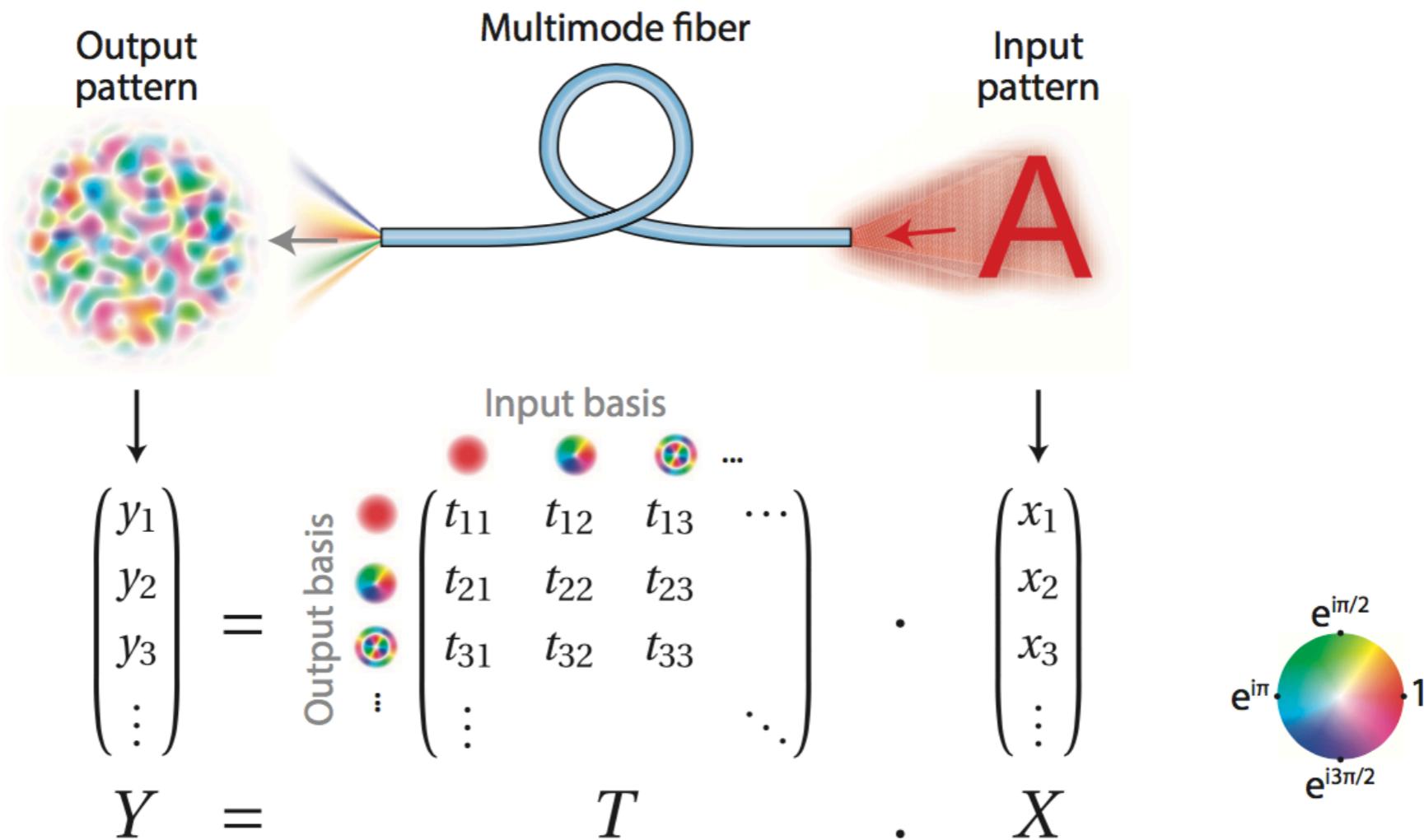
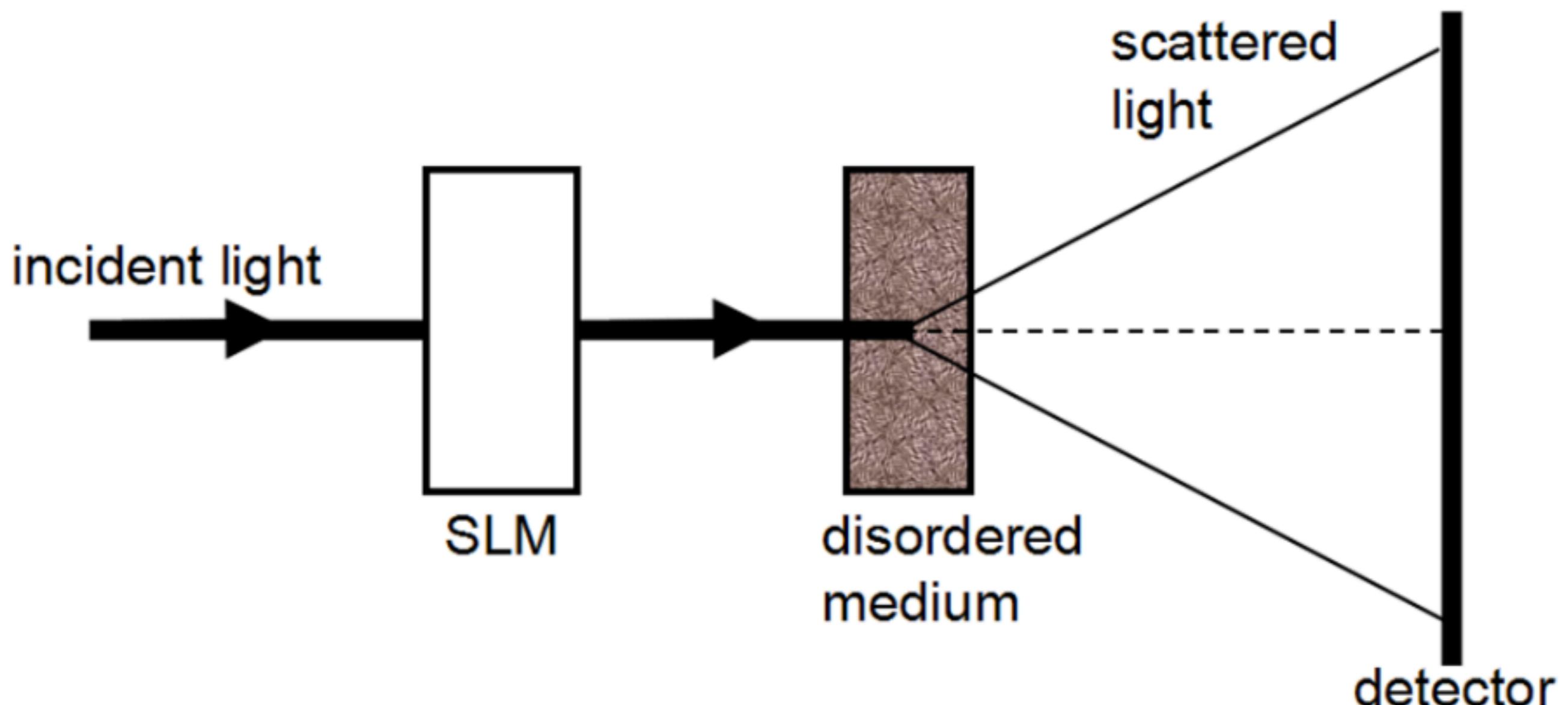


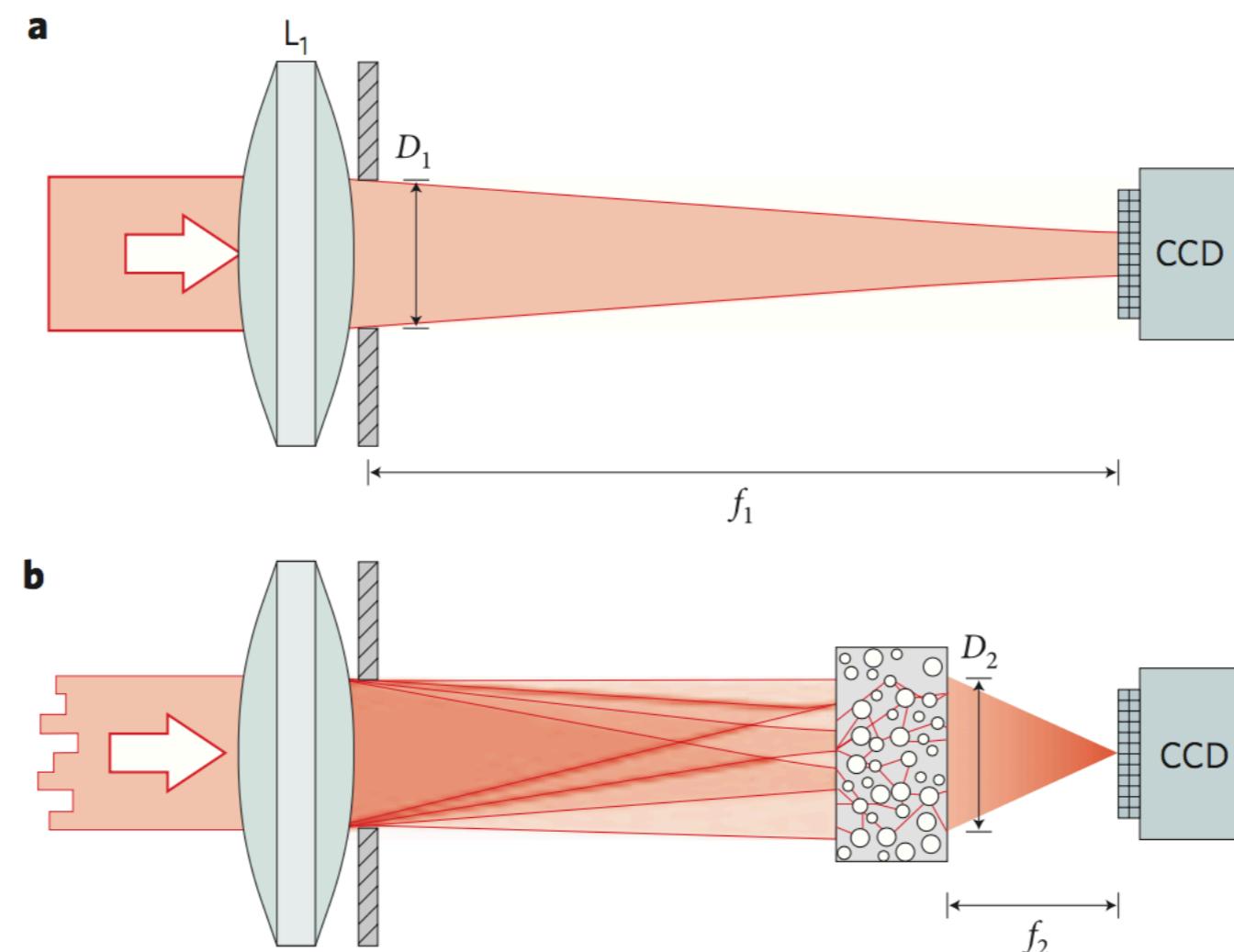
Figure 3.1: The transmission matrix model for multimode fibers. The input pattern X is injected into the fiber, where it undergoes a transformation T , and exits as the output pattern Y . X and Y are vectors of complex coefficients describing the optical fields at input and output respectively in a suitable basis. In this illustration, the input and output bases consist of the fiber modes, but other choices are possible too. T is the transmission matrix relating the input coefficients X and the output coefficients Y .

Transmission Matrix - Disordered media



Exploiting disorder for perfect focusing

I. M. Vellekoop^{1†}, A. Lagendijk² and A. P. Mosk^{1*}



Transmission Matrix - Disordered media

PRL 104, 100601 (2010)

 Selected for a **Viewpoint** in *Physics*
PHYSICAL REVIEW LETTERS

week ending
12 MARCH 2010



Measuring the Transmission Matrix in Optics: An Approach to the Study and Control of Light Propagation in Disordered Media

S. M. Popoff, G. Lerosey, R. Carminati, M. Fink, A. C. Boccara, and S. Gigan

Institut Langevin, ESPCI ParisTech, CNRS UMR 7587, ESPCI, 10 rue Vauquelin, 75005 Paris, France

(Received 27 October 2009; revised manuscript received 11 January 2010; published 8 March 2010)

We introduce a method to experimentally measure the monochromatic transmission matrix of a complex medium in optics. This method is based on a spatial phase modulator together with a full-field interferometric measurement on a camera. We determine the transmission matrix of a thick random scattering sample. We show that this matrix exhibits statistical properties in good agreement with random matrix theory and allows light focusing and imaging through the random medium. This method might give important insight into the mesoscopic properties of a complex medium.

Transmission Matrix - Disordered media

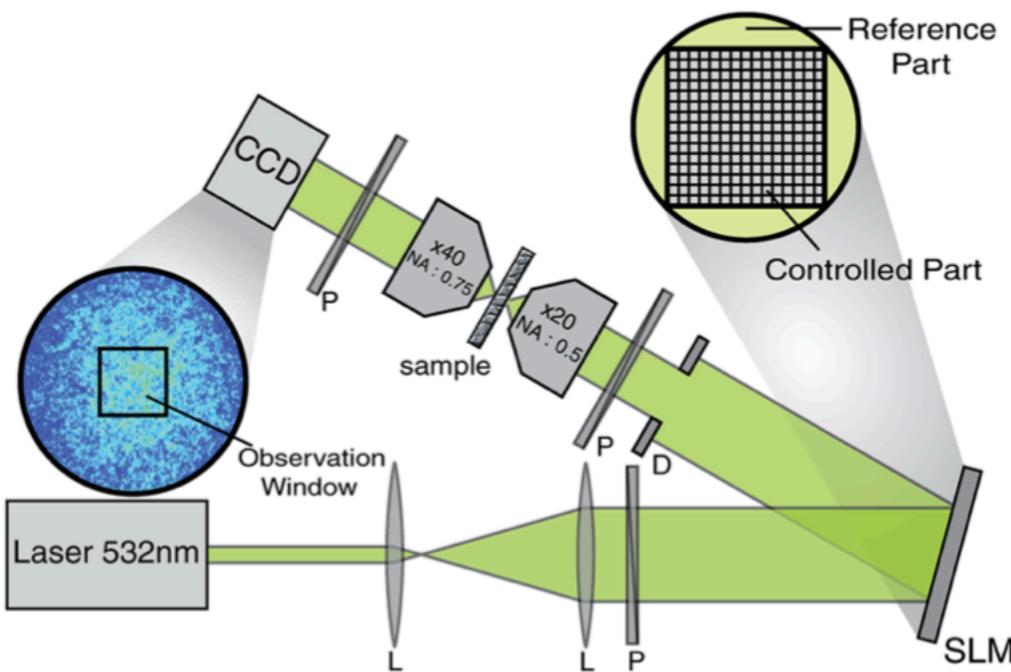


FIG. 1 (color online). Schematic of the apparatus. The laser is expanded and reflected off a SLM. The phase-modulated beam is focused on the multiple-scattering sample and the output intensity speckle pattern is imaged by a CCD camera: lens (L), polarizer (P), diaphragm (D).

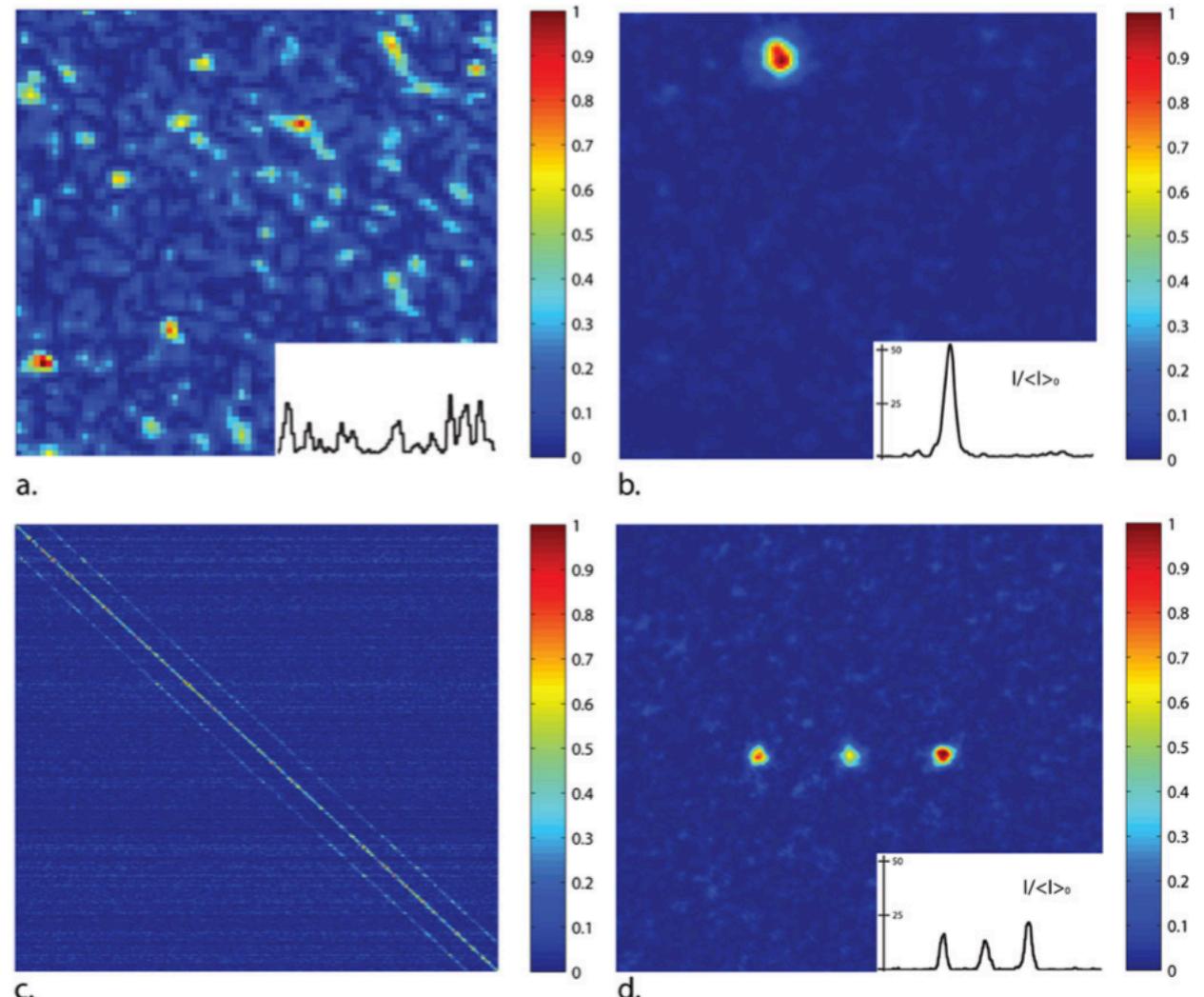


FIG. 2 (color online). Experimental results of focusing. (a) Initial aspect of the output speckle. (b) We measure the TM for 256 controlled segments and use it to perform phase conjugation. (c) Norm of the focusing operator $O_{\text{norm}}^{\text{foc}}$. (d) Example of focusing on several points. (The insets show intensity profiles along one direction.)

The information age in optics: Measuring the transmission matrix

Elbert G. van Putten and **Allard P. Mosk**

Complex Photonic Systems, Faculty of Science and Technology and MESA+ Institute for Nanotechnology, University of Twente, P.O. Box 217, 7500 AE Enschede, The Netherlands

March 8, 2010 • *Physics* 3, 22

The transmission of light through a disordered medium is described in microscopic detail by a high-dimensional matrix. Researchers have now measured this transmission matrix directly, providing a new approach to control light propagation.

Transmission Matrix - Disordered media

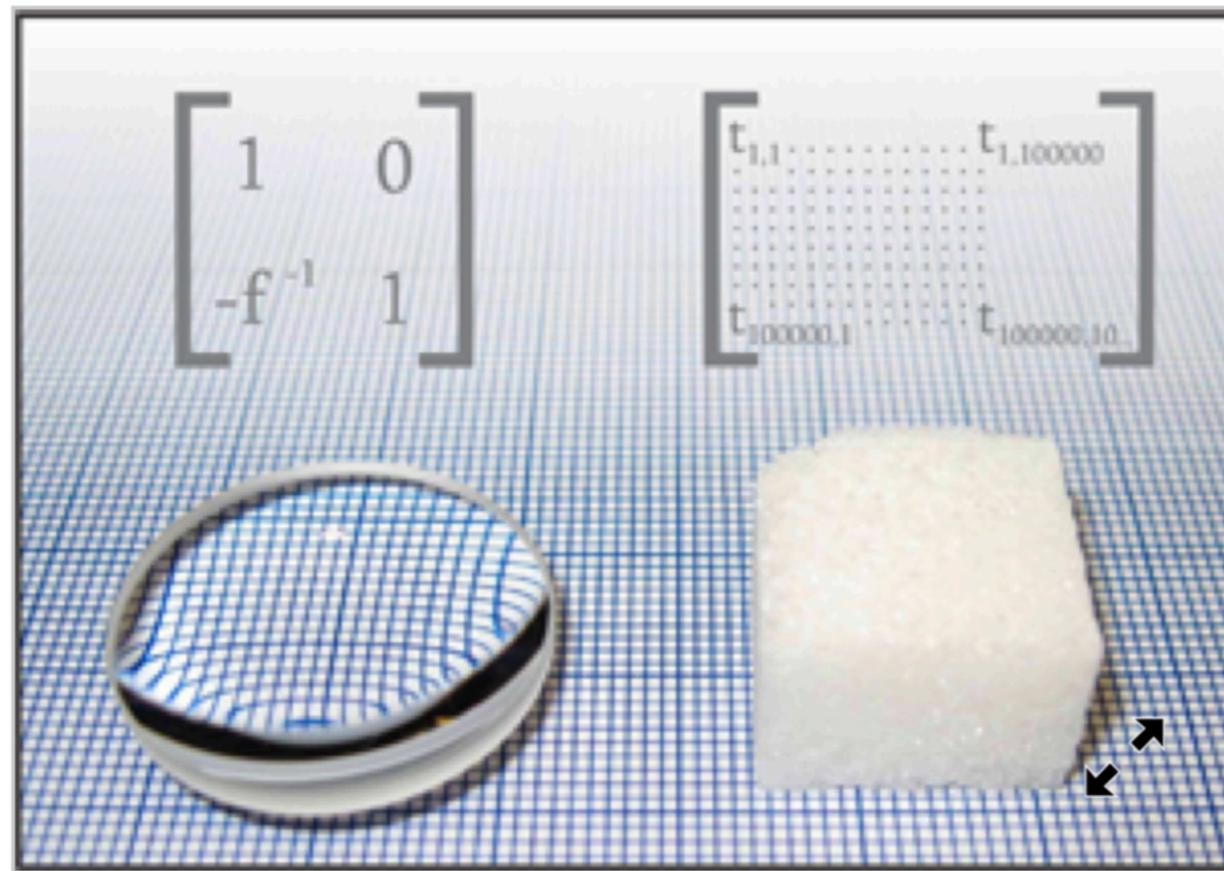


Figure 1: Two optical elements fully characterized by their transmission matrix, which relates the incident wave front to the transmitted one. In the case of a thin lens, the transformation of the wave front is described by a 2×2 matrix operating on a vector describing the wave front curvature [27]. For more complex elements such as a sugar cube the transmission matrix operates in a basis of transversal modes, which is very large. Full knowledge of the transmission matrix enables disordered materials to focus light as lenses.

ARTICLES

Optical phase conjugation for turbidity suppression in biological samples

ZAHID YAQOOB^{1†}, DEMETRI PSALTIS^{1,2}, MICHAEL S. FELD³ AND CHANGHUEI YANG^{1*}

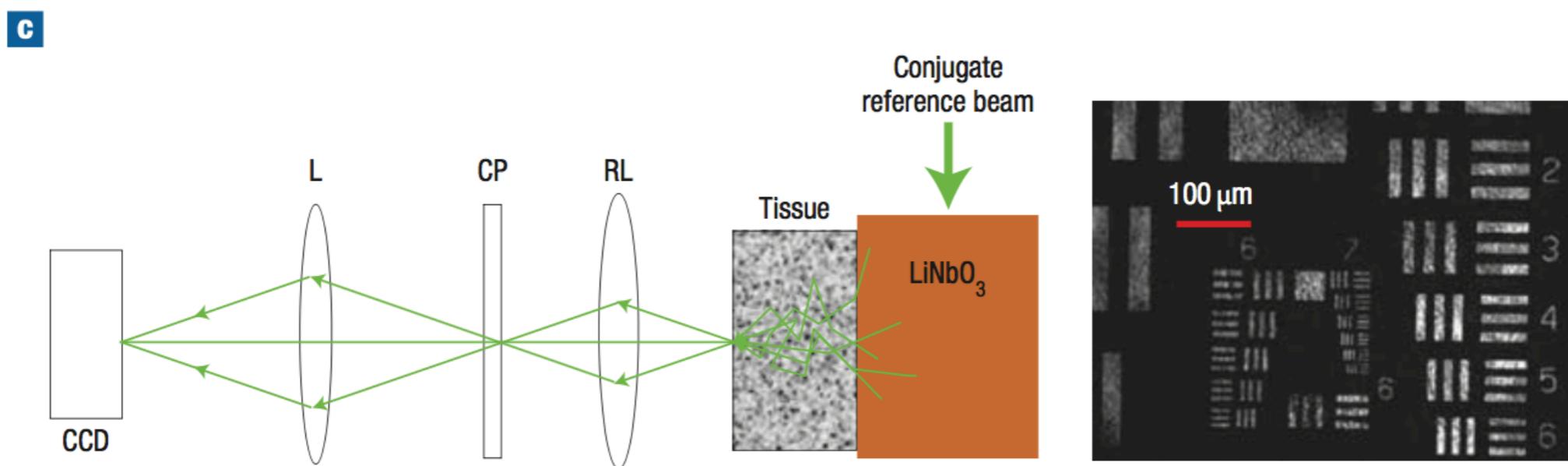
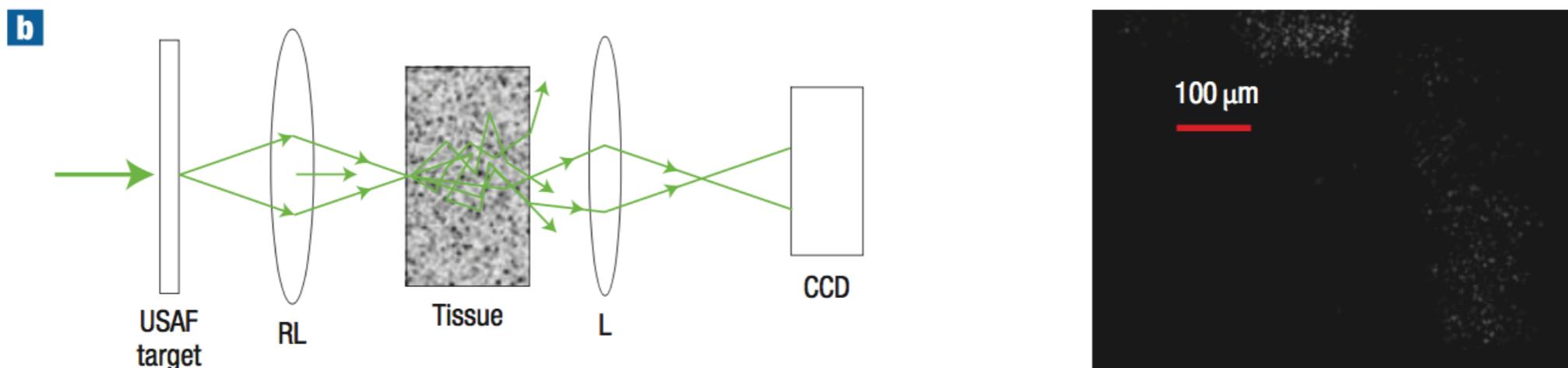
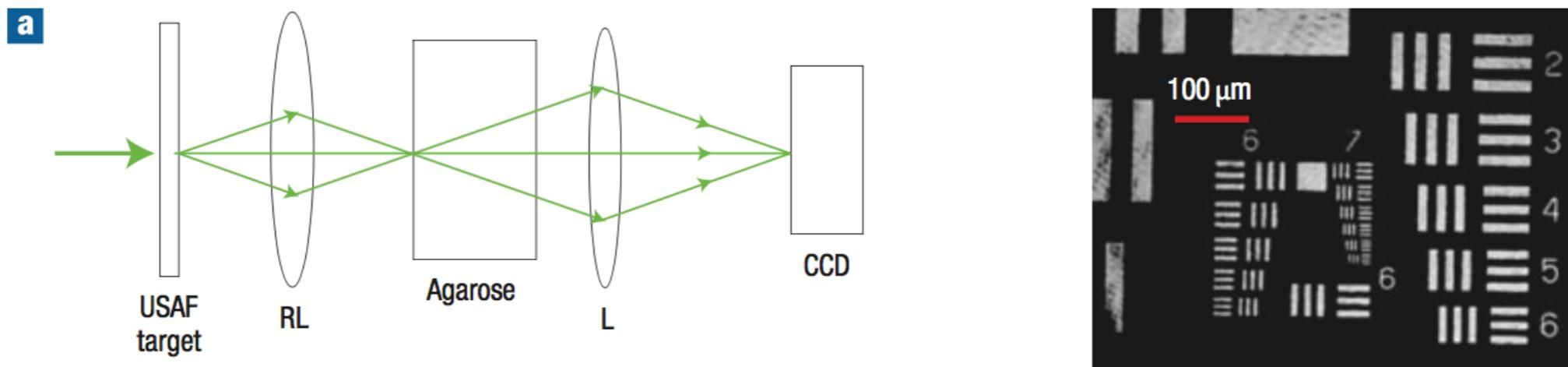
¹Electrical Engineering, California Institute of Technology, 1200 East California Boulevard, Pasadena, California 91125, USA

²School of Engineering, Ecole Polytechnique Fédérale de Lausanne, CH-1015 Lausanne, Switzerland

³Massachusetts Institute of Technology, G. R. Harrison Spectroscopy Laboratory, 77 Massachusetts Avenue 6-207, Cambridge, Massachusetts 02139, USA

†Present address: Massachusetts Institute of Technology, G. R. Harrison Spectroscopy Laboratory, 77 Massachusetts Avenue 6-208, Cambridge, Massachusetts 02139, USA

Transmission Matrix - Disordered media



TRANSMISSION MATRIX - PROPERTIES

Transmission Matrix

$$\begin{pmatrix} E_1^{\text{out}} \\ E_2^{\text{out}} \\ E_3^{\text{out}} \\ \vdots \\ E_M^{\text{out}} \end{pmatrix} = \begin{pmatrix} t_{11} & t_{12} & t_{13} & \cdots & t_{1N} \\ t_{21} & t_{22} & t_{23} & \cdots & t_{2N} \\ t_{31} & t_{32} & t_{33} & \cdots & t_{3N} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ t_{M1} & t_{M2} & t_{M3} & \cdots & t_{MN} \end{pmatrix} \begin{pmatrix} E_1^{\text{in}} \\ E_2^{\text{in}} \\ E_3^{\text{in}} \\ \vdots \\ E_N^{\text{in}} \end{pmatrix}.$$

Output amplitudes Transmission matrix Input amplitudes

Linear Transformation

$$E_m^{\text{out}} = \sum_{n=1}^N t_{mn} E_n^{\text{in}},$$

Matrices and Linear Transformation

Theorem of Linear Algebra

Let $T : \mathbb{R}^n \mapsto \mathbb{R}^m$ be a linear transformation. Then the matrix A satisfying $T(\vec{x}) = A\vec{x}$ is given by

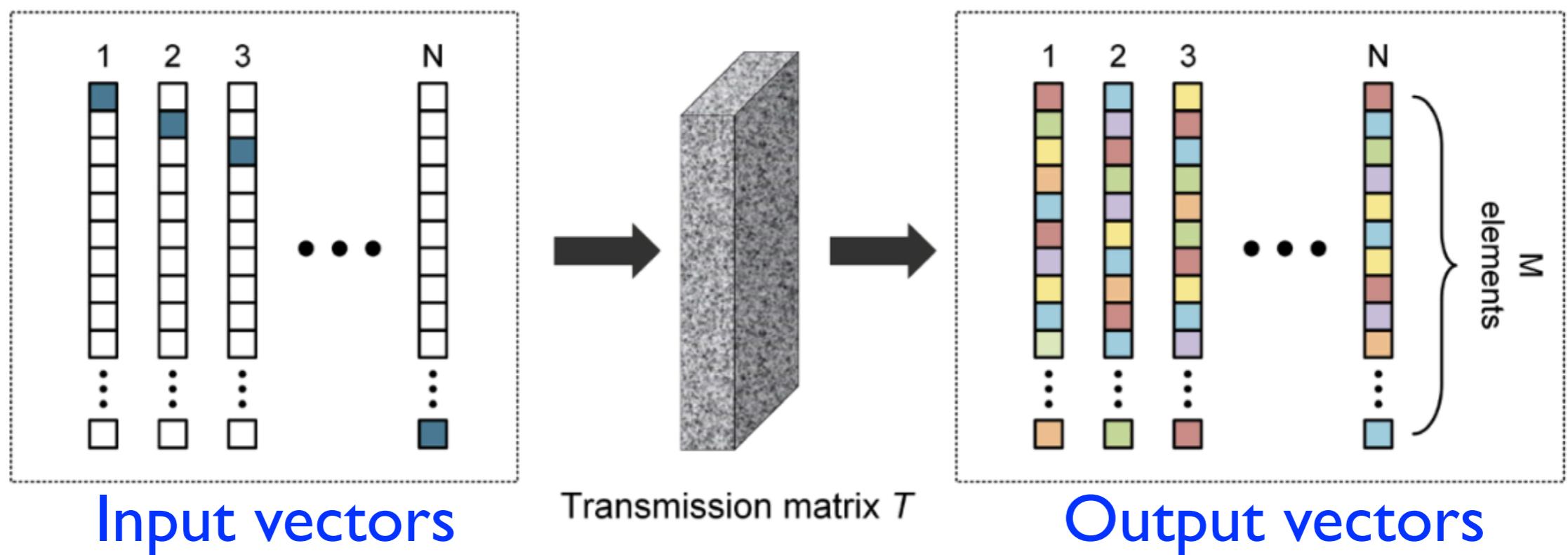
$$A = \begin{bmatrix} | & & | \\ T(\vec{e}_1) & \dots & T(\vec{e}_n) \\ | & & | \end{bmatrix}$$

where \vec{e}_i is the i^{th} column of I_n , and then $T(\vec{e}_i)$ is the i^{th} column of A .

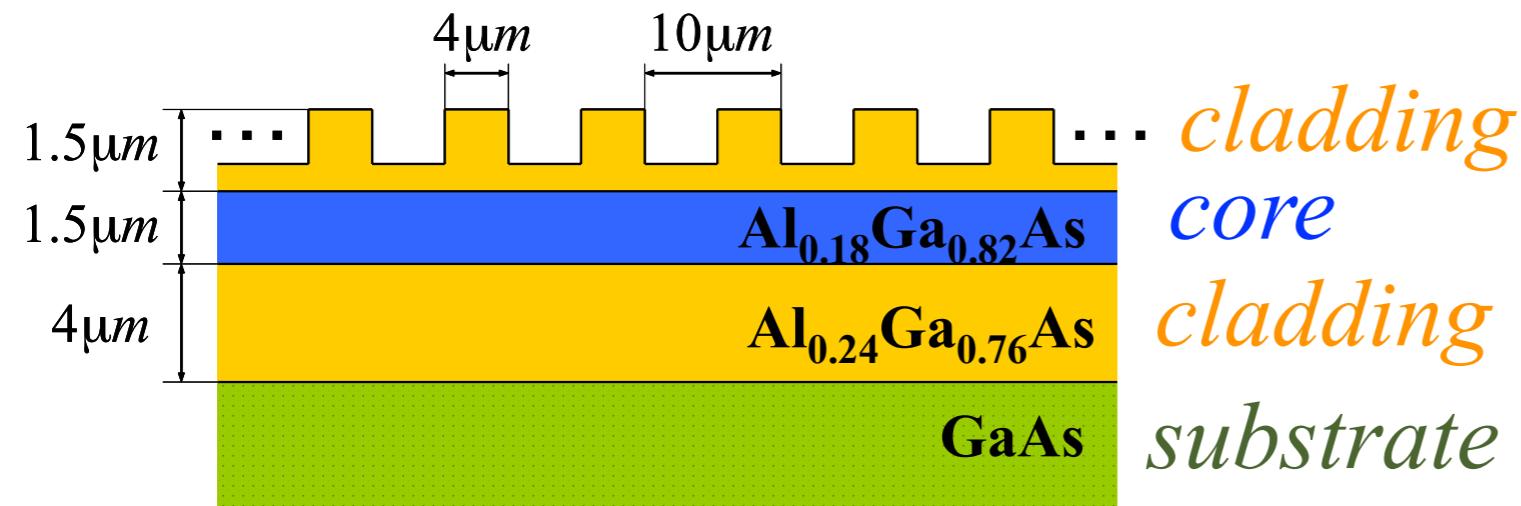


Every linear transformation is related to a matrix, given the bases of the two vector spaces that the linear map is defined

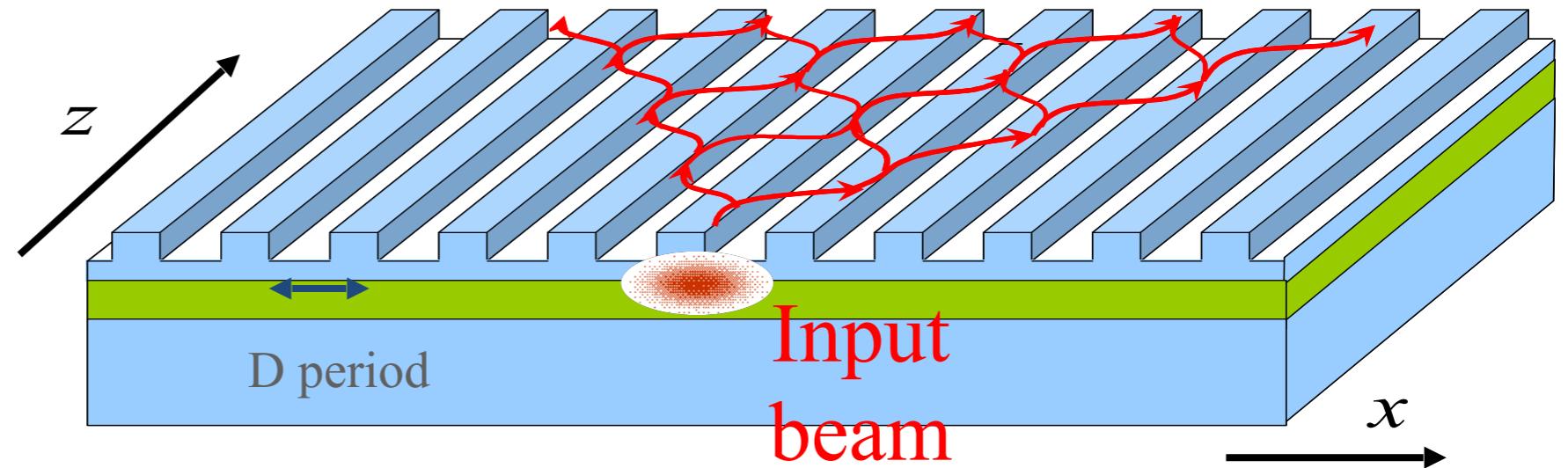
Transmission Matrix - Linear Transformation



Simplest Example - Waveguide Array



cladding
core
cladding
substrate



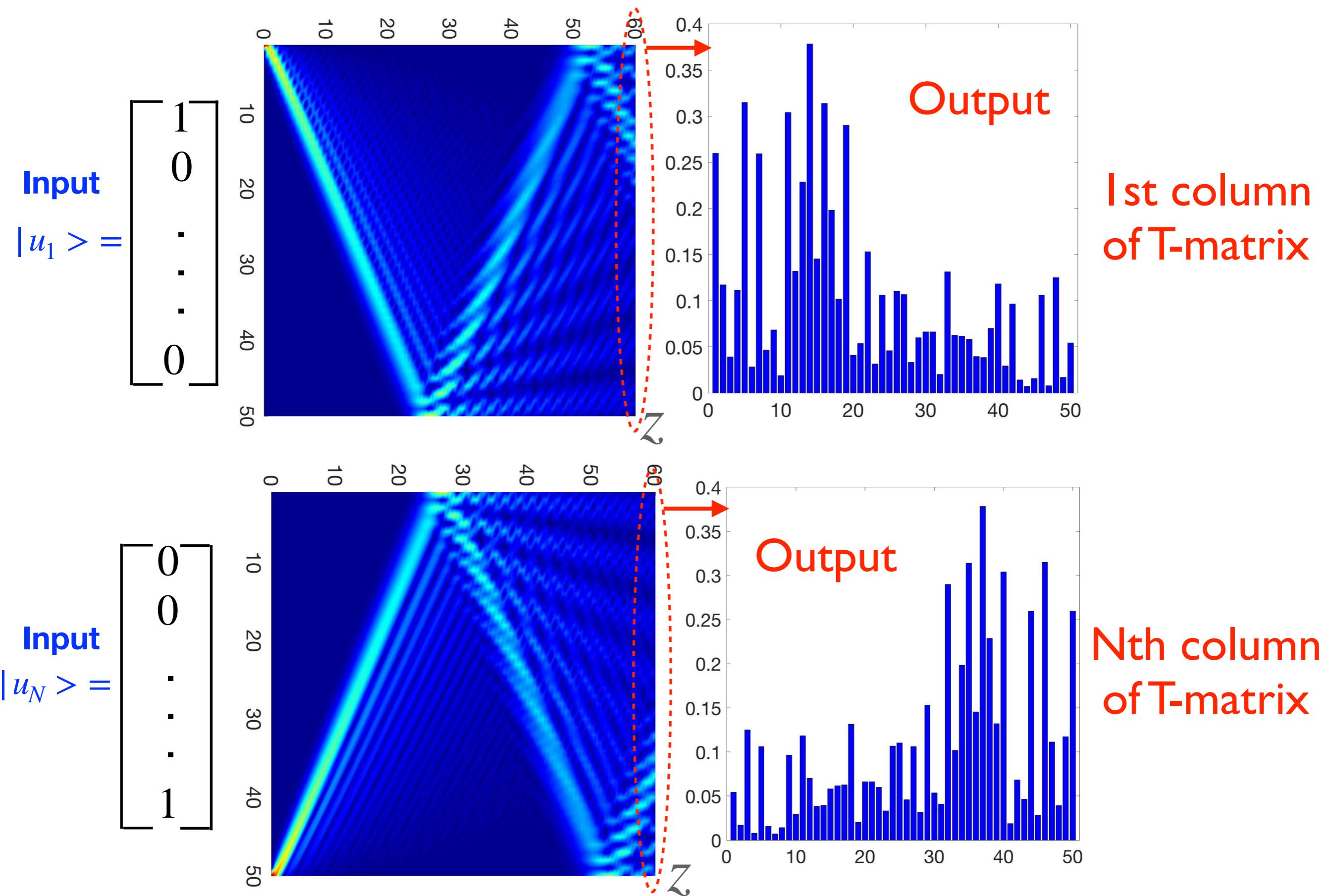
$$i \frac{dc_n}{dz} + \kappa(c_{n+1} + c_{n-1}) + V_n c_n = 0$$

coupling

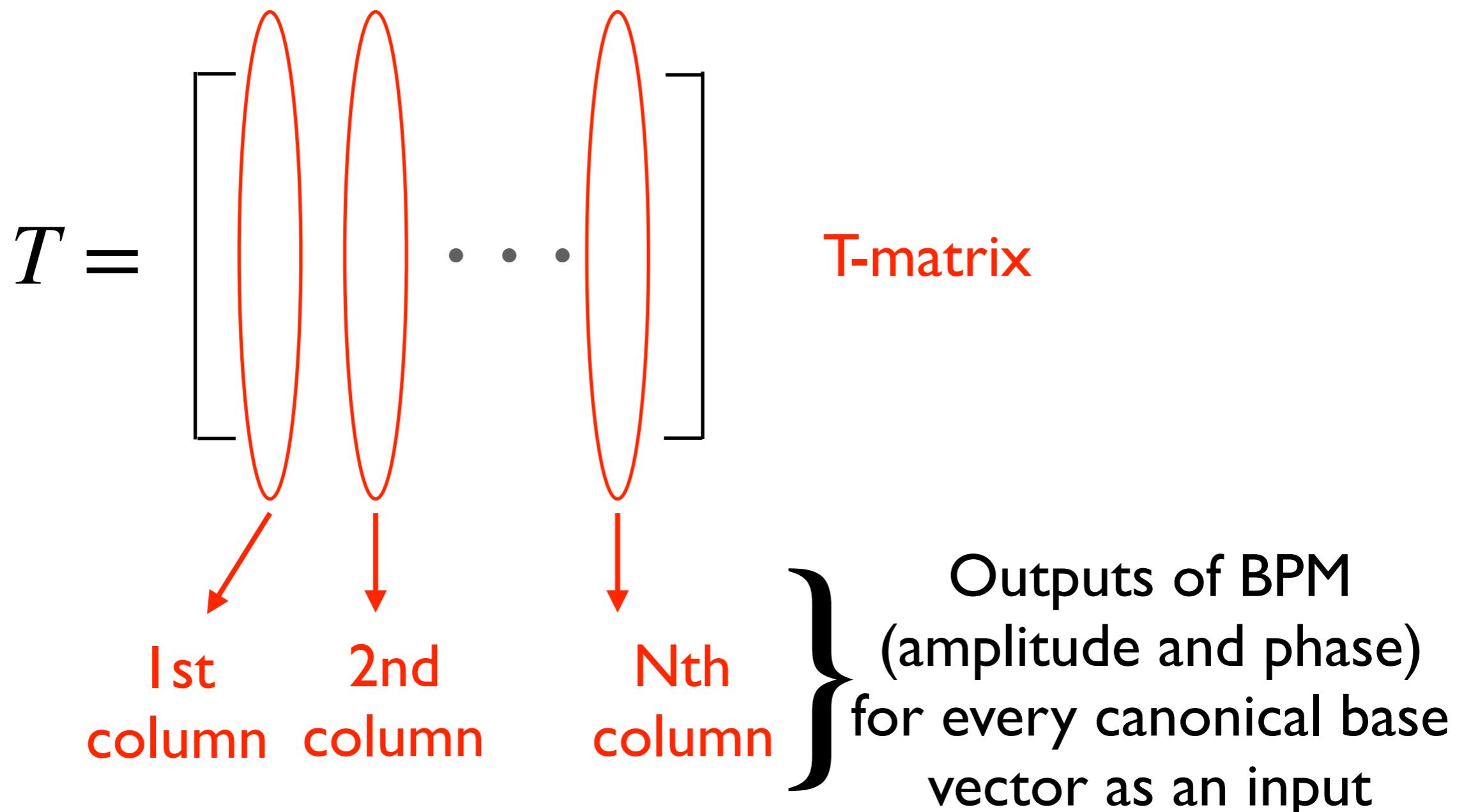
Discrete Diffraction Equation

Coupled mode Theory (tight-bind approximation)

Constructing the transmission matrix



Constructing the transmission matrix

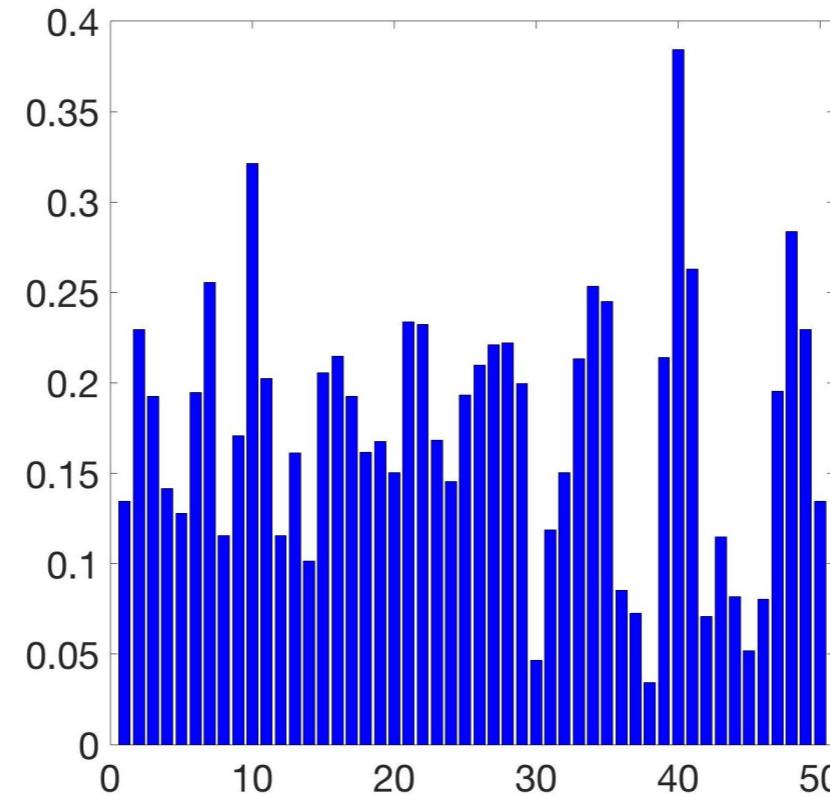
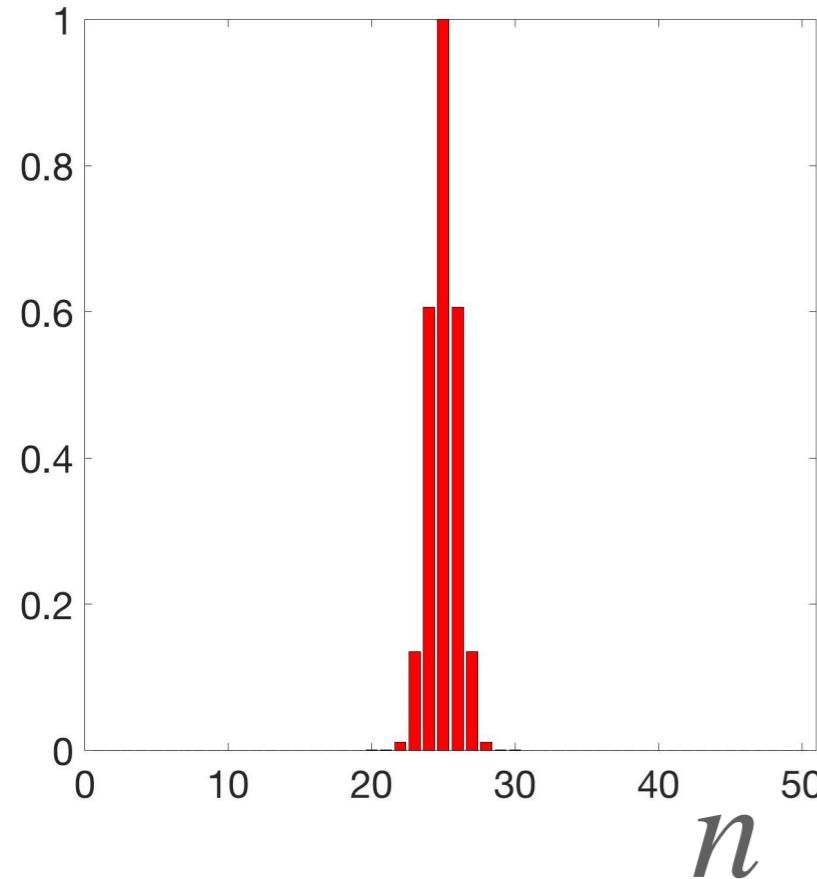


Focusing by inversion

$$T|\psi_{in}\rangle = |\psi_{out}\rangle \Rightarrow$$

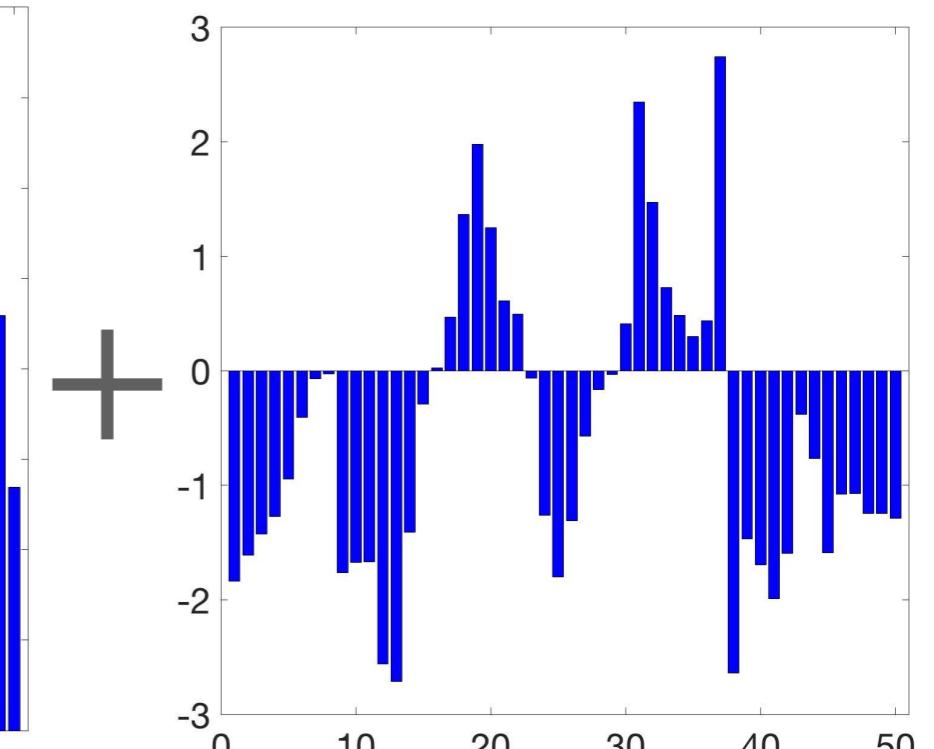
$$T^{-1}|\psi_{out}\rangle = |\psi_{in}\rangle$$

Desired Focused spot



Amplitude

Calculate input



Phase

Focusing by inversion

