

Solution for exercise 9 – Part I:

Design and force balancing of a flexure-based parallel stage

Context:

A flexure-base parallel stage is subjected to fluctuating linear accelerations along the x -axis due to motion of its base. To be insensitive to this specific perturbation, we wish to force balance the mechanism.

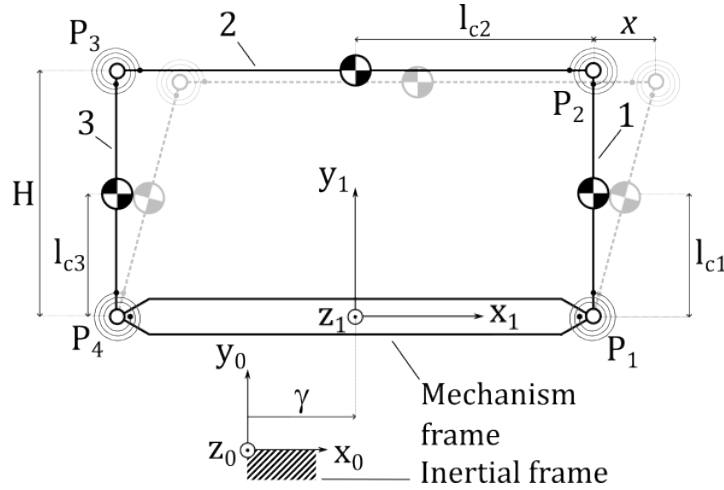


Fig. 1: Flexure-based parallel stage Pseudo-Rigid-Body-Model before its force balancing.

Figure 1 shows the Pseudo-Rigid-Body-Model (PRBM) of the studied flexure-based parallel stage before its force balancing. The mechanism is shown in its rest position (black lines) and in its deformed position (transparent black lines). The mechanism has one degree-of-freedom and is composed of three moving rigid bodies: solids **1**, **2** and **3** articulated in P_1 , P_2 , P_3 and P_4 .

Each solid $i = 1 \dots 3$, has:

- a mass m_i ,
- a center of mass (COM) c_i (represented by the symbol \odot) located by a distance l_{ci} ,
- an inertia tensor $\bar{\bar{J}}_i$ written at c_i .

Each pivot $P_{i=1\dots4}$ has an angular stiffness k_i that represents the intrinsic stiffness of the future flexure-based implementation.

First, the parallel stage mechanism is actuated by a voice coil whose permanent magnet is fixed to the mechanism's base and whose coil is fixed to solid **2**. The voice coil applies a force F at c_2 along the x -axis. The parameter x is used to locate the displacement of solid **2** relative to its base.

Last, the linear displacement that corresponds to the motion of the mechanism's base is represented by the parameter γ .

Part I: Force balancing of the flexure-based parallel stage

In order to force balance the mechanism, we propose to redistribute its mass by lowering the COM of solids 1 and 3 thanks to two identical supplementary masses m_{bal} . These two masses are located by the distance l_{bal} .

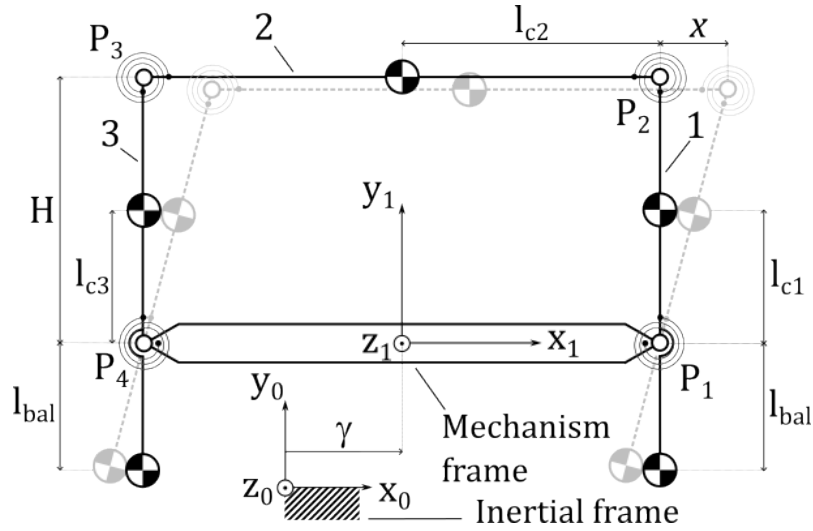


Fig. 2: Flexure-based parallel stage Pseudo-Rigid-Body-Model after its force balancing.

Numerical Values:

- $H = 20e-3$ [m]
- $m_1 = 1.7138708e-02$ [kg]
- $m_2 = 6.8122300e-02$ [kg]
- $m_3 = m_1$
- $m_{bal} = 3.5565044e-02$ [kg]
- $l_{c1} = 1.0411598e-2$ [m]
- l_{c2} : not needed
- $l_{c3} = l_{c1}$
- l_{bal} : to be found
- $\bar{J}_1 = \begin{bmatrix} 5.3559891e-6 & 0 & 0 \\ 0 & 1.9478338e-6 & 0 \\ 0 & 0 & 5.3569824e-6 \end{bmatrix}$ [kg. m²]
- \bar{J}_2 : not needed
- $\bar{J}_3 = \bar{J}_1$
- $\bar{J}_{bal} = \begin{bmatrix} 2.0034975e-6 & 0 & 0 \\ 0 & 2.5379149e-6 & 0 \\ 0 & 0 & 1.1271682e-6 \end{bmatrix}$ [kg. m²]

1. Using the Euler-Lagrange formalism, calculate the equation of motion of the parallel stage as a function of F , x and γ , assuming small deformations around the rest position. The equation of motion is of the form:

$$F = m_{eq}\ddot{x} + k_{eq}x + \tau_{coupling}\dot{\gamma}$$

Reminder: The Euler-Lagrange equation that needs to be solved is

$$F = \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}} \right) - \frac{\partial L}{\partial x}$$

Where $L = K - V$ is the Lagrangian; K and V are the kinetic and potential energies of the system, respectively.

SOLUTION:

We introduce the parameter θ to simplify the writing of the equations. Parameter θ , considered positive **anticlockwise**, represents the angle of solid 1 (respectively, solid 3) in its deformed position relative to its vertical rest position.

Using parameter θ , the kinetic and potential energies of the total mechanism can be expressed as

$$K = \frac{1}{2}m_{1,tot}\dot{c}_{1,tot}^2 + \frac{1}{2}m_2\dot{c}_2^2 + \frac{1}{2}m_{3,tot}\dot{c}_{3,tot}^2 + \frac{1}{2}J_{1,tot,zz}\dot{\theta}^2 + \frac{1}{2}J_{3,tot,zz}\dot{\theta}^2$$

$$V = \frac{1}{2}(k_1 + k_2 + k_3 + k_4)\theta^2$$

Where $m_{1,tot}$ and $\dot{c}_{1,tot}$ are the mass and the COM linear velocity vector of solid 1 equipped with its balancing mass; \dot{c}_2 is the COM linear velocity vector of solid 2; $m_{3,tot}$ and $\dot{c}_{3,tot}$ are the mass and the COM linear velocity vector of solid 3 equipped with its balancing mass; $\dot{\theta}$ is the angular velocity vector of solids 1 and 3; $J_{1,tot,zz}$ and $J_{3,tot,zz}$ are the moment of inertia at COM of solids 1 and 3 equipped with their respective balancing mass.

With

$$m_{1,tot} = m_1 + m_{bal} \quad \text{and} \quad m_{3,tot} = m_3 + m_{bal}$$

$$\dot{c}_{1,tot} = (\dot{\gamma} - l_{c1,tot} \cos \theta \dot{\theta})x_0 - l_{c1,tot} \sin \theta \dot{\theta}y_0$$

$$\dot{c}_2 = (\dot{\gamma} - H \cos \theta \dot{\theta})x_0 - H \sin \theta \dot{\theta}y_0$$

$$\dot{c}_{3,tot} = (\dot{\gamma} - l_{c3,tot} \cos \theta \dot{\theta})x_0 - l_{c3,tot} \sin \theta \dot{\theta}y_0$$

$$l_{c1,tot} = \frac{m_{bal}l_{bal} - m_1l_{c1}}{m_{1,tot}} \quad \text{and} \quad l_{c3,tot} = \frac{m_{bal}l_{bal} - m_3l_{c3}}{m_{3,tot}}$$

$$J_{1,tot,zz} = J_{1,zz} + J_{bal,zz} + (l_{c1,tot} + l_{c1})^2 m_1 + (l_{bal} - l_{c1,tot})^2 m_{bal}$$

$$J_{3,tot,zz} = J_{3,zz} + J_{bal,zz} + (l_{c3,tot} + l_{c3})^2 m_3 + (l_{bal} - l_{c3,tot})^2 m_{bal}$$

The kinetic energy simplifies and can then be written as

$$\begin{aligned} K = & \frac{1}{2} m_{1,tot} (\dot{\gamma}^2 + l_{c1,tot}^2 \dot{\theta}^2 - 2l_{c1,tot} \dot{\gamma} \dot{\theta} \cos \theta) \\ & + \frac{1}{2} m_2 (\dot{\gamma}^2 + H^2 \dot{\theta}^2 + 2H \dot{\gamma} \dot{\theta} \cos \theta) \\ & + \frac{1}{2} m_{3,tot} (\dot{\gamma}^2 + l_{c3,tot}^2 \dot{\theta}^2 - 2l_{c3,tot} \dot{\gamma} \dot{\theta} \cos \theta) + \frac{1}{2} J_{1,tot,zz} \dot{\theta}^2 \\ & + \frac{1}{2} J_{3,tot,zz} \dot{\theta}^2 \end{aligned}$$

We also have the following simplifications:

$$\begin{aligned} m_1 = m_3, \quad l_{c1} = l_{c3}, \quad J_{1,zz} = J_{3,zz} \quad \text{and} \quad k_1 = k_2 = k_3 = k_4. \\ x = H \sin \theta \approx H\theta \quad \text{and} \quad \cos \theta \approx 1 \end{aligned}$$

Using these simplifications, we can rewrite K and V so that:

$$\begin{aligned} K \approx & \frac{1}{2} (2m_{1,tot} + m_2) \dot{\gamma}^2 + \frac{1}{2} \frac{2J_{1,tot,zz} + 2m_{1,tot}l_{c1,tot}^2 + m_2H^2}{H^2} \dot{\theta}^2 \\ & + \frac{m_2H - 2m_{1,tot}l_{c1}}{H} \dot{\gamma} \dot{\theta} \end{aligned}$$

$$V \approx 2 \frac{k_1}{H^2} x^2$$

We then find

$$F = \frac{2J_{1,tot,zz} + 2m_{1,tot}l_{c1,tot}^2 + m_2H^2}{H^2} \ddot{x} + 4 \frac{k_1}{H^2} x + \frac{m_2H - 2m_{1,tot}l_{c1,tot}}{H} \ddot{\gamma}$$

So that

$$m_{eq} = \frac{2J_{1,tot,zz} + 2m_{1,tot}l_{c1,tot}^2 + m_2H^2}{H^2}$$

$$k_{eq} = 4 \frac{k_1}{H^2}$$

$$\tau_{coupling} = \frac{m_2 H - 2m_{1,tot} l_{c1,tot}}{H}$$

2. Identify the dynamic coupling term $\tau_{coupling}$ that couples the mechanism motion to its base acceleration \ddot{y} .

SOLUTION:

$$\tau_{coupling} = \frac{m_2 H - 2m_{1,tot} l_{c1,tot}}{H}$$

3. Find the length l_{bal} that cancels $\tau_{coupling}$.

SOLUTION:

$$l_{bal} = \frac{2m_1 l_{c1} + m_2 H}{2m_{bal}} = 24.1716 \text{ [mm]}$$

4. Compute the linear equivalent masses $m_{eq,init}$ and $m_{eq,tot}$ of the parallel stage before and after it is forced balanced.

SOLUTION:

$$m_{eq,init} = \frac{2(J_{1,zz} + m_1 l_{c1}^2) + m_2 H^2}{H^2} \approx 104.2 \text{ [g]}$$

$$m_{eq,tot} = \frac{2(J_{1,tot,zz} + m_{1,tot} l_{c1,tot}^2) + m_2 H^2}{H^2} \approx 213.7 \text{ [g]}$$

5. Considering $F = 0$, write the analytical transfer functions $T_{init}(\omega)$ and $T_{final}(\omega)$ of the parallel stage before and after it is forced balanced:

$$T(\omega) = \frac{X(\omega)}{(j\omega)^2 \Gamma(\omega)}.$$

SOLUTION:

$$T(\omega) = \frac{\tau_{coupling}}{k_{eq}} \frac{1}{\left(\frac{\omega}{\omega_0}\right)^2 - 1}$$

With

$$\omega_0^2 = \frac{k_{eq}}{m_{eq}}$$