

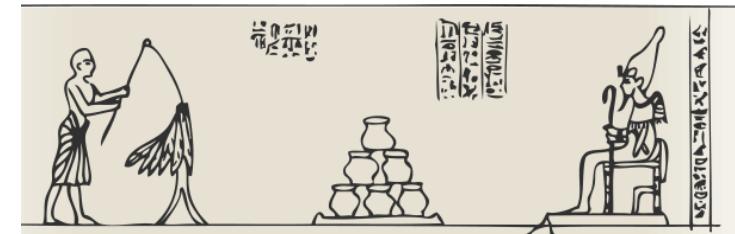
Heat Pump Systems

Summary W2

Prof. J. Schiffmann

Heating & Cooling in Historical Context

- Heat through combustion is simple and cheap
- Artificial cooling is more challenging to achieve
- Industrial vapor compression heat pump cycles only beginning of 20th century
- Today heat pumps are spread within various applications ranging from domestic to industrial applications



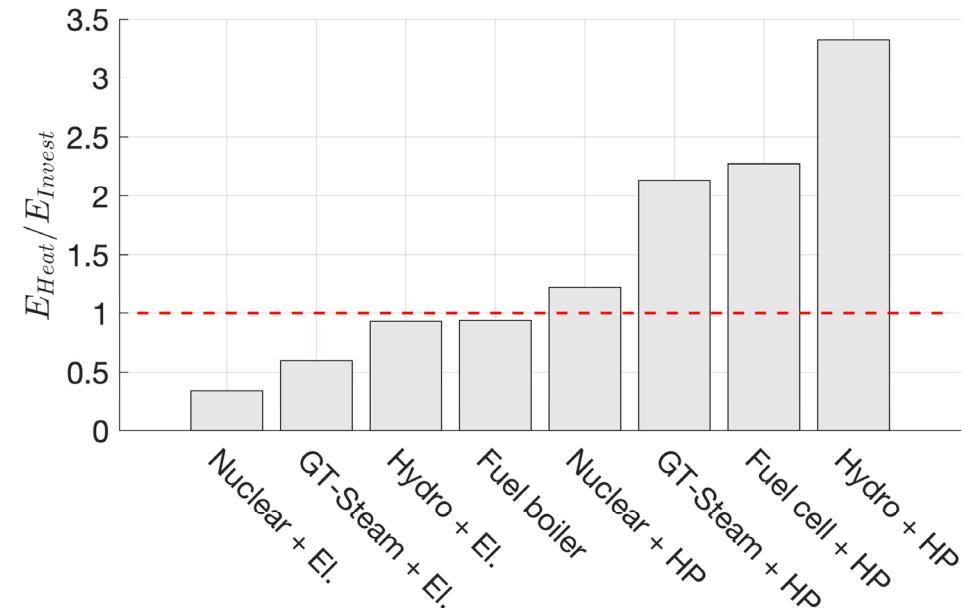
<http://www.carel-japan.com/high-efficiency-solutions/evaporative-cooling/>



[americanhistory.si.edu blog ice-harvesting-electric-refrigeration](http://americanhistory.si.edu/blog/ice-harvesting-electric-refrigeration)

Domestic HVAC

- 27% of primary energy consumption relates to domestic HVAC
- Heat pumps play key role in reducing energy consumption and CO₂ emissions



What is a Heat Pump Thermodynamically?

- Bithermal thermodynamic cycle working in anti-clockwise direction
- Work is invested to drive the cycle, which absorbs and supplies heat at different temperatures

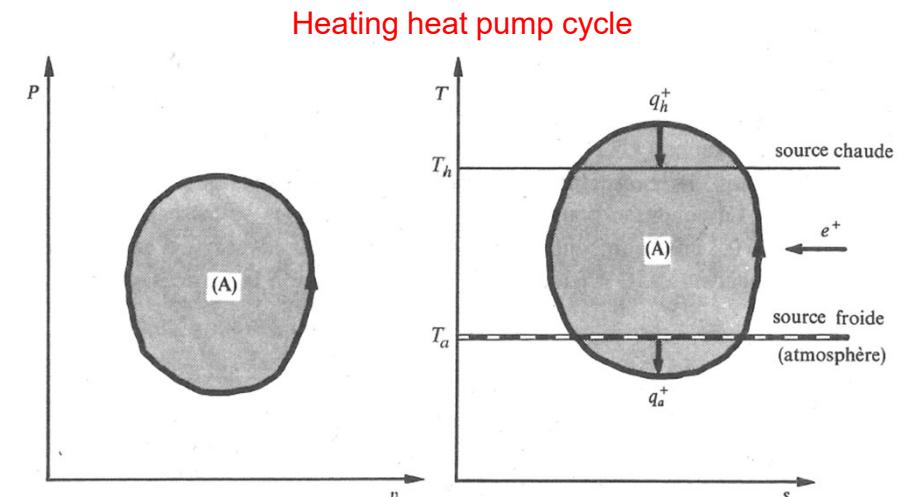
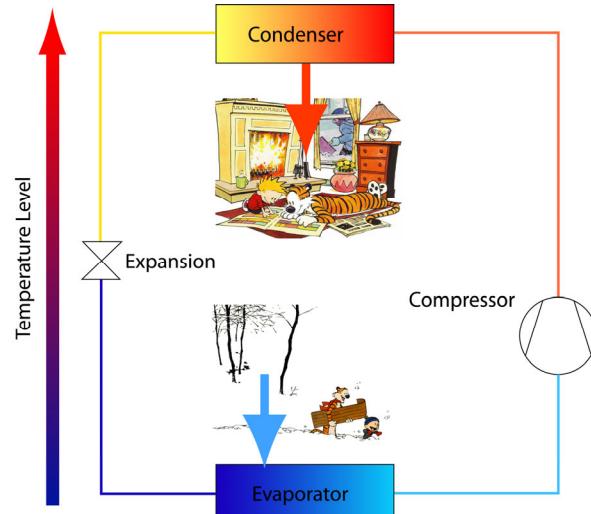


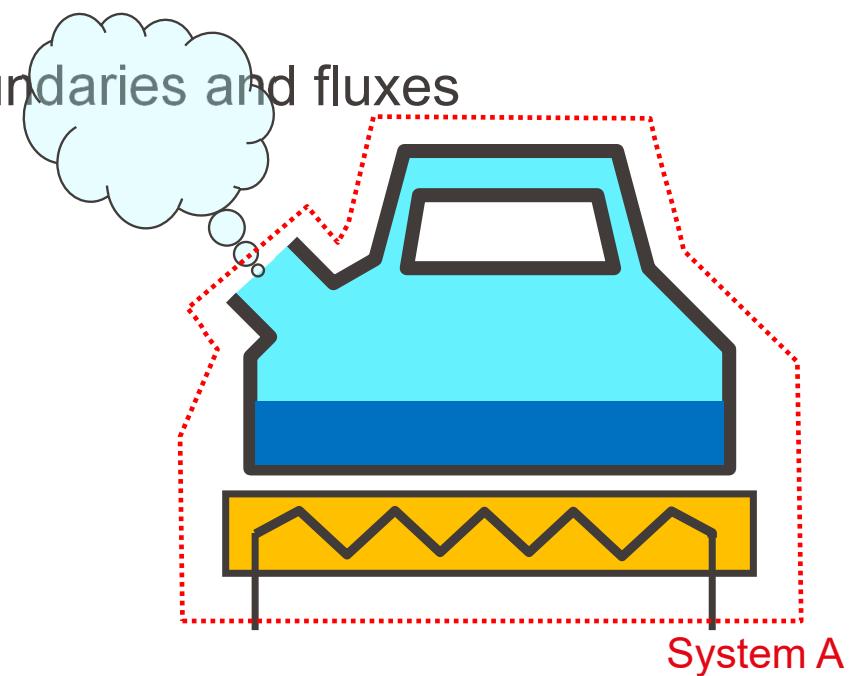
Fig 13.18 Favrat

Systematic Approach in Thermodynamics

1. What is known? → Sketch with known properties

2. What is problem? → Define objectives of analysis

3. Define the system → Identify system boundaries and fluxes



Systematic Approach in Thermodynamics (cont.)

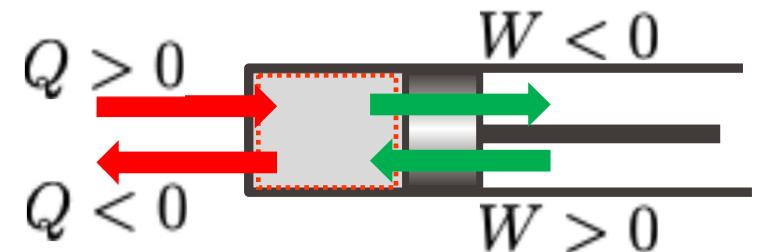
1. What is known? → Sketch with known properties
2. What is problem? → Define objectives of analysis
3. Define the system → Identify system boundaries and fluxes
4. Define assumptions → Identify suitable simplifying assumptions
5. Thermodynamic analysis → Apply physical laws
6. Discussion → Critical analysis of results & assumptions

Heat Pump Systems

Thermodynamics Crash Course
First Law for Closed Systems

Introduction

- Energy is fundamental quantity in thermodynamics
- Energy can only be stored, transformed, or transferred
- Energy cannot be destroyed or produced
- In closed system, energy can only be transferred through work and heat
- Reference is system
 - If work/heat is provided from surrounding to system $\rightarrow W > 0, Q > 0$
 - If system provides work/heat to surrounding $\rightarrow W < 0, Q < 0$

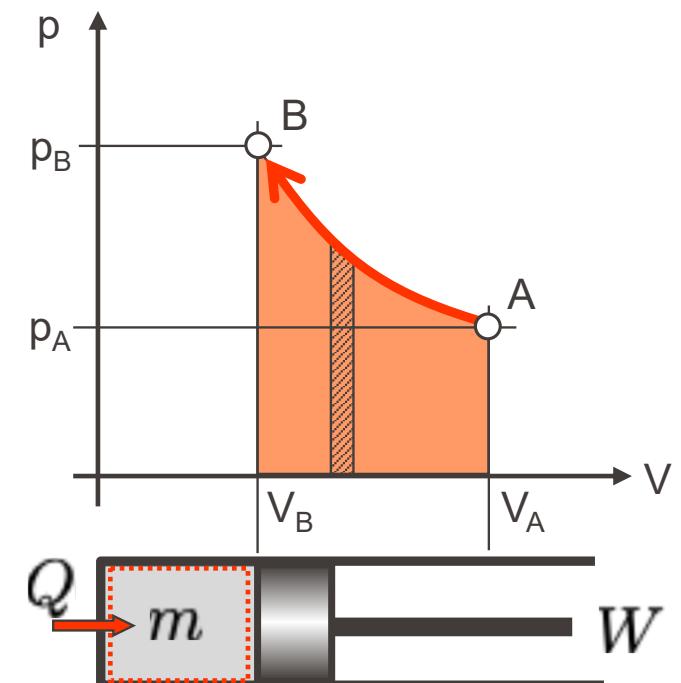


Work

- Consider closed system where piston moves from A to B

$$W = \int_A^B \delta W = - \int_{x_A}^{x_B} pAdx = - \int_{V_A}^{V_B} pdV$$

- Work depends on evolution of p vs. V
→ work depends on process details
- Surface under transformation line in pV-diagram represents work
- Work is no thermodynamic state property

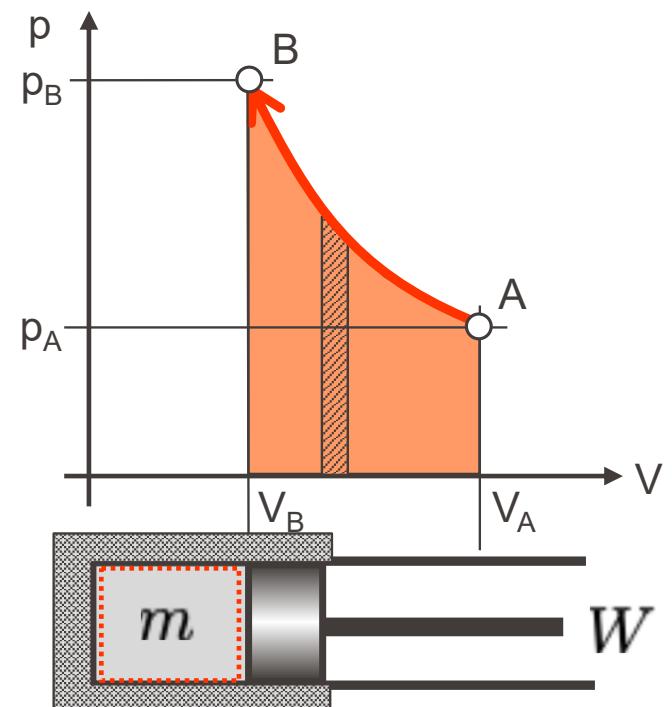


Work

- Consider closed and adiabatic system where piston moves from A to B
- Only possible energy transfer is via work
- Work on adiabatic system between fixed states depends only on start and end states

$$W = E_B - E_A$$

- Total energy E is thermodynamic state property



Total Energy E

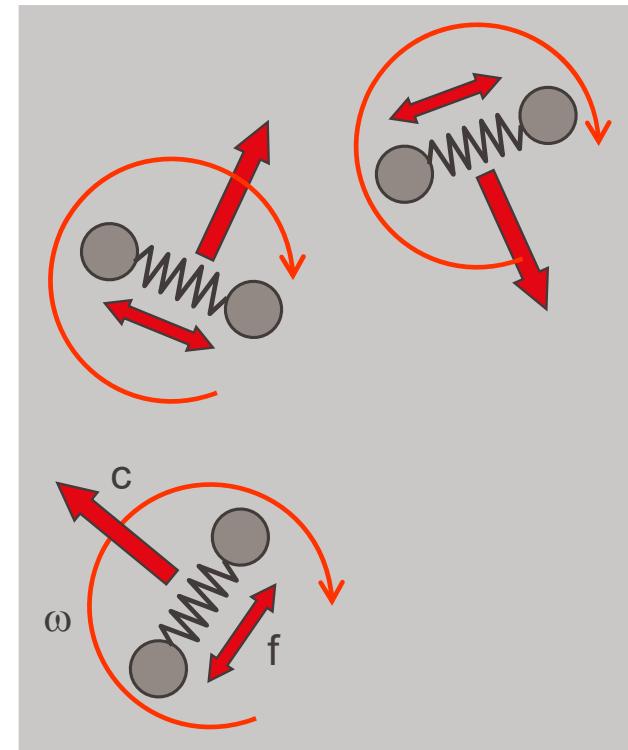
- Total energy E includes all possible forms of energy in system
 - Kinetic energy KE
 - Potential energy PE
 - Internal energy U
- All energy changes that are not kinetic or potential are summarized as internal energy

$$E_B - E_A = (KE_B - KE_A) + (PE_B - PE_A) + (U_B - U_A)$$

- Adiabatic work is transformed into kinetic, potential and internal energy

Internal Energy U

- Internal energy of fluid corresponds to kinetic energy related to molecular motion, chemical bonds, intramolecular forces
- Temperature plays important role
- No motion at absolute zero



Non-Adiabatic Transformation

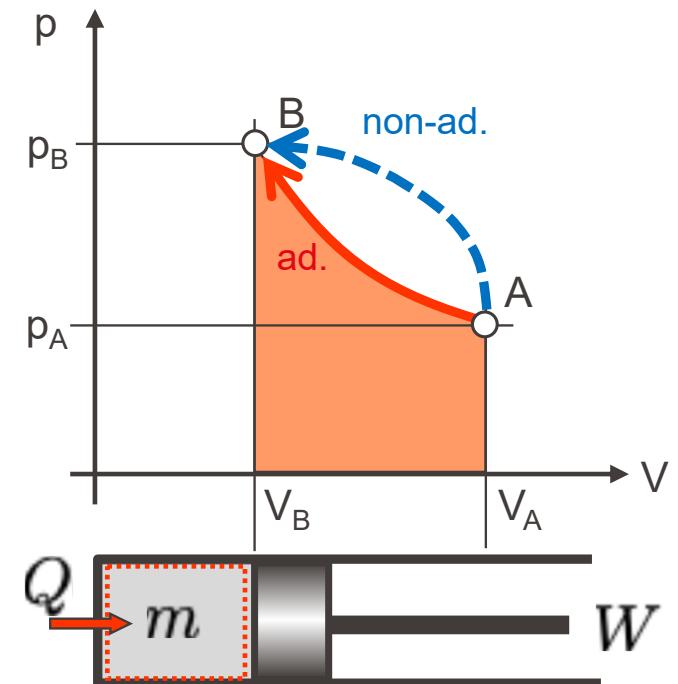
- Consider closed system where piston moves from A to B

$$W_{ad.} = E_B - E_A$$

$$W_{non-ad.} \neq E_B - E_A$$

- In non-adiabatic system, change of E cannot be solely explained by work
- Total energy change result of work and heat

$$Q + W_{non-ad.} = E_B - E_A$$



First Law in Closed System

- In general process from A to B, energy is transferred to achieve $E_B - E_A$
- Energy transfer occurs across system boundary through
 - Work W
 - Heat Q
- First law becomes

$$\underbrace{E_B - E_A}_{\text{State}} = \underbrace{W_{AB} + Q_{AB}}_{\text{Process}}$$

- Change of total energy in system corresponds to net transfer of work and heat across system boundary

Equivalent Formulations of First Law in Closed System

- Balance between end and start states

$$E_B - E_A = W_{AB} + Q_{AB}$$

- Differential balance

$$dE = \delta W + \delta Q$$

- Power balance

$$\frac{dE}{dt} = \dot{W} + \dot{Q}$$

where $E = KE + PE + U = m \left(\frac{w^2}{2} + gz + u \right)$

Heat

- Like work, heat is not a thermodynamic state, depends on process
- Heat flux

$$\dot{q} = \frac{\dot{Q}}{A} \quad \dot{Q} = \int \dot{q} dA$$

- Heat transfer mechanisms: conduction and radiation
- Energy transfer from a body to a fluid referred to as convection, i.e. combined effects of conduction and bulk motion of fluid

Conduction

- Heat transfer mechanism in solids and liquids at rest
- Heat transferred by activity at molecular scale
- Governed by Fourier's law

$$\dot{q} = \frac{\dot{Q}}{A} = -\lambda \frac{dT}{dx}$$

- Typical values for conductivity

	λ [W/mK]
Gases	0.01 – 0.2
Liquids	0.1 - 1
Solids	1 - 450

Radiation

- Radiation

- No medium required (\rightarrow vacuum, space)
- Governed by Boltzmann-equation

$$\dot{q} = \epsilon \sigma T^4$$

$$\sigma = 5.67 \cdot 10^{-8} \left[\frac{W}{m^2 K^4} \right]$$

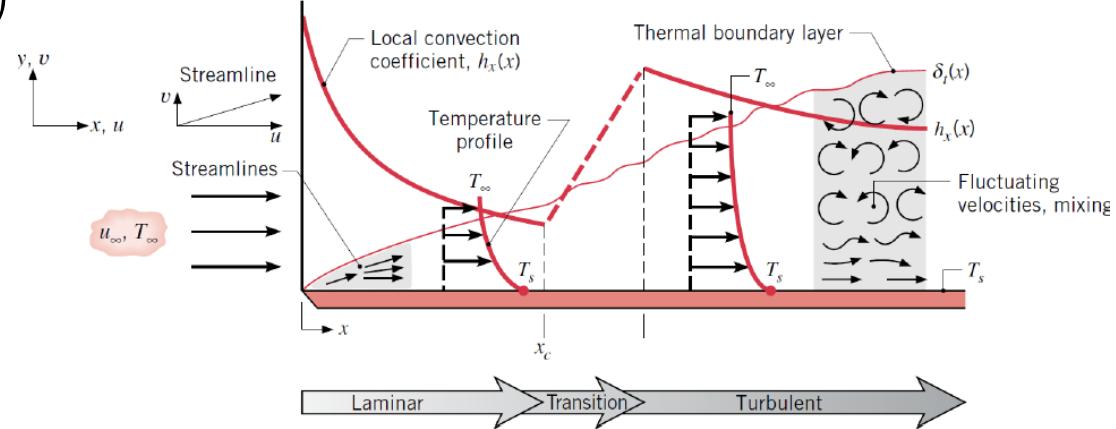
- Emissivity $\epsilon \in [0, 1]$
- Emissivity depends on surface shape, material properties, surface finish, orientation

Convection

- Heat transfer activated by fluid motion at interface with solid
- Coupling between fluid motion and fluid conduction
- Challenging to calculate → correlations for specific configurations
- Governed by Newton's law

$$\dot{q} = \frac{\dot{Q}}{A} = \alpha (T_{Wall} - T_{Fluid})$$

- Laminar regime → low flux
- Turbulence enhances energy exchange between layers



Convection Correlations

- Correlations for forced convection

$$Nu = \frac{\alpha L}{\lambda} = f_{forced}(Re, Pr)$$

$$Re = \frac{vL}{\nu}$$

- Correlations for natural convection

$$Nu = \frac{\alpha L}{\lambda} = f_{nat}(Gr, Pr)$$

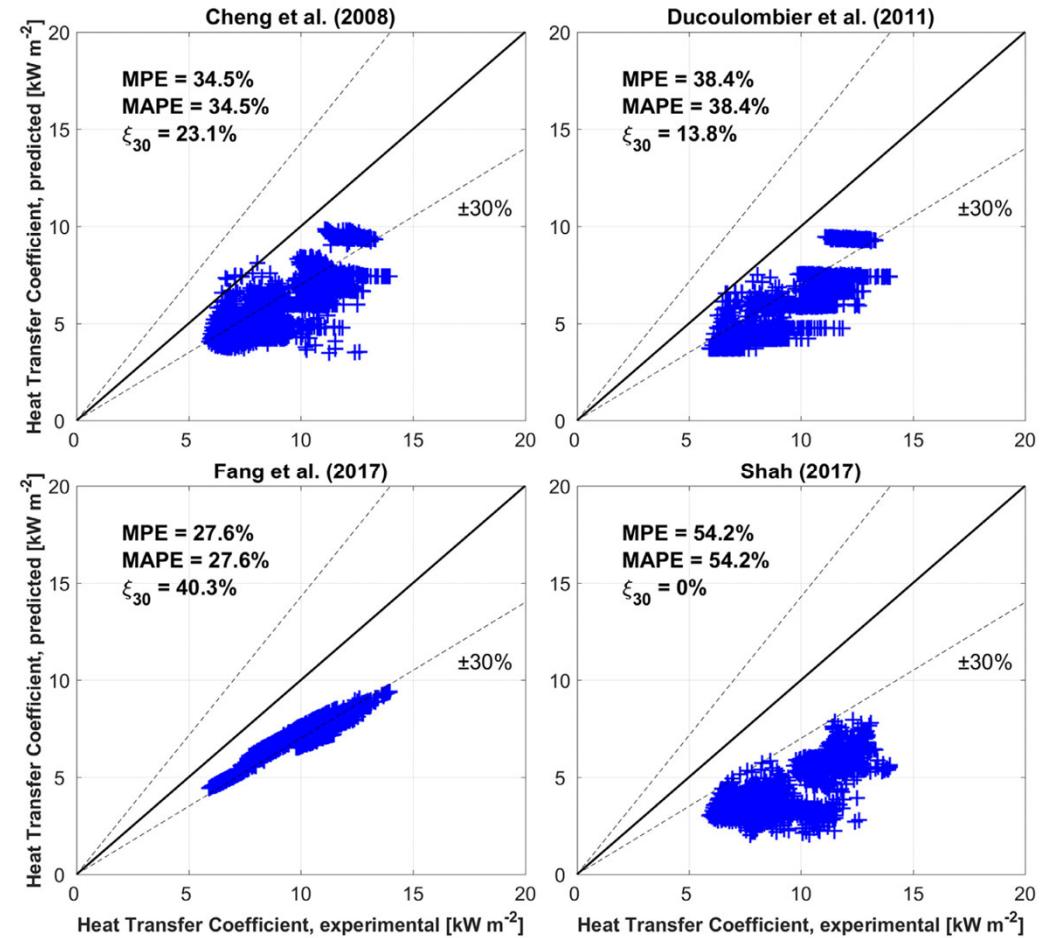
$$Gr = \frac{L^3 g \beta \Delta T}{\nu^2}$$

- Fluid properties condensed into *Pr*-number

$$Pr = \frac{\nu \rho c_p}{\lambda}$$

Performance of Typical Correlations

- Available correlations usually not very accurate
- Heat transfer coefficient highly dependent on fluid, mass-flux, geometry



Heat Pump Systems

Thermodynamics Crash Course
Cycles

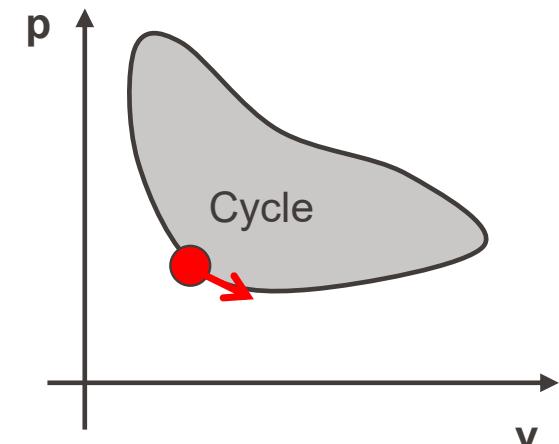
Cycle Characteristics

- Cyclic process returns to starting point periodically
- Energy balance of cycle

$$\Delta E_C = W_C + Q_C$$



Net work and net heat during
one period of the cycle process



Cycle Characteristics

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- Energy balance of cycle

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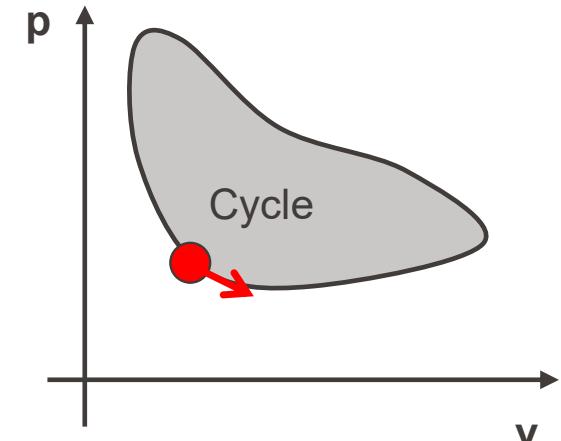
Net work and net heat during
one period of the cycle process

- End corresponds to starting point

$$\Delta E_C = 0$$

$$W_C = -Q_C = - \oint pdV$$

- In cycle, net heat corresponds to net work \rightarrow true for all cycles



Power Cycles

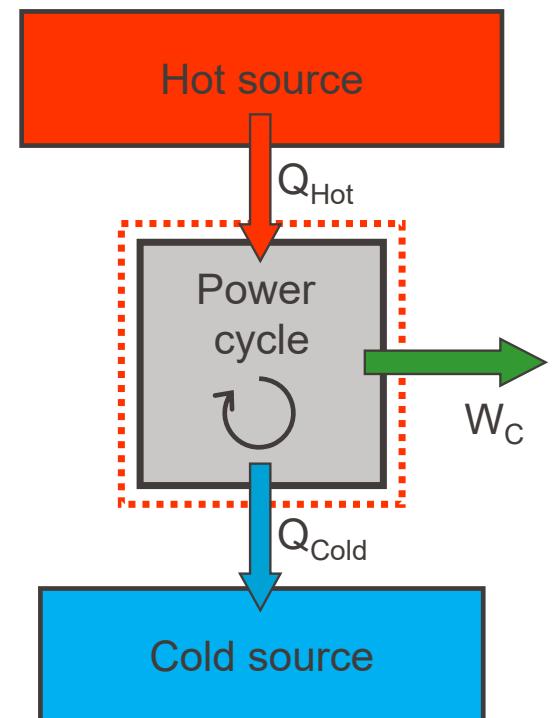
- Power cycle absorbs heat from hot source, transforms part of it into work, and delivers remaining heat to cold source
- Power cycles operate in clockwise direction
- Energy balance

$$W_C = Q_{Hot} + Q_{Cold}$$

$$W_C = Q_{Hot}^+ - Q_{Cold}^-$$



Notation to get positive values
for heat and work

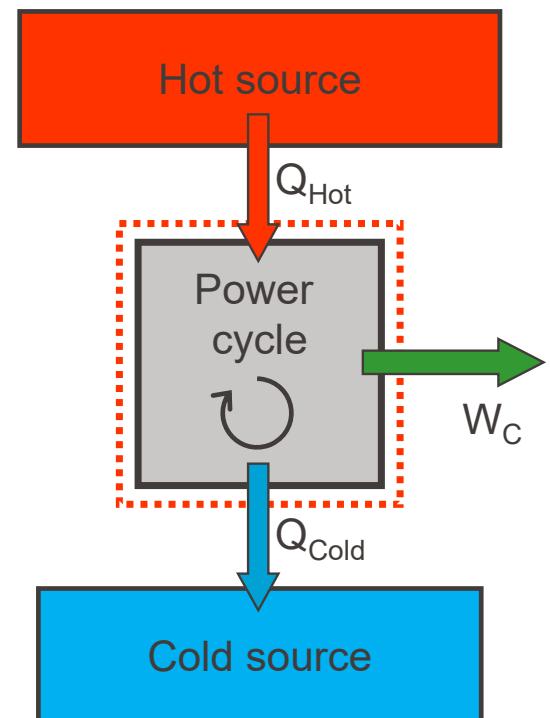


Power Cycle Efficiency

- Efficiency is ratio between yield and investment

$$\eta_{th} = \frac{W_C}{Q_{Hot}^+} = 1 - \frac{Q_{Cold}^-}{Q_{Hot}^+}$$

- Due 1st law and $Q_{Cold}^- > 0$ thermal efficiency $\eta_{th} < 1$



Typical thermal efficiencies

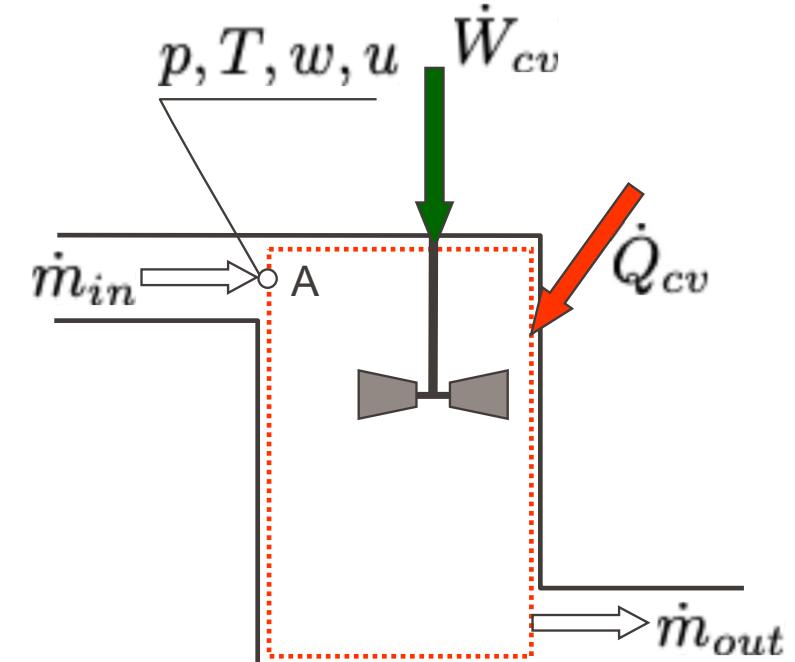
▪ Steam locomotives	0.12
▪ Nuclear power station	0.3 – 0.34
▪ Coal plant	0.35 – 0.48
▪ Gas turbine engines	0.3 – 0.42
▪ Combined cycles	0.62
▪ Internal combustion engines	0.35
▪ Large diesel engines	0.58

Heat Pump Systems

Thermodynamics Crash Course
First Law for Open Systems

Open Systems

- Characterized by mass fluxes across system boundary
 - Energy balance
 - Mass balance
- Possible energy transfer
 - Work
 - Heat
 - Convective contribution through incoming and outgoing mass-flows



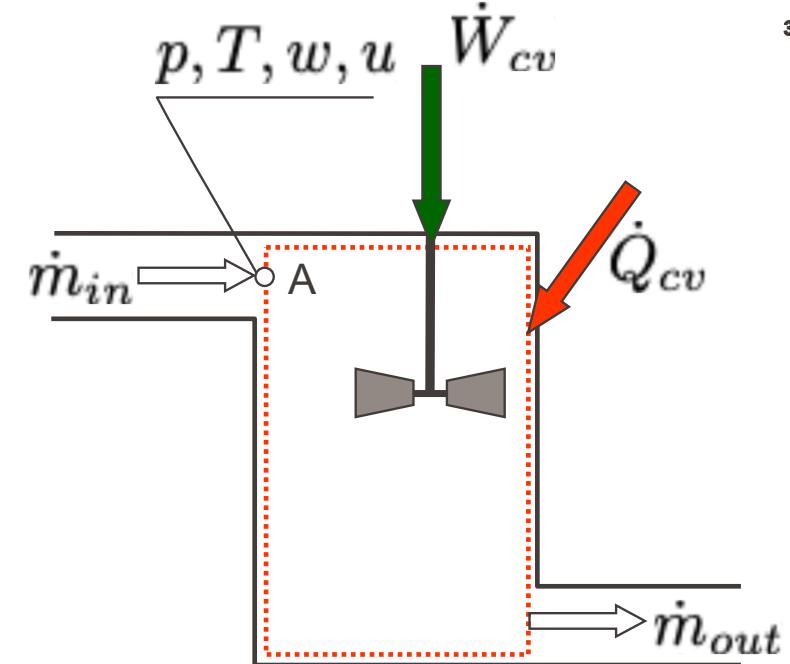
Open Systems: Mass Balance

- Mass balance

$$\frac{dm_{cv}}{dt} = \sum \dot{m}_{in} - \sum \dot{m}_{out}$$

- Mass in control volume

$$m_{cv} = \int_V \rho dV \text{ where } \rho(\vec{x}, t)$$



Open Systems: Mass Balance

- Mass balance

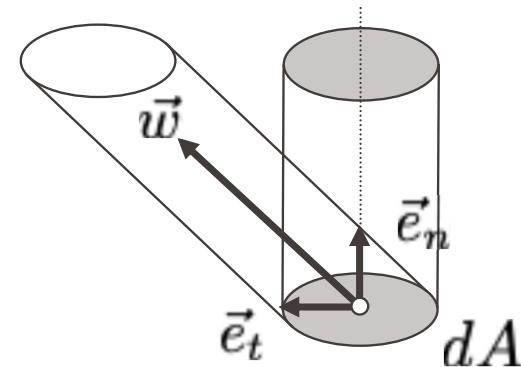
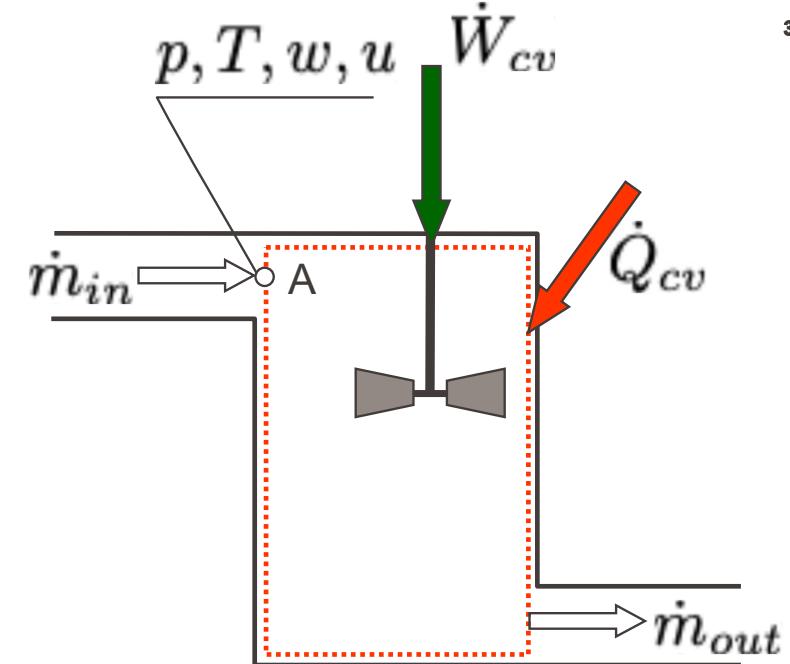
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- Mass in control volume

$$m_{cv} = \int_V \rho dV \text{ where } \rho(\vec{x}, t)$$

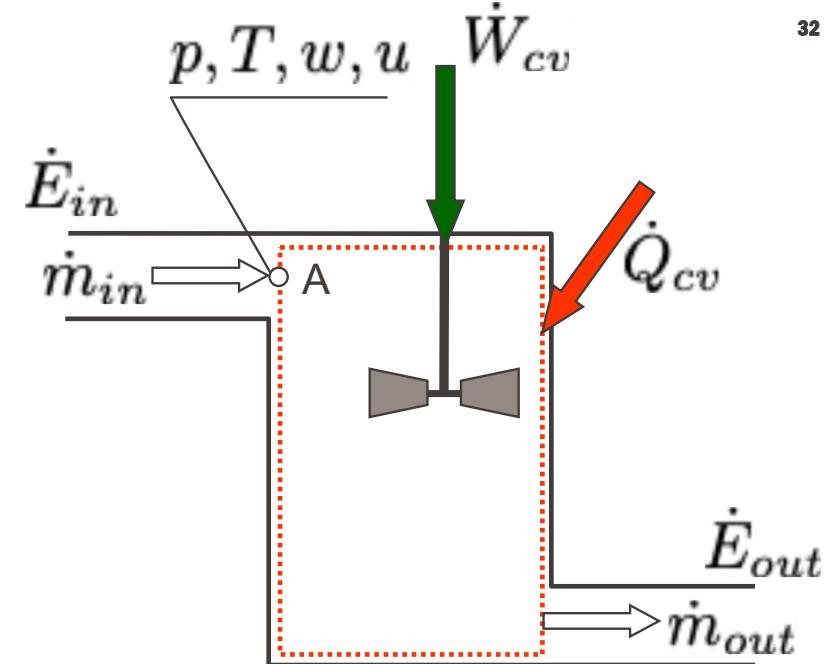
- Mass flow

$$\dot{m} = \int_A \rho \vec{w} \vec{e}_n dA$$



Open Systems: Energy Balance

- Incoming and outgoing mass-flows add and retrieve total energy from system
- Convective flow across system boundary composed of internal, kinetic, and potential energy



$$\frac{dE_{cv}}{dt} = \dot{W} + \dot{Q} + \underbrace{\dot{m}_{in} \left(u_{in} + \frac{w_{in}^2}{2} + gz_{in} \right)}_{\dot{E}_{in}} - \underbrace{\dot{m}_{out} \left(u_{out} + \frac{w_{out}^2}{2} + gz_{out} \right)}_{\dot{E}_{out}}$$

Open Systems: Transfer Power

- Transfer to mass across system boundary requires work that is transferred to and from system

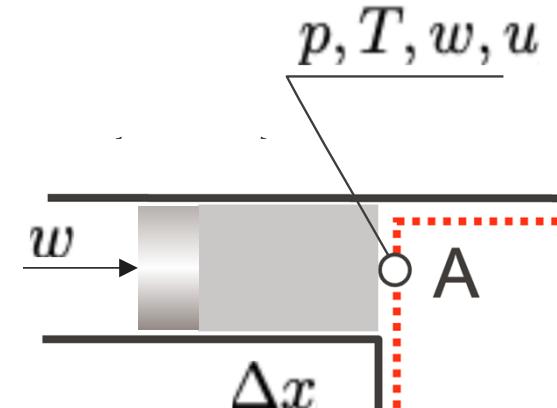
- Transfer power: $\dot{W} = pAw$

- Expression for total work on system:

$$\dot{W} = \dot{W}_{cv} + p_{in}A_{in}w_{in} - p_{out}A_{out}w_{out}$$

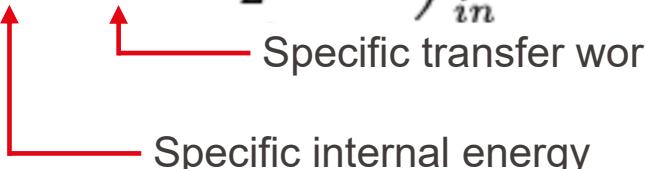
$$\dot{W} = \dot{W}_{cv} + \dot{m}_{in}p_{in}v_{in} - \dot{m}_{out}p_{out}v_{out}$$

where \dot{W}_{cv} corresponds to work other than transfer work delivered to system



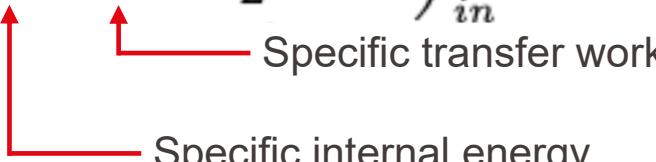
Open Systems: Energy Balance

- Combining energy balance and expression for total work on system

$$\frac{dE_{cv}}{dt} = \dot{W}_{cv} + \dot{Q}_{cv} + \dot{m}_{in} \left(u + pv + \frac{w^2}{2} + gz \right)_{in} - \dot{m}_{out} \left(u + pv + \frac{w^2}{2} + gz \right)_{out}$$


Open Systems: Energy Balance

- Combining energy balance and expression for total work on system

$$\frac{dE_{cv}}{dt} = \dot{W}_{cv} + \dot{Q}_{cv} + \dot{m}_{in} \left(u + pv + \frac{w^2}{2} + gz \right)_{in} - \dot{m}_{out} \left(u + pv + \frac{w^2}{2} + gz \right)_{out}$$


Specific transfer work

Specific internal energy

- With definition of enthalpy h as a new state variable: $h = u + pv$

$$\frac{dE_{cv}}{dt} = \dot{W}_{cv} + \dot{Q}_{cv} + \dot{m}_{in} \left(h + \frac{w^2}{2} + gz \right)_{in} - \dot{m}_{out} \left(h + \frac{w^2}{2} + gz \right)_{out}$$

Open System Balances

- Energy balance:

$$\frac{dE_{cv}}{dt} = \dot{W}_{cv} + \dot{Q}_{cv} + \sum \dot{m}_{in} \left(h + \frac{w^2}{2} + gz \right)_{in} - \sum \dot{m}_{out} \left(h + \frac{w^2}{2} + gz \right)_{out}$$

Diagram illustrating the components of the energy balance equation:

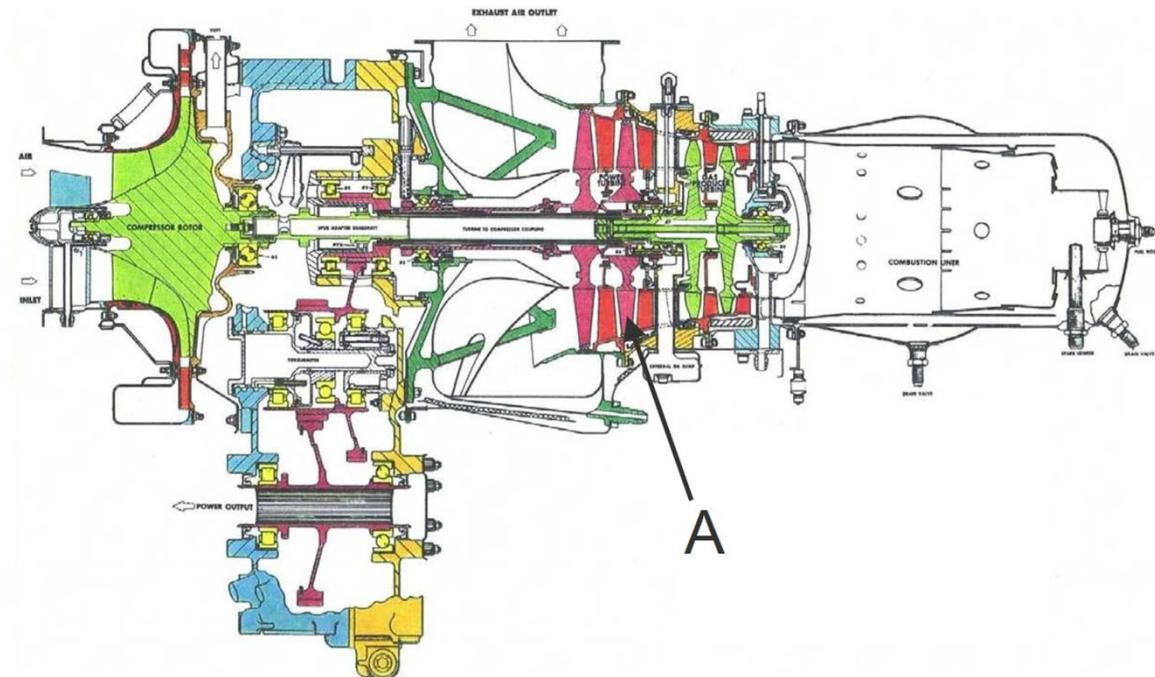
- Net work: \dot{W}_{cv}
- Net heat: \dot{Q}_{cv}
- Net convected power: $\sum \dot{m}_{in} \left(h + \frac{w^2}{2} + gz \right)_{in} - \sum \dot{m}_{out} \left(h + \frac{w^2}{2} + gz \right)_{out}$
- Specific potential energy: gz
- Specific kinetic energy: $\frac{w^2}{2}$
- Specific internal + transfer energy: h

- Mass balance:

$$\frac{dm_{cv}}{dt} = \sum \dot{m}_{in} - \sum \dot{m}_{out}$$

Example: Helicopter Engine

- Mass balance
- Energy balance



Heat Pump Systems

Thermodynamics Crash Course
Second Law

Natural Systems: Observations

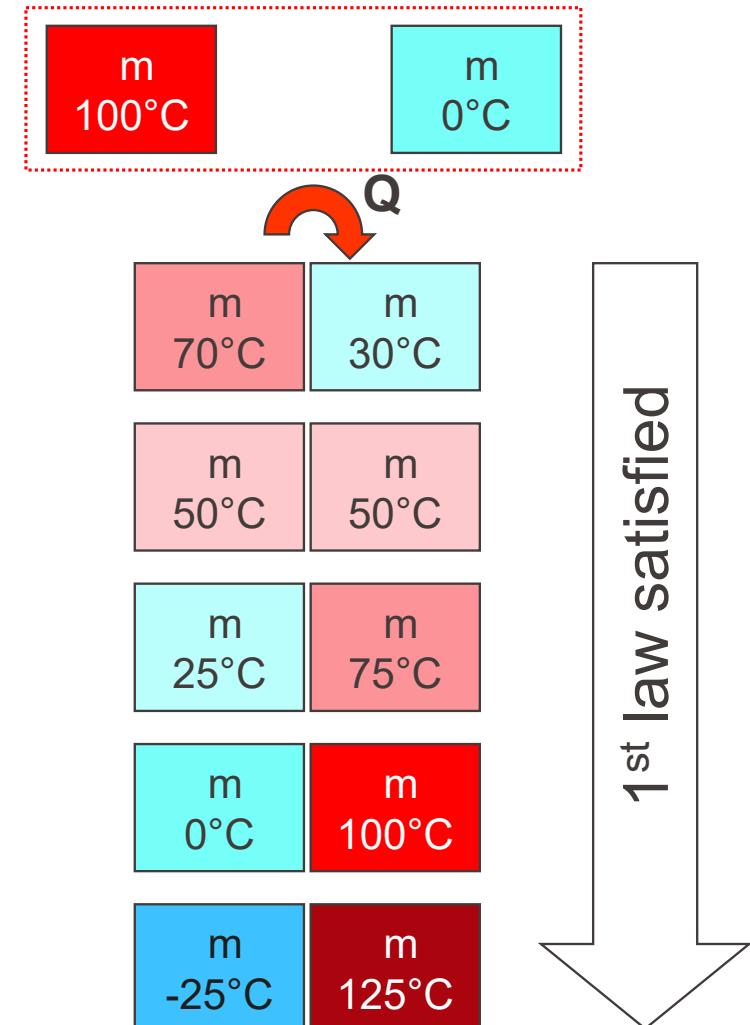
- Systems left on their own are subject to natural balancing processes until state of equilibrium is reached
- Spontaneous balancing processes only take place if there is a natural tension / imbalance / potential
- Examples of natural balancing processes
 - Body at higher temperature naturally cools down until reaching ambient temperature
 - Pressurized air bottle naturally leaks until ambient pressure is reached
 - Suspended mass naturally falls to floor

Natural Systems: Observations

- Natural balancing processes only in one direction
- Reversal is possible only by investing energy from outside system
- Problem: 1st law does not prevent reverse processes
- Description of observations requires additional law → 2nd law

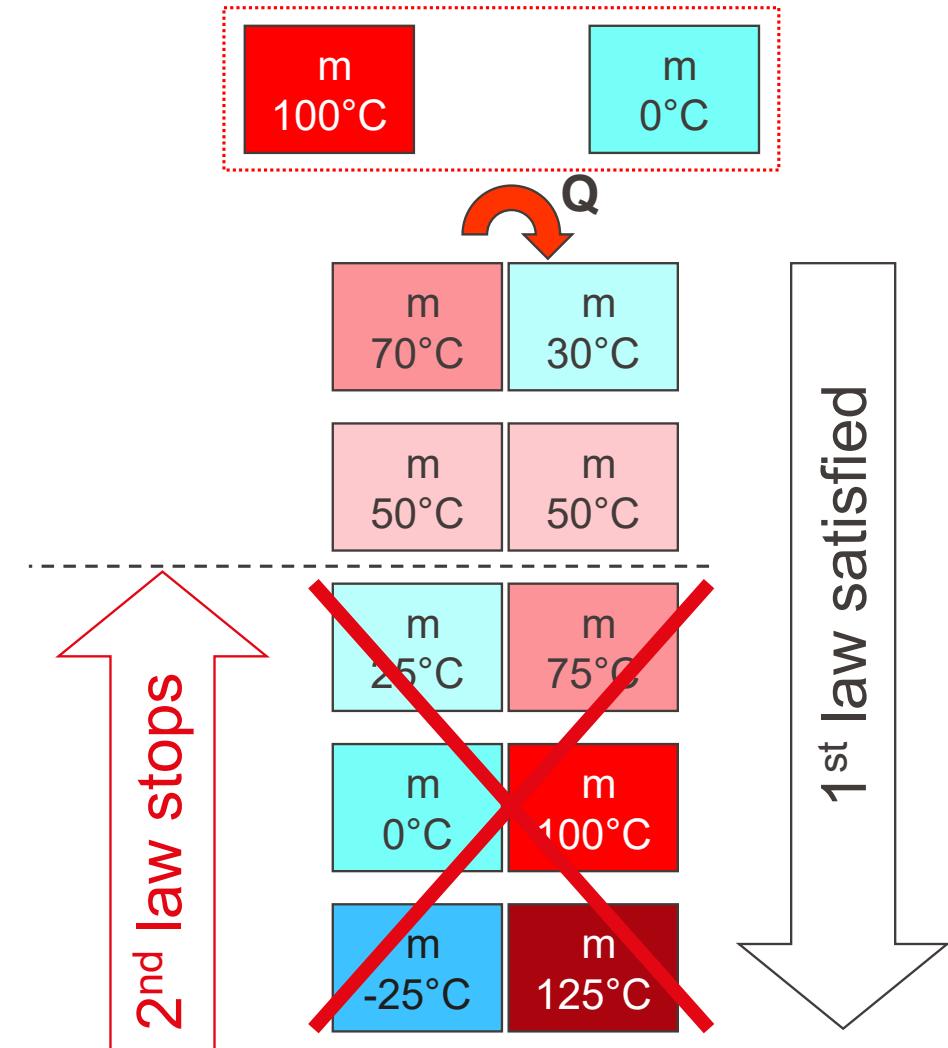
Experiment

- Closed system contains two equivalent masses, one at 0°C , one at 100°C
- 1st law violates observed natural processes



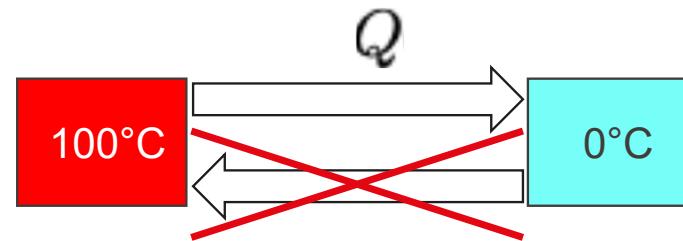
Experiment

- Closed system contains two equivalent masses, one at 0°C , one at 100°C
- 1st law violates observed natural processes
- 2nd law prevents violation



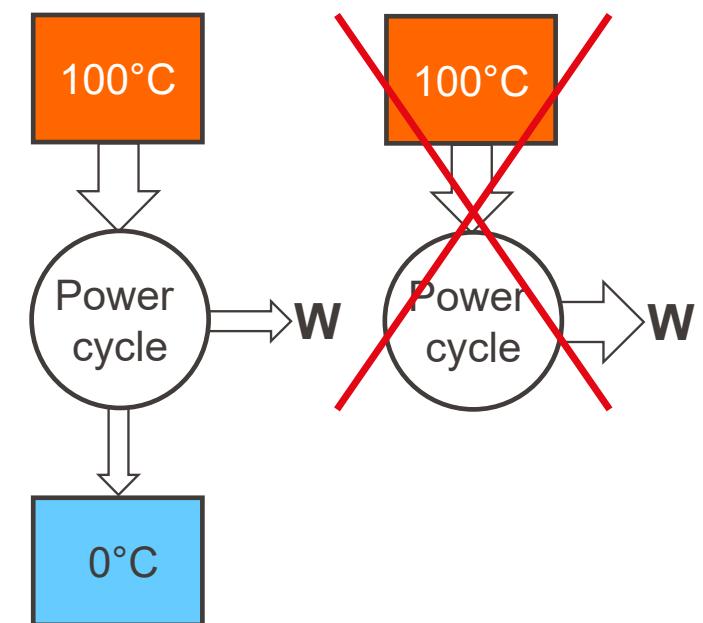
Formulations of 2nd Law

- By Clausius (1854): There is no change of state whose only result is the transfer of heat from a body at a lower temperature to a body at a higher temperature
 - Heat does not flow naturally from a low-temperature reservoir to a higher one



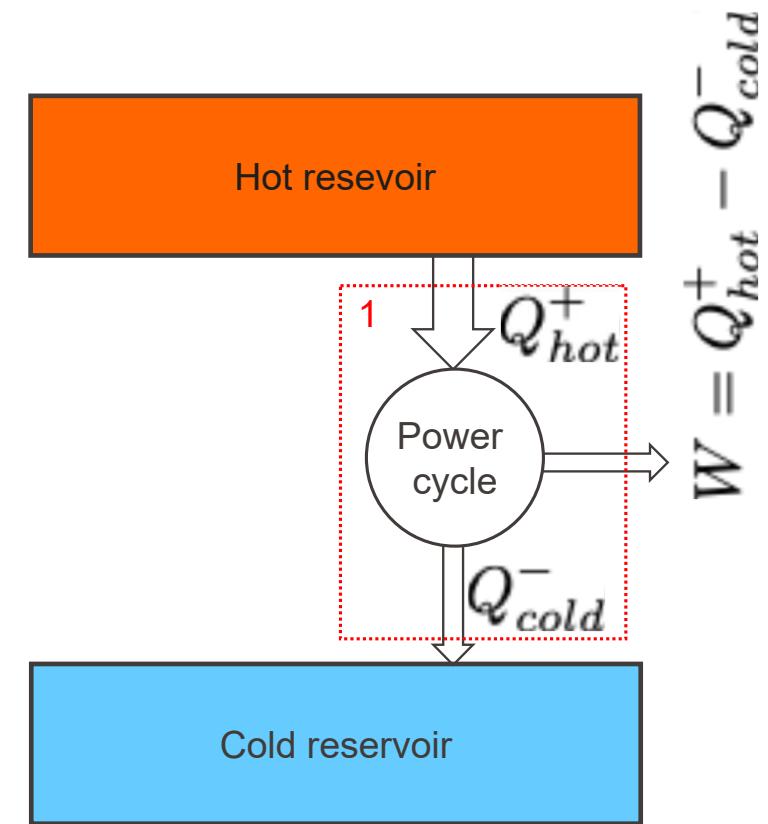
Formulations of 2nd Law

- Kelvin-Planck (1848/1926): It is impossible to construct a device which, operating in a cycle, will produce no other effect than the extraction of heat from a reservoir and the performance of equivalent amount of work
 - Power cycles requires hot and cold source
 - Thermal efficiency < 1



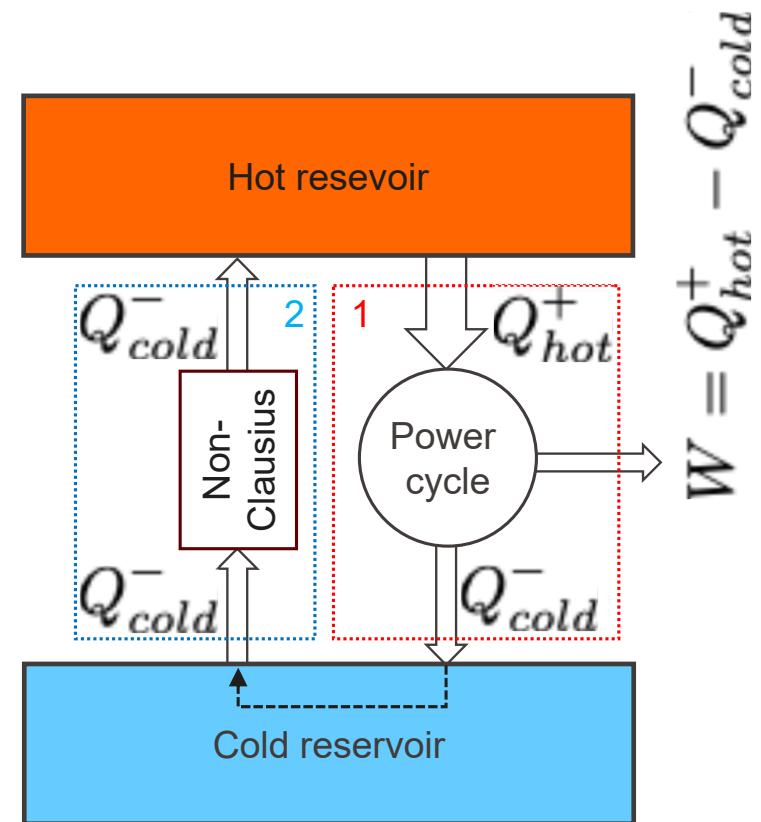
Equivalence of 2nd Law Formulations

- System 1 composed of power cycle according to Kelvin-Planck



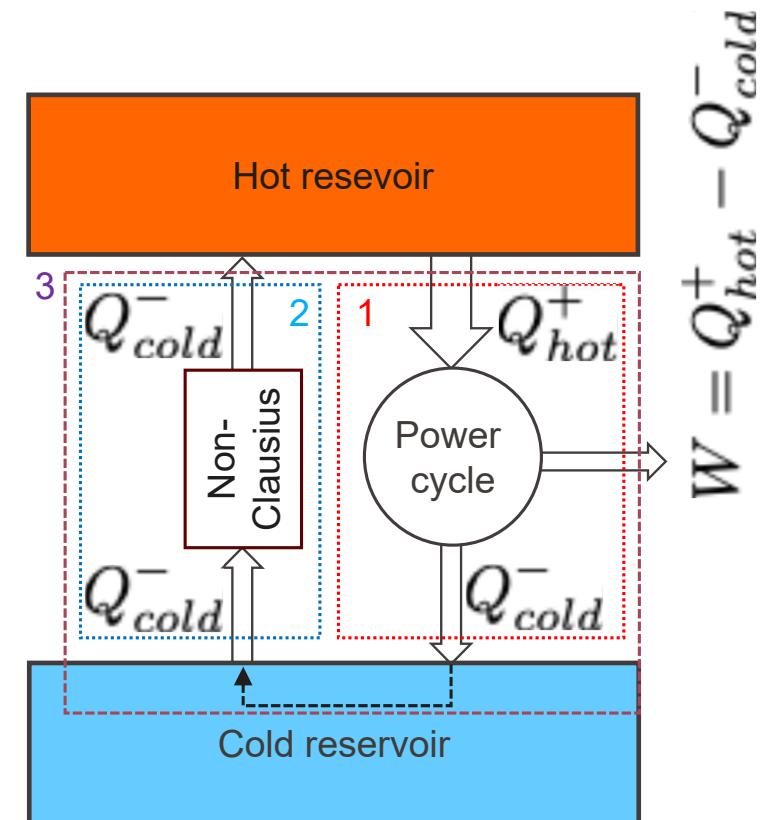
Equivalence of 2nd Law Formulations

- System 1 composed of power cycle according to Kelvin-Planck
- System 2 transferring heat opposed to Clausius
- Heat rejected by power cycle flows back to hot reservoir



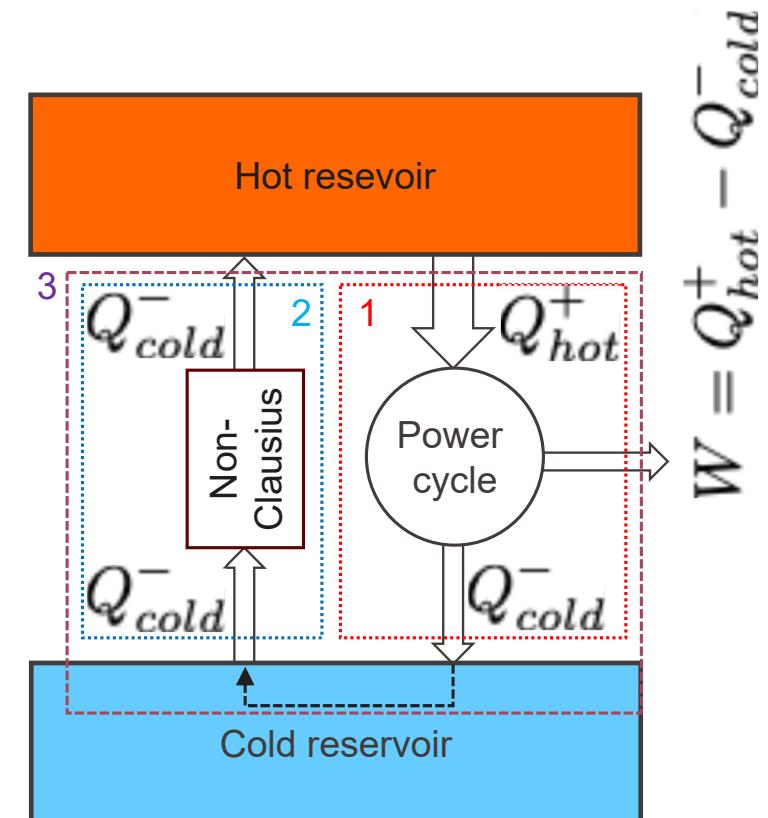
Equivalence of 2nd Law Formulations

- System 1 composed of power cycle according to Kelvin-Planck
- System 2 transferring heat opposed to Clausius
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- Combined system 3 draws net heat $Q_{hot}^+ - Q_{cold}^-$ and transform it into W without rejecting heat to cold reservoir
→ violates Kelvin-Planck formulation



Equivalence of 2nd Law Formulations

- System 1 composed of power cycle according to Kelvin-Planck
- System 2 transferring heat opposed to Clausius
- Heat rejected by power cycle flows back to hot reservoir
- Combined system 3 draws net heat $Q_{hot}^+ - Q_{cold}^-$ and transform it into W without rejecting heat to cold reservoir
→ violates Kelvin-Planck formulation
- Violation of Kelvin-Planck requires violation of Clausius → equivalence



Concept of Reversibility

- Knowing 1st and 2nd law that indicate what is possible, we want to know what a perfect machine looks like
- A perfect machine is reversible
- A process is called reversible if a system can be returned to its initial state without causing any changes in the environment
- Processes that take place in one direction only, are irreversible
 - Cooling of a cup of tea, emptying of a pressurized bottle, dissipation, falling mass, plastic deformation, spontaneous chemical reaction, mixing, ...

Concept of Reversibility

- Irreversibility is a dissipation of energy, a loss of potential work
→ engineers want to avoid them
- Reversible processes are hypothetical and represent processes without any losses
- Reversible processes are not possible practically, but they represent the limits of what is possible while satisfying 1st and 2nd law
- Reversible processes are useful as references to measure the thermodynamic performance of real processes

Carnot Principles

- The thermal efficiency of an irreversible power cycle is always lower than that of a reversible cycle between the same thermal reservoirs
- All reversible power cycles between the same thermal reservoirs have the same thermal efficiency
- The efficiency of a reversible machine is independent of the process, the components, and the working fluid

Thermal Efficiency

- Reversible power cycles between same thermal reservoirs have same thermal efficiency → efficiency only function of reservoir temperatures

$$\eta_{th-rev.} = g(T_{hot}, T_{cold})$$

- Definition of thermal efficiency $\eta_{th} = 1 - \frac{Q_{Cold}^-}{Q_{Hot}^+}$

- Consequently $\left[\frac{Q_{Cold}^-}{Q_{Hot}^+} \right]_{rev.} = f(T_{hot}, T_{cold})$

Thermal Efficiency

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- Definition of thermal efficiency $\eta_{th} = 1 - \frac{Q_{Cold}^-}{Q_{Hot}^+}$

- Consequently $\left[\frac{Q_{Cold}^-}{Q_{Hot}^+} \right]_{rev.} = f(T_{hot}, T_{cold}) = \frac{T_{cold}}{T_{hot}}$

 Kelvins approach: ratio of reversibly transferred heats equals ratio of reservoir temperatures

Thermal Efficiency

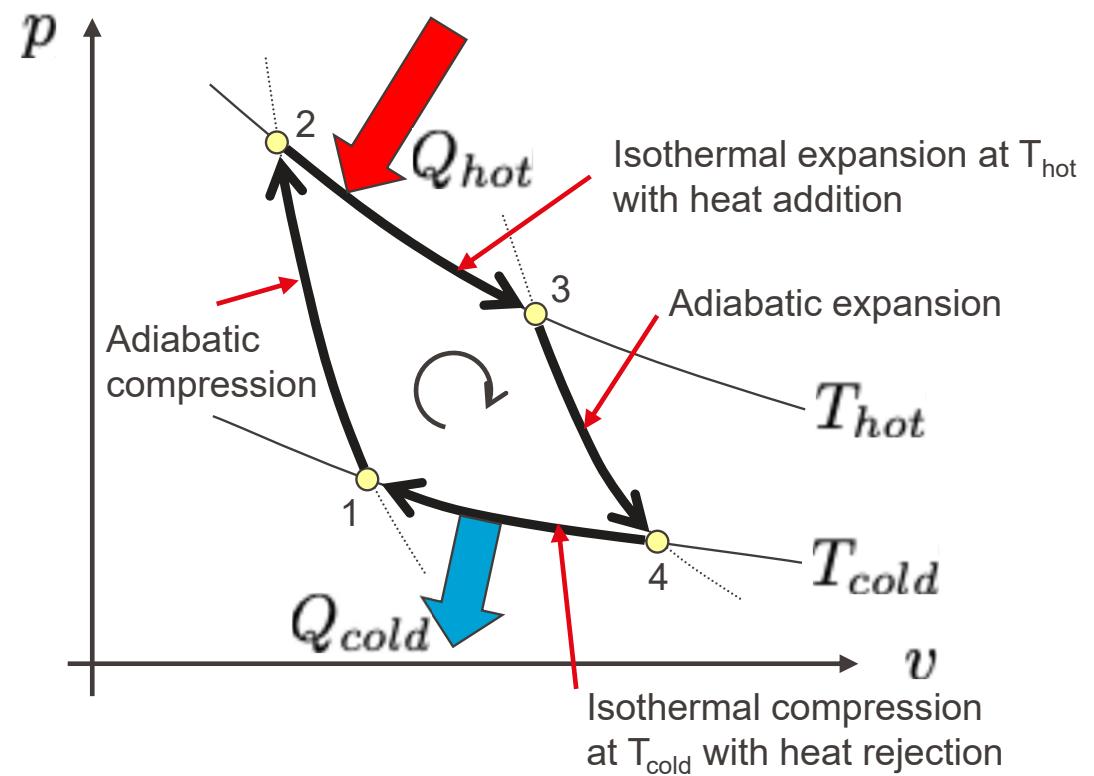
- Through Kelvins definition of the function, thermal efficiency of reversible cycle expressed as:

$$\eta_{th-rev.} = 1 - \frac{T_{cold}}{T_{hot}} = \eta_c$$

- Carnot cycle is one famous reversible power cycle
- Thermal efficiency of reversible cycle is called Carnot-efficiency
- Carnot efficiency is reference for assessing performance of power cycles

Carnot Cycle

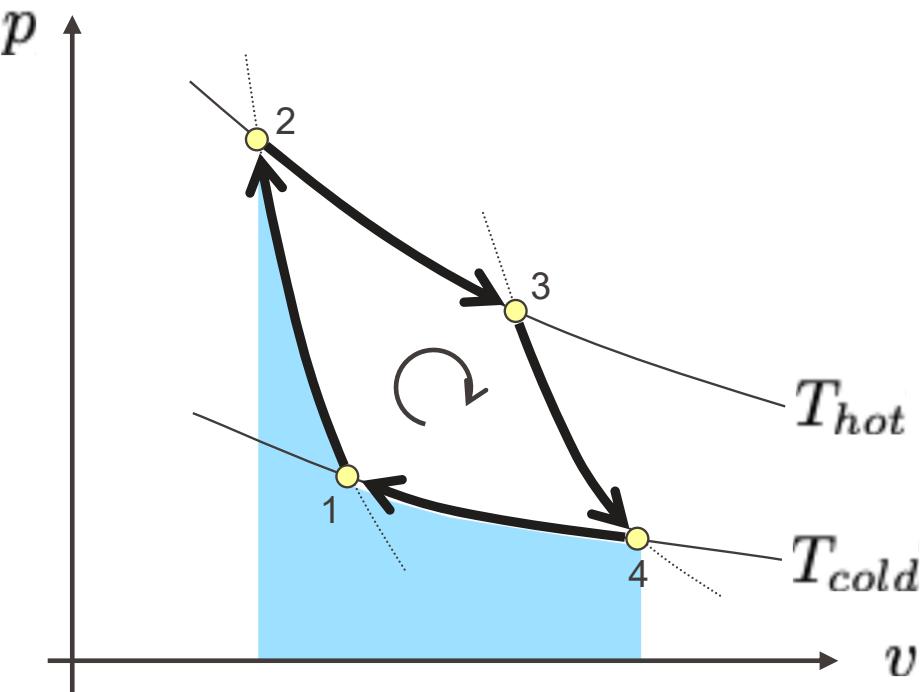
- Cycle composed of two reversible adiabatic and two reversible isothermal processes



Carnot Cycle

- Compression work

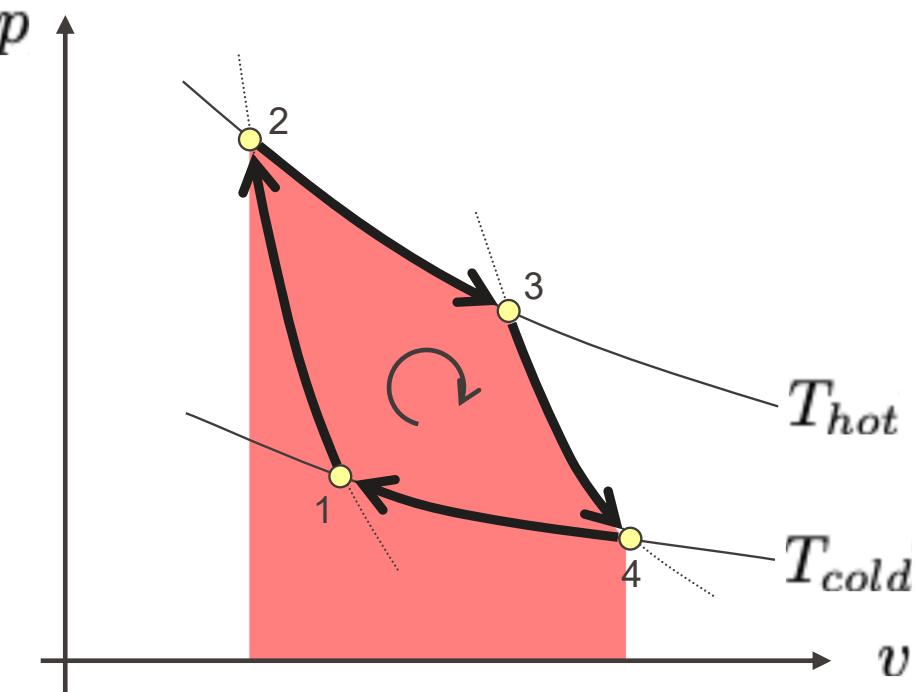
$$W_{42} = - \int_4^2 pdV > 0$$



Carnot Cycle

- Expansion work

$$W_{234} = - \int_2^4 p dV < 0$$



Carnot Cycle

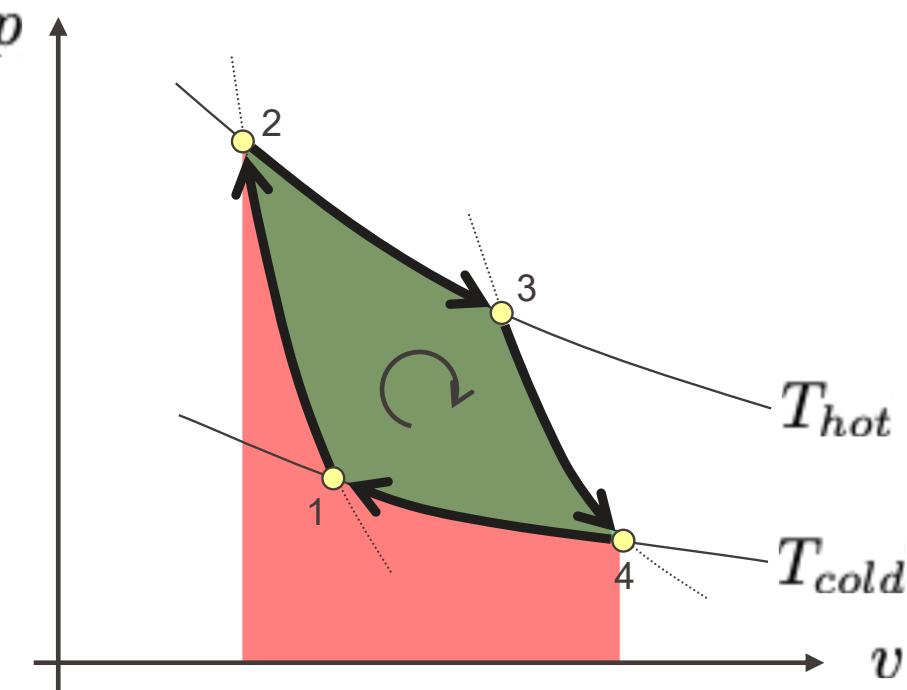
- Net work

$$W_{234} = - \int_2^4 p dV < 0$$

$$W_C = W_{234} - W_{42}$$

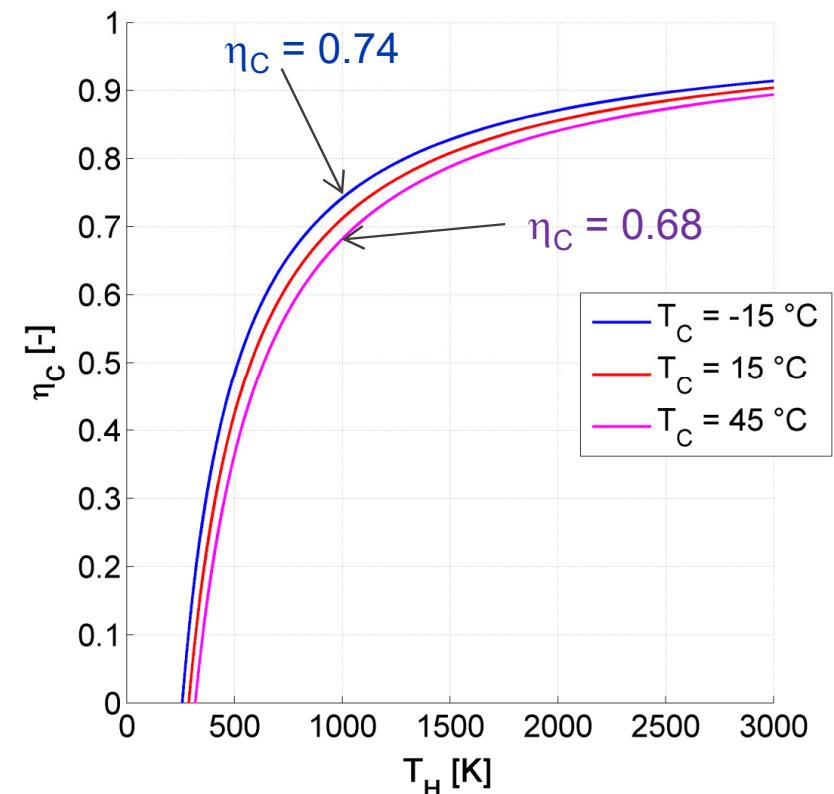
$$W_C = \eta_{th-rev.} Q_{hot}$$

$$W_C = \left[1 - \frac{T_{cold}}{T_{hot}} \right] Q_{hot}$$



Thermal Efficiency

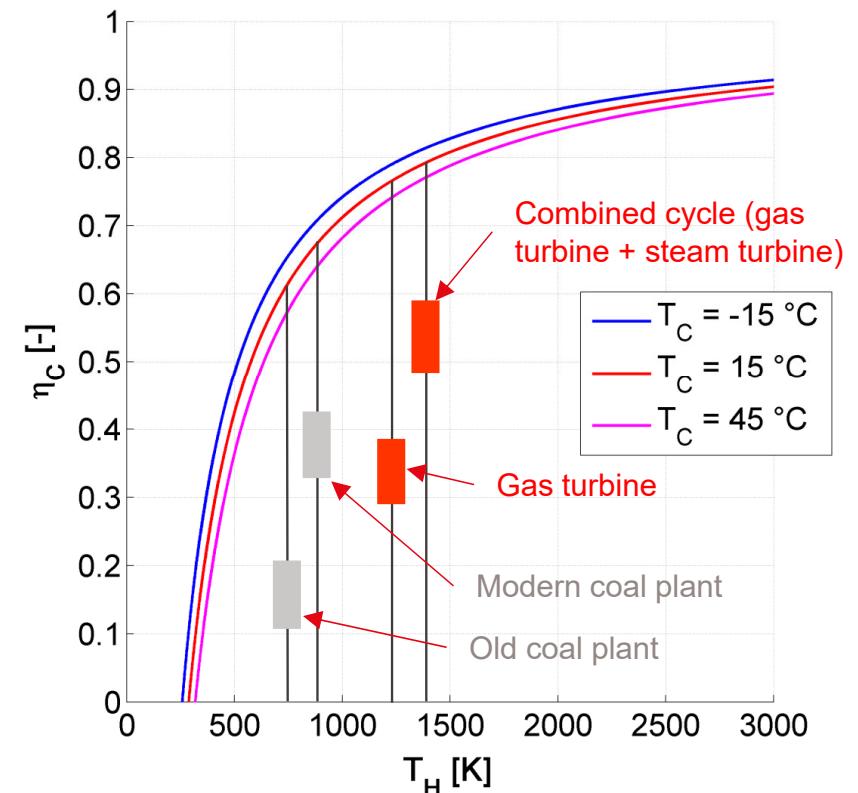
- Maximum thermal efficiency limited by Carnot efficiency
- Carnot efficiency rises with temperature difference between reservoirs
- Maximizing hot source temperature is key drive to improve thermal performance



$$\eta_c = 1 - \frac{T_{cold}}{T_{hot}}$$

Typical Thermal Efficiencies

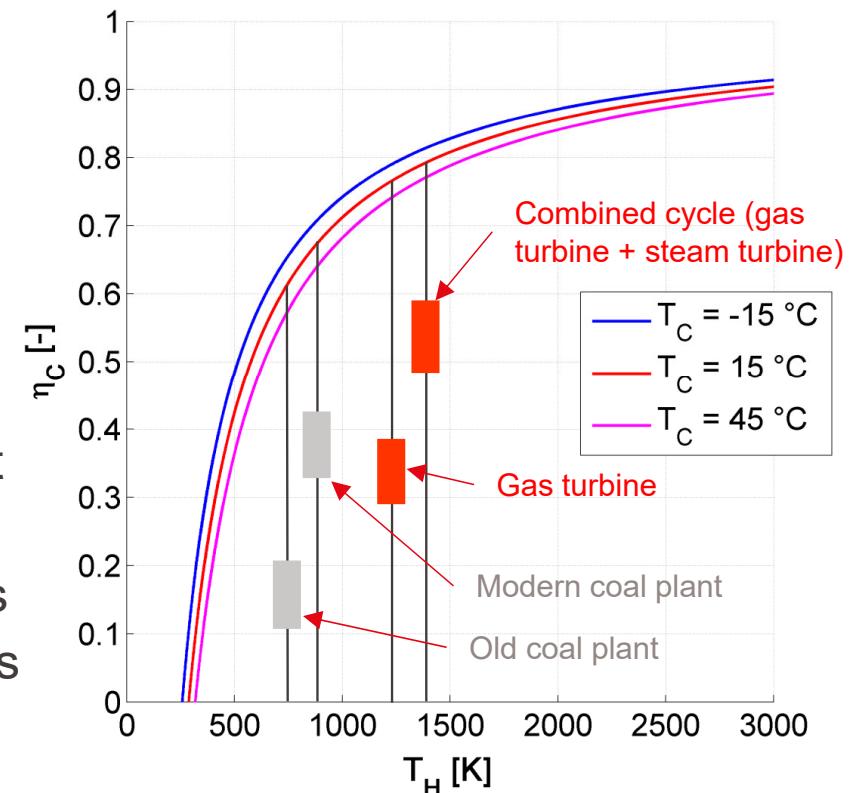
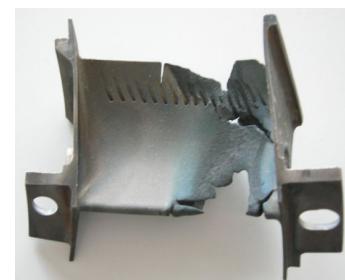
- Real life thermal efficiency always lower than Carnot efficiency



$$\eta_c = 1 - \frac{T_{cold}}{T_{hot}}$$

Typical Thermal Efficiencies

- Real life thermal efficiency always lower than Carnot efficiency
- Limitations stem from irreversibility and practical limitations
 - Temperature drop required to transfer heat
 - Friction, secondary flows
 - Pressure gradient needed to transfer mass
 - Limited temperature resistance of materials



$$\eta_c = 1 - \frac{T_{cold}}{T_{hot}}$$

Outlook for W3

- Thermodynamics crash course
 - Entropy, entropy balance in closed/open systems
 - Isentropic efficiencies
 - Concept of exergy
 - Exergy efficiency

- Theory questions
- Reversible, adiabatic expansion
- Polytropic expansion
- Maximum reversible work output of a finite source
- Filling empty gas tank
- Reversible cycle efficiencies