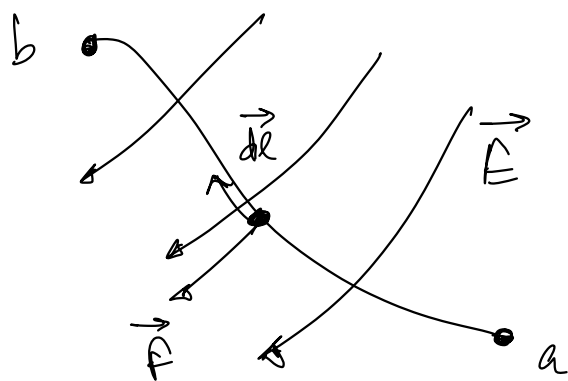


# Moving charge in an electric field



$$\vec{E} = -\nabla\phi \rightarrow \text{potential}$$

$$\vec{F} = z \vec{E}$$

charge

$$\delta W = \vec{F} \cdot d\vec{l}$$

$$\delta W = \vec{f} \cdot d\vec{l}$$

force we are applying to move the object

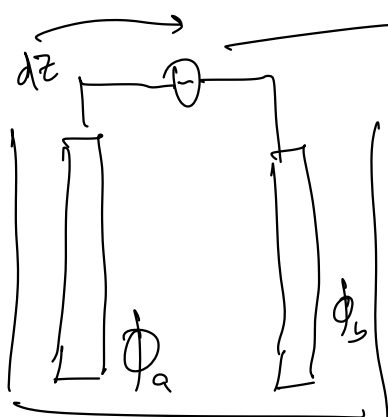
"work done on the charge against the action of the field"

$$\delta W = z \nabla\phi \cdot d\vec{l}$$

$$= z \left( \frac{\partial\phi}{\partial x} dx + \frac{\partial\phi}{\partial y} dy + \frac{\partial\phi}{\partial z} dz \right)$$

$$W = \int_a^b z \nabla\phi \cdot d\vec{l} = z (\phi_b - \phi_a) \Rightarrow \boxed{W = \Delta\phi z}$$

battery

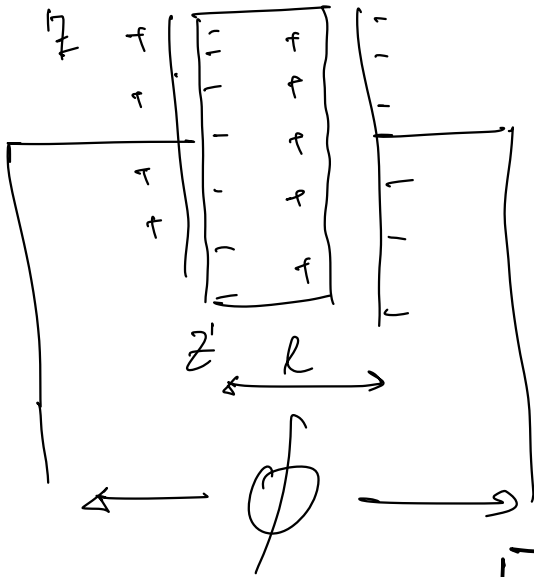


high resistance  
(trickle of charge)

$$\delta W = \Delta\phi dz$$

$\mathcal{E}$  electromotive force

## \* work to polarize a material (eg a ferroelectric material)



$\vec{E} \rightarrow$  electric field  
 $\vec{P} = \epsilon_0 \chi_e \vec{E}$  material property.  
 $\rightarrow$  induced polarization in the material (per unit volume)

$$\delta W = \vec{E} d\vec{P}$$

where  $\vec{P} = V\vec{P}$

$\rightarrow$  volume of the material.  
 $\rightarrow$  dipole moment

## \* work to magnetize a material

$$\delta W = \mu_0 H dM \rightarrow \text{total magnetic moment}$$

permeability of free space

applied magnetic field

## SUMMARY:

| <u>form of work</u>                  | <u><math>\delta W</math></u>          | <u>state Variables</u>              |
|--------------------------------------|---------------------------------------|-------------------------------------|
| lengthen a rod                       | $F dx$                                | $F \quad l$                         |
| hydrostatic pressure                 | $-p dV$                               | $p \quad V$                         |
| general strain energy                | $\sum_i \sum_j V_{ij} d\epsilon_{ij}$ | $\sigma_{ij} \quad V \epsilon_{ij}$ |
| moving charge through<br>a potential | $\phi dQ$                             | $\phi \quad Q$                      |
| polarization                         | $\vec{E} d\vec{P}$                    | $\vec{E} \quad \vec{P}$             |
| magnetization                        | $\mu_0 \vec{H} d\vec{M}$              | $\vec{H} \quad \vec{M}$             |

GENERAL

$$\delta W = \underbrace{\gamma}_{\substack{\text{intensive} \\ \text{(force)}}} \cdot d \underbrace{X}_{\substack{\text{extensive} \\ \text{(displacement)}}}$$

→ reversible work expressed in terms of state variables of the system

$\gamma$  &  $X$  are CONJUGATE VARIABLES

# Exact & Inexact differentials:

$d \rightarrow$  exact

$\delta \rightarrow$  inexact

$\vec{F} \cdot d\vec{r} \rightarrow$  exact?

exact = differential of a function

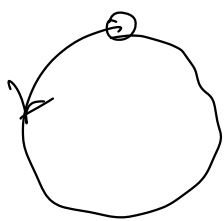
$$\vec{F} = \vec{\nabla} \phi(r)$$

$$F_x = \left( \frac{\partial \phi}{\partial x} \right)_{y,z} \quad F_y = \left( \frac{\partial \phi}{\partial y} \right)_{x,z} \quad F_z = \left( \frac{\partial \phi}{\partial z} \right)_{x,y}$$

Consequence

$$\int_{r_1}^{r_2} \vec{F} \cdot d\vec{r} = \int_{r_1}^{r_2} \vec{\nabla} \phi(r) \cdot d\vec{r} = \phi(\vec{r}_2) - \phi(\vec{r}_1)$$

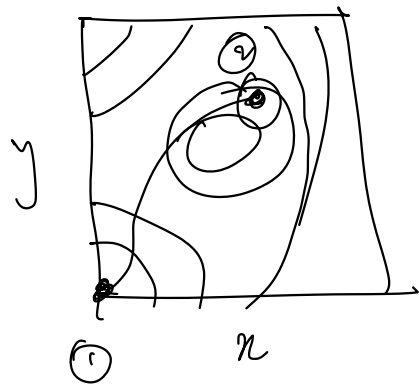
$\vec{F} \cdot d\vec{r} = \text{exact} \rightarrow F$  is "conservative"



$$\oint \vec{F} \cdot d\vec{r} = 0 \rightarrow \text{no net energy transfer}$$

EXAMPLE:

Calculate energy transfer

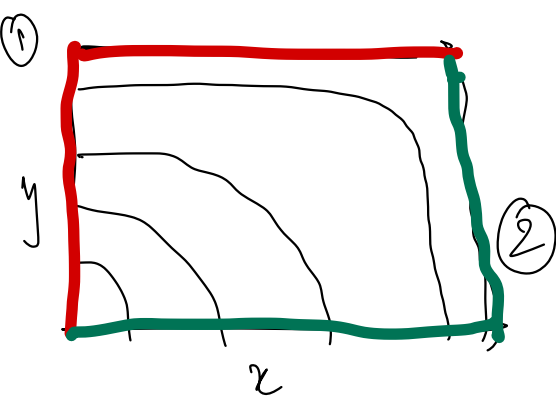


$$E_{\text{potential in field}} = mgh(x, y)$$

$$W = \int_{(1)}^{(2)} \vec{F} \cdot d\vec{r} = \int_{(1)}^{(2)} -\vec{\nabla} E \cdot d\vec{r}$$

$$W = -mg(h_2 - h_1)$$

What makes a differential inexact?



(I) Missing forces

$$dW = \int f_x dx$$

$$W(\text{path 1}) = 0 + 0$$

$\uparrow$   $\uparrow$   
 $dx=0$   $f_x=0$

$$W(\text{path 2}) \neq 0$$

$$\delta W = \int f_x dx \rightarrow \text{not exact}$$

not complete / inexact

(II) Dissipative forces (eg friction)

$$\vec{F} = -\vec{\nabla} E - f \hat{e}_r$$

$\hookrightarrow$  Unit vector along the path

$$\delta W = \int_1^2 \vec{f} \cdot d\vec{r} = - \int_1^2 \vec{\nabla} E \cdot d\vec{r} - \int_1^2 f \hat{e}_r \cdot d\vec{r}$$

$$= -mg(h_2 - h_1) - f L \quad \text{length of the path}$$

The force is not conservative:

$$\oint \delta W = -fL \neq 0$$

In all these "inexact" differentials we are missing additional contributions  $\Rightarrow$  "HEAT"

$$dU = \delta W + \delta q$$

$\nearrow$  exact differential       $\nwarrow$  inexact differential       $\nwarrow$  inexact differential

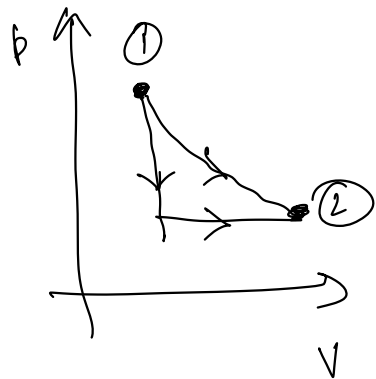
$\Rightarrow$  THERMO HAS 3 LAWS THAT INTRODUCE NEW STATE FUNCTIONS

- \* 0<sup>th</sup> law  $\rightarrow$  T (temperature)
- \* 1<sup>st</sup> law  $\rightarrow$  U (internal energy)
- \* 2<sup>nd</sup> law  $\rightarrow$  S (entropy)

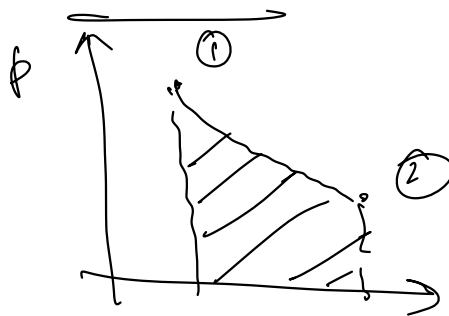
# FIRST LAW OF THERMODYNAMICS

- introduces internal energy state function
- equivalence of work and heat

\* Work depends on path : (eg) ideal gas  $\rightarrow$  reversible changes of state

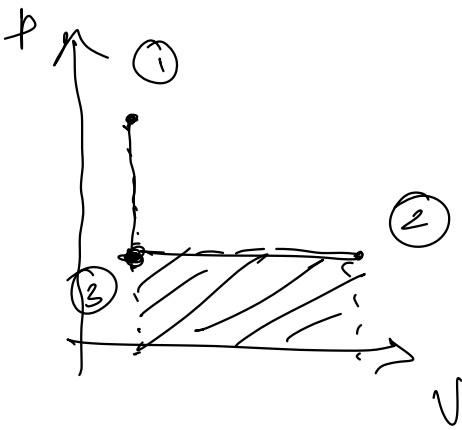


isothermal path



$$W = \int_1^2 -p dV$$

$$= \int_1^2 -\frac{RT}{V} dV = -RT \log \left( \frac{V_2}{V_1} \right)$$



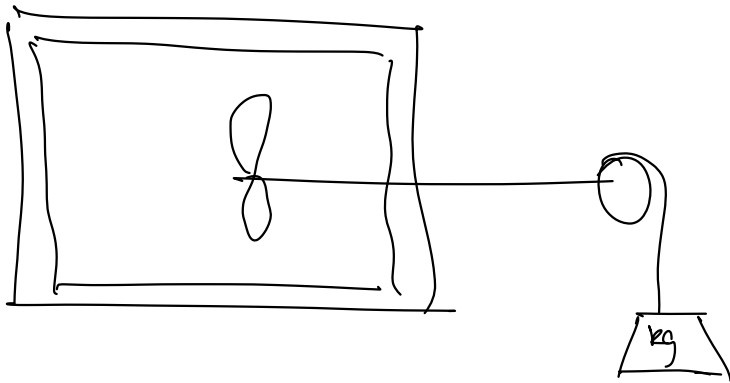
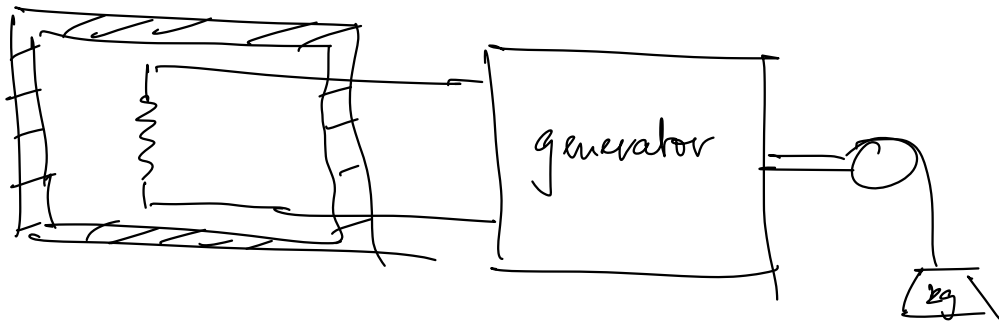
$$W = \underbrace{\int_1^3 -p dV}_0 + \underbrace{\int_3^2 -p dV}_{-(V_2 - V_3)p}$$

# JOULE'S STATEMENT OF FIRST LAW:

If an adiabatic system is caused to change from a prescribed initial to a prescribed final state, the work expended is the same for all paths connecting the two states.

$$W_{1 \rightarrow 2}^{\text{adiabatic}} = U_2 - U_1 = \Delta U$$

\* Work performed adiabatically is independent of the process



same amount  
of work connects  
the same two  
states  $(p_i, T_i)$

$\longrightarrow (p_f, T_f)$



1<sup>st</sup> law: \* system caused to change from an initial state to a final state adiabatically

$\Rightarrow$  Work done on the system is the same for all adiabatic paths connecting the two states  
 $\Rightarrow$  IMPLIES STATE FUNCTIONS

INTERNAL ENERGY:

$$W_{i \rightarrow f} (\text{adiabatic}) = U_f - U_i = \Delta U$$

\* general non-adiabatic transformation

$$\Delta U \neq W$$

$$\Delta U = W + Q \longrightarrow \text{valid for any process}$$

$$\boxed{dU = \delta W + \delta Q} \longrightarrow \text{differential form}$$

exact differential      inexact differential

reversible:

$$\boxed{dU = \sum_i \gamma_i dx_i + \delta Q}$$

# SECOND LAW OF THERMODYNAMICS

Kelvin-Planck Statements  $\longleftrightarrow$  Clausius Statements  $\longleftrightarrow$  Carothedory Statement

All equivalent  $\Rightarrow$  proofs can be found in any standard thermo book

$$dU = \delta Q + \sum_i \gamma_i dx_i$$

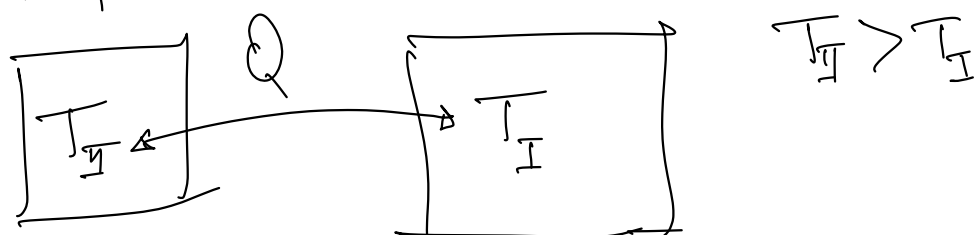
~~~~~

↑  
??

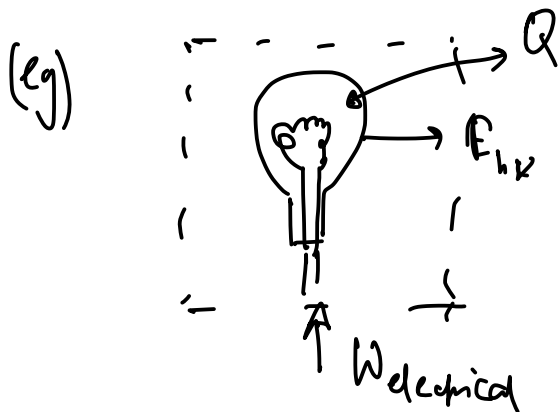
expressed in terms of state variables of the system

What do we do about the "heat" term?

NOTICE: The 1<sup>st</sup> law does not make any statements about the direction of the process.



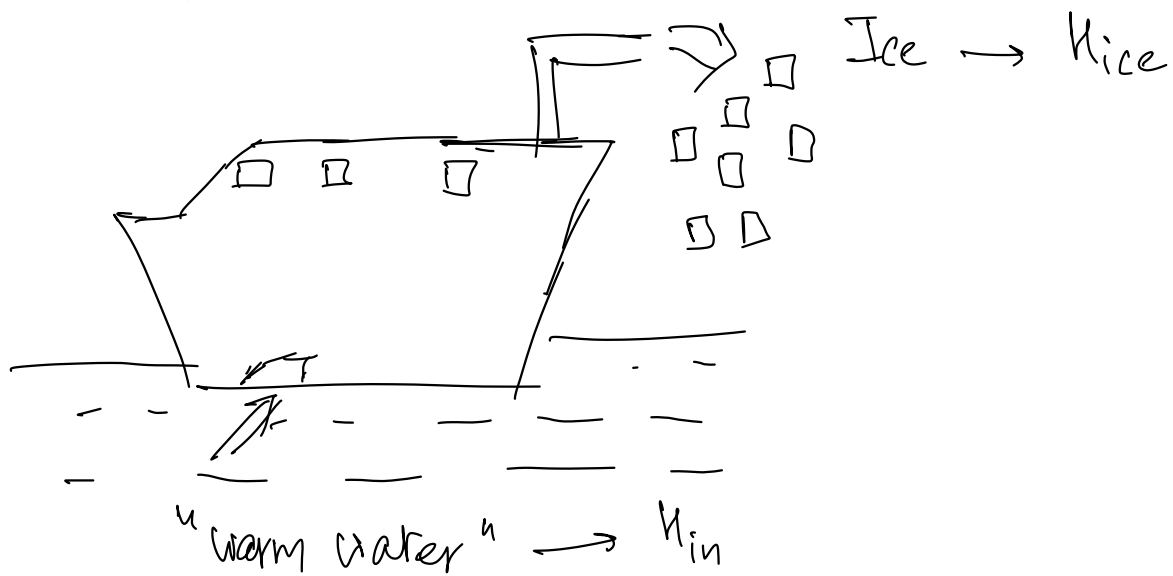
Based on the 1<sup>st</sup> law of thermo heat can flow from  $II \rightarrow I$  &  $I \rightarrow II$



No reason (from the 1<sup>st</sup> law) that we can't run this process in reverse.

# \* POWERING THE BOATS IN LAC LEMAN WITH LAKE WATER :

Simplon / LaSuisse : CGN boats  $\rightarrow$  850 people !!



$\Rightarrow$  1500 horsepower (actually 1400)  $\rightarrow$  LOTS OF ENERGY : Cars are usually  $\approx$  100-200 hp

$$\Rightarrow 1120 \text{ kJ/s}$$

$$W = -1120 \text{ kJ/s} = (H_{\text{ice}} - H_{\text{in}}) \dot{m}$$

$$H_{\text{ice}} - H_{\text{in}} \approx -334 \text{ J/g}$$

$$\Rightarrow \boxed{\dot{m} \approx 3.35 \text{ kg/s}}$$