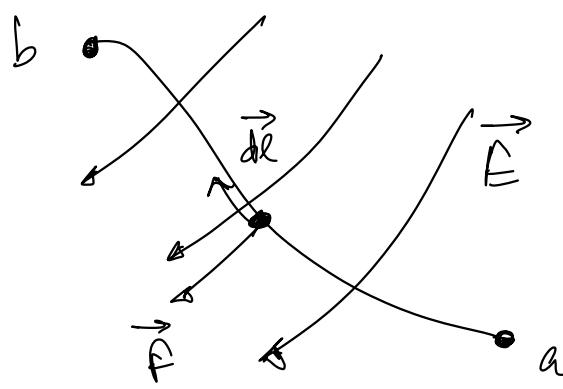


Moving charge in an electric field



$$\vec{E} = -\nabla\phi \rightarrow \text{potential}$$

$$\vec{F} = z \vec{E}$$

charge

$$\delta W = \vec{F} \cdot \vec{dl}$$

$$\delta W = \vec{F} \cdot \vec{dl}$$

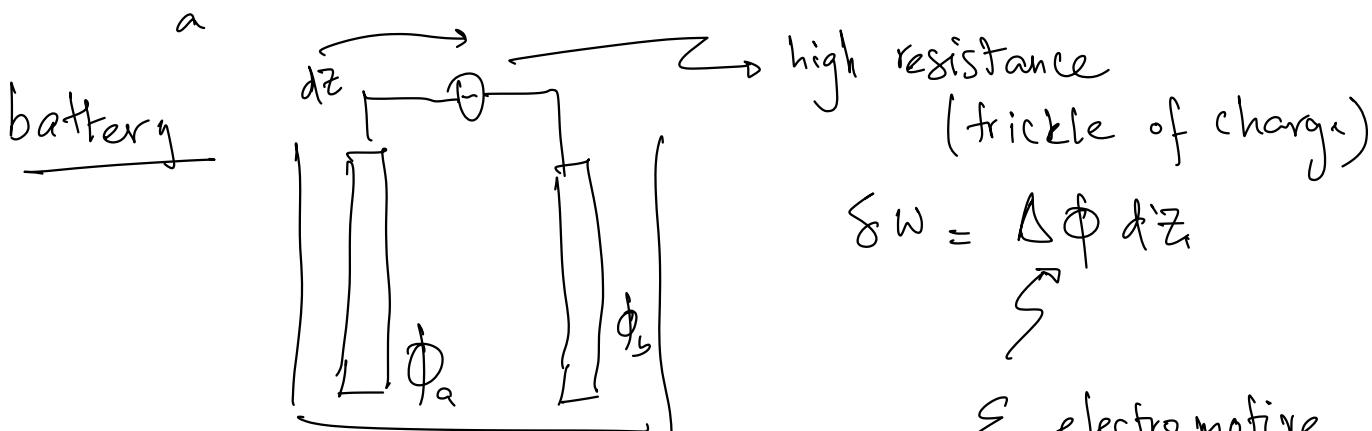
force we are applying to move the object

"Work done on the charge against the action of the field"

$$\delta W = z \nabla\phi \cdot \vec{dl}$$

$$= z \left(\frac{\partial\phi}{\partial x} dx + \frac{\partial\phi}{\partial y} dy + \frac{\partial\phi}{\partial z} dz \right)$$

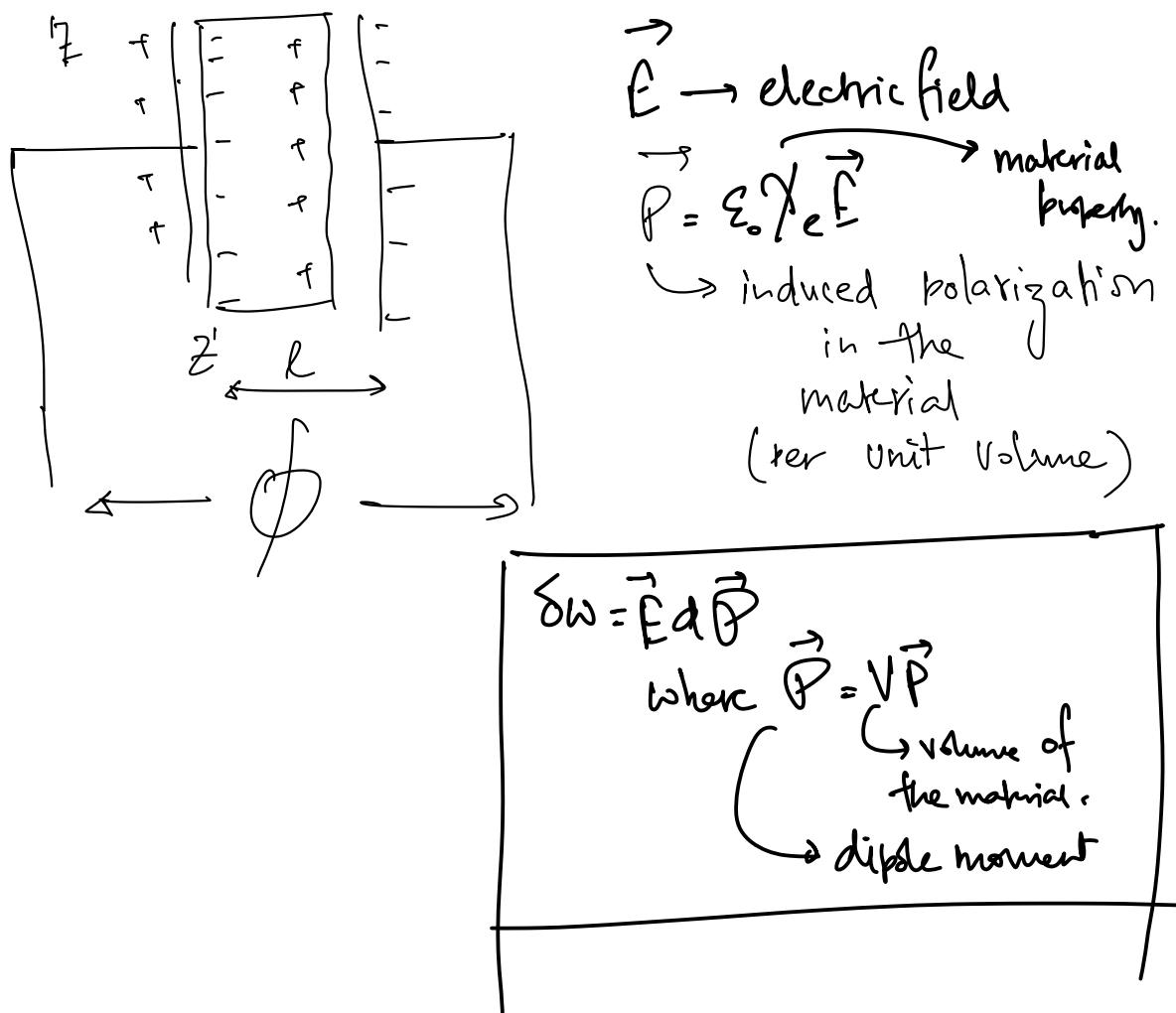
$$W = \int_a^b z \nabla\phi \cdot \vec{dl} = z (\phi_b - \phi_a) \Rightarrow [W = \Delta\phi z]$$



$$\delta W = \Delta\phi dz$$

E electromotive force

* Work to polarize a material (eg a ferroelectric material)



* Work to magnetize a material

$$\delta W = \mu_0 H dM \rightarrow \text{total magnetic moment}$$

permeability of free space applied magnetic field

SUMMARY:

<u>form of work</u>	<u>δW</u>	<u>state variables</u>
lengthen a rod	$\int F dl$	$F \quad l$
hydrostatic pressure	$-\int p dV$	$p \quad V$
general strain energy	$\sum_i \sum_j V_0 \epsilon_{ij} d\epsilon_{ij}$	$\epsilon_{ij} \quad V_0 \epsilon_{ij}$
moving charge through a potential	$\int \phi d\vec{r}_i$	$\phi \quad \vec{r}_i$
polarization	$\int \vec{E} d\vec{P}$	$\vec{E} \quad \vec{P}$
magnetization	$\int \vec{M}_0 \vec{H} d\vec{M}$	$\vec{H} \quad \vec{M}$

GENERAL

$$\delta W = \gamma \cdot dX$$

γ (intensive)
 dX (extensive)

reversible work expressed in terms of state variables of the system

of the system

γ & X are CONJUGATE VARIABLES

Exact & Inexact differentials:

\oint → exact

δ → inexact

$\vec{F} \cdot d\vec{r}$ → exact?

exact = differential of a function

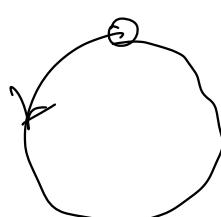
$$\vec{F}_z = \vec{\nabla} \phi(r)$$

$$F_x = \left(\frac{\partial \phi}{\partial x} \right)_{y,z} \quad F_y = \left(\frac{\partial \phi}{\partial y} \right)_{x,z} \quad F_z = \left(\frac{\partial \phi}{\partial z} \right)_{x,y}$$

Consequence

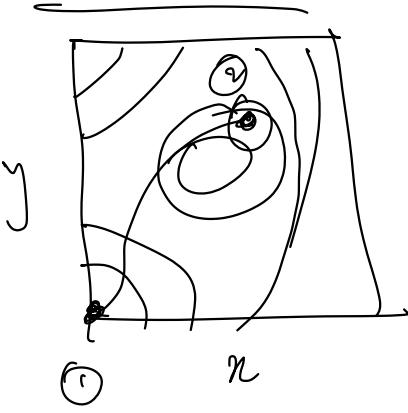
$$\int_{\vec{r}_1}^{\vec{r}_2} \vec{F} \cdot d\vec{r} = \int_{\vec{r}_1}^{\vec{r}_2} \vec{\nabla} \phi(r) \cdot d\vec{r} = \phi(\vec{r}_2) - \phi(\vec{r}_1)$$

$\vec{F} \cdot d\vec{r}$ = exact → F is "conservative"



$\oint \vec{F} \cdot d\vec{r} = 0 \rightarrow$ no net energy transfer

EXAMPLE:

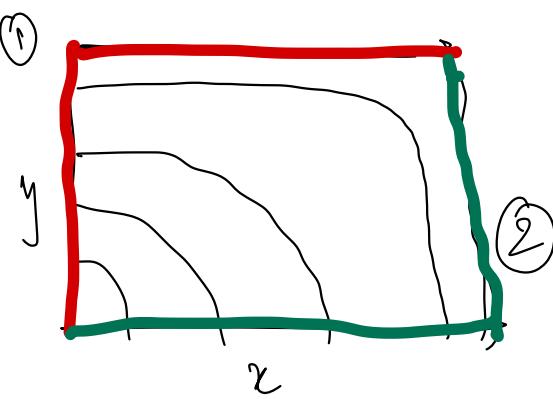


Calculate energy transfer

$$\text{potential in field} = mg h(x, y)$$
$$W = \int \vec{F} \cdot d\vec{r} = \int -\vec{\nabla} E \cdot d\vec{r}$$

$$W_2 - mg(h_2 - h_1)$$

What makes a differential inexact?



① Missing forces

$$dW = \int f_x dx$$

$$W(\text{path 1}) = 0 + 0$$

\uparrow \uparrow

$$dx = 0 \quad f_x = 0$$

$$W(\text{path 2}) \neq 0$$

$$\delta W = \int f_x dx \rightarrow \text{not exact}$$

\uparrow

$$\text{not complete / inexact}$$

② Dissipative forces (e.g. friction)

$$\vec{F} = -\vec{\nabla} E - f \hat{e}_r$$

\curvearrowright Unit vector along the path

$$\delta W = \int \vec{f} \cdot d\vec{r} = - \int \vec{F} \cdot d\vec{r} - \int f \hat{e}_r \cdot d\vec{r}$$

$$= -mg(h_2 - h_1) - f L \xrightarrow{\text{length of the path}}$$

The force is not conservative:

$$\oint \delta W = -fL \neq 0$$

In all these "inexact" differentials we are missing additional contributions \Rightarrow "HEAT"

$$dU = \delta W + \delta q$$

\swarrow \swarrow \swarrow \swarrow

exact differential inexact inexact
 differential

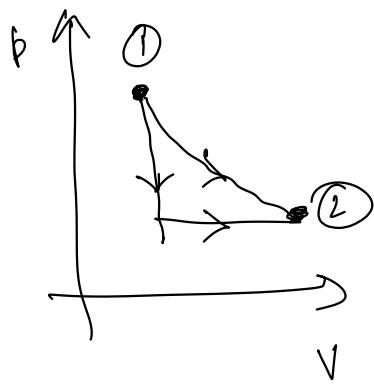
\Rightarrow THERMO HAS 3-LAWS THAT INTRODUCE NEW STATE FUNCTIONS

- * 0th law $\rightarrow T$ (temperature)
- * 1st law $\rightarrow U$ (internal energy)
- * 2nd law $\rightarrow S$ (entropy)

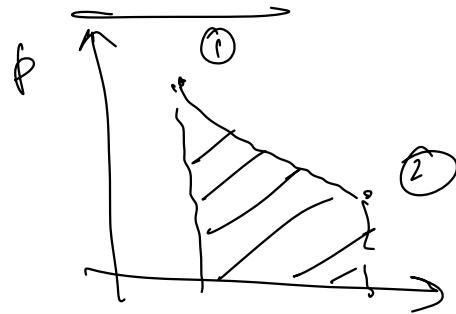
FIRST LAW OF THERMODYNAMICS

- introduces internal energy state function
- equivalence of work and heat

* Work depends on path : (eg) ideal gas \rightarrow reversible changes of state

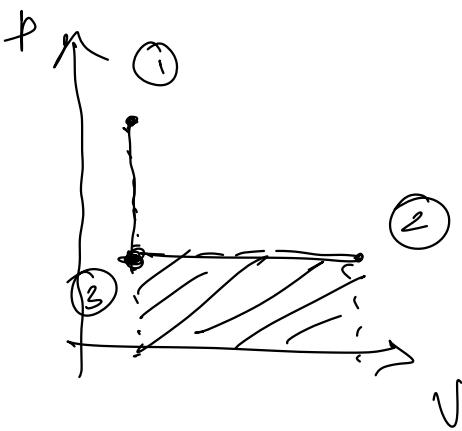


isothermal path



$$W = \int_1^2 -P dV$$

$$= \int_1^2 -\frac{RT}{V} dV = -RT \ln \left(\frac{V_2}{V_1} \right)$$



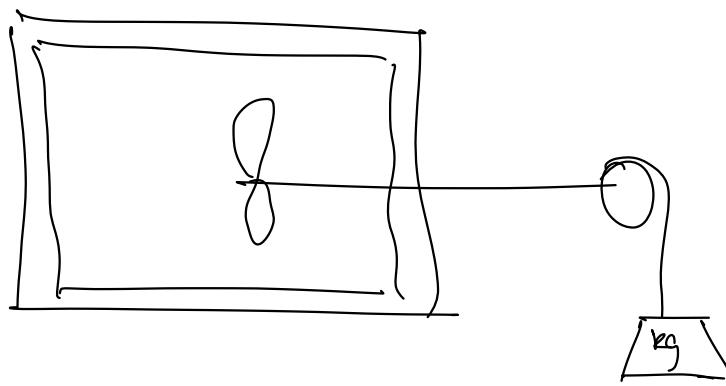
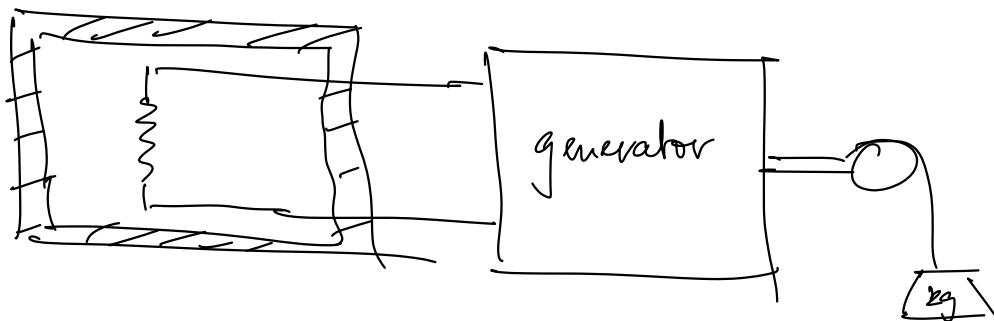
$$W = -P dV + \int_2^3 -P dV$$
$$= - (V_2 - V_3) P$$

JOULE'S STATEMENT OF FIRST LAW:

If an adiabatic system is caused to change from a prescribed initial to a prescribed final state, the work expended is the same for all paths connecting the two states.

$$W_{1 \rightarrow 2}^{\text{adiabatic}} = U_2 - U_1 = \Delta U$$

* Work performed adiabatically is independent of the process



same amount
of work connects
the same two
states (p_i, T_i)

$\rightarrow (p_f, T_f)$

1st law: * system caused to change from an initial state to a final state adiabatically

⇒ Work done on the system is the same for all adiabatic paths connecting the two states
⇒ IMPLIES STATE FUNCTION

INTERNAL ENERGY:

$$W_{i \rightarrow f} (\text{adiabatic}) = U_f - U_i = \Delta U$$

* General non-adiabatic transformation

$$\Delta U \neq W$$

$$\Delta U = W + Q \longrightarrow \text{Valid for any process}$$

$$\boxed{dU = \delta W + \delta Q} \rightarrow \text{differential form}$$

exact differential inexact differential

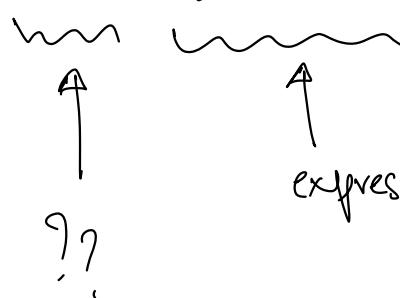
Reversible:

$$\boxed{\int dU = \sum_i \gamma_i dX_i + \delta Q}$$

SECOND LAW OF THERMODYNAMICS

Kelvin-Planck Statements \leftrightarrow Clausius Statements \leftrightarrow Caratheodory Statement
All equivalent \Rightarrow proofs can be found in any standard thermo book

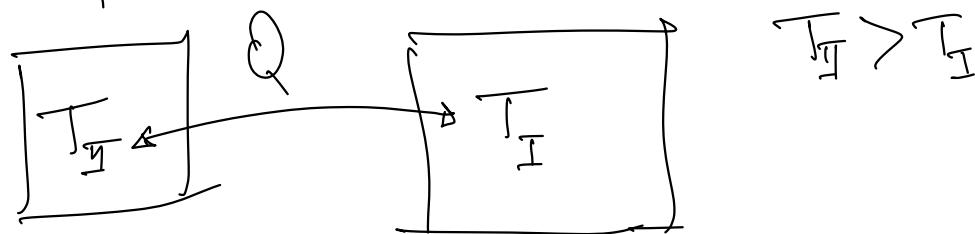
$$dU = \delta Q + \sum_i Y_i dx_i$$



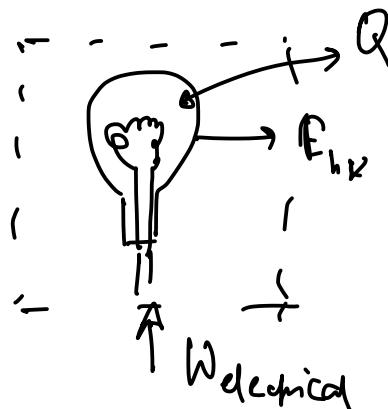
expressed in terms of state variables of the system

What do we do about the "heat" term?

Notice: The 1st law does not make any statements about the direction of the process.



Based on the 1st law of thermo heat can flow from $II \rightarrow I$ & $I \rightarrow II$

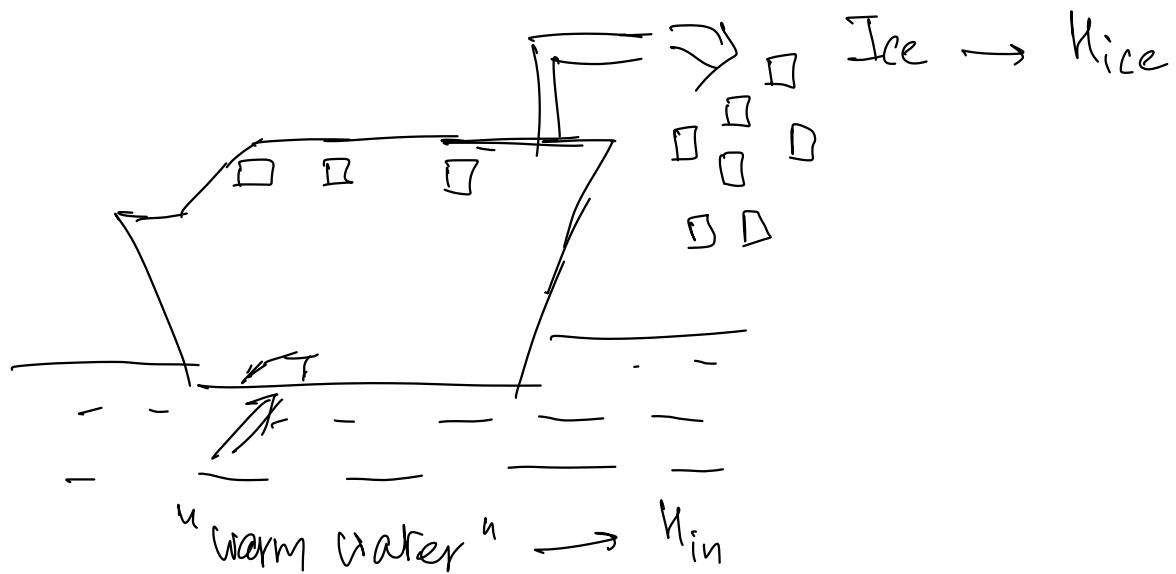


Notice (from the 1st law) that we can't run this process in reverse.

(eg)

* POWERING THE BOATS IN LAC LÉMAN WITH LARE WATER :

Simplon | LaSuisse : CGN boats \rightarrow 850 people !!



\approx 1500 horsepower (actually 1400) \rightarrow LOTS OF ENERGY : cars are usually \approx 100-200 hp

$$W = -1120 \text{ kJ/s} = (H_{ice} - H_{in}) \dot{m}$$

$$H_{ice} - H_{in} \approx -334 \text{ J/g}$$

$$\Rightarrow \dot{m} \approx 3.35 \text{ kg/s}$$