

3.2 Finite element approximation

$$\begin{cases} -\operatorname{div}(\varepsilon \nabla u) + \underline{b} \cdot \nabla u = f & \Omega \\ \vartheta(u) = g & \partial\Omega \end{cases}$$

$$u = u_0 + u_g \quad u_0 \in H_0^1(\Omega) \quad u_g \text{ lifting.}$$

$$V_n^g = \{ u \in H^1(\Omega) \mid \vartheta(u) = \mathcal{I}_n g, \quad u|_K \in \mathcal{P}^r(K) \quad \forall K \in \mathcal{T}_h \}$$

$$\|u - u_n\|_{H^1} \leq \underbrace{M/d}_{\text{inf}} \inf_{v_n \in V_n^g} \|u - v_n\|_{H^1}$$

$$\frac{M}{d} \sim Pe = \frac{|b|}{\varepsilon}$$

$$r=1 \quad \inf_{v_n} \|u - v_n\|_{H^1} \leq \underline{h}. \quad \cancel{\|u\|_{H^2}^2}$$

$$\text{if } Pe h = \frac{|b| h}{2\varepsilon} > 1 \Rightarrow Pe > \frac{2}{h}$$

$$\Rightarrow \|u - u_n\|_{H^1} \lesssim Pe \cdot h \Rightarrow \text{the error is bounded by a constant that is not small.}$$

Artificial viscosity

$$a(u, v) = \int_{\Omega} \varepsilon \nabla u \cdot \nabla v + (\underline{b} \cdot \nabla u) v$$

$$a_n(u, v) = \int_{\Omega} \left(\varepsilon + \gamma \frac{|b|h}{2} \right) \nabla u \cdot \nabla v + (\underline{b} \cdot \nabla u) v$$

γ = a parameter that I need to choose.

let's try an error estimate:

$$a(u, v) = \int_{\Omega} f v$$

$$a_n(u_n, v_n) = \int_{\Omega} f v_n$$

• non-consistent method

$$a_n(u, v_n) \neq \int_{\Omega} f v_n$$

a_n

$$a_n(u, u) \geq \alpha_n \|u\|_{H^1}^2$$

$$a_n(u, v) \leq M_n \|u\|_{H^1} \cdot \|v\|_{H^1}$$

$$\alpha_n \|u_n - w_n\|_{H^1}^2 \leq a_n(u_n - w_n, u_n - w_n) =$$

$$\underbrace{a_n(u_n - u, u_n - w_n)}_I + \underbrace{a(u - w_n, u_n - w_n)}_{II}$$

$$II \leq \underbrace{\uparrow}_{\text{continuity of } a_n} M_n \cdot \|u - w_n\|_{H^1} \cdot \|u_n - w_n\|_{H^1}$$

$$I = a_n(u_n, u_n - w_n) - a_n(u, u_n - w_n) = \int_{\Omega} f(u_n - w_n) - a_n(u, u_n - w_n)$$

$$= \int_{\Omega} f(u_n - w_n) - \int_{\Omega} \varepsilon \nabla u \nabla (u_n - w_n) - \gamma \frac{|b|h}{2} \int_{\Omega} \nabla u \nabla (u_n - w_n)$$

$u = \text{solution of the continuous problem!}$

$$- \int_{\Omega} \underline{b} \cdot \nabla u (u_n - w_n)$$

$$I = - \gamma \frac{|b|h}{2} \int_{\Omega} \nabla u \nabla (u_n - w_n)$$

$$I \leq \|u_n - w_n\|_{H^1} \cdot \sup_{v_n \in V_n^0} \frac{- \gamma \frac{|b|h}{2} \int_{\Omega} \nabla u \cdot \nabla v_n}{\|v_n\|_{H^1}}$$

$$\leq \|u_n - w_n\|_{H^1} \cdot \frac{\gamma |b|h}{2} \cdot \|u\|_{H^1}$$

$$\alpha_n \|u_n - w_n\|_{H^1}^2 \leq I + II$$

$$\leq \gamma \frac{|b|h}{2} \|u\|_{H^1} \cdot \|u_n - w_n\| + M_n \|u - w_n\| \cdot \|u_n - w_n\|$$

where $w_n \in V_n^g$ is any function in the space.

$$\|u - u_n\|_{H^1} \leq \|u - w_n\|_{H^1} + \|u_n - w_n\|_{H^1}$$

$$\|u - u_n\|_{H^1} \leq \left(1 + \frac{M_n}{\alpha_n}\right) \|u - w_n\|_{H^1} + \frac{\gamma |b|h}{2} \|u\|_{H^1}$$

$$M_n \leq |b|$$

$$\frac{M_n}{\alpha_n} = \frac{2|b|}{2\varepsilon + \gamma |b|h} \lesssim \frac{|b|}{\gamma |b|h} \leq \frac{1}{\gamma} \cdot \frac{1}{h}$$

$$\alpha_n = \varepsilon + \gamma \frac{|b|h}{2}$$

$$\|u - u_n\| \lesssim \left(1 + \frac{1}{\gamma} \cdot \frac{1}{h}\right) \inf_{w_n \in V_n^g} \|u - w_n\|_{H^1} + \frac{1}{\alpha} \frac{\gamma |b|h}{2} \|u\|_{H^1}$$

Streamline diffusion $a_n(u, v) = a(u, v) + \frac{\gamma h}{|b|} \int_{\Omega} (\underline{b} \cdot \nabla u) (\underline{b} \cdot \nabla v)$

Fully consistent methods

SUPG = streamline - Petrov Galerkin

GALS = Galerkin - least squares

$$-\varepsilon \Delta u + \underline{b} \cdot \nabla u = f$$

$$\bullet \quad \underline{b} \cdot \nabla v$$

$$\int (-\varepsilon \Delta u + \underline{b} \cdot \nabla u) \underline{b} \cdot \nabla v = \int f \underline{b} \cdot \nabla v$$

formally.

SUPG

$$a_n(u_h, v_h) = a(u_h, v_h) +$$

$$\sum_{k \in T_h} \delta_k \int_k (-\varepsilon \Delta u_h + \underline{b} \cdot \nabla u_h) \cdot \underline{b} \cdot \nabla v_h$$

$$\langle f_h, v_h \rangle = \int_{\Omega} f \cdot v_h + \sum_{K \in \mathcal{T}_h} \delta_K \int_K f \underline{b} \cdot \nabla v_h$$

SUPG is $a_h(u_h, v_h) = \langle f_h, v_h \rangle$.

- SUPG is consistent

$$\begin{aligned} a_h(u, v_h) &= \boxed{a(u, v_h)} + \sum_K \delta_K \int_K \boxed{(-\varepsilon \Delta u + \underline{b} \cdot \nabla u) \cdot \underline{b}} \nabla v_h \\ &= \boxed{\int f v_h} + \sum_K \delta_K \int_K \boxed{f} \underline{b} \cdot \nabla v_h \end{aligned}$$

Consistency $\rightarrow$ Galerkin orthogonality. $\rightarrow$

$$\|u - u_h\|_{H^1} \leq \underbrace{\|M\|_{\alpha}}_{a_h(\cdot, \cdot)} \|u - I_h u\|_{H^1}$$

$$\delta_K = \frac{h_K}{\|\underline{b}\|} \cdot \gamma_K$$

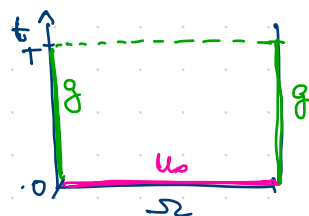
γ_K = depends on the degree,
on the shape of K.

CHAPTER 4

Parabolic problems -

Heat equation.

$$\begin{cases} \frac{\partial u}{\partial t} - \operatorname{div}(A \nabla u) = f & \Omega \times (0, T) \\ u(x, 0) = u_0(x) & \Omega \\ \gamma(u(x, t)) = g(t) & \partial\Omega \end{cases}$$



propagation of the heat in a medium.

$$u(x, t) : \Omega \times (0, T) \rightarrow \mathbb{R}$$

A is the diffusion matrix, A is positive definite.

4.1 Weak formulation & wellposedness.

$$\begin{aligned} L^2(0, T; L^2(\Omega)) &= \{ t \mapsto u(x, t) \in L^2(\Omega) \mid t \mapsto \|u\|_{L^2(\Omega)} \in L^2(0, T) \} \\ &= \{ u(x, t) : \int_0^T \int_{\Omega} |u|^2 < +\infty \}. \end{aligned}$$

$$\begin{aligned} L^2(0, T; H^1(\Omega)) &= \{ t \mapsto u(x, t) \in H^1(\Omega) \mid t \mapsto \|u(x, t)\|_{H^1(\Omega)} \in L^2(0, T) \} \\ &= \{ u(x, t) : \int_0^T \|u\|_{H^1(\Omega)}^2 < +\infty \}. \end{aligned}$$

$$L^\infty(0, T; L^2(\Omega)) = \{ u(x, t) : \max_{a.e. t \in (0, T)} \|u\|_{L^2(\Omega)} < +\infty \}$$

$$C^0(0, T; L^2(\Omega)) = \{ u(x, t) : \max_{t \in (0, T)} \|u\|_{L^2(\Omega)} < +\infty \}$$

Theorem 1 If $f \in L^2(0, T; L^2(\Omega))$, $g = 0$

and $u_0 \in L^2(\Omega)$, then $\exists! u$

$$u \in L^2(0, T; H^1(\Omega)) \left[\text{and } \frac{\partial u}{\partial t} \in L^2(0, T, (H_0^1(\Omega))') \right]$$

theorem 2

Indeed $u \in C^0(0, T; L^2(\Omega))$

theorem 3

$f \in L^2(0, T; L^2(\Omega))$, $g = 0$ and

$u_0 \in H_0^1(\Omega) \Rightarrow$

$u \in L^\infty(0, T; H^1(\Omega)) \quad \frac{\partial u}{\partial t} \in L^2(0, T; L^2(\Omega))$

Sketch of the proofs of theorem 1 & 3.

$H_0^1(\Omega) \quad V_N = \text{span} \{ \varphi_1, \dots, \varphi_N \}$

$\{ \varphi_1, \varphi_2, \dots \}$ is basis for $H_0^1(\Omega)$.

$$\int_{\Omega} \left(\frac{\partial u}{\partial t} - \text{div}(A \nabla u) \right) \cdot v = f \cdot v, \quad u \in H_0^1(\Omega)$$

$$\int_{\Omega} \frac{\partial u}{\partial t} \cdot v + \int_{\Omega} A \nabla u \cdot \nabla v = \int_{\Omega} f v, \quad v \in H_0^1(\Omega). \quad (1)$$

Galerkin method

$u_N \in V_N$ such that

$$\int_{\Omega} \frac{\partial u_N}{\partial t} v_N + \int_{\Omega} A \nabla u_N \cdot \nabla v_N = \int_{\Omega} f v_N$$

$$u_N = \sum_i u_N^i(t) \varphi_i \quad \underline{u}_N = [u_N^1 \dots u_N^N]$$

$v_N = \sum v_N^e \varphi_e$
test function.

$$\frac{\partial u_N}{\partial t} = \sum_i (u_N^i(t))' \varphi_i$$

$$\int_{\Omega} \frac{\partial u_N}{\partial t} \cdot v_N = \sum_i \sum_e (u_N^i(t))' \overbrace{v_e \int_{\Omega} \varphi_i \varphi_e}^{\text{Mie}} = \underline{\tilde{M}} \dot{\underline{u}}_N$$

$$\underline{u}_N = [u_N^1 \dots u_N^N]$$

$$\underline{v} = [v_N^1 \dots v_N^N]$$

$$\begin{cases} M \dot{\underline{u}}_N + A \underline{u}_N = \underline{f}_N \\ \underline{u}_N(0) = \underline{u}_N^0 \end{cases}$$

$$\underline{f}_N^i = \int f \varphi_i$$

$\underline{u}_N^0$ = coefficients of $u_0(x)$ in the basis.

system of ODEs.

M is symmetric and positive definite $\rightarrow \exists! \underline{u}_N$
 A is positive definite.

$\Rightarrow u_N(x, t)$ exists as Galerkin solution of (1)

the solution u of (HE) is the "limit" of u_N .

We go back to Galerkin

$$\int_{\Omega} \frac{\partial u_N}{\partial t} v_N + \int_{\Omega} A \nabla u_N \nabla v_N = \int_{\Omega} f v_N \quad v_N = u_N$$

$$\int_0^T \left[\frac{1}{2} \int_{\Omega} \frac{\partial}{\partial t} |u_N|^2 + \int_{\Omega} A |\nabla u_N|^2 \right] = \int_{\Omega} f u_N$$

$$\frac{1}{2} \|u_N(t)\|_{L^2(\Omega)}^2 - \frac{1}{2} \|u_N(0)\|_{L^2(\Omega)}^2 + \int_0^T \int_{\Omega} A |\nabla u_N|^2 = \int_0^T \int_{\Omega} f \cdot u_N$$

• coercity $\int_0^T \alpha \|u_N\|_{H_0^1(\Omega)}^2 \leq \int_0^T \int_{\Omega} A |\nabla u_N|^2$

$$\frac{1}{2} \|u_N(t)\|_{L^2(\Omega)}^2 + \alpha \int_0^T \|u_N\|_{H_0^1(\Omega)}^2 \leq \int_0^T \int_{\Omega} f u_N + \frac{1}{2} \|u_N(0)\|_{L^2(\Omega)}^2$$

the sequence u_N is bounded in $L^\infty(0, T; L^2(\Omega))$
and in $L^2(0, T; H^1_0(\Omega)) \dots$

$$\begin{aligned} \int_0^T \int_{\Omega} f u_N &\leq \int_0^T \|f\|_{L^2(\Omega)} \cdot \|u_N\|_{L^2(\Omega)} \\ &\leq \|f\|_{L^2(L^2)} \cdot \|u_N\|_{L^2(L^2)} \\ &\leq \|f\|_{L^2(L^2)} \cdot \|u_N\|_{L^2(H^1_0)} \\ &\leq \frac{1}{2\alpha} \|f\|_{L^2(L^2)}^2 + \frac{\alpha}{2} \|u_N\|_{L^2(H^1_0)}^2 \end{aligned}$$

$$\left[\begin{aligned} 2ab &\leq a^2 + b^2 \quad \leftarrow (a-b)^2 = a^2 + b^2 - 2ab \geq 0 \\ 2(\gamma a)\left(\frac{b}{\gamma}\right) &\leq \gamma^2 a^2 + \frac{b^2}{\gamma^2} \end{aligned} \right]$$

$$\frac{1}{2} \|u_N(T)\|_{L^2(\Omega)}^2 + \alpha \int_0^T \|u_N\|_{H^1_0(\Omega)}^2 \leq \int_0^T \int_{\Omega} f u_N + \frac{1}{2} \|u_N(0)\|_{L^2(\Omega)}^2$$

$$\frac{1}{2} \|u_N(T)\|_{L^2(\Omega)}^2 + \frac{\alpha}{2} \int_0^T \|u_N\|_{H^1_0}^2 \leq \frac{1}{2\alpha} \int_0^T \|f\|_{L^2(\Omega)}^2 + \frac{1}{2} \|u_N(0)\|_{L^2}^2$$

u_N is bounded $\uparrow$

(2)

initial condition

- there is a mathematical theory that guarantees that u_N admits a limit u a subsequence of

- Such a limit is a solution of the heat equation.

Let's continue with the bounds.

$$\int_{\Omega} \frac{\partial u_N}{\partial t} \cdot v_N + \int_{\Omega} A \nabla u_N \cdot \nabla v_N = \int f v_N \quad v_N = \frac{\partial u_N}{\partial t}$$

$$\int_{\Omega} \left| \frac{\partial u_N}{\partial t} \right|^2 + \underbrace{\int_{\Omega} A \nabla u_N \cdot \nabla \frac{\partial u_N}{\partial t}}_{\frac{1}{2} \frac{\partial}{\partial t} \int_{\Omega} A |\nabla u_N|^2} = \int f \frac{\partial u_N}{\partial t}$$

$$\begin{aligned} \bullet \int_0^T \int_{\Omega} \left| \frac{\partial u_N}{\partial t} \right|^2 + \frac{1}{2} \int_{\Omega} A |\nabla u_N(T)|^2 - \frac{1}{2} \int_{\Omega} A |\nabla u_N(0)|^2 \\ = \int_0^T \int_{\Omega} f \frac{\partial u_N}{\partial t} \end{aligned}$$

$$\bullet \int_0^T \int_{\Omega} \left| \frac{\partial u_N}{\partial t} \right|^2 + \frac{1}{2} \alpha \|u_N(T)\|_{H_0^1}^2 \leq \|f\|_{L^2(L^2)} \cdot \left\| \frac{\partial u_N}{\partial t} \right\|_{L^2(L^2)} + \frac{1}{2} M \cdot \|u_N(0)\|_{H_0^1}^2$$

$$\bullet \int_0^T \int_{\Omega} \left| \frac{\partial u_N}{\partial t} \right|^2 + \alpha \|u_N(T)\|_{H_0^1}^2 \leq \|f\|_{L^2(L^2)} + M \|u_N(0)\|_{H_0^1}^2$$

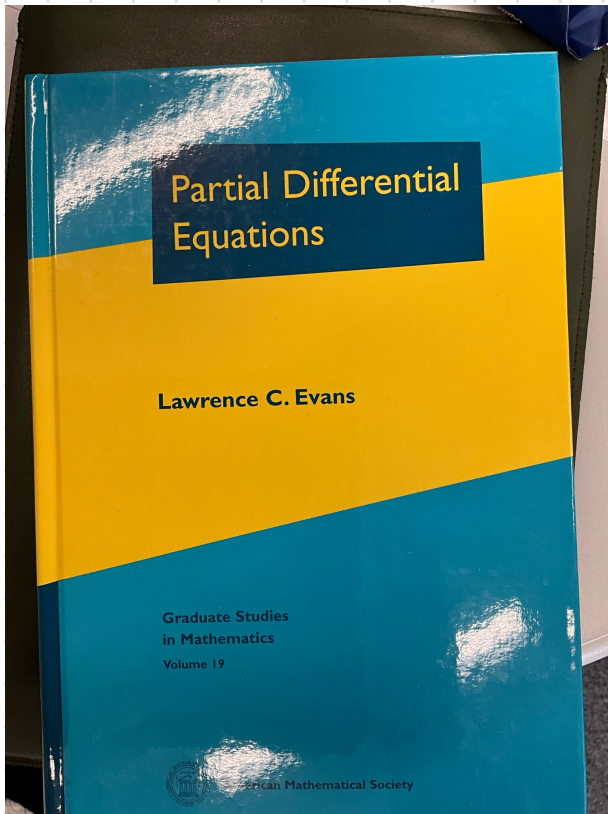
(3)

if $u_N(0) \in H_0^1(\Omega)$, $u_N(0)$ is u_0 , the initial cond.

$$\Rightarrow \frac{\partial u_N}{\partial t} \in L^2(0, T; L^2(\Omega)), \quad u_N \in L^\infty(0, T; H_0^1(\Omega))$$

→ these regularities extend to the limit, i.e., to the solution u of (HE).





Read here if you are interested
in the fully detailed mater. proof!
CHAPTER 7 :-)