

Enroluation: tablette, une seule page.

examen, 2 h 15
poids

10 à 15 questions QCM, 1 ou 2 exercices à rédiger
 $\frac{2}{3}$
ou $\frac{1}{2}$

$\frac{1}{3}$
 $\frac{1}{2}$
transf. de Laplace + thm des résidus

formules à savoir (pas les tables des transformées)

Séries de Fourier: chaleur, ondes

Transformées de Fourier: $\longrightarrow$

$\xrightarrow{x}$ ou $(0, L)$
 $\xrightarrow{1}$ $x \in \mathbb{R}$

$$\frac{\partial u}{\partial x} - \frac{\partial^2 u}{\partial x^2} = f$$

L'eq. des ondes et transf de Fourier (pas dans le livre)

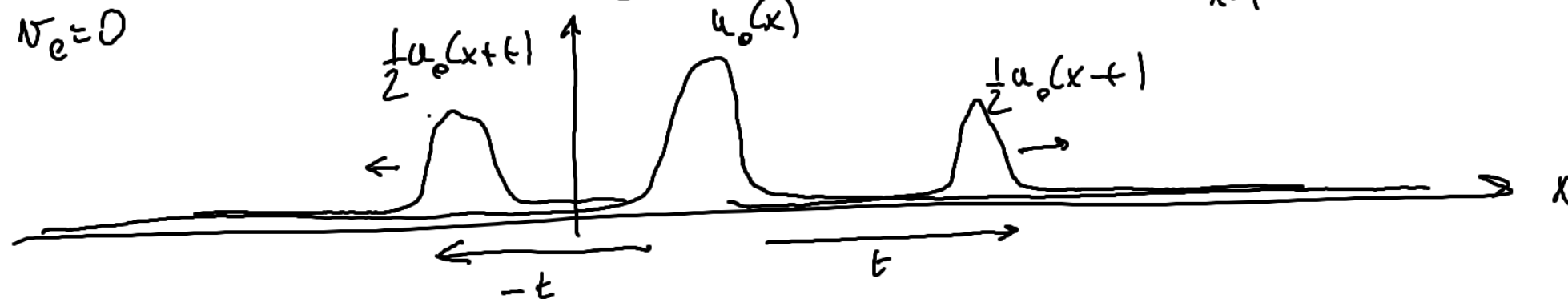
Etant donné $f(x,t)$ $x \in \mathbb{R}, t > 0$, $u_0(x), v_0(x)$ on cherche $u(x,t)$ tq

$$\begin{cases} \frac{\partial^2 u}{\partial t^2}(x,t) - \frac{\partial^2 u}{\partial x^2}(x,t) = f(x,t) & x \in \mathbb{R} \quad t > 0 \quad \text{eq. ondes} \\ u(x,0) = u_0(x) \\ \frac{\partial u}{\partial t}(x,0) = v_0(x) \end{cases}$$

Remarques: • Pas de conditions aux limites mais $\int_{-\infty}^{+\infty} |u_0(x)| dx < +\infty$ et $\int_{-\infty}^{+\infty} |v_0(x)| dx < +\infty$
 • Si $f(x,t) = 0$ la sol est donnée par

$$u(x,t) = \frac{1}{2} (u_0(x-t) + u_0(x+t)) + \frac{1}{2} \int_{x-t}^{x+t} v_0(y) dy \quad \text{Duhamel}$$

$v_0 = 0$



On pose $\hat{u}(\alpha, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} u(x, t) e^{-i\alpha x} dx$ $\alpha \in \mathbb{R}$ t paramètre.

idem pour $\hat{f}(\alpha, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f(x, t) e^{-i\alpha x} dx$ et $\hat{u}_0(\alpha) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} u_0(x) e^{-i\alpha x} dx$, $\hat{v}_0(\alpha) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} v_0(x) e^{-i\alpha x} dx$

$\hat{u}(\alpha, t)$ inconnu, $\hat{f}(\alpha, t)$, $\hat{u}_0(\alpha)$, $\hat{v}_0(\alpha)$ connus

On multiplie l'éq des ondes par $e^{-i\alpha x}$ $\alpha \in \mathbb{R}$ et on intègre entre $x = -\infty$ et $x = +\infty$:

$$\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} \left(\frac{\partial^2 u}{\partial t^2}(x, t) - \frac{\partial^2 u}{\partial x^2}(x, t) \right) e^{-i\alpha x} dx = \underbrace{\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f(x, t) e^{-i\alpha x} dx}_{\hat{f}(\alpha, t)}$$

$$\frac{\partial^2}{\partial t^2} \hat{u}(\alpha, t) - (i\alpha)^2 \hat{u}(\alpha, t) = \hat{f}(\alpha, t)$$

Condition initiale $\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} u(x, 0) e^{-i\alpha x} dx = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} u_0(x) e^{-i\alpha x} dx : \hat{u}(\alpha, 0) = \hat{u}_0(\alpha)$

de même on obtient

$$\hat{v}(\alpha, 0) = \hat{v}_0(\alpha)$$

$$\begin{aligned} f: \mathbb{R} &\rightarrow \mathbb{R} \\ x &\mapsto f(x) \\ \mathcal{F}(f')(\alpha) &= i\alpha \mathcal{F}(f)(\alpha) \\ \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f'(x) e^{-i\alpha x} dx \\ &= i\alpha \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f(x) e^{-i\alpha x} dx \end{aligned}$$

On obtient donc, pour chaque $\alpha \in \mathbb{R}$, l'eq. différentielle :

$$\begin{cases} \frac{\partial^2}{\partial t^2} \hat{u}(\alpha, t) + \alpha^2 \hat{u}(\alpha, t) = f(\alpha, t) & t > 0 \\ \hat{u}(\alpha, 0) = \hat{u}_0(\alpha) \\ \frac{\partial}{\partial t} \hat{u}(\alpha, 0) = \hat{v}_0(\alpha) \end{cases}$$

α est un paramètre.

$$\begin{cases} y''(t) + \alpha^2 y(t) = 0 \\ y(0) = 1 \\ y'(0) = 0 \\ y(t) = \cos(\alpha t) \end{cases}$$

On résout cette eq. diff en utilisant la tr. de Laplace :

$$\hat{u}(\alpha, t) = \hat{u}_0(\alpha) \cos(\alpha t) + \frac{\hat{v}_0(\alpha)}{\alpha} \sin(\alpha t) + \int_0^t f(\alpha, s) \sin(\alpha(t-s)) ds$$

On a la formule suivante : pour $a \in \mathbb{R}$

$$\mathcal{F}(f(x+a))(\alpha) = e^{i\alpha a} \mathcal{F}(f)(\alpha)$$

En effet : $\mathcal{F}(f(x+a))(\alpha) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f(x+a) e^{-i\alpha x} dx$

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} f(y) e^{-i\alpha(y-a)} dy = e^{i\alpha a} \mathcal{F}(f)(\alpha)$$

$$x+a=y \quad dx=dy \quad x=y-a$$

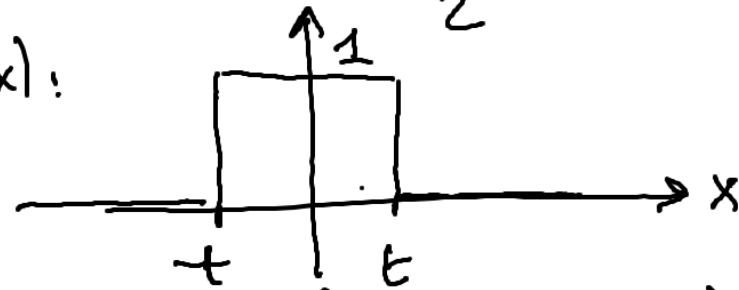
$$\mathcal{F}(f)(\alpha) \cdot \mathcal{F}(g)(\alpha) = \frac{1}{\sqrt{2\pi}} \mathcal{F}(f+g)(\alpha)$$

On a: $\hat{u}_0(\alpha) \cos(\alpha t) = \hat{u}_0(\alpha) \frac{e^{i\alpha t} + e^{-i\alpha t}}{2} = \mathcal{F}(u_0)(\alpha) \frac{e^{i\alpha t} + e^{-i\alpha t}}{2}$

$$= \frac{\mathcal{F}(u_0(x+t))(\alpha) + \mathcal{F}(u_0(x-t))(\alpha)}{2}$$

Transformée de Fourier inverse de $\hat{u}_0(\alpha) \cos \alpha t$: $\frac{u_0(x+t) + u_0(x-t)}{2}$

On a: $\mathcal{F}(g_t(x))(\alpha) = \sqrt{\frac{2}{\pi}} \frac{\sin(\alpha t)}{\alpha}$ où $g_t(x)$:



c.f tables

Donc $\hat{u}_0(\alpha) \frac{\sin(\alpha t)}{\alpha} = \mathcal{F}(u_0)(\alpha) \cdot \mathcal{F}\left(\sqrt{\frac{\pi}{2}} g_t(x)\right) = \frac{1}{\sqrt{2\pi}} \mathcal{F}\left(u_0 * \sqrt{\frac{\pi}{2}} g_t(x)\right)$

avec $(u_0 * \sqrt{\frac{\pi}{2}} g_t)(x) = \int_{-\infty}^{+\infty} u_0(y) \underbrace{\sqrt{\frac{\pi}{2}} g_t(x-y)}_{1 \text{ si } -t \leq x-y \leq t} dy$ (def de *: $(f * g)(x) = \int_{-\infty}^{+\infty} f(y) g(x-y) dy = \int_{-\infty}^{+\infty} f(x-y) g(y) dy$)

$$= \sqrt{\frac{\pi}{2}} \int_{x-t}^{x+t} u_0(y) dy$$

$$\left[\hat{u}_0(\alpha) \frac{\sin(\alpha t)}{\alpha} = \mathcal{F}\left(\frac{1}{2} \int_{x-t}^{x+t} u_0(y) dy \right) \right]$$

$$\begin{aligned}\hat{u}(\alpha, t) &= \hat{u}_0(\alpha) \cos(\alpha t) + \frac{\hat{v}_0(\alpha)}{\alpha} \sin(\alpha t) + \int_0^t \hat{f}(\alpha, s) \sin(\alpha(t-s)) ds \\ &= \mathcal{F}\left(\frac{u_0(x-t) + u_0(x+t)}{2}\right) + \mathcal{F}\left(\frac{1}{2} \int_{x-t}^{x+t} v_0(y) dy\right) + \int_0^t \hat{f}(\alpha, s) \sin(\alpha(t-s)) ds\end{aligned}$$

Transformée de Fourier inverse (formule d'inversion) $\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} \underbrace{\int_0^t \hat{f}(\alpha, s) \sin(\alpha(t-s)) ds}_{\text{}} e^{i\alpha x} d\alpha$

Solution $u(x, t) = \frac{1}{2}(u_0(x-t) + u_0(x+t)) + \frac{1}{2} \int_{x-t}^{x+t} v_0(y) dy +$