

Analysis III - 203(d)

Winter Semester 2024

Session Extra: December 29, 2025

Exercise 1 Recall the ReLU function

$$\text{ReLU}(x) = \max(0, x). \quad (1)$$

A shallow neural network with one output neuron, m internal neurons, and n input neurons has the general form

$$f(x_1, \dots, x_n) = \sum_{k=1}^m \text{ReLU} \left(\sum_{i=1}^n A_{ki} x_i + b_k \right) \quad (2)$$

where A_{ki} are weights and where b_k is a shift parameter.

- Most training algorithms require the gradient of this network. Compute the derivatives $\partial_i f$
- There is interest in training algorithms that use the Hessian matrix. Compute the partial derivatives $\partial_{ij}^2 f$.

Solution 1 We begin with the first derivatives. Obviously,

$$\partial_i f(x_1, \dots, x_n) = \sum_{k=1}^m \partial_i \text{ReLU} \left(\sum_{i=1}^n A_{ki} x_i + b_k \right) \quad (3)$$

The ReLU function can only be differentiated in the sense of distributions. Recalling the Heaviside function,

$$H(x) = \begin{cases} 1 & x > 0 \\ 0 & x \leq 0 \end{cases}, \quad (4)$$

it is now possible to write the derivative as

$$\partial_i f(x_1, \dots, x_n) = \sum_{k=1}^m H \left(\sum_{i=1}^n A_{ki} x_i + b_k \right) A_{ki} \quad (5)$$

$$= \sum_{k=1}^m A_{ki} H \left(\sum_{i=1}^n A_{ki} x_i + b_k \right). \quad (6)$$

This computes the first derivatives.¹

We continue with the second derivatives. Recall that the Heaviside function has a derivative in the sense of distributions, which is the Dirac Delta. Hence

$$\partial_{ji}^2 f(x_1, \dots, x_n) = \sum_{k=1}^m A_{ki} \partial_j H \left(\sum_{i=1}^n A_{ki} x_i + b_k \right) \quad (7)$$

$$= \sum_{k=1}^m A_{ki} A_{kj} \delta_0 \left(\sum_{i=1}^n A_{ki} x_i + b_k \right). \quad (8)$$

This computes the second derivatives.²

Exercise 2 In standard models of elasticity, a long straight beam of elastic material, such as wood or metal, can be modeled as a one-dimensional interval. When it is subject to an outside force f , such as gravity, then the deformation from the base is modeled by the beam equation

$$u''''(x) = f(x) \quad (9)$$

Here, the fourth derivative u'''' can be interpreted as the curvature of a curvature and f describes the direction (upwards, downwards) and magnitude of the force.

So-called non-local interactions are modeled via a convolutional term $k \star u$, where k indicates how parts of a beam are influenced by neighboring parts. With that in mind, we consider a generalized beam equation

$$u''''(x) + cu(x) + (k \star u)(x) = f(x). \quad (10)$$

This is a so-called integro-differential equation.

Compute the Fourier transform of this differential equation for general source terms, and write it down for the particular example

$$k(x) = e^{-|y|}, \quad f(x) = e^{-y^2}. \quad (11)$$

You are not expected to solve this equation.

Solution 2 The Fourier transform of this equation reads

$$(i\alpha)^4 \hat{u}(\alpha) + c\hat{u}(\alpha) + \sqrt{2\pi} \hat{k}(\alpha) \hat{u}(\alpha) = \hat{f}(\alpha). \quad (12)$$

¹Remark: the derivative does not have a meaningful value if one of the arguments of the Heaviside function is zero (or close to zero within the range of rounding errors). When training a neural network via gradient descent, this is “justified” by the assumption that these arguments being close to zero is very unlikely to happen in practice.

²Remark: the situation here is even worse than with the first derivatives. The Dirac Delta is zero everywhere except at the origin, and it equals a pointmass at the origin. Correspondingly, second-order training algorithms, such as, e.g., Newton’s method, are not well-defined for such a network. This is a possible incentive to replace ReLU by other activation functions, such as $S(x) = \ln(1 + e^x)$.

We simplify this to:

$$\alpha^4 \hat{u}(\alpha) + c \hat{u}(\alpha) + \sqrt{2\pi} \hat{k}(\alpha) \hat{u}(\alpha) = \hat{f}(\alpha). \quad (13)$$

We isolate $\hat{u}(\alpha)$, giving us:

$$\hat{u}(\alpha) = \frac{\hat{f}(\alpha)}{\alpha^4 + c + \sqrt{2\pi} \hat{k}(\alpha)}. \quad (14)$$

In the particular case of

$$k(x) = e^{-|y|}, \quad f(x) = e^{-y^2}, \quad (15)$$

one finds

$$\hat{k}(x) = \sqrt{\frac{2}{\pi}} \frac{1}{1 + \alpha^2}, \quad \hat{f}(x) = \frac{1}{\sqrt{2}} e^{-\alpha^2/4}. \quad (16)$$

One can simplify

$$\hat{u}(\alpha) = \frac{\hat{f}(\alpha)}{\alpha^4 + c + \sqrt{2\pi} \hat{k}(\alpha)} = \frac{\hat{f}(\alpha)}{\alpha^4 + c + \frac{2}{1+\alpha^2}} = \underbrace{\frac{1 + \alpha^2}{(1 + \alpha^2)(\alpha^4 + c) + 2}}_{=:\hat{h}(\alpha)} \hat{f}(\alpha). \quad (17)$$

The solution of the ODE can be found if we know the inverse Fourier transform $h(x)$ of the factor $\hat{h}(\alpha)$. Then the product of $\hat{h}(\alpha)\hat{f}(\alpha)$ can be transformed into the convolution $\sqrt{2\pi}h \star f$. Finding the inverse Fourier transform of such a factor will be possible with techniques of complex analysis, to be discussed next semester.

Exercise 3 Solve the integro-differential equation

$$9u(x) + 2 \int_{-\infty}^{+\infty} (u''(t) - 4u(t)) e^{-2|x-t|} dt = \frac{1}{x^2 + 1}.$$

Solution 3 We use the Fourier transform. First, we transform the equation. Beginning with the source term, we use the convolution theorem and transform the derivatives:

$$\mathfrak{F} \left[\frac{1}{x^2 + 1} \right] = 9\mathfrak{F}[u] + 2\mathfrak{F} \left[\int_{-\infty}^{+\infty} (u''(t) - 4u(t)) e^{-2|x-t|} dt \right] \quad (18)$$

$$= 9\mathfrak{F}[u] + 2\sqrt{2\pi} \mathfrak{F}[u'' - 4u] \cdot \mathfrak{F}[e^{-2|x|}] \quad (19)$$

$$= 9\mathfrak{F}[u] + 2(-\alpha^2 \mathfrak{F}[u] - 4\mathfrak{F}[u]) \cdot \mathfrak{F}[e^{-2|x|}]. \quad (20)$$

For convenience, we use the \hat{u} notation for the Fourier transform and use the Fourier transform table:

$$\sqrt{\frac{\pi}{2}}e^{-|\alpha|} = 9\hat{u}(\alpha) - 2\sqrt{2\pi}(\alpha^2\hat{u}(\alpha) + 4\hat{u}(\alpha)) \cdot 2\sqrt{\frac{2}{\pi}}\frac{1}{4+\alpha^2} \quad (21)$$

$$= 9\hat{u}(\alpha) - 2\sqrt{2\pi}\hat{u}(\alpha)(\alpha^2 + 4) \cdot 2\sqrt{\frac{2}{\pi}}\frac{1}{4+\alpha^2} \quad (22)$$

$$= 9\hat{u}(\alpha) - 2\sqrt{2\pi}\hat{u}(\alpha) \cdot 2\sqrt{\frac{2}{\pi}} \quad (23)$$

$$= 9\hat{u}(\alpha) - 8\hat{u}(\alpha) = \hat{u}(\alpha). \quad (24)$$

We transform back and obtain the solution

$$u(x) = \frac{1}{1+x^2}. \quad (25)$$

Exercise 4 Solve the integral equation

$$u(t) + \lambda \int_0^{+\infty} e^{-|y|} u(t-y) dy = e^{-|t|}.$$

Solution 4 We observe that this is a convolutional integral equation

$$u(t) + \lambda(K \star u)(t) = f(t)$$

with

$$K(t) = \begin{cases} e^{-t} & t > 0 \\ 0 & \text{otherwise,} \end{cases} \quad f(t) = e^{-|t|}.$$

We look up the Fourier transforms:

$$\hat{K}(\alpha) = \frac{1}{\sqrt{2\pi}} \frac{1}{1+i\alpha}, \quad \hat{f}(\alpha) = \sqrt{\frac{2}{\pi}} \frac{1}{1+\alpha^2}$$

We Fourier transform the equation,

$$\hat{u}(\alpha) + \lambda\sqrt{2\pi}\hat{K}(\alpha)\hat{u}(\alpha) = \hat{f}(\alpha),$$

and isolating \hat{u} we find

$$\hat{u}(\alpha) = \frac{\hat{f}(\alpha)}{1 + \lambda\sqrt{2\pi}\hat{K}(\alpha)},$$

Let us simplify this. Using $1 + \alpha^2 = 1 - i^2\alpha^2 = (1+i\alpha)(1-i\alpha)$, we obtain

$$\hat{u}(\alpha) = \frac{\hat{f}(\alpha)}{1 + \lambda\frac{1}{1+i\alpha}}$$

$$= \frac{\hat{f}(\alpha)(1+i\alpha)}{(1+i\alpha)+\lambda} = \sqrt{\frac{2}{\pi}} \frac{\frac{1+i\alpha}{1+\alpha^2}}{(1+\lambda)+i\alpha} = \sqrt{\frac{2}{\pi}} \frac{1}{1-i\alpha} \cdot \frac{1}{(1+\lambda)+i\alpha}.$$

We already know that

$$\hat{v}(\alpha) = \frac{1}{(1+\lambda)+i\alpha} \implies v(t) = \sqrt{2\pi} \begin{cases} e^{-(1+\lambda)t} & t > 0 \\ 0 & \text{otherwise.} \end{cases}$$

Using the modulation identity, we find

$$\hat{w}(\alpha) = \frac{1}{1-i\alpha} \implies w(t) = \sqrt{2\pi} \begin{cases} e^t & t < 0 \\ 0 & \text{otherwise.} \end{cases}$$

Hence

$$\hat{u}(\alpha) = \sqrt{\frac{2}{\pi}} \hat{v}(\alpha) \hat{w}(\alpha)$$

which means

$$u(t) = \sqrt{\frac{2}{\pi}} \frac{1}{\sqrt{2\pi}} (v \star w)(t)$$

We calculate

$$u(t) = \sqrt{\frac{2}{\pi}} \frac{1}{\sqrt{2\pi}} \int_0^{+\infty} \sqrt{2\pi} e^{-(1+\lambda)y} w(t-y) dy = \sqrt{\frac{2}{\pi}} \int_0^{+\infty} e^{-(1+\lambda)y} w(t-y) dy.$$

The function w vanishes for $t-y > 0$, that is, $t > y$. Hence we only need to consider the case $t < y$. For negative $t < 0$, the condition $t < y$ is already implied by $0 < y$, and thus

$$u(t) = \sqrt{\frac{2}{\pi}} \sqrt{2\pi} \int_0^{\infty} e^{-(1+\lambda)y} e^{t-y} dy \quad (26)$$

$$= 2e^t \int_0^{\infty} e^{-(2+\lambda)y} dy = 2e^t \int_0^{\infty} e^{-(2+\lambda)y} dy = 2e^t \frac{1}{2+\lambda}. \quad (27)$$

For non-negative $t \geq 0$, we further simplify

$$\begin{aligned} u(t) &= \sqrt{\frac{2}{\pi}} \sqrt{2\pi} \int_t^{\infty} e^{-(1+\lambda)y} e^{t-y} dy \\ &= 2e^t \int_t^{\infty} e^{-(2+\lambda)y} dy = 2 \frac{e^t}{-(2+\lambda)} \left(0 - e^{-(2+\lambda)t} \right) = 2 \frac{e^t}{2+\lambda} e^{-(2+\lambda)t} = 2 \frac{e^{-(1+\lambda)t}}{2+\lambda}. \end{aligned}$$

For completeness only, we check that this satisfies the equation. When $t < 0$, then

$$2 \frac{e^t}{2+\lambda} + \lambda \int_0^{+\infty} e^{-y} 2 \frac{e^{t-y}}{2+\lambda} dy \quad (28)$$

$$= \frac{2}{2+\lambda}e^t + \frac{2\lambda}{2+\lambda}e^t \int_0^{+\infty} e^{-2y} dy = \frac{2}{2+\lambda}e^t + \frac{\lambda}{2+\lambda}e^t = e^t = e^{-|t|}. \quad (29)$$

When $t > 0$, we split the integral to switch between the two different regimes of u . We find

$$\begin{aligned} & u(t) + \lambda \int_0^t e^{-y} u(t-y) dy + \lambda \int_t^\infty e^{-y} u(t-y) dy \\ &= u(t) + \lambda \int_0^t e^{-y} 2 \frac{e^{-(1+\lambda)(t-y)}}{2+\lambda} dy + \lambda \int_t^\infty e^{-y} 2e^{t-y} \frac{1}{2+\lambda} dy \\ &= u(t) + \frac{2}{2+\lambda} e^{-(1+\lambda)t} \lambda \int_0^t e^{-y} e^{(1+\lambda)y} dy + \frac{2}{2+\lambda} e^t \lambda \int_t^\infty e^{-y} e^{-y} dy \\ &= u(t) + u(t) \lambda \int_0^t e^{\lambda y} dy + \frac{2}{2+\lambda} e^t \lambda \int_t^\infty e^{-2y} dy \\ &= u(t) + u(t) (e^{\lambda t} - 1) + \frac{2\lambda}{2+\lambda} e^t \frac{e^{-2t}}{2} \\ &= u(t) + u(t) (e^{\lambda t} - 1) + \frac{\lambda}{2+\lambda} e^{-t}. \end{aligned}$$

We plug in the definition of $u(t)$ for $t \geq 0$ and observe

$$\begin{aligned} & u(t) + u(t) (e^{\lambda t} - 1) + \frac{\lambda}{2+\lambda} e^{-t} \\ &= \frac{2}{2+\lambda} e^{-(1+\lambda)t} + \frac{2}{2+\lambda} e^{-(1+\lambda)t} (e^{\lambda t} - 1) + \frac{\lambda}{2+\lambda} e^{-t} \\ &= \frac{2}{2+\lambda} e^{-(1+\lambda)t} + \frac{2}{2+\lambda} (e^{-t} - e^{-(1+\lambda)t}) + \frac{\lambda}{2+\lambda} e^{-t} = e^{-t}. \end{aligned}$$

This completes the verification of the solution.