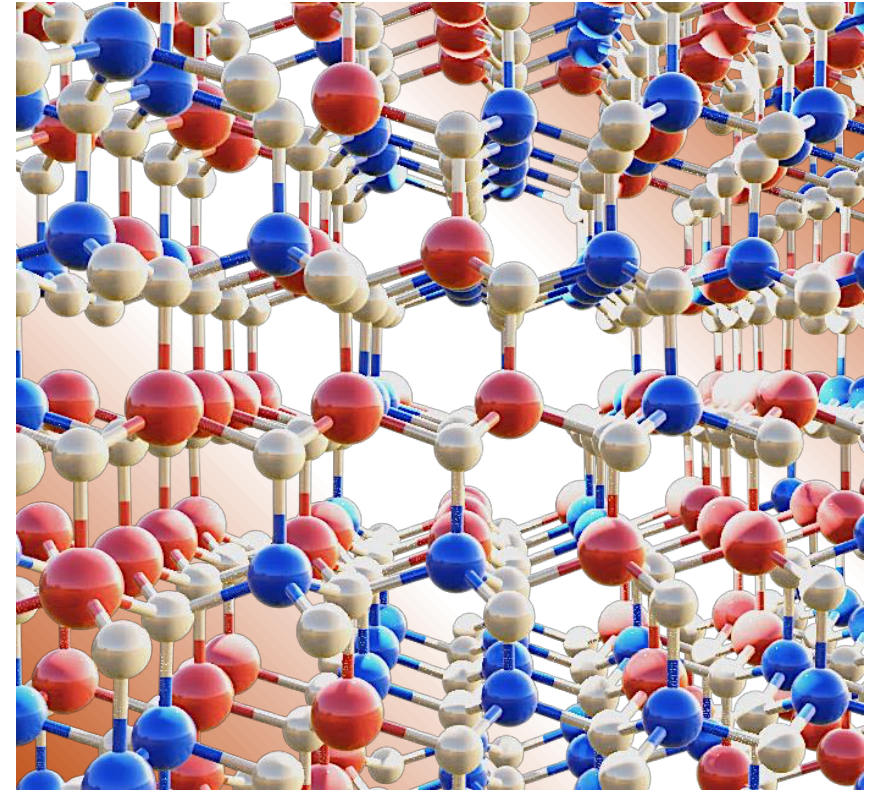


# The physics of nitride semiconductors

## Lecture 1: Overview of quantum mechanics, and introduction to band theory



Emmanouil (Manos) Kioupakis<sup>1,2</sup>

<sup>1</sup> Professor, Materials Science and Engineering, University of Michigan

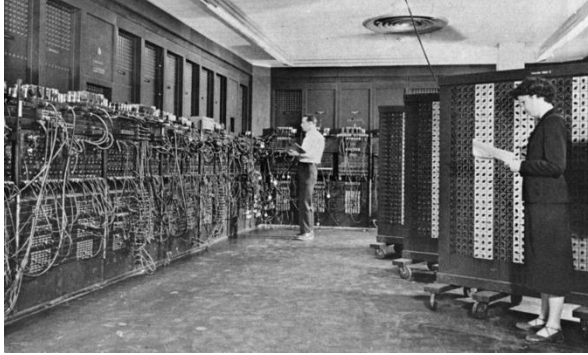
<sup>2</sup> Visiting Professor, IMX and IEM, STI, EPFL



EPFL



# Semiconductors and their applications



ENIAC: first programmable, electronic, general-purpose digital computer, 1945-1955. Based on vacuum tubes. ~ 500 floating-point operations per second (Flops)



Efficient light emitters

Solar cells



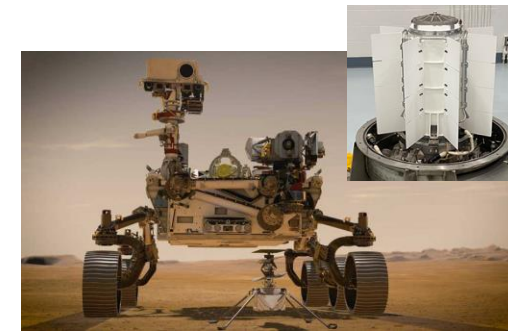
Indoor agriculture



Today's smartphones: ~ 100 **G**Flops



Today's state of the art supercomputers:  
~4 **P**Flops



Thermoelectric electricity production from heat



# Nitrides: important for energy efficient lighting



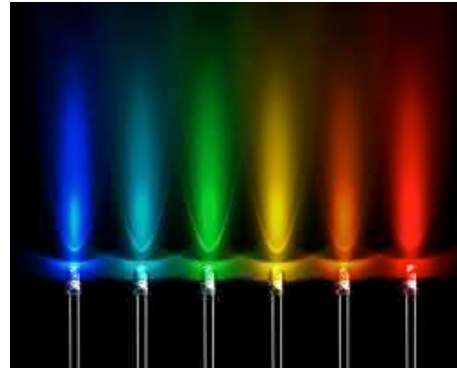
Incandescent



Fluorescent



LED



- To **manufacture modern nitride electronic and optoelectronic devices**, we need to understand the behavior of electrons in nitrides and how they interact with external inputs (voltage, light, heat, pressure).
- Also, we need to **understand the interactions of light and electrons with materials** in order to develop modern characterization tools to study materials at the atomic scale.
- The behavior of electrons at the atomic scale and the interactions of light with materials cannot be described with the laws of classical mechanics (i.e., Newton's laws).
- To understand the electrons in electronic materials and devices, we need to apply the principles of **quantum mechanics**



# Wave properties of matter

## Classical physics:

Particles: localized, indivisible

Waves: extended, divisible

	Particle	Wave
Electron	✓	✗
Light	✗	✓

## Quantum Physics:

	Particle	Wave
Electron	✓	?
Light	?	✓

# Wave properties of matter

## Classical physics:

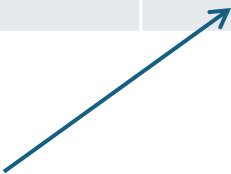
Particles: localized, indivisible

Waves: extended, divisible

	Particle	Wave
Electron	✓	✗
Light	✗	✓

## Quantum Physics:

	Particle	Wave
Electron	✓	?
Light	✓	✓



Evidence: Photoelectric effect, blackbody radiation, Compton scattering, X-rays

Photon energy  $E = hf$

Photon momentum  $p = h/\lambda$

$h$  = Planck's constant =  $6.626 \times 10^{-34}$  J s

# Wave properties of matter

## Classical physics:

Particles: localized, indivisible

Waves: extended, divisible

	Particle	Wave
Electron	✓	✗
Light	✗	✓

**Electrons also have wave-like properties!**

Wavelength  $\lambda = h/p$ , frequency  $f = E/h$

## Quantum Physics:

	Particle	Wave
Electron	✓	✓
Light	✓	✓

Evidence: Photoelectric effect, blackbody radiation, Compton scattering, X-rays

Photon energy  $E = hf$

Photon momentum  $p = h/\lambda$

$h$  = Planck's constant =  $6.626 \times 10^{-34}$  J s

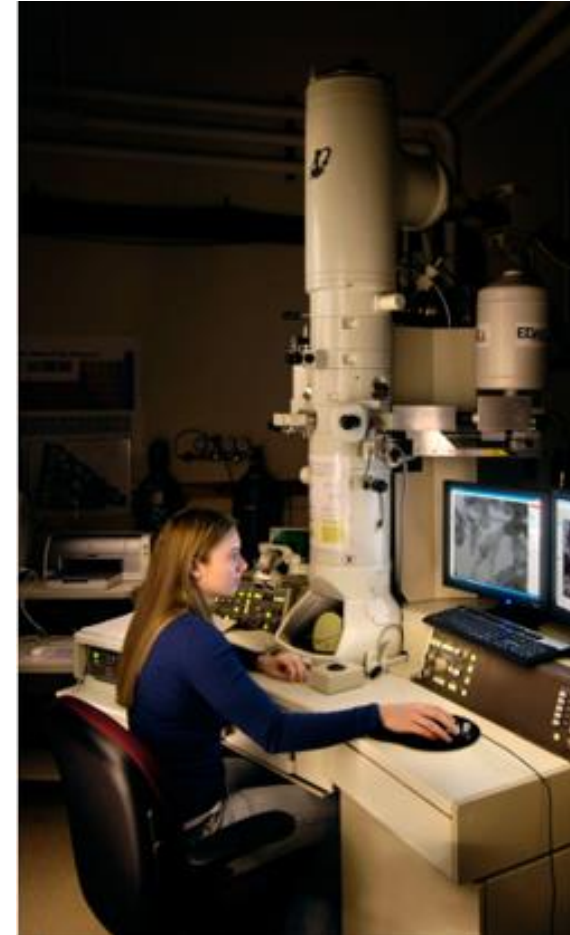
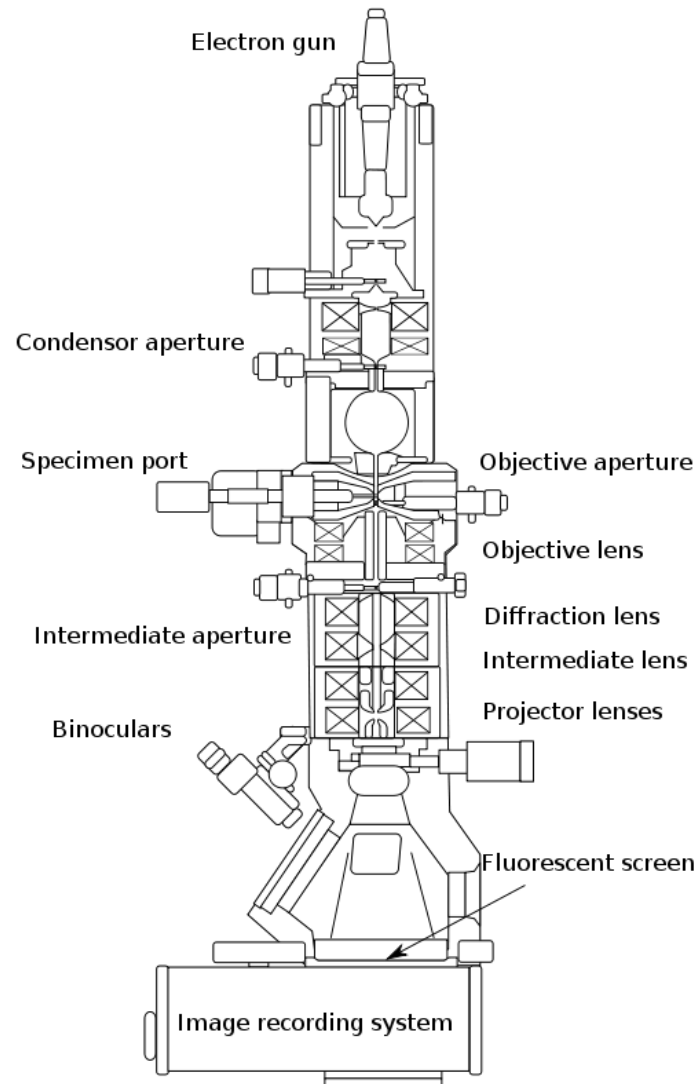


# Proof and application: the electron microscope

Transmission electron microscope (TEM)

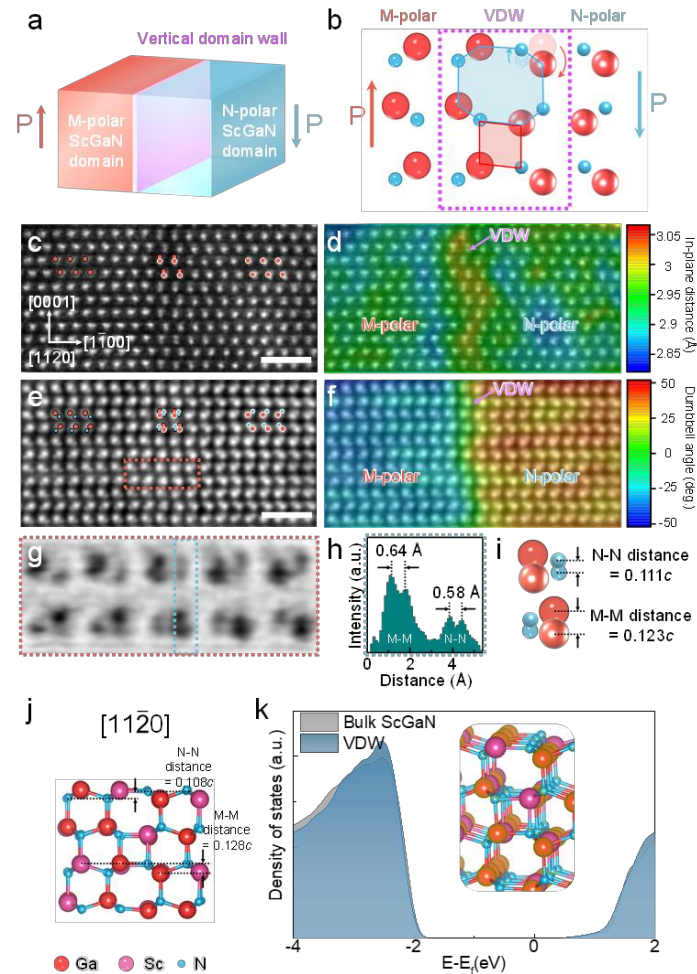
Uses electron waves to image materials

Wavelength of electrons shorter than wavelength of light, can image materials with sub-atomic resolution

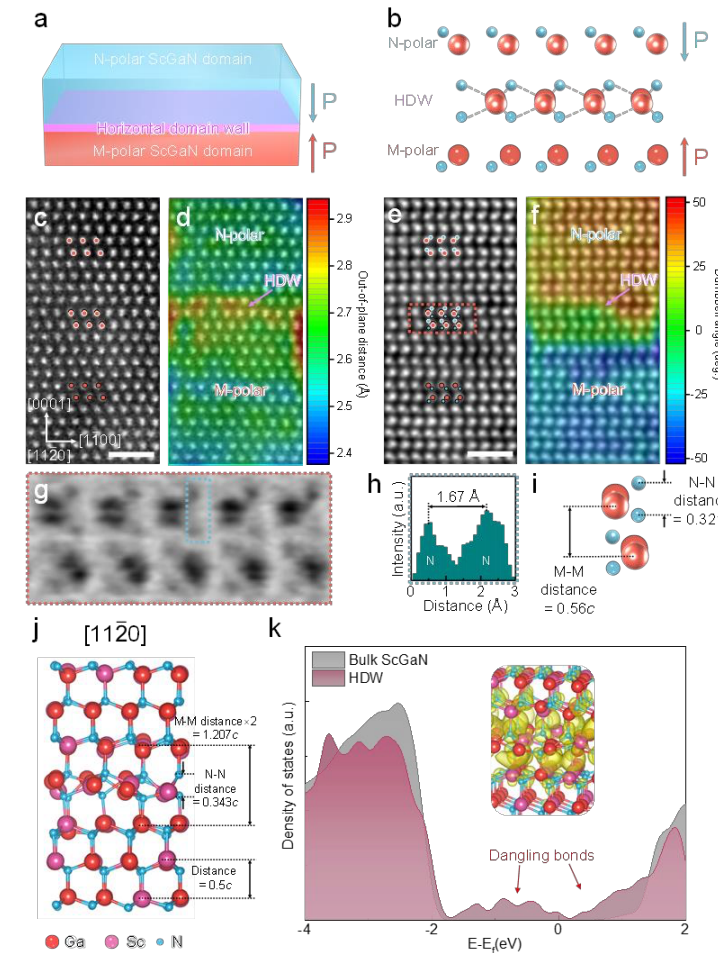


# Application: seeing domain walls in ferroelectric ScGaN

Vertical charge-neutral domain walls



Horizontal charged domain walls, stabilized by dangling bonds



Images from D. Wang, ..., E. Kioupakis, Z. Mi, “Electric-Field-Induced Domain Walls in Wurtzite Ferroelectrics”, in press (<https://arxiv.org/abs/2312.08645v2>)

# Particle waves

The electron has energy  $E$  and momentum  $p$ . Its wavelength  $\lambda$  and frequency  $f$  are given by:

$$E = hf \rightarrow f = E/h$$

$$p = \frac{h}{\lambda} \rightarrow \lambda = \frac{h}{p}$$

$\Psi$  = **wave function** = wave associated with motion of electron  
(just like the electric field is for photons)

$$P(x, t) = |\Psi(x, t)|^2$$

electron at (x,t)

probabilities

= probability of finding an

→ Quantum mechanics describes

$$\int_{-\infty}^{\infty} P(x, t) dx = \int_{-\infty}^{\infty} |\Psi(x, t)|^2 dx = 1$$

**Normalization  
condition**



# Quantum Mechanics

It is a rigorous formalism to describe the motion of probability waves (wave functions) associated with the motion of particles (e.g., electrons) under the influence of an external potential (e.g., atomic nuclei or other electrons).

Also, it is a recipe of how to derive all measurable quantities of a particle (position, momentum, energy, interaction with light, etc.) from the knowledge of the wave function.

Key finding (1926): **Schrödinger's equation**, the equation that determines the wave function given the external potential.



$$i\hbar \frac{\partial}{\partial t} \Psi(x, t) = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \Psi(x, t) + V(x, t) \Psi(x, t)$$

# Schrödinger's equation

$$i\hbar \frac{\partial}{\partial t} \Psi(x, t) = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \Psi(x, t) + V(x, t) \Psi(x, t)$$

What is  $V(x,t)$ ? A: the potential energy that electrons experience (e.g., attraction by the nuclei, electric field in electronic devices, oscillating electric field in optical devices)

What is  $i$  ? A: imaginary number  $i^2 = -1$ . Wave functions are complex! Cannot be measured

To fully determine  $\Psi(x,t)$  we also need:

- *Boundary conditions*: Value of  $\Psi$  at boundary (determined by the problem we need to solve)
- *Initial condition*: Value of  $\Psi$  at  $t=0$

Once we find  $\Psi(x,t)$ , we can determine all measurable properties of the system.



# Time-independent Schrödinger equation

If  $V(x,t)=V(x)$  only (independent of  $t$ ) then write  $\Psi$  as:

$$\Psi(x, t) = \psi(x) f(t)$$

Separate time-*dependent* Schrödinger into two equations:

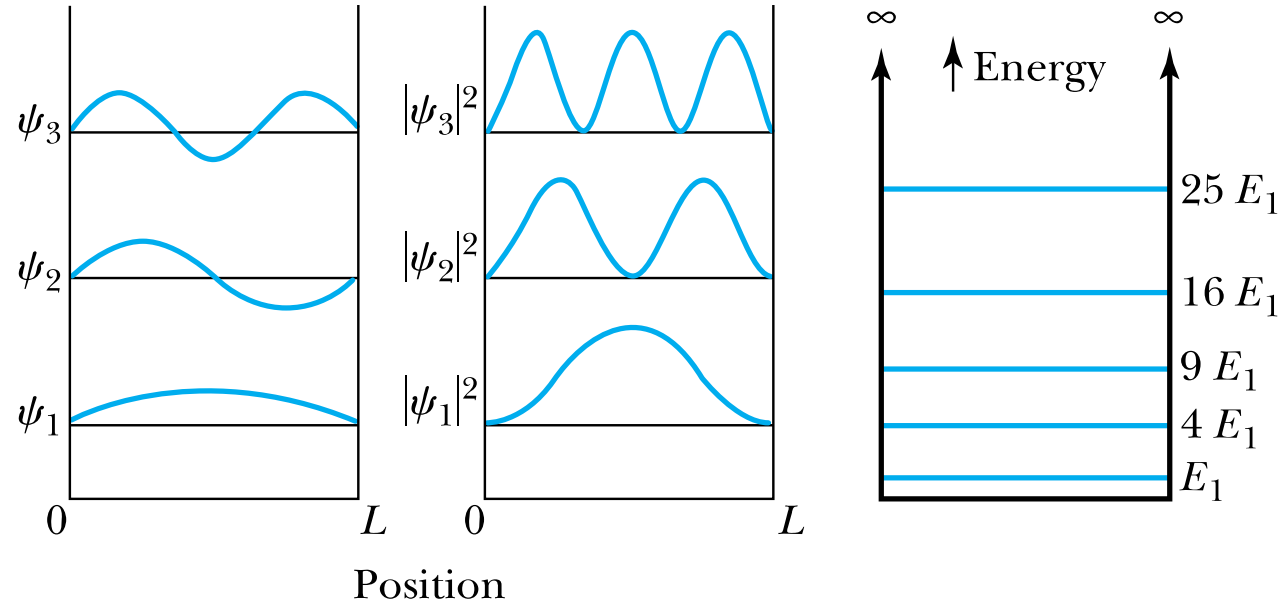
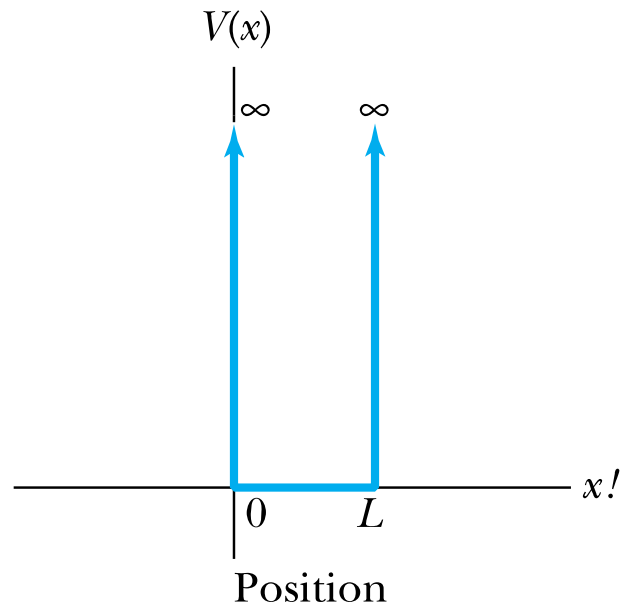
$$\left\{ \begin{array}{l} f(t) = e^{-i\omega t} \\ \boxed{-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + V(x)\psi(x) = E\psi(x)} \end{array} \right.$$

**Time-independent Schrödinger equation (+ boundary conditions).**

**Energy eigenvalue problem**, solve to find energy and  $\psi(x)$ .

# Example 1: particle in an infinite square well

$$V(x) = \begin{cases} \infty & x \leq 0, x \geq L \\ 0 & 0 < x < L \end{cases}$$

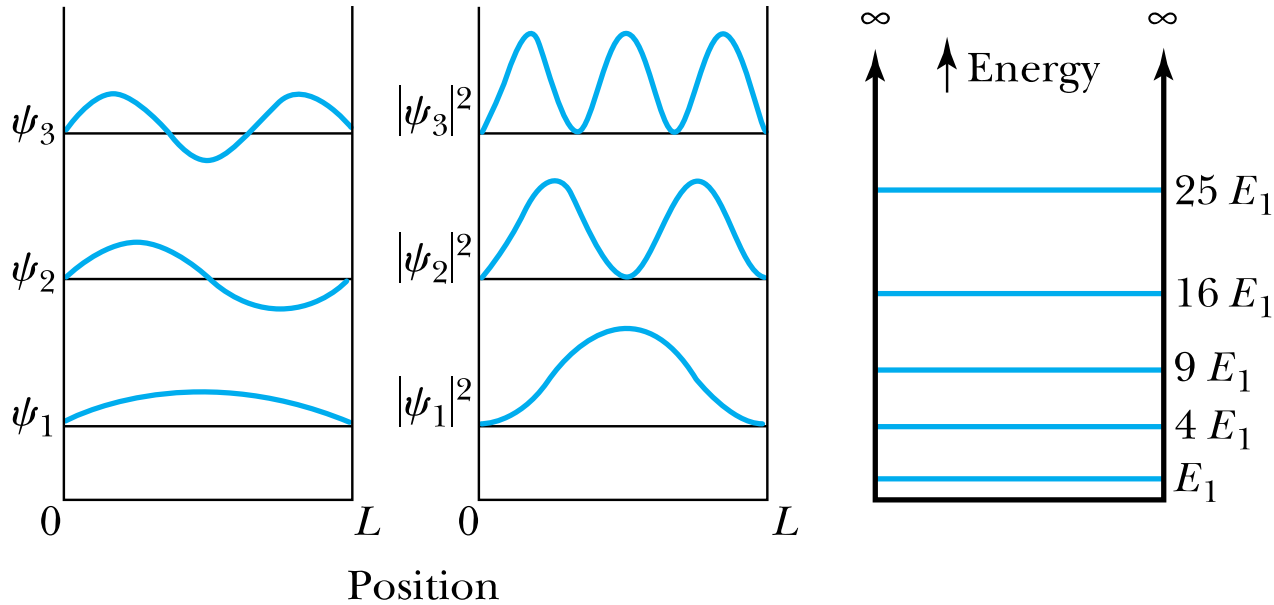


$$\psi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right) \quad (n = 1, 2, 3, \dots)$$

$$E_n = n^2 \frac{\pi^2 \hbar^2}{2mL^2} \quad (n = 1, 2, 3, \dots)$$

**Figure 6.2** Infinite square-well potential. The potential is  $V = \infty$  everywhere except the region  $0 < x < L$ , where  $V = 0$ .

# Some key observations:



$$\psi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right) \quad (n = 1, 2, 3, \dots)$$

$$E_n = n^2 \frac{\pi^2 \hbar^2}{2mL^2} \quad (n = 1, 2, 3, \dots)$$

- Ground state energy ( $n=1$ ) is not zero → **zero point motion**. Electron is not at rest even at  $T = 0$  K.
- Ground-state energy (which is kinetic energy only) increases as  $+1/L^2$  with decreasing  $L$  → Increases more than attraction by the nucleus in atoms ( $\sim -1/L$ ), therefore **electrons cannot fall into the nucleus**.
- Wave functions are oscillatory. **Ground state has no nodes** (no zeros). One more node (zero of wave function) for each increasing value of  $n$ .
- Wave functions are **orthogonal** to each other:
 
$$\int_{-\infty}^{\infty} \psi_m^*(x) \psi_n(x) dx \quad \begin{array}{l} \text{If } m=n \rightarrow \text{Integral} = 1 \\ \text{If } m \neq n \rightarrow \text{Integral} = 0 \end{array}$$
- Wave functions are **complete**: any wave function can be written as a linear combination of solutions to Schrödinger.

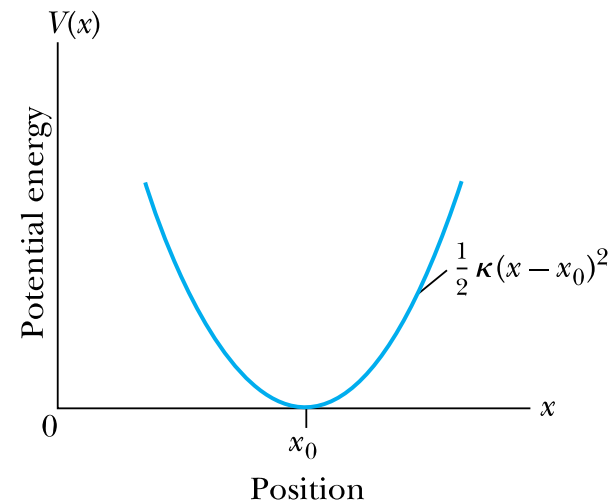
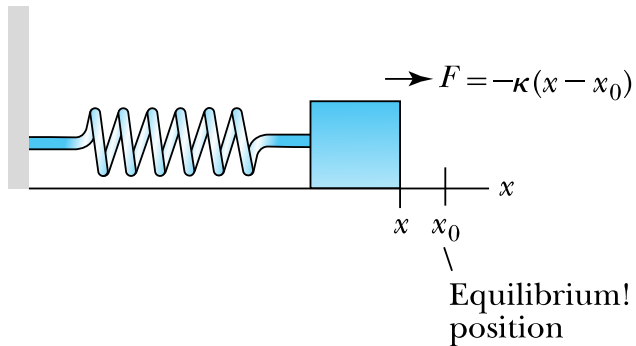
$$f(x) = \sum_{n=1}^{\infty} c_n \psi_n(x)$$

$$c_n = \int_{-\infty}^{\infty} \psi_n^*(x) f(x) dx$$



# Example 2: Quantum harmonic oscillator

$$V(x) = \frac{1}{2}kx^2 = \frac{1}{2}m\omega^2x^2, \quad \omega = \sqrt{\frac{k}{m}}$$

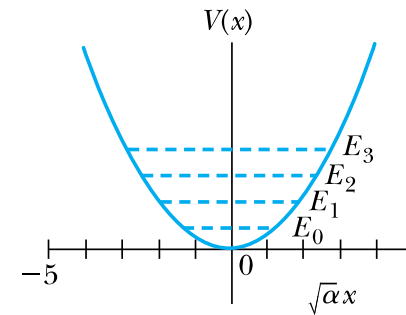


Solution: 
$$E_n = \left(n + \frac{1}{2}\right) \hbar\omega$$

$$= \left(n + \frac{1}{2}\right) \hbar\sqrt{\frac{k}{m}}, \quad n = 0, 1, 2, 3, \dots$$

Wave functions

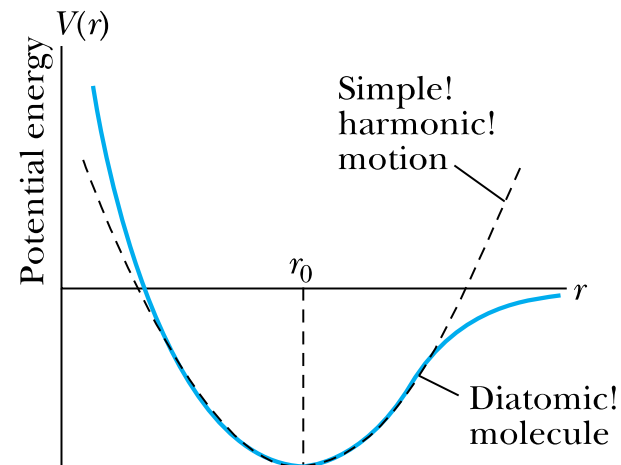
$$\begin{aligned}\psi_3(x) &= \left(\frac{\alpha}{\pi}\right)^{1/4} \frac{1}{\sqrt{3}} (\sqrt{\alpha}x) (2\alpha x^2 - 3) e^{-\alpha x^2/2} \\ \psi_2(x) &= \left(\frac{\alpha}{\pi}\right)^{1/4} \frac{1}{\sqrt{2}} (2\alpha x^2 - 1) e^{-\alpha x^2/2} \\ \psi_1(x) &= \left(\frac{\alpha}{\pi}\right)^{1/4} \sqrt{2\alpha} x e^{-\alpha x^2/2} \\ \psi_0(x) &= \left(\frac{\alpha}{\pi}\right)^{1/4} e^{-\alpha x^2/2}\end{aligned}$$



Why is it important?

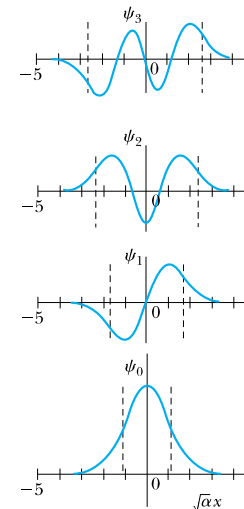
**Every** potential is harmonic near a minimum.

E.g., atomic vibrations around equilibrium position.



Simple!  
harmonic!  
motion

Diatomic!  
molecule



Energy levels are equally spaced.

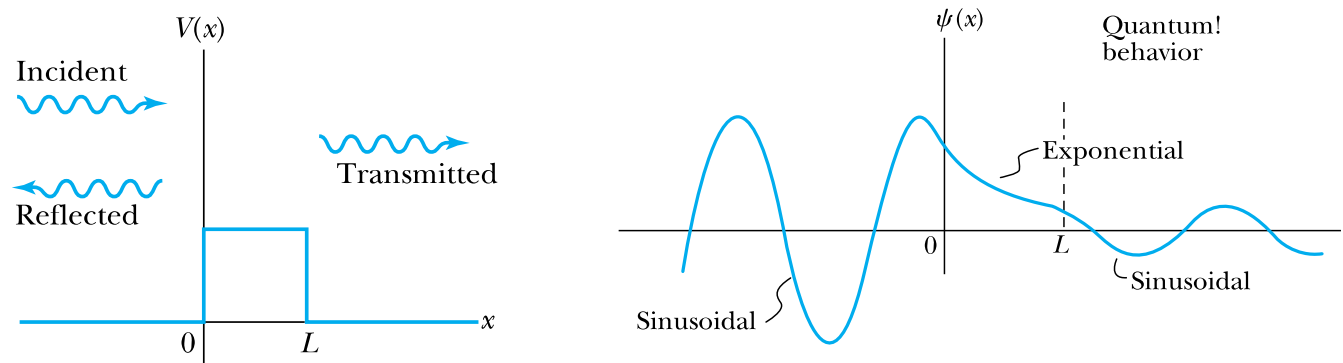
Quantum of vibrational energy  
= **phonon**

$$\Delta E = \hbar\omega = hf$$

Energy is absorbed and emitted by vibrating atoms in discrete amounts.

# Example 3: Quantum tunneling

- Even if  $E < V_0$ , there is a probability that the particle can make it across the barrier → **Tunneling (transmission probability  $T \neq 0$ )**
- Even if  $E > V_0$ , there is a probability that the particle will be reflected at the barrier (**transmission probability  $T \neq 1$** )

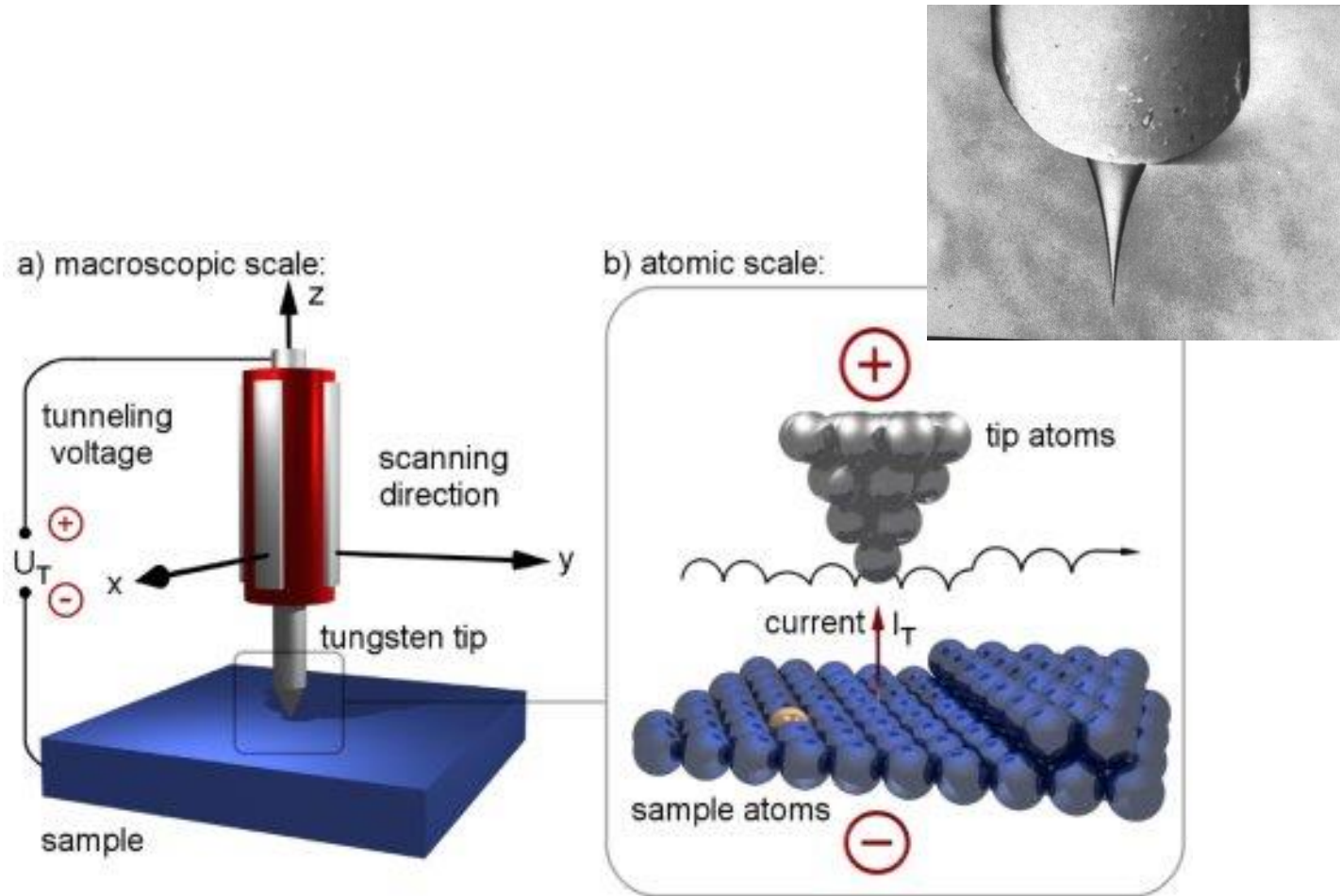


What is the probability  $T$  that a particle with energy  $E$  will be transmitted across a barrier of width  $L$  and height  $V_0$ ?

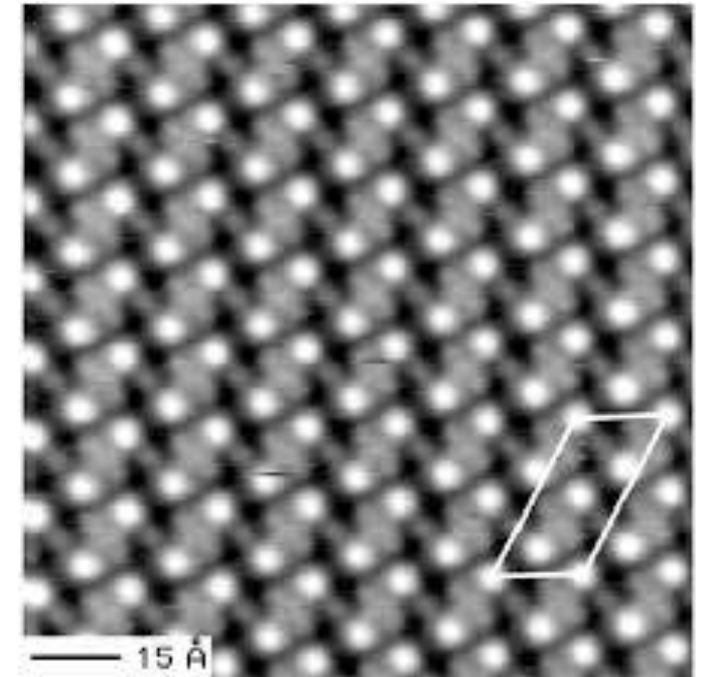
Answer:

$$\begin{aligned} \text{For } E < V_0: T &= \left[ 1 + \frac{V_0^2 \sinh^2(\kappa L)}{4E(V_0 - E)} \right]^{-1}, \quad \kappa \equiv \sqrt{\frac{2m(V_0 - E)}{\hbar^2}} \\ \text{For } E > V_0: T &= \left[ 1 + \frac{V_0^2 \sin^2(k_2 L)}{4E(E - V_0)} \right]^{-1}, \quad k_2 \equiv \sqrt{\frac{2m(E - V_0)}{\hbar^2}} \end{aligned}$$

# Application: scanning tunneling microscope



Surface reconstruction of GaN



Smith et al., MRS Online Proceedings Library, **482**, 363 (1997)

# Schrödinger's equation for the hydrogen atom

Electron under the influence of the attractive potential by the proton:

$$V(r) = -\frac{e^2}{4\pi\epsilon_0 r}$$

3D Schrödinger's equation:

$$\psi = \psi(x, y, z)$$

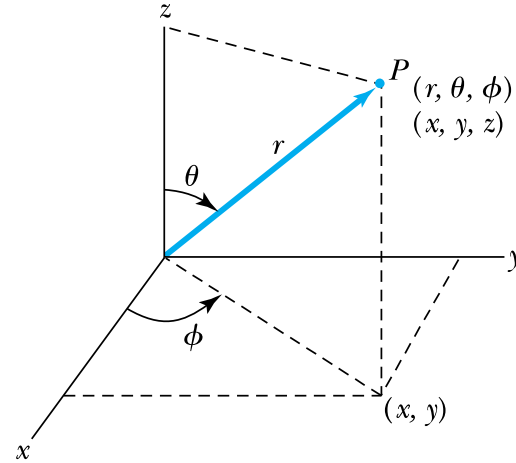
$$-\frac{\hbar^2}{2m} \left( \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} \right) + V\psi = E\psi$$

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

$$-\frac{\hbar^2}{2m} \nabla^2 \psi + V\psi = E\psi$$

# Schrödinger's equation in spherical coordinates

$$\psi = \psi(r, \theta, \phi)$$



$$r = \sqrt{x^2 + y^2 + z^2}$$

$$\theta = \cos^{-1} \frac{z}{r} \quad (\text{Polar angle})$$

$$\phi = \tan^{-1} \frac{y}{x} \quad (\text{Azimuthal angle})$$

$$x = r \sin \theta \cos \phi$$

$$y = r \sin \theta \sin \phi$$

$$z = r \cos \theta$$

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \psi}{\partial \phi^2} +$$

$$+ \frac{2\mu}{\hbar^2} (E - V) \psi = 0$$

(7.3)



# Three equations, one for each coordinate

$$\psi(r, \theta, \phi) = R(r)f(\theta)g(\phi)$$

## Radial equation

$$\frac{1}{r^2} \frac{d}{dr} \left( r^2 \frac{dR}{dr} \right) + \frac{2\mu}{\hbar^2} \left[ E - V - \frac{\hbar^2}{2\mu} \frac{\ell(\ell + 1)}{r^2} \right] R = 0$$

$$n = 1, 2, 3, 4, \dots$$

$$l = 0, 1, 2, 3, \dots, n - 1$$

$$m_l = -l, -l + 1, \dots, -1, 0, 1, \dots, l - 1, l$$

## Angular equation

$$\frac{1}{\sin \theta} \frac{d}{d\theta} \left( \sin \theta \frac{df}{d\theta} \right) + \left[ \ell(\ell + 1) - \frac{m_\ell^2}{\sin^2 \theta} \right] f = 0$$

Also: spin quantum number  $m_s = +1/2, -1/2$

## Azimuthal equation

$$\frac{d^2 g}{d\phi^2} = -m_\ell^2 g$$

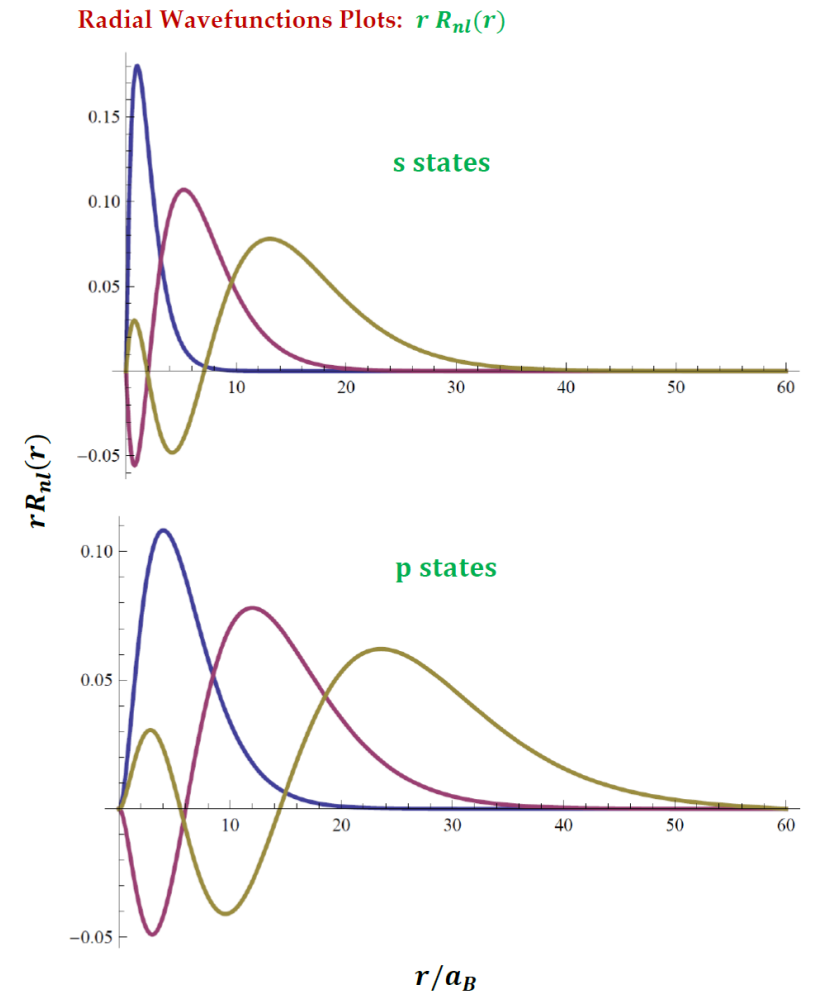
Only radial equation depends on the potential  $V(r)$

→ Angular and azimuthal equations same for all atoms

# Solutions to radial equation

**Table 7.1** Hydrogen Atom Radial Wave Functions

$n$	$\ell$	$R_{n\ell}(r)$
1	0	$\frac{2}{(a_0)^{3/2}} e^{-r/a_0}$
2	0	$\left(2 - \frac{r}{a_0}\right) \frac{e^{-r/2a_0}}{(2a_0)^{3/2}}$
2	1	$\frac{r}{a_0} \frac{e^{-r/2a_0}}{\sqrt{3}(2a_0)^{3/2}}$
3	0	$\frac{1}{(a_0)^{3/2}} \frac{2}{81\sqrt{3}} \left(27 - 18\frac{r}{a_0} + 2\frac{r^2}{a_0^2}\right) e^{-r/3a_0}$
3	1	$\frac{1}{(a_0)^{3/2}} \frac{4}{81\sqrt{6}} \left(6 - \frac{r}{a_0}\right) \frac{r}{a_0} e^{-r/3a_0}$
3	2	$\frac{1}{(a_0)^{3/2}} \frac{4}{81\sqrt{30}} \frac{r^2}{a_0^2} e^{-r/3a_0}$



# Solutions to angular equation

**Table 7.2** Normalized Spherical Harmonics  $Y(\theta, \phi)$

$\ell$	$m_\ell$	$Y_{\ell m_\ell}$
0	0	$\frac{1}{2\sqrt{\pi}}$
1	0	$\frac{1}{2}\sqrt{\frac{3}{\pi}} \cos \theta$
1	$\pm 1$	$\mp \frac{1}{2}\sqrt{\frac{3}{2\pi}} \sin \theta e^{\pm i\phi}$
2	0	$\frac{1}{4}\sqrt{\frac{5}{\pi}} (3 \cos^2 \theta - 1)$
2	$\pm 1$	$\mp \frac{1}{2}\sqrt{\frac{15}{2\pi}} \sin \theta \cos \theta e^{\pm i\phi}$
2	$\pm 2$	$\frac{1}{4}\sqrt{\frac{15}{2\pi}} \sin^2 \theta e^{\pm 2i\phi}$
3	0	$\frac{1}{4}\sqrt{\frac{7}{\pi}} (5 \cos^3 \theta - 3 \cos \theta)$
3	$\pm 1$	$\mp \frac{1}{8}\sqrt{\frac{21}{\pi}} \sin \theta (5 \cos^2 \theta - 1) e^{\pm i\phi}$
3	$\pm 2$	$\frac{1}{4}\sqrt{\frac{105}{2\pi}} \sin^2 \theta \cos \theta e^{\pm 2i\phi}$
3	$\pm 3$	$\mp \frac{1}{8}\sqrt{\frac{35}{\pi}} \sin^3 \theta e^{\pm 3i\phi}$



# The Orbitron gallery of atomic orbitals



1s



2s



3s



4s



5s



6s



7s



# Quantum mechanics of multi-electron atoms

Schrödinger's equation for Helium atom:

2 electrons, one set of coordinates ( $r_1, r_2$ ) for each electron

$\psi = \psi(\vec{r}_1, \vec{r}_2)$       Wave function of 2 electrons

$$\hat{H}\psi = E\psi$$

$$\hat{H} = \underbrace{-\frac{\hbar^2}{2m}\nabla_1^2}_{\text{K.E. of e \#1}} - \underbrace{\frac{\hbar^2}{2m}\nabla_2^2}_{\text{K.E. of e \#2}} - \underbrace{\frac{1}{4\pi\epsilon_0}\frac{2e^2}{r_1}}_{\text{P.E. of e \#1}} - \underbrace{\frac{1}{4\pi\epsilon_0}\frac{2e^2}{r_2}}_{\text{P.E. of e \#2}} + \underbrace{\frac{1}{4\pi\epsilon_0}\frac{e^2}{|\vec{r}_1 - \vec{r}_2|}}_{\text{e-e interaction}}$$

Impossible to solve analytically due to interaction term.

We can solve it numerically using computers.

However, we can learn a lot by analyzing the wave functions of H

# Wave functions of multielectron atoms

The wave functions of multielectron atoms have the same functional form as the wave functions of the hydrogen atom:

$$\psi_{n,l,m_l}(r, \theta, \phi) = R_{n,l}(r)Y_{l,m_l}(\theta, \phi)$$

The possible values of the integers  $n, l, m_l$  are **exactly the same**.

The spherical harmonics  $Y_{lm_l}$  are **exactly the same** as in the case of hydrogen (and in general for any spherically symmetric potential).

The radial wave functions  $R(r)$  and the energies of the orbitals differ from H. They depend on the atomic number of each element and they are determined by the interplay of electron-nucleus and electron-electron interactions. (we will see examples later)



# Pauli exclusion principle

To understand multielectron atoms, we need one more quantum principle: **Pauli exclusion principle**.

**No two electrons can have the same set of quantum numbers.**

For atoms, **no two electrons can have the same set of**  $(n, l, m_l, m_s)$

→ Explains organization of electrons in the energy levels of atoms.

The principle is more general: applies to all particles with **half-integer spin = fermions**.

As opposed to particles with **integer spin = bosons** (not subject to Pauli exclusion).

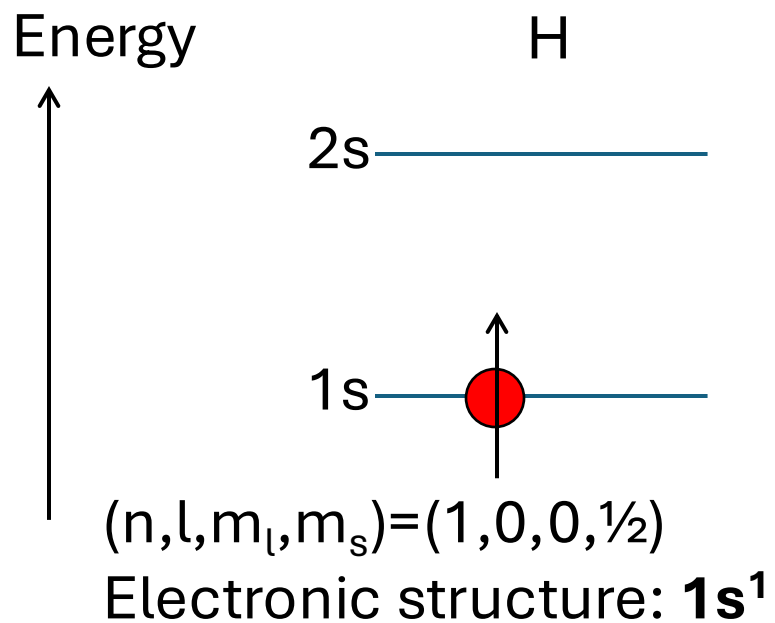
Examples of fermions and bosons?

Fermions: electron, proton, neutron, ...

Bosons: photon, He<sup>4</sup> nucleus, Higgs, ...

Origin: **indistinguishability** of quantum particles (see Gibbs paradox)

# Electronic structure of H and He



Labeling atomic shells:

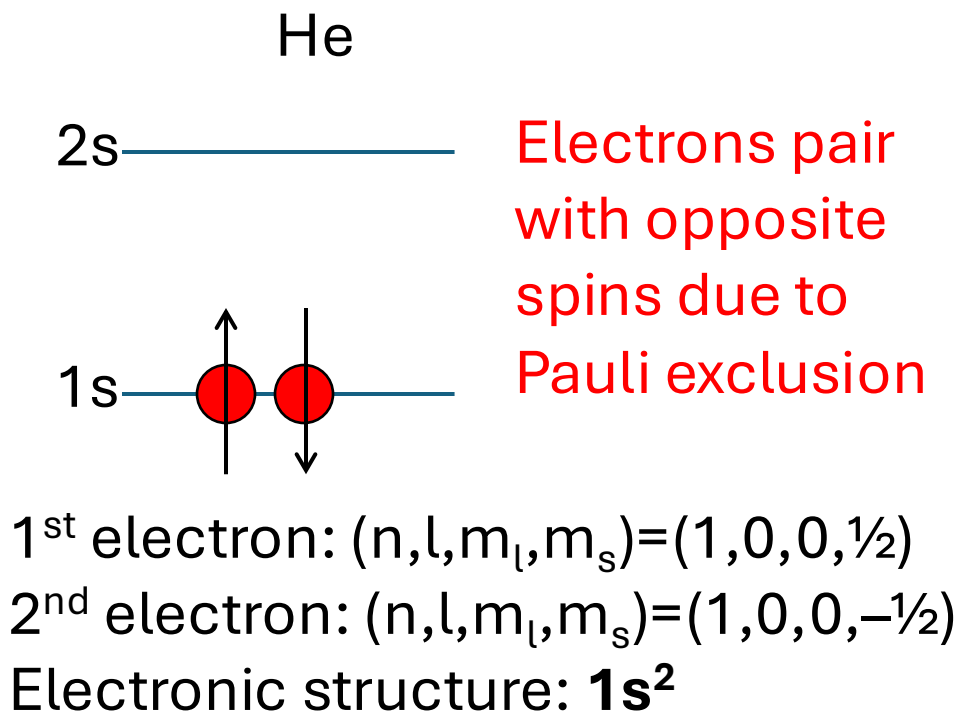
$n=1 \rightarrow K$

$n=2 \rightarrow L$

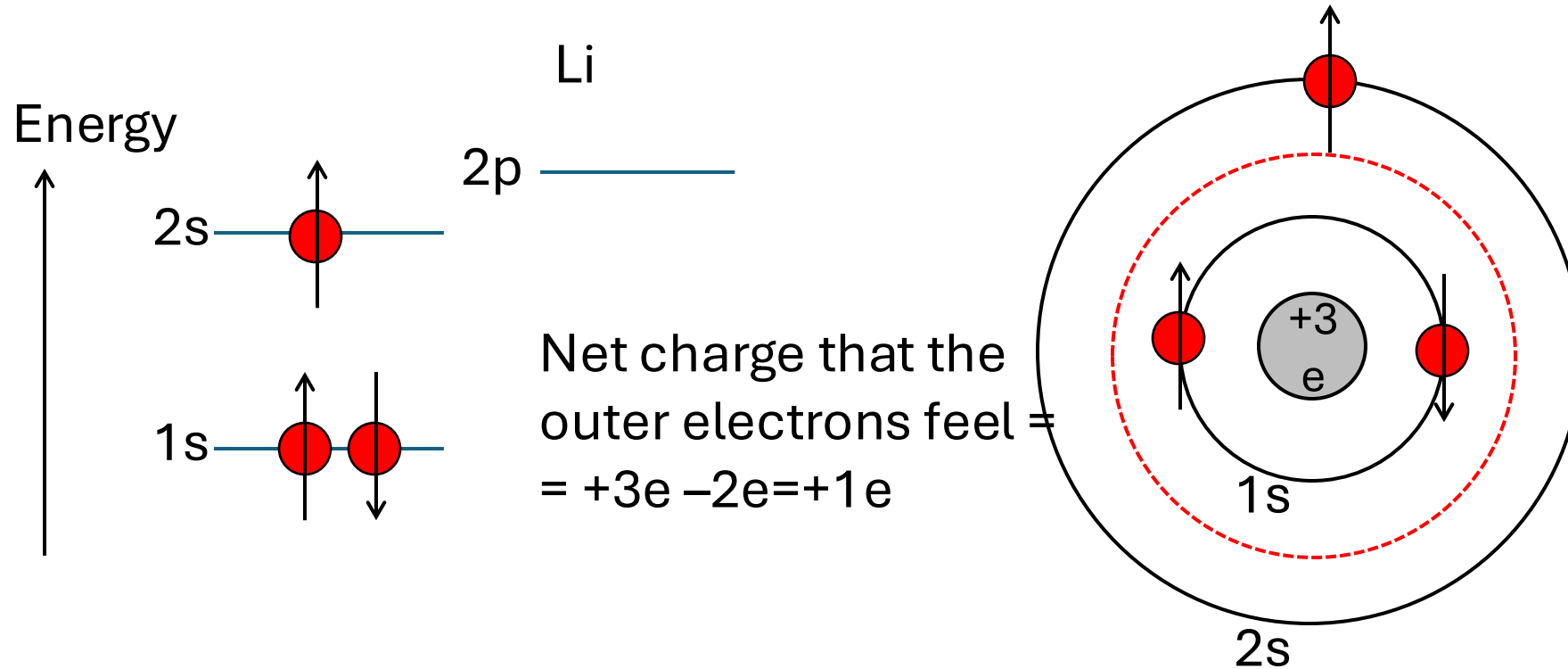
$n=3 \rightarrow M$

$n=4 \rightarrow N$

...



# Electronic structure of Li



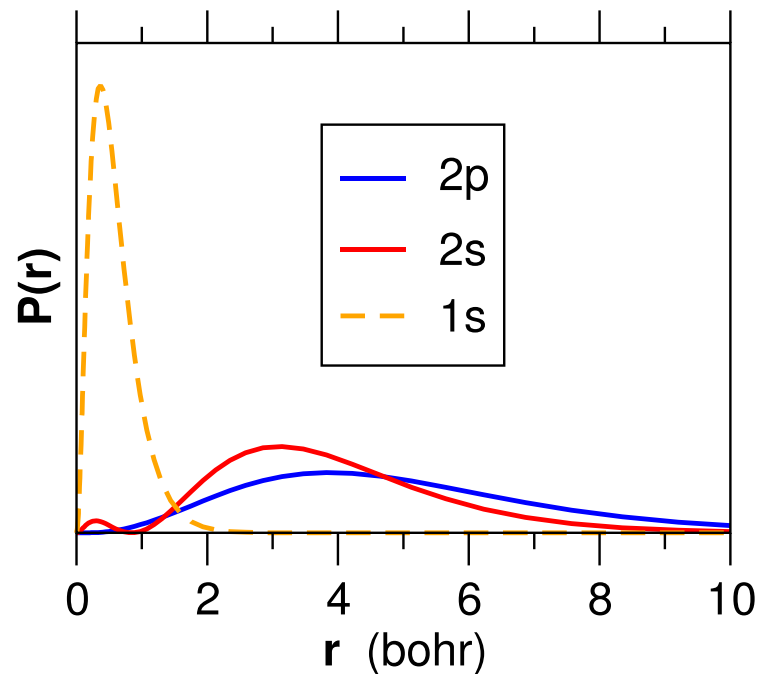
In H, the 2s and 2p states have the same energy

An electron in H is attracted by the bare Coulomb potential of the nucleus.

An outer electron in Li (or other multi-electron atoms) feels a reduced (“screened”) attractive potential due to the repulsion by the inner electrons.

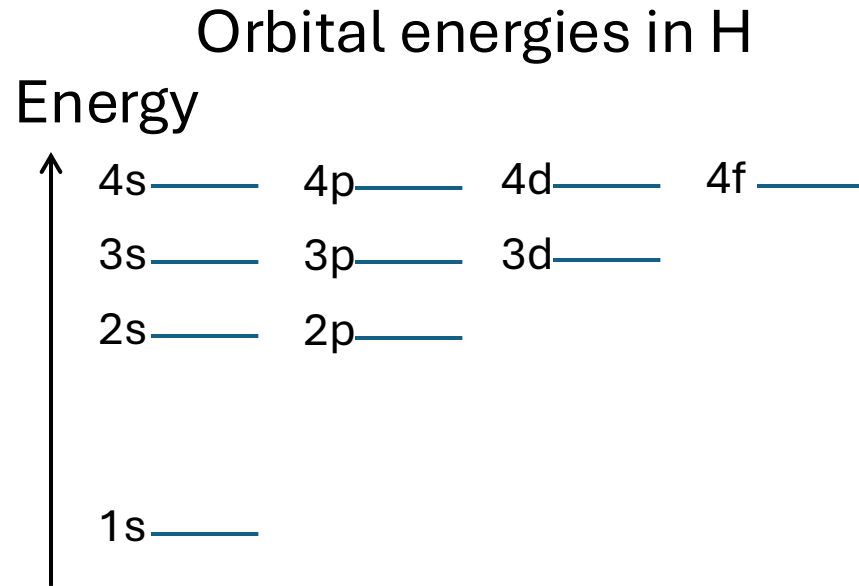
# Radial distribution function in Li

Calculated data for the radial probability distribution function  $P(r)$ :

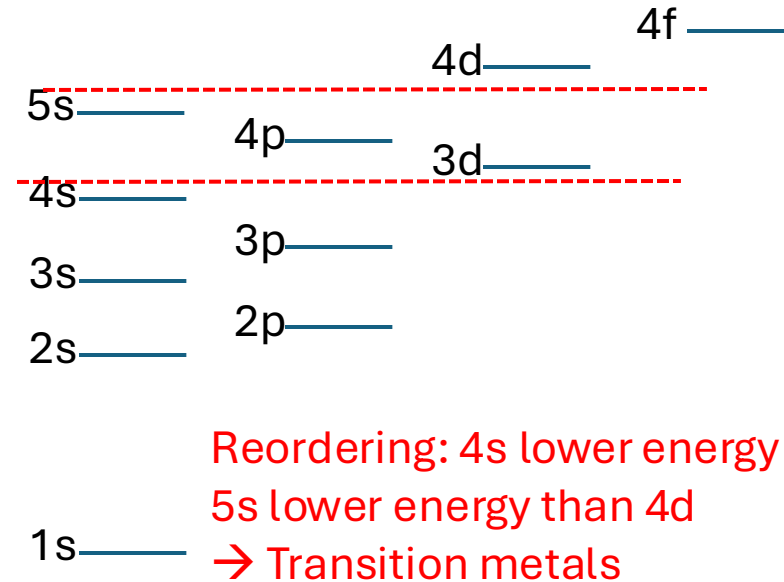


The electron in the 2s state penetrates more through the cloud of 1s electrons and feels a stronger attraction by the nucleus. 2s electrons have lower potential energy.

# Orbital energies in multielectron atoms



Orbital energies in multielectron atoms



Filling the states: How many electrons in each subshell?

s:  $l=0 \rightarrow m_l = \{0\} \times 2$  (for spin)  $\rightarrow$  2 electrons

p:  $l=1 \rightarrow m_l = \{-1, 0, +1\} \times 2$  (for spin)  $\rightarrow$  6 electrons

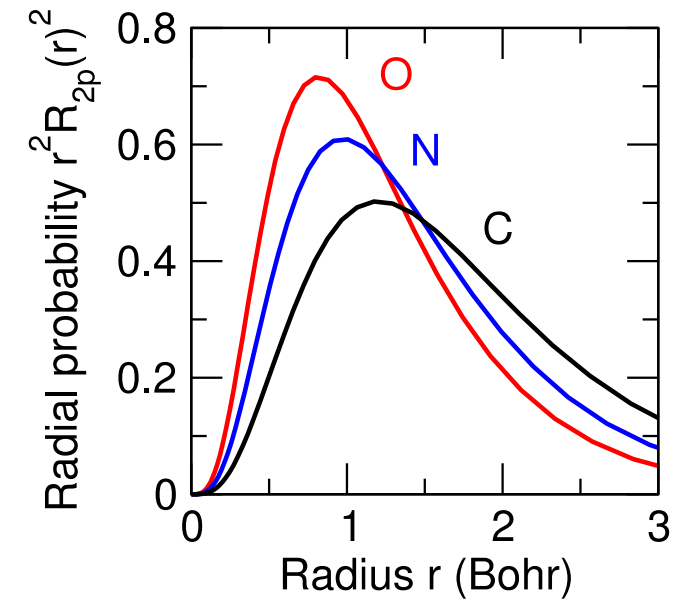
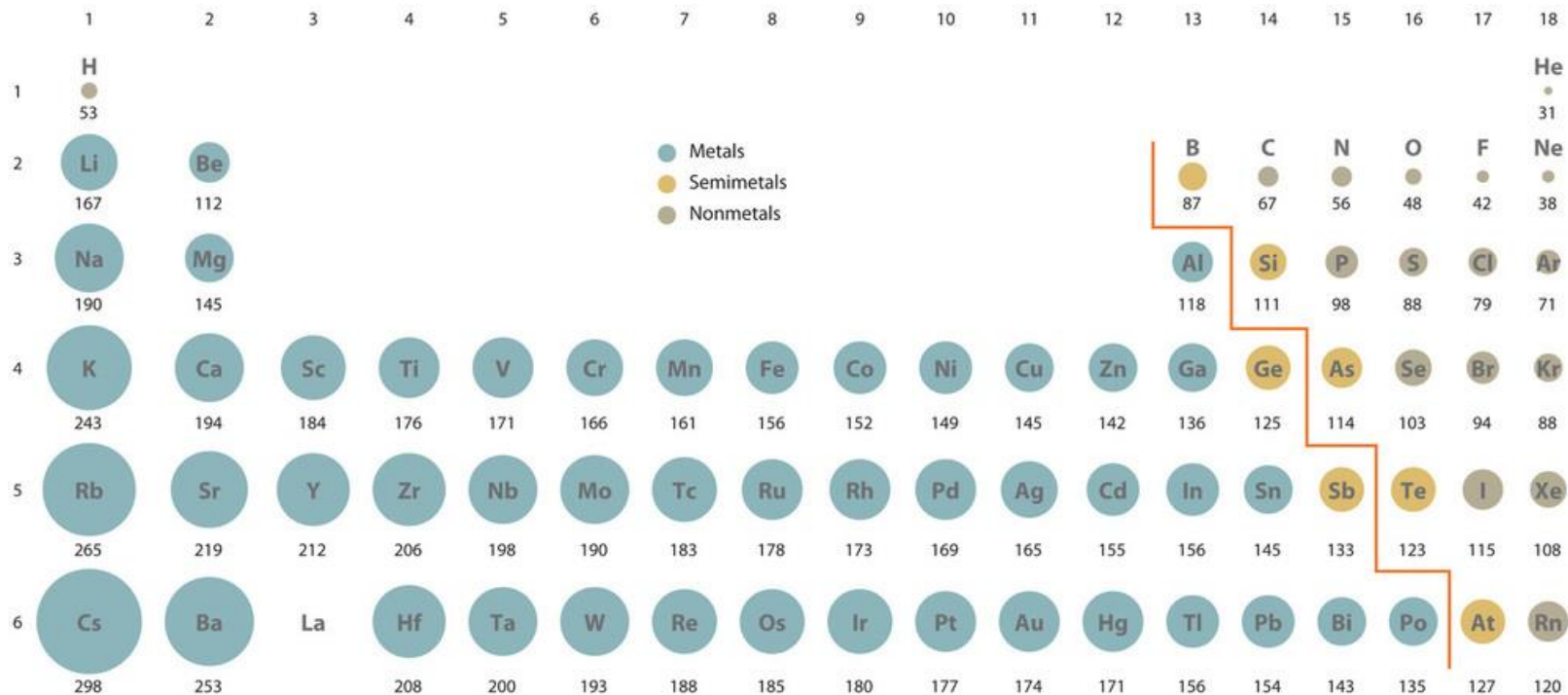
d:  $l=2 \rightarrow m_l = \{-2, -1, 0, +1, +2\} \times 2$  (for spin)  $\rightarrow$  10 electrons

f:  $l=3 \rightarrow m_l = \{-3, -2, -1, 0, +1, +2, +3\} \times 2$  (for spin)  $\rightarrow$  14 electrons





# Size of the atoms



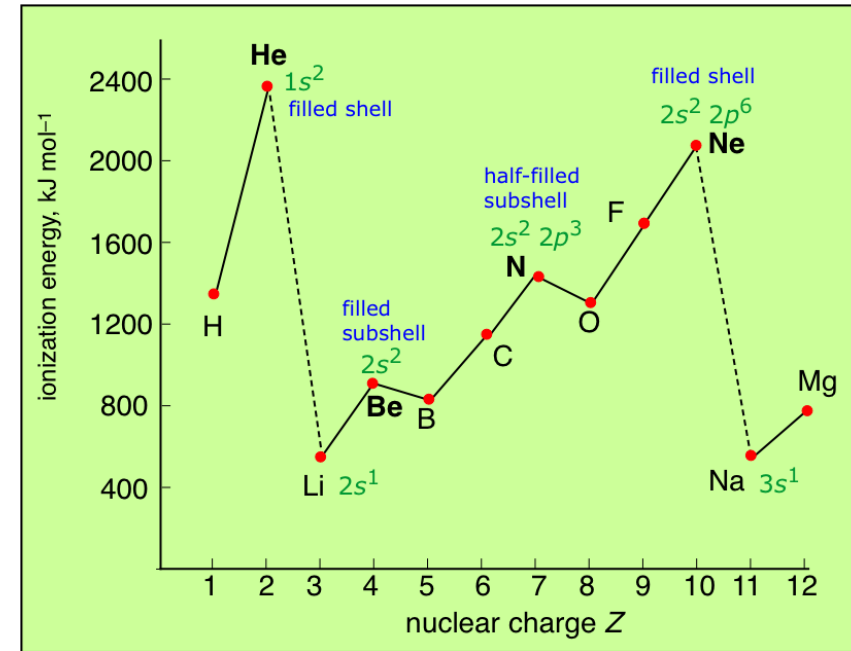
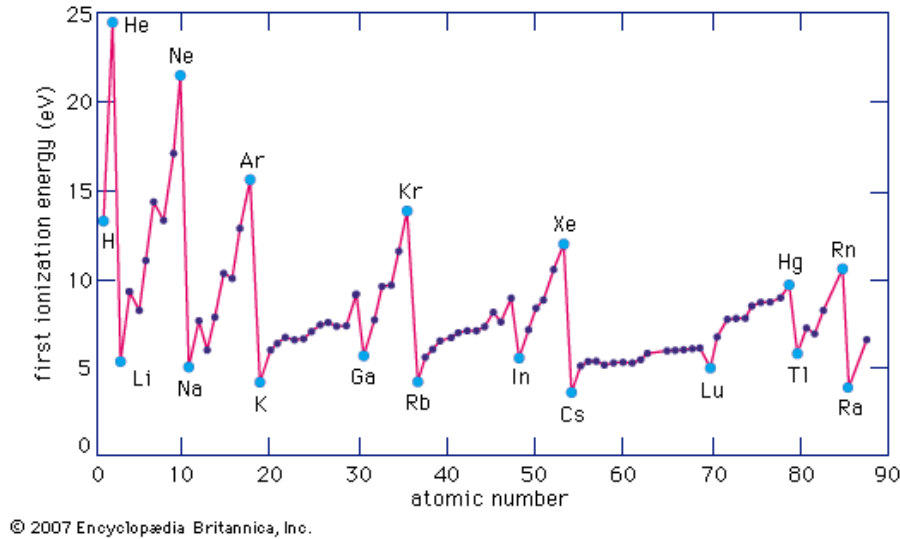
Why does the size of the atoms decrease from left to right?

A: increasing  $Z \rightarrow$  stronger attraction to nucleus.

Why does the size increase from top to bottom?

A: outer shells are more weakly bound to nucleus (screening by inner electrons)

# Ionization energy




Ionization energy = energy needed to remove outermost electron  
 = measure of stability of e in atoms.  
 What trends do you see?

- Large for inert gases, low for alkalis and alkaline earths
- Li → Ne: overall increase for increasing  $Z$
- Be → B: occupation of p orbitals. N → O: double occupation of p orbital

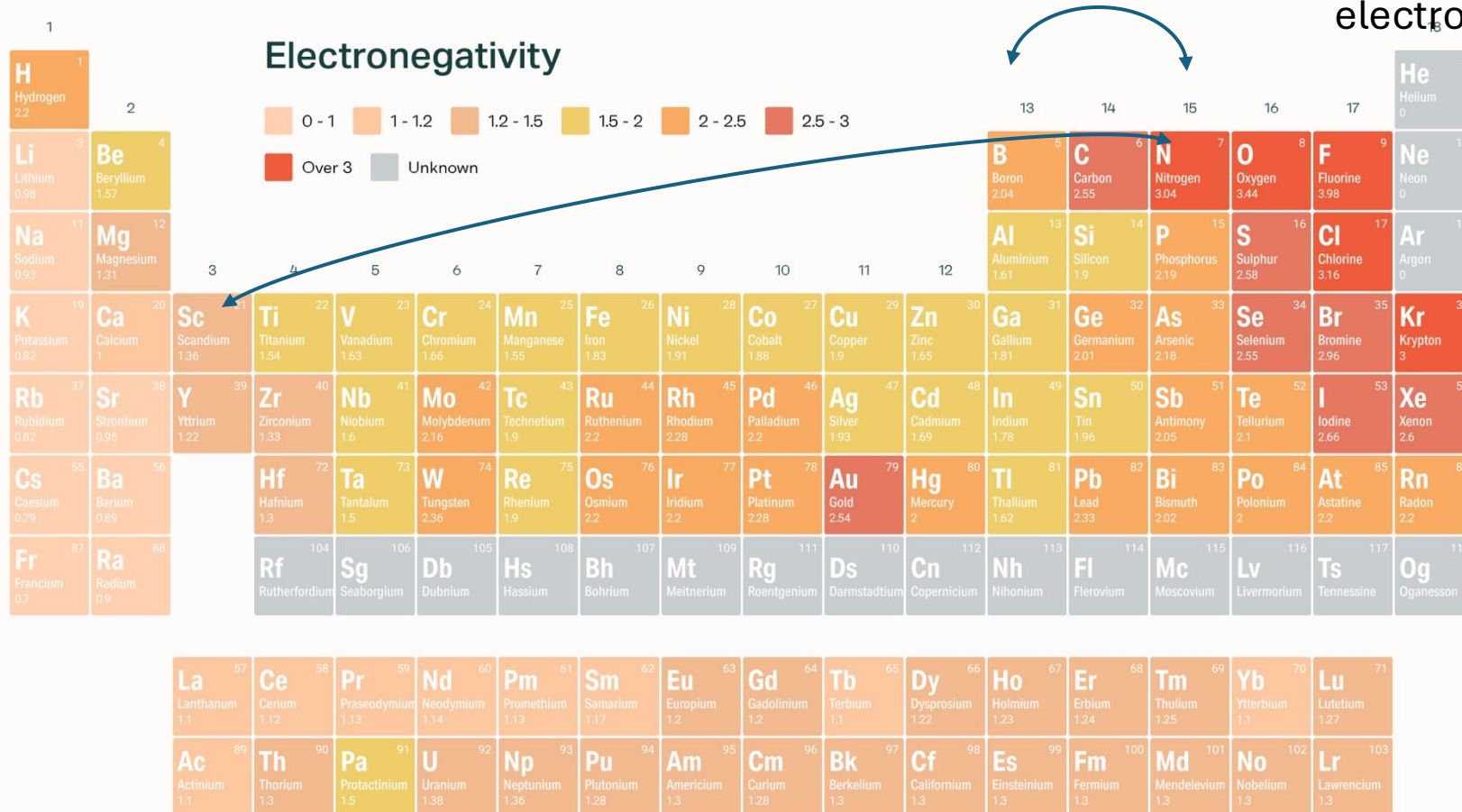
# Electronegativity = tendency to attract shared electrons

B, Ga, Al, In, Sc, Y, La: metals, electropositive, tend to give electrons (cations).

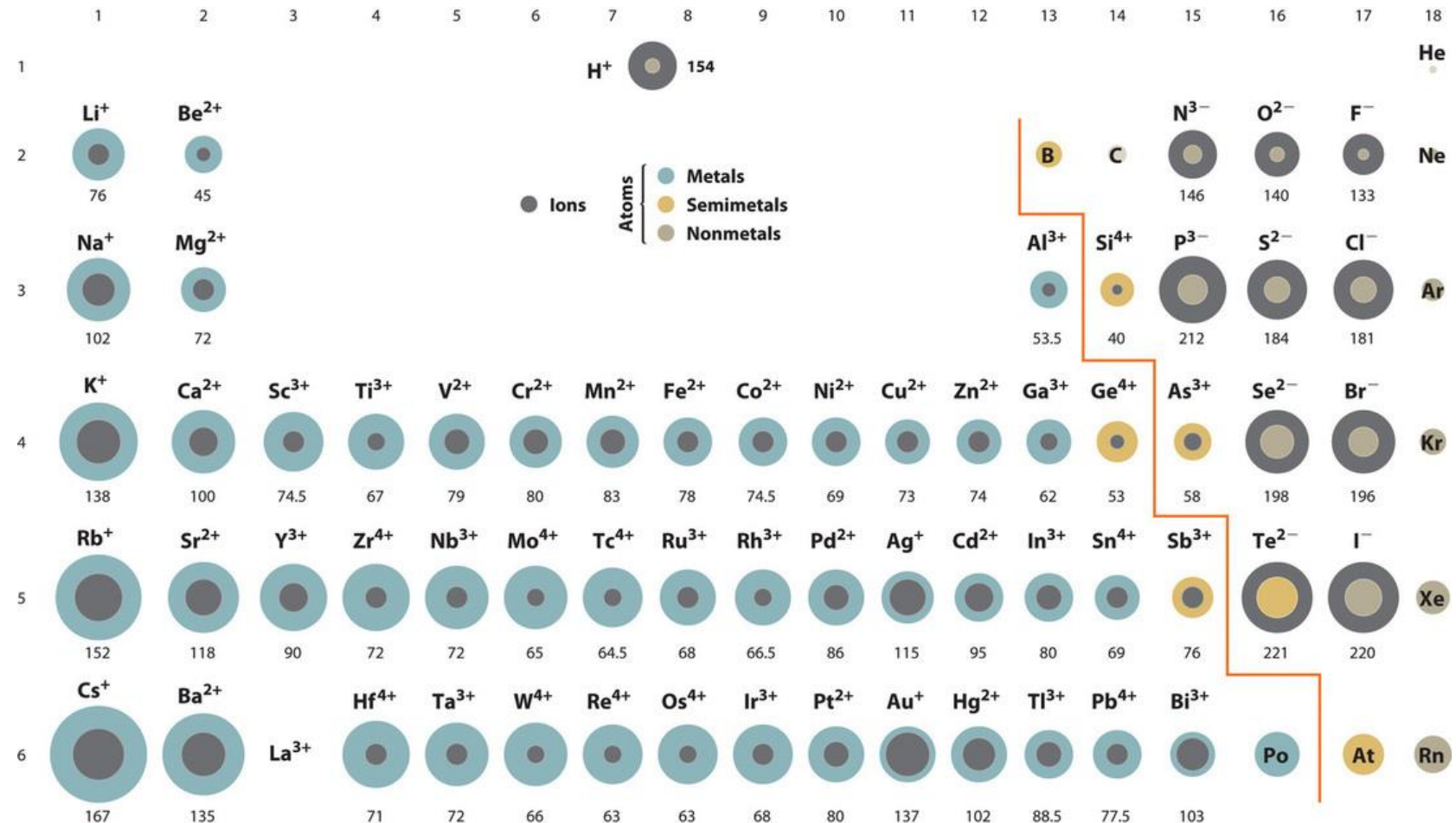
N: nonmetal, electronegative, tends to attract electrons (anion).

**Main group-III elements:** small electronegativity difference with Nitrogen  
  $\rightarrow$  Primarily **covalent** bonding

**Transition metal group-III elements:**  
large electronegativity difference with  
Nitrogen → **ionic** bonding



# Size of ions



In ceramic materials (e.g., compound semiconductors) the elements have a partial ionic character (e.g.,  $\text{Ga}^{3+}\text{N}^{3-}$ ). Need to consider size of ions when electrons are removed from (cations) or added to (anions) the outer shell.

E.g.,  $\text{Ga}^{3+}$ : large size difference between size of atom and size of ion. Dense packing of atoms enables strong orbital overlap. Explains high mobility of Ga-based semiconductors.



# Molecules and Chemical Bonding

Molecule = collection of nuclei + electrons. Need to describe coupled motion of electrons and nuclei.

Two basic approximations:

- 1) **Born-Oppenheimer** or **Adiabatic** approximation: since electrons are much lighter and move much faster than nuclei ( $m_e/M_p \sim 1/1836$ ), they respond instantaneously to the nuclear motion.
  - Solve electronic motion assuming that the nuclei are static
  - Solve nuclear motion assuming that electrons are always in their ground state

Good approximation in general, corrections needed for lightest elements.

- 2) **Classical Nuclei** approximation: electrons need quantum mechanics, but the motion of the nuclei can be described by classical physics.



# Molecules and Chemical Bonding

In molecules, the valence electrons do not stay in their atomic orbitals. Instead, they join together to form bonds.

Quantum principles still apply: discrete energy levels, electrons described by wave functions.

Wave functions in atoms =  
atomic orbitals (e.g.,  $1s$ ,  $2p_x$ , ...)

Wave functions in molecules =  
**molecular orbitals** (e.g.,  $\sigma$ ,  $\pi$ , ...)

Need to solve Schrödinger's equation  
for molecule. Hard to do, need  
approximations.



<http://magneticcarpet.deviantart.com/art/Molecular-orbital-102480755>



# Linear Combination of Atomic Orbitals

Near an atom, molecular wave functions look like the atomic wave functions of the atom.

Assumption: the molecular wave function can be approximated by a **Linear Combination of Atomic Orbitals (LCAO method)**

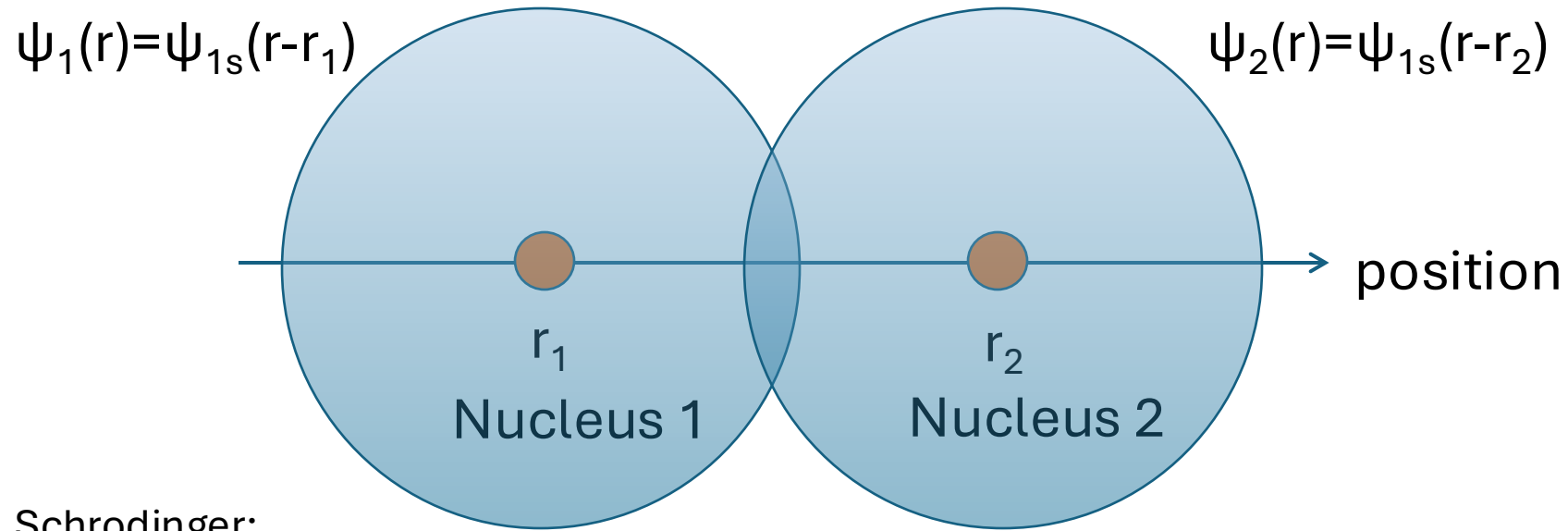
$$\psi_{\text{molecule}}(\vec{r}) = c_1 \psi_{\text{atom1}}(\vec{r} - \vec{r}_1) + c_2 \psi_{\text{atom2}}(\vec{r} - \vec{r}_2)$$

$c_1, c_2$  = numerical prefactors, chosen to satisfy Schrodinger's equation.

Physical meaning of  $c_1, c_2$ : Measure how much each of  $\psi_1, \psi_2$  contribute to the molecular wave function, depend on the electronegativity of the atoms.

$|c_1|^2$  = probability that the electron in the molecular orbital is found in state  $\psi_1$

# Example: LCAO for the H<sub>2</sub> molecule



Schrodinger:

$$H\psi_{\text{molecule}}(\vec{r}) = E_{\text{molecule}}\psi_{\text{molecule}}(\vec{r})$$

$$H = -\frac{\hbar^2}{2m}\nabla^2 + V_{\text{molecule}}(\vec{r})$$

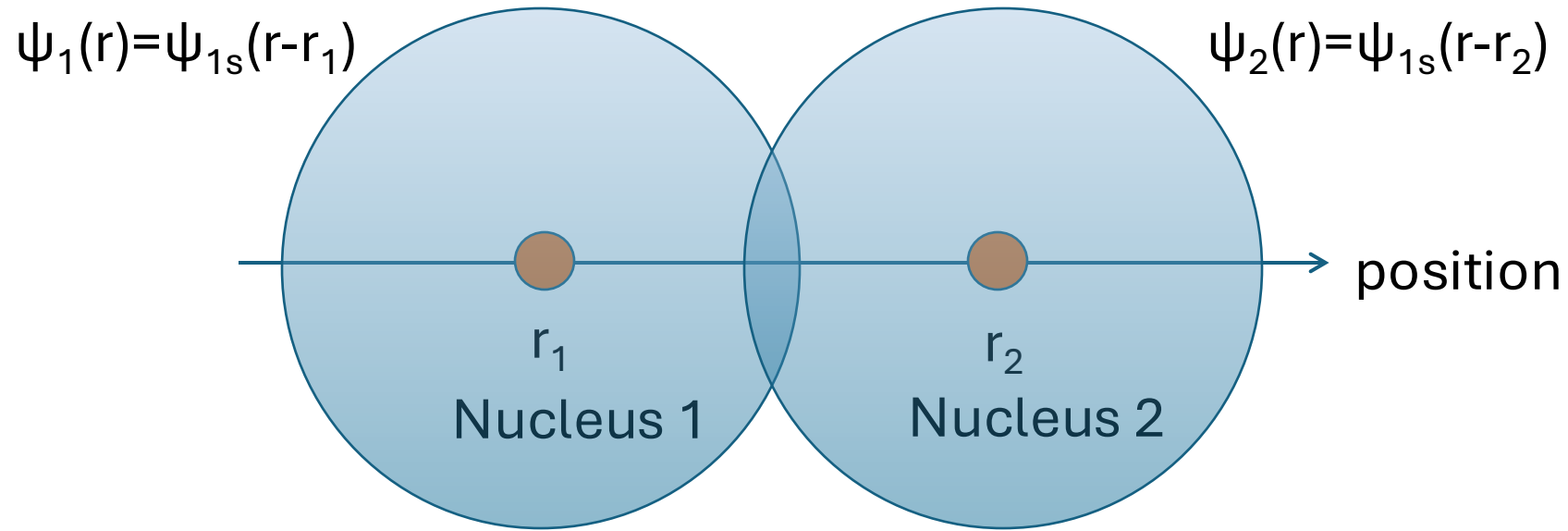
Potential:

$$V_{\text{molecule}}(\vec{r}) = V_{\text{atom1}}(\vec{r} - \vec{r}_1) + V_{\text{atom2}}(\vec{r} - \vec{r}_2)$$

Wave function:

$$\begin{aligned}\psi_{\text{molecule}}(\vec{r}) &= c_1\psi_{\text{atom1}}(\vec{r} - \vec{r}_1) + c_2\psi_{\text{atom2}}(\vec{r} - \vec{r}_2) \\ &= c_1\psi_1(\vec{r}) + c_2\psi_2(\vec{r})\end{aligned}$$

# Example: LCAO for the H<sub>2</sub> molecule

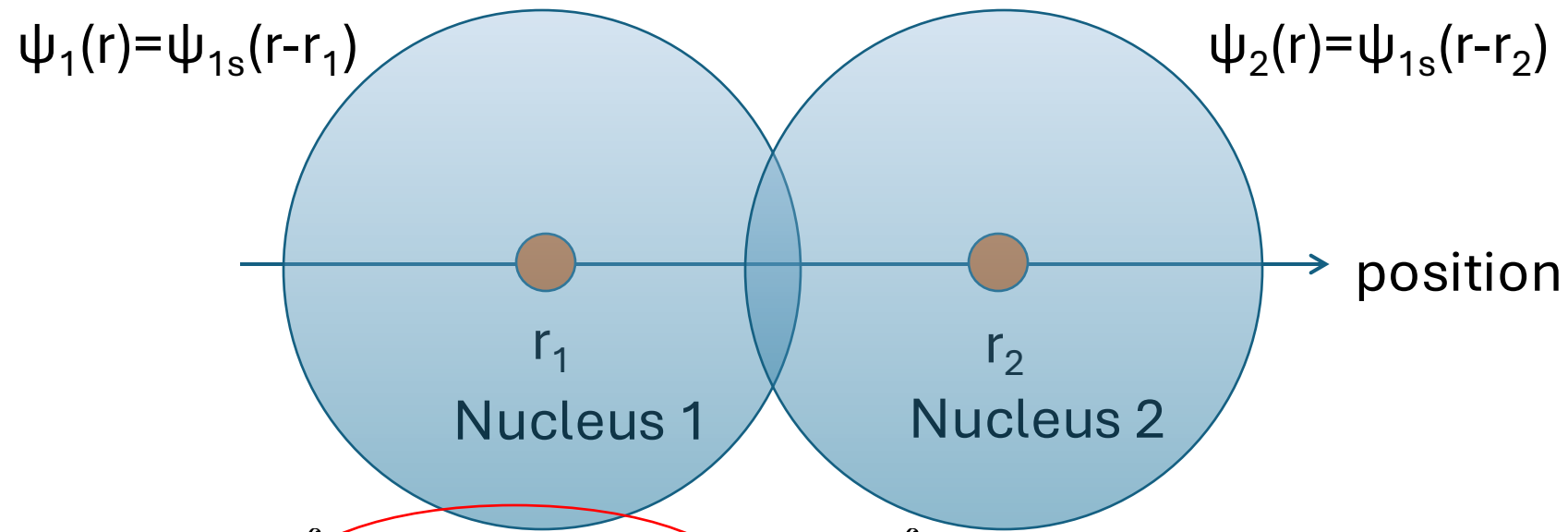


$$c_1(H\psi_1) + c_2(H\psi_2) = Ec_1\psi_1 + Ec_2\psi_2$$

Multiply both sides by  $\psi_1^*$  and integrate over  $r$ :

$$\begin{aligned} c_1 \int \psi_1^*(\vec{r}) H \psi_1(\vec{r}) d\vec{r} + c_2 \int \psi_1^*(\vec{r}) H \psi_2(\vec{r}) d\vec{r} = \\ = Ec_1 \int \psi_1^*(\vec{r}) \psi_1(\vec{r}) d\vec{r} + Ec_2 \int \psi_1^*(\vec{r}) \psi_2(\vec{r}) d\vec{r} \end{aligned}$$

# Physical meaning of the integrals

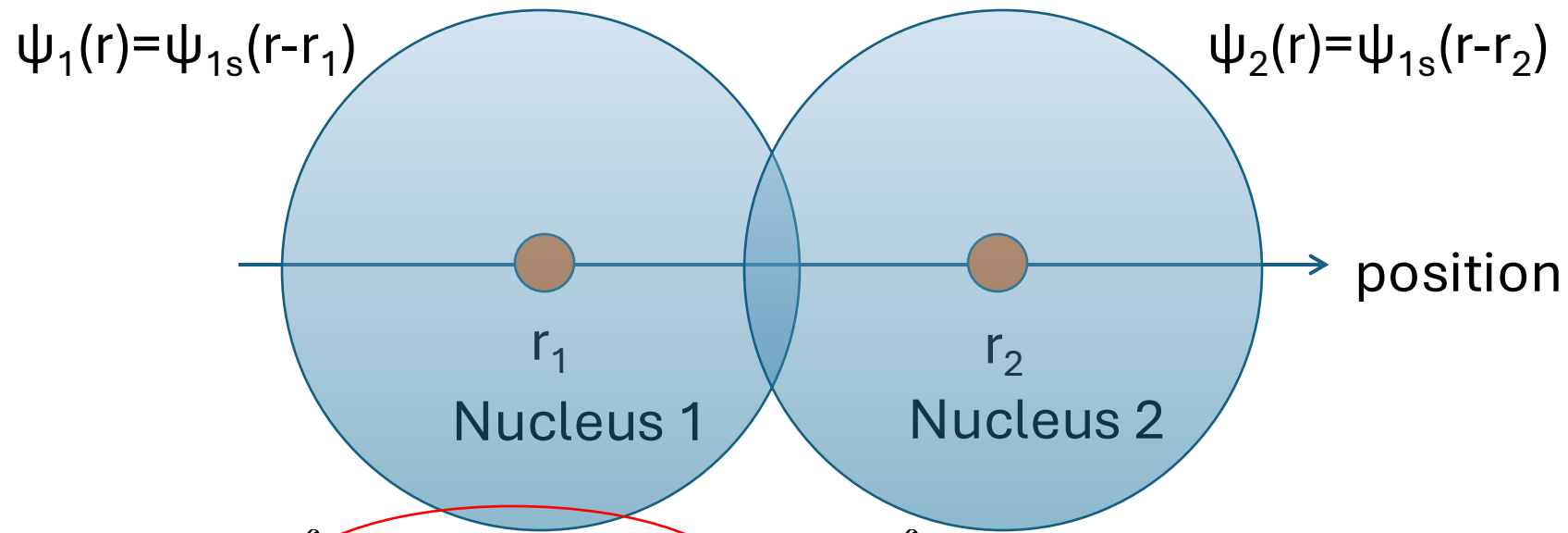


$$c_1 \int \psi_1^*(\vec{r}) H \psi_1(\vec{r}) d\vec{r} + c_2 \int \psi_1^*(\vec{r}) H \psi_2(\vec{r}) d\vec{r} =$$
$$= E c_1 \int \psi_1^*(\vec{r}) \psi_1(\vec{r}) d\vec{r} + E c_2 \int \psi_1^*(\vec{r}) \psi_2(\vec{r}) d\vec{r}$$

$$H = -\frac{\hbar^2}{2m} \nabla^2 + V_{\text{molecule}}(\vec{r})$$

$$V_{\text{molecule}}(\vec{r}) = V_{\text{atom1}}(\vec{r} - \vec{r}_1) + V_{\text{atom2}}(\vec{r} - \vec{r}_2)$$

# Physical meaning of the integrals



$$c_1 \int \psi_1^*(\vec{r}) H \psi_1(\vec{r}) d\vec{r} + c_2 \int \psi_1^*(\vec{r}) H \psi_2(\vec{r}) d\vec{r} =$$

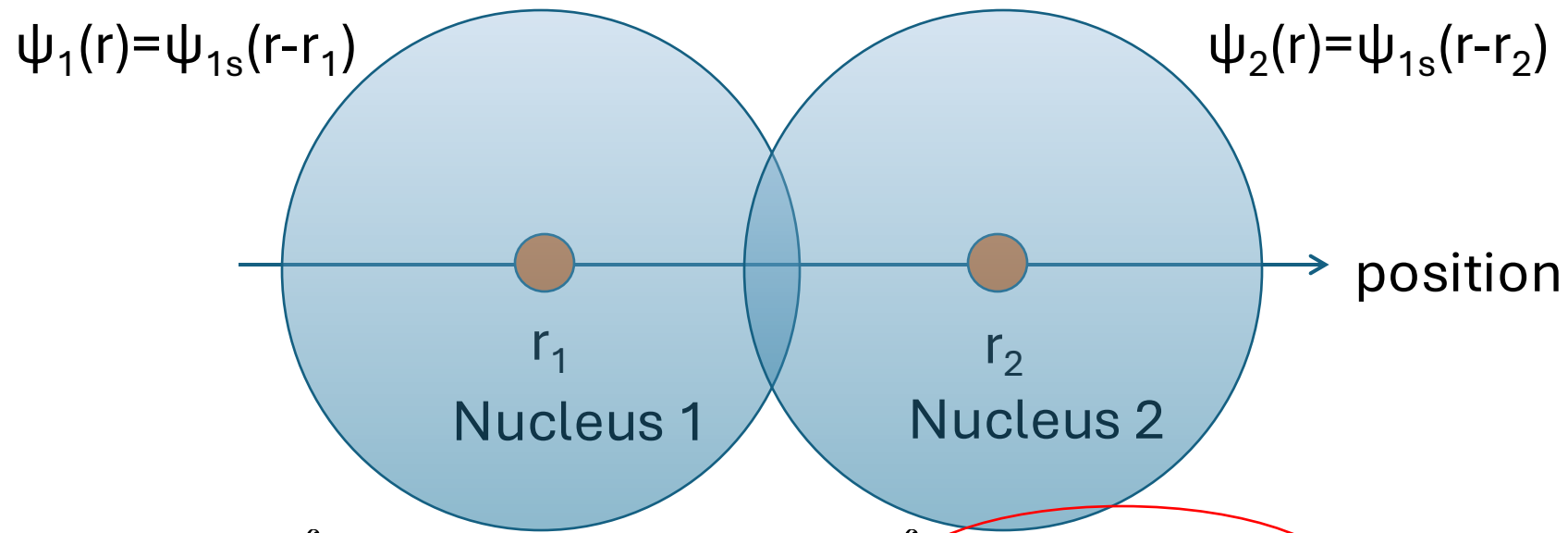
$$= E c_1 \int \psi_1^*(\vec{r}) \psi_1(\vec{r}) d\vec{r} + E c_2 \int \psi_1^*(\vec{r}) \psi_2(\vec{r}) d\vec{r}$$

In the neighborhood of atom 1 (where  $\psi_1$  is large):  $V_{\text{molecule}} \cong V_{\text{atom1}}$

$$\Rightarrow \int \psi_1^*(\vec{r}) H \psi_1(\vec{r}) d\vec{r} \simeq \int \psi_1^*(\vec{r}) H_{\text{atom1}} \psi_1(\vec{r}) d\vec{r} = E_{1s} \equiv \epsilon_1$$

= Energy of the atomic orbital  $\psi_1$  of atom 1 = energy of 1s orbital of Hydrogen atom

# Physical meaning of the integrals



$$c_1 \int \psi_1^*(\vec{r}) H \psi_1(\vec{r}) d\vec{r} + c_2 \int \psi_1^*(\vec{r}) H \psi_2(\vec{r}) d\vec{r} =$$

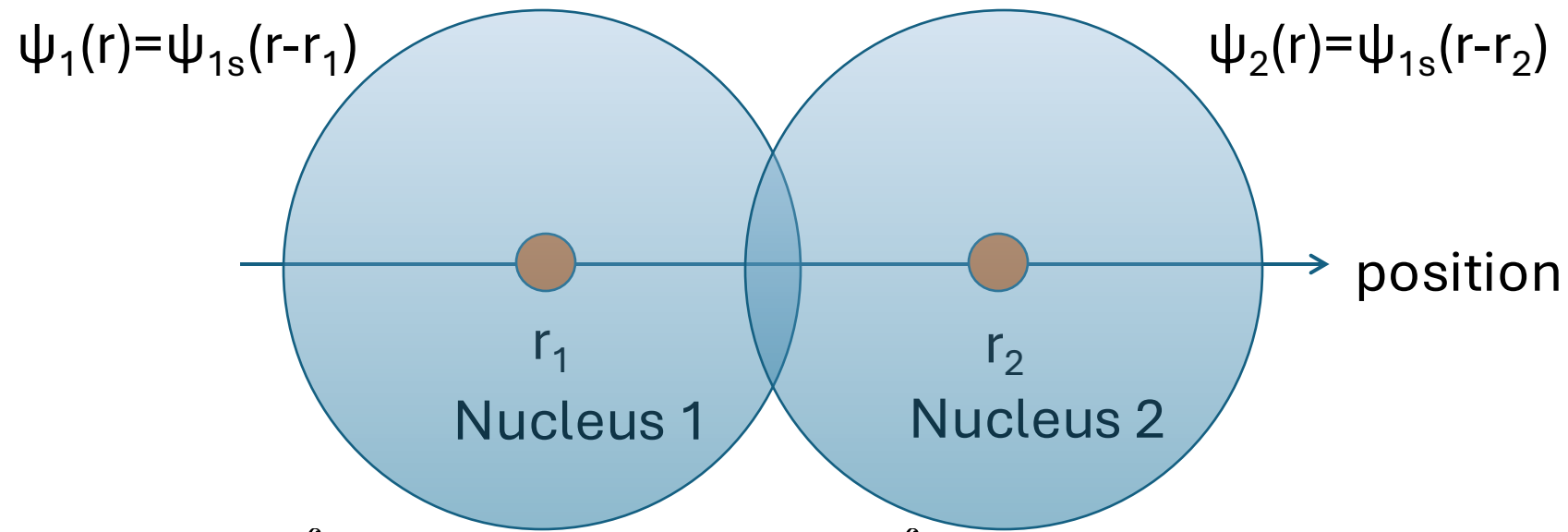
$$= E c_1 \int \psi_1^*(\vec{r}) \psi_1(\vec{r}) d\vec{r} + E c_2 \int \psi_1^*(\vec{r}) \psi_2(\vec{r}) d\vec{r}$$

**Interatomic matrix element:**

$$\int \psi_1^*(\vec{r}) H \psi_2(\vec{r}) d\vec{r} = -V_s$$

= a negative number. It is related to the probability that an electron can transition from atomic orbital  $\psi_1$  to  $\psi_2$  = relates to **sharing** of electrons between atoms.

# Physical meaning of the integrals

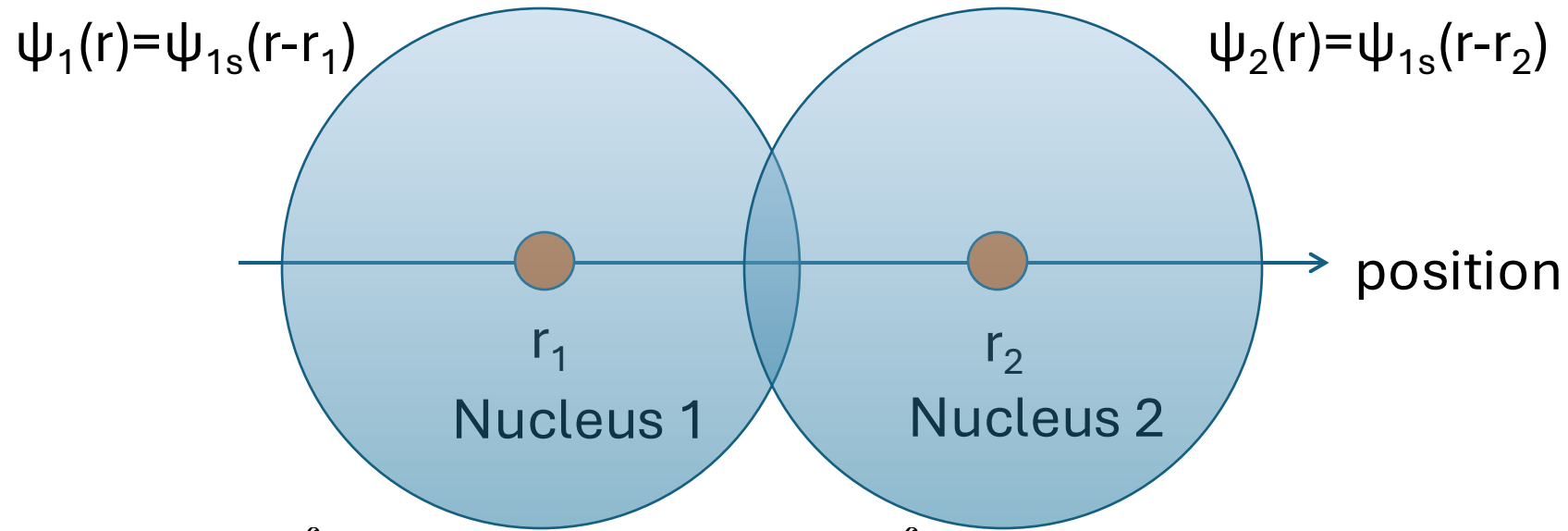


$$\begin{aligned} c_1 \int \psi_1^*(\vec{r}) H \psi_1(\vec{r}) d\vec{r} + c_2 \int \psi_1^*(\vec{r}) H \psi_2(\vec{r}) d\vec{r} = \\ = Ec_1 \int \psi_1^*(\vec{r}) \psi_1(\vec{r}) d\vec{r} + Ec_2 \int \psi_1^*(\vec{r}) \psi_2(\vec{r}) d\vec{r} \end{aligned}$$

**Question: what is the value of this integral?**

**Answer: = 1**

# Physical meaning of the integrals



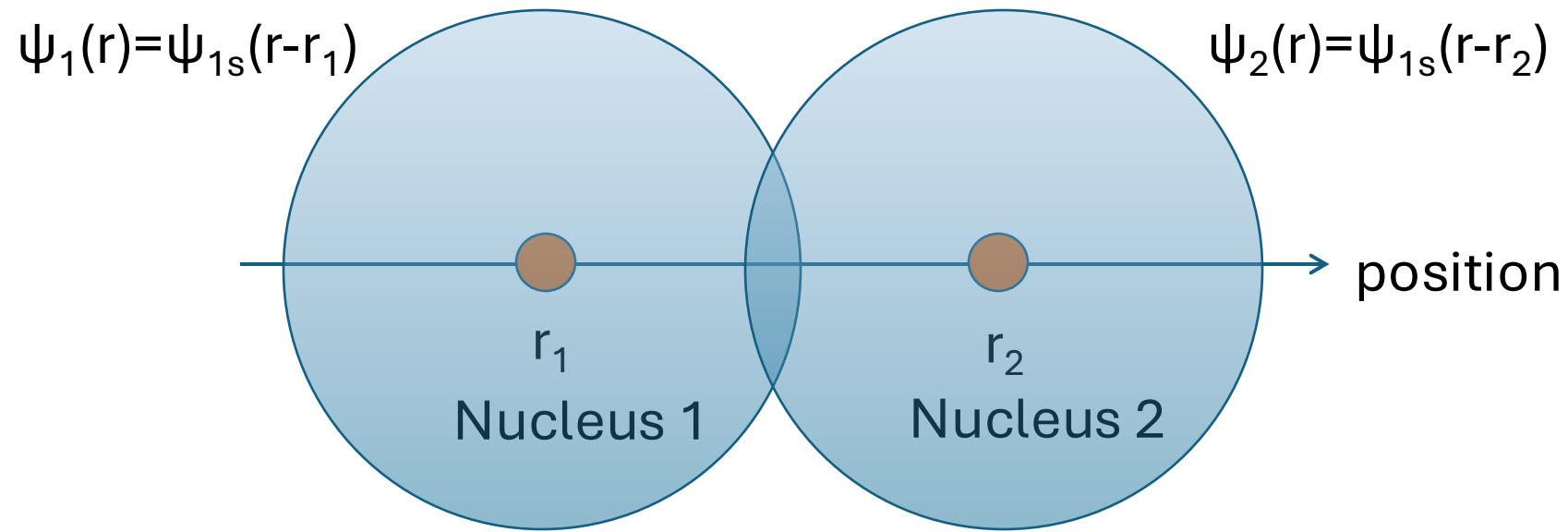
$$\begin{aligned} c_1 \int \psi_1^*(\vec{r}) H \psi_1(\vec{r}) d\vec{r} + c_2 \int \psi_1^*(\vec{r}) H \psi_2(\vec{r}) d\vec{r} = \\ = E c_1 \int \psi_1^*(\vec{r}) \psi_1(\vec{r}) d\vec{r} + E c_2 \int \psi_1^*(\vec{r}) \psi_2(\vec{r}) d\vec{r} \end{aligned}$$

**Question: what is the value of this integral?**

**Answer: overlap integral**  $\int \psi_1^*(\vec{r}) \psi_2(\vec{r}) d\vec{r} = S \simeq 0$



# Physical meaning of the integrals



$$\Rightarrow c_1 \epsilon_1 - c_2 V_s = E c_1$$

One equation with 3 unknowns. Question: What to do next?

Answer: go back to Schrodinger equation, multiply by  $\psi_2$  and integrate

# Question: and now what?

$$\begin{array}{l} c_1 \epsilon_1 - c_2 V_s = E c_1 \\ -c_1 V_s + c_2 \epsilon_2 = E c_2 \end{array}$$

For  $H_2$  :  $\epsilon_1 = \epsilon_2 = E_{1s}$

Hint: how many unknowns and how many equations?

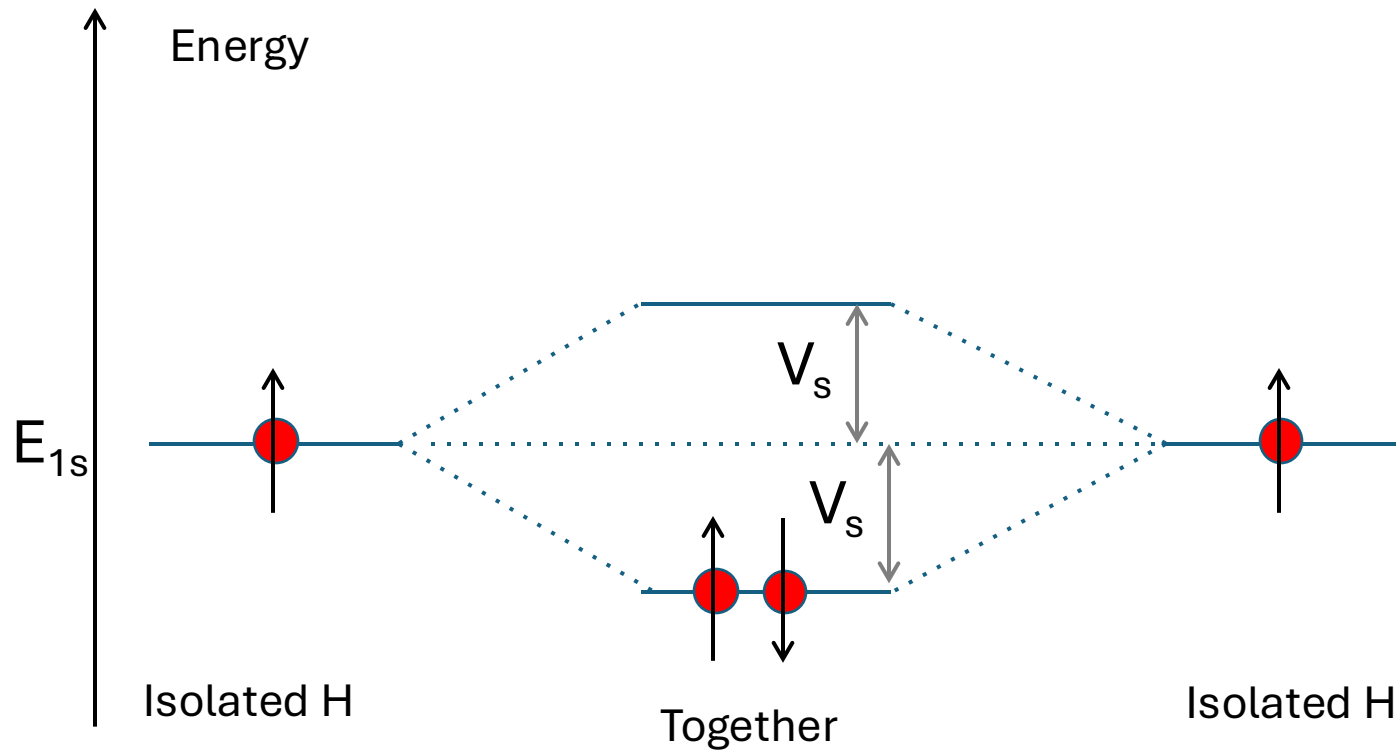
Answer: determinant of coefficients must be zero. Otherwise, the system of the two equations has a unique trivial solution:

$c_1 = c_2 = 0$

$$\begin{array}{l} c_1 (\epsilon_1 - E) - c_2 V_s = 0 \\ -c_1 V_s + c_2 (\epsilon_2 - E) = 0 \end{array}$$

Algebra ...

# Energy levels of H<sub>2</sub> molecular orbitals



Question: Does the energy increase or decrease when the two H atoms are brought together? What can you say about bond formation?

Answer: Energy is lowered by  $2V_s$ , the molecule has a lower energy than individual atoms → **chemical bond** forms between H atoms.

# The wave functions of the H<sub>2</sub> molecule (molecular orbitals)

What is the wave function for  $E_{\text{bonding}} = E_{1s} - V_s$ ?

And what is the wave function for  $E_{\text{antibonding}} = E_{1s} + V_s$ ?

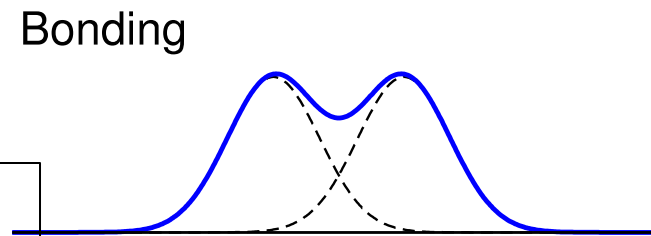
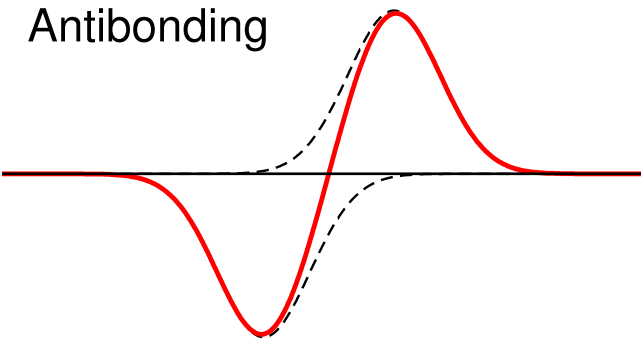
Answer: go back to the set of equations and set the energy equal to the bonding or the antibonding orbital energy, find the ratio of  $c_1/c_2$ . Then, impose normalization  $c_1^2 + c_2^2 = 1$

For the **bonding** orbital:  
( $\sigma$  orbital, no node)

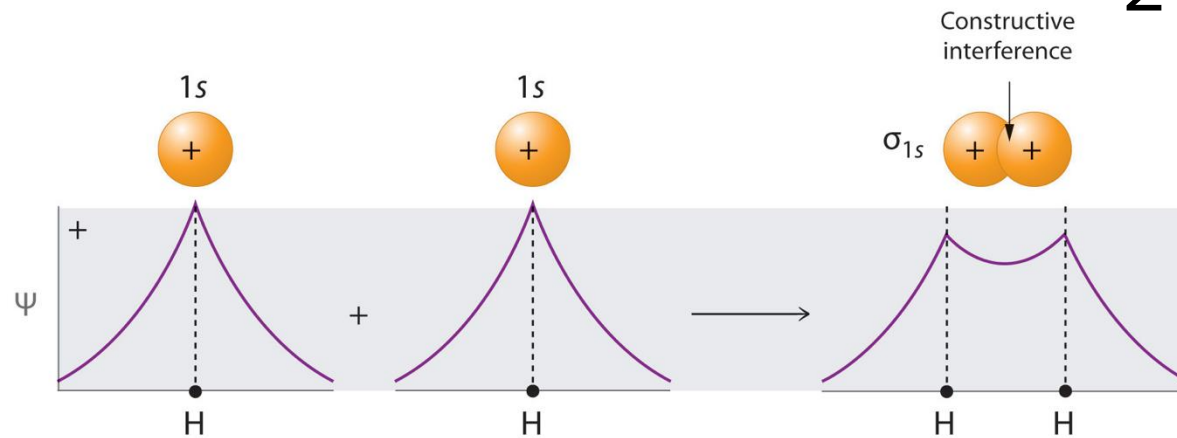
$$c_1 = c_2 = \frac{1}{\sqrt{2}}$$

For the **antibonding** orbital:  
( $\sigma^*$  orbital, with node)

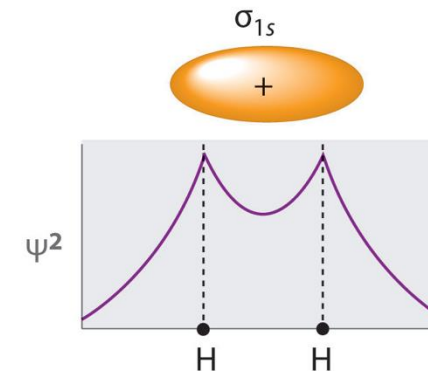
$$c_1 = -c_2 = \frac{1}{\sqrt{2}}$$



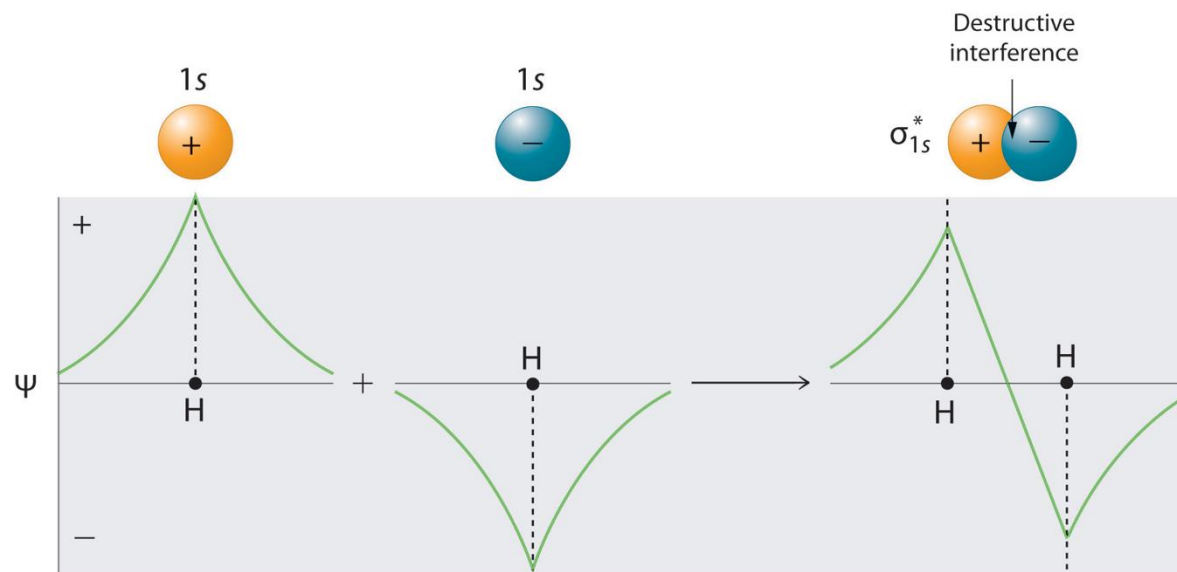
# Molecular orbitals of H<sub>2</sub> molecule



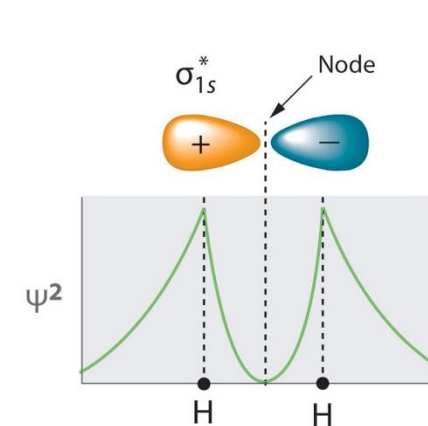
(a) Wave functions combined for  $\sigma_{1s}$



(b) Bonding probability density

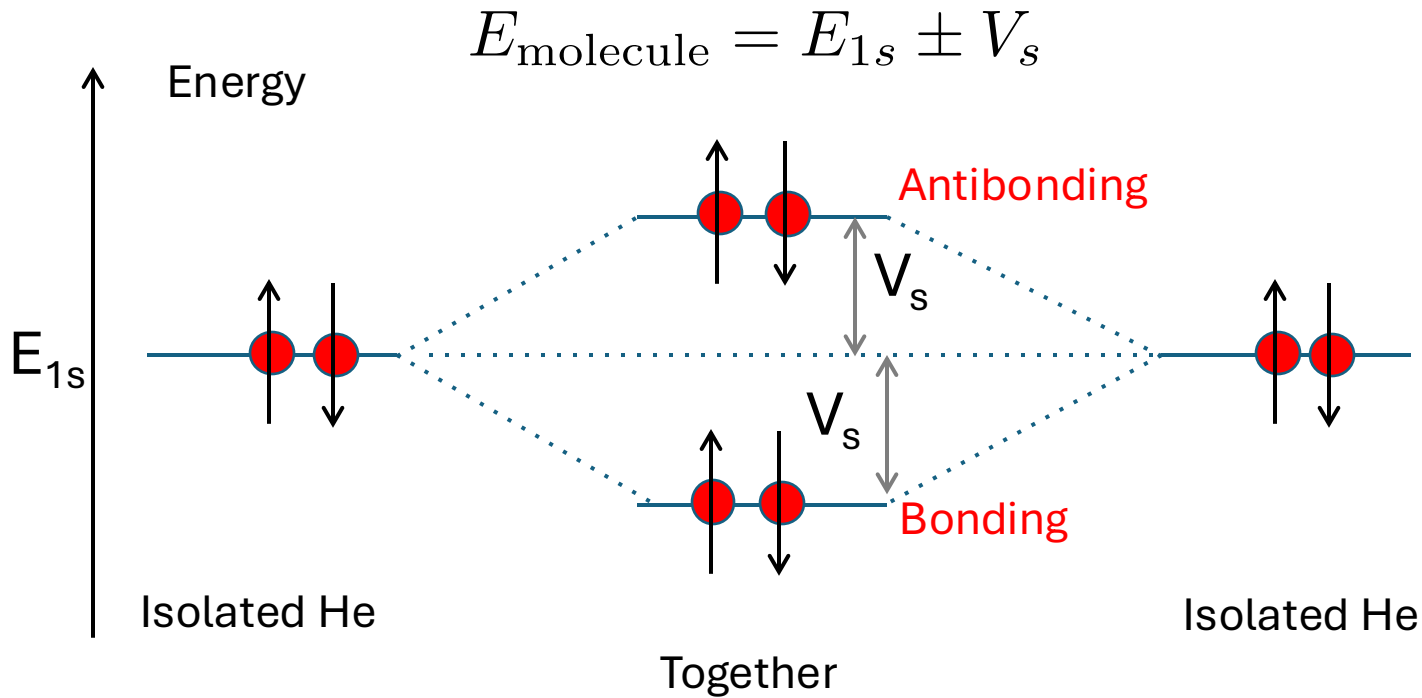


(c) Wave functions combined for  $\sigma_{1s}^*$



(d) Antibonding probability density

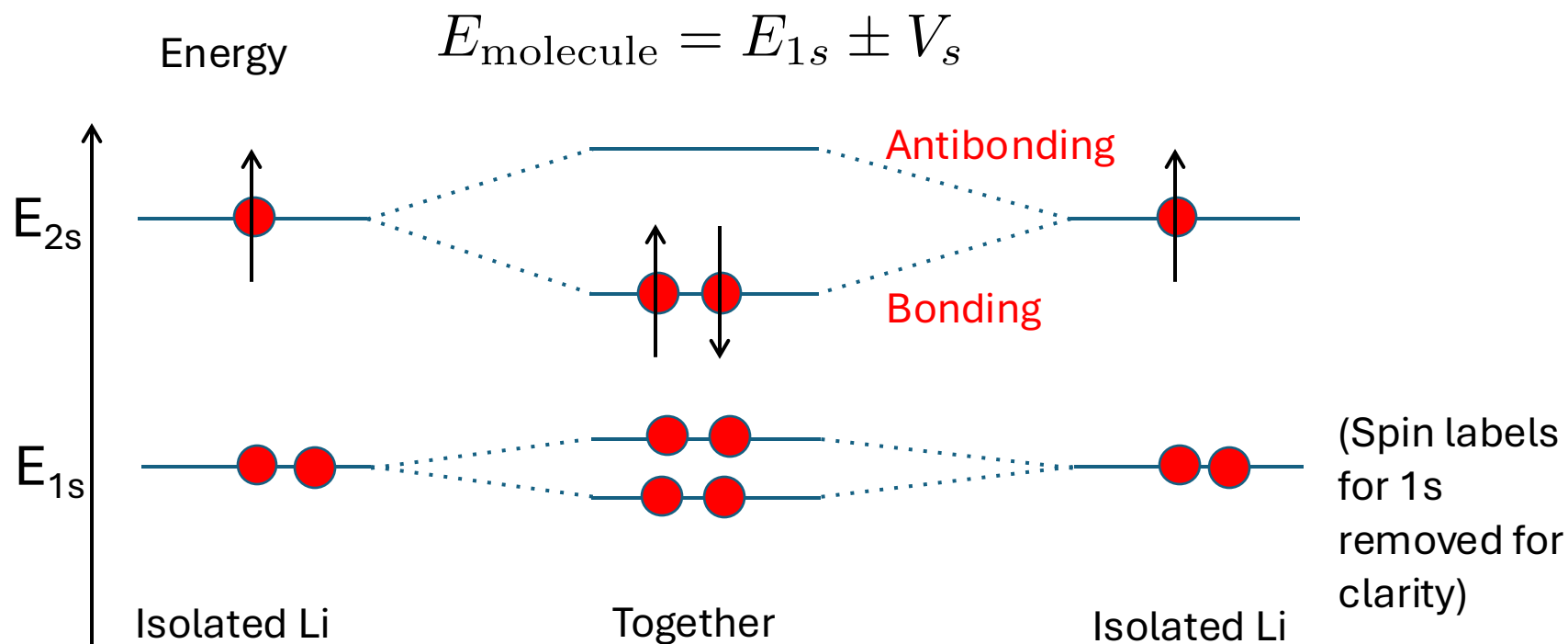
Follow the same steps for the He<sub>2</sub> molecule. Does the He<sub>2</sub> molecule bind? Why or why not?



Answer: The four electrons occupy both the bonding and the antibonding orbitals. The net energy gain is zero. → No bond formation.  
(in reality: weak van der Waals bond, but broken by zero point motion: Helium remains liquid even at 0 K.)

# In-class exercise: Li<sub>2</sub>

Follow the same steps for the Li<sub>2</sub> molecule. Does the Li<sub>2</sub> molecule bind or not? If yes, which electrons contribute to bonding?



Answer: the filled 1s electrons (**core electrons**) do not contribute to bonding. The **valence 2s electrons** lower the energy and bind Li<sub>2</sub>.  $V_s$  is larger for 2s electrons than 1s electrons. Why?

Answer: 2s orbital is larger, 2s electrons more easily shared than 1s.



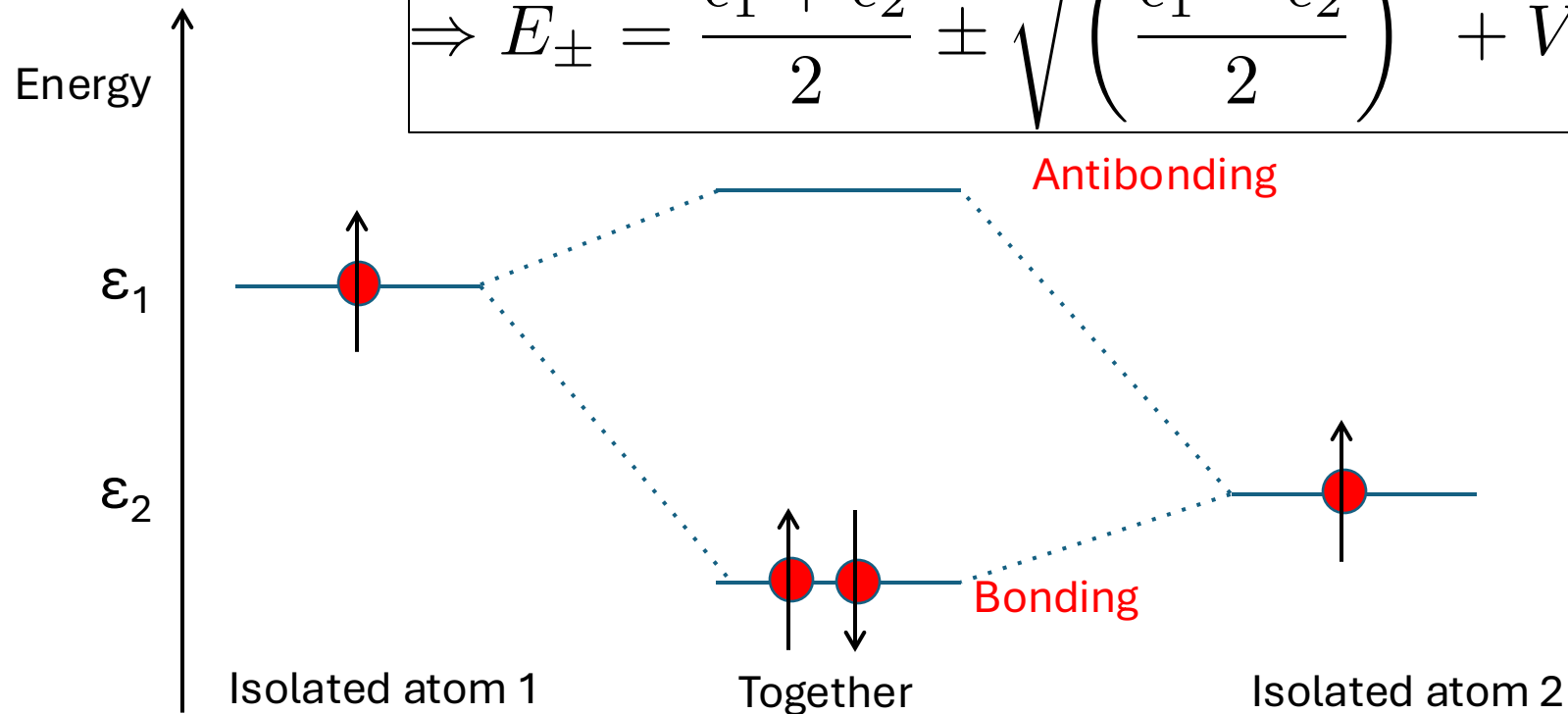
# What if the two atoms are different?

In-class: what are the energies of the molecular orbitals if  $\epsilon_1 \neq \epsilon_2$ ?

Answer: go back to the eigenvalue problem (determinant) and solve for  $\epsilon_1 \neq \epsilon_2$ .

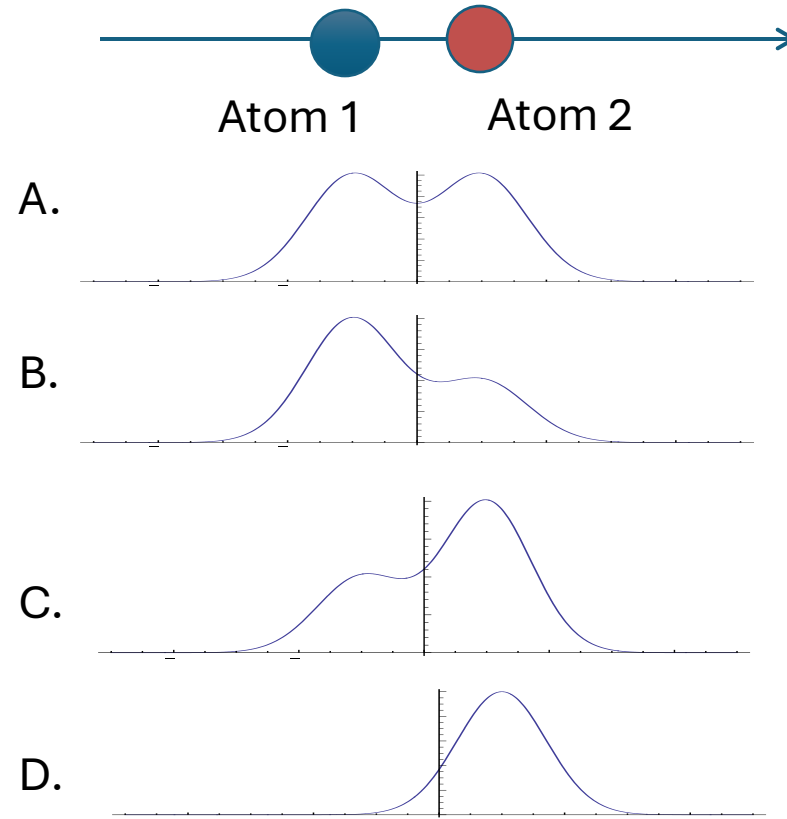
$$\text{Det} = 0 \Rightarrow (\epsilon_1 - E)(\epsilon_2 - E) - V_s^2 = 0$$

$$\Rightarrow E_{\pm} = \frac{\epsilon_1 + \epsilon_2}{2} \pm \sqrt{\left(\frac{\epsilon_1 - \epsilon_2}{2}\right)^2 + V_s^2}$$



# Question

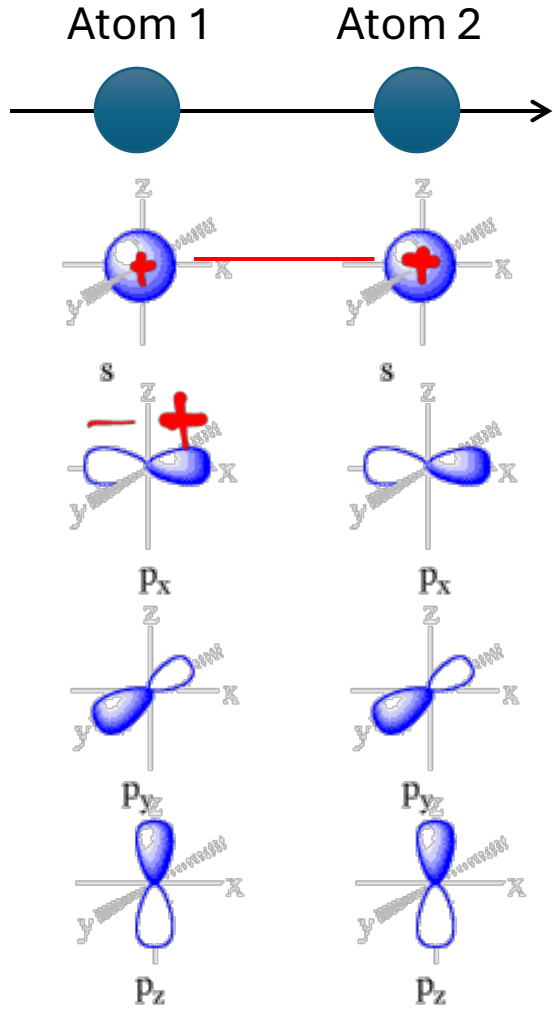
If  $\epsilon_1 > \epsilon_2$ , what is the shape of the bonding molecular orbital?



Answer: **C**: Atom 2 is more electronegative ( $\epsilon_2 < \epsilon_1$ , therefore the potential well is deeper). There is a higher probability that the electron will be nearer atom 2. However, there is still some probability that the electron will be shared between both atoms.

# Diatomic molecules with p orbitals

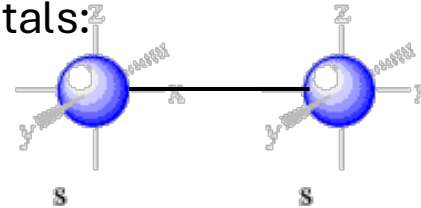
The p valence electrons (e.g., C, N, O) also participate in bonding



Key quantities: **interatomic matrix elements** for each pair of atomic orbitals

$$\int \psi_1^*(\vec{r}) H \psi_2(\vec{r}) d\vec{r}$$

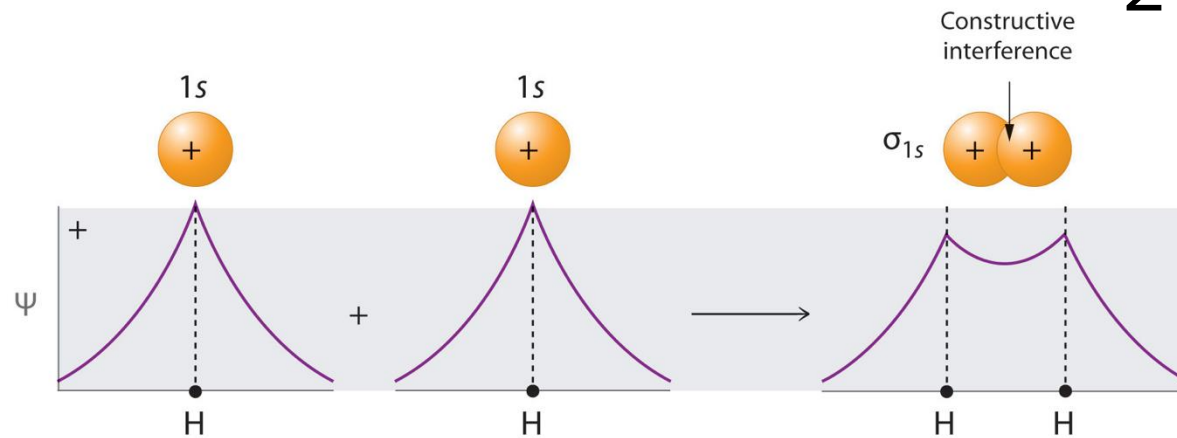
For a pair of s orbitals:



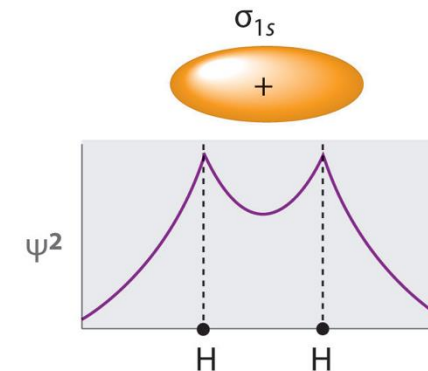
$$\int \psi_1^*(\vec{r}) H \psi_2(\vec{r}) d\vec{r} = -V_s = V_{ss\sigma}$$

Two atomic s orbitals combine to form  $\sigma$  and  $\sigma^*$  molecular orbitals

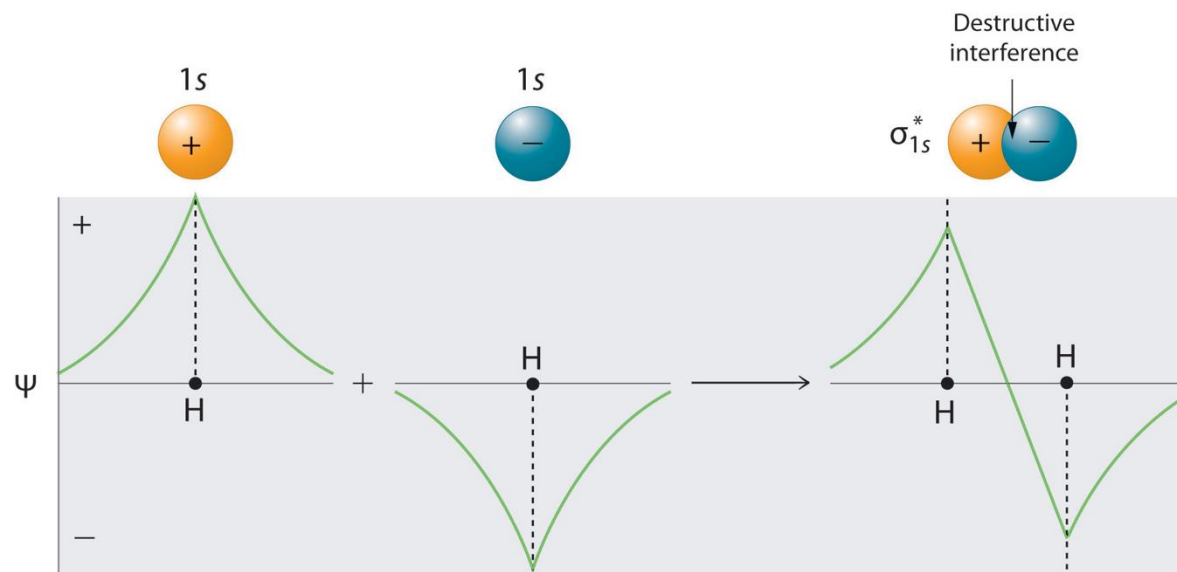
# Molecular orbitals of H<sub>2</sub> molecule



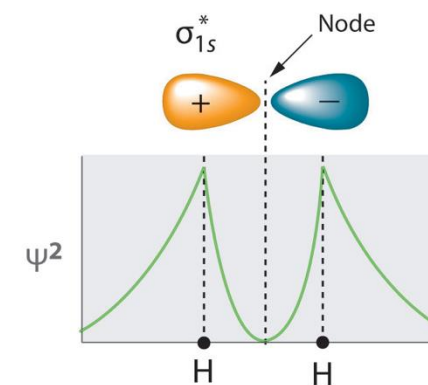
(a) Wave functions combined for  $\sigma_{1s}$



(b) Bonding probability density



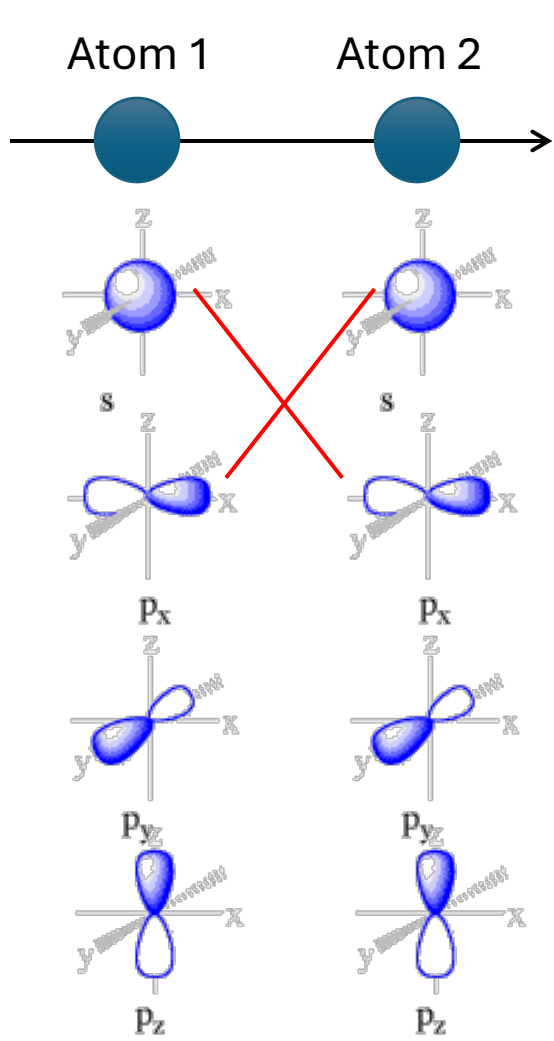
(c) Wave functions combined for  $\sigma_{1s}^*$



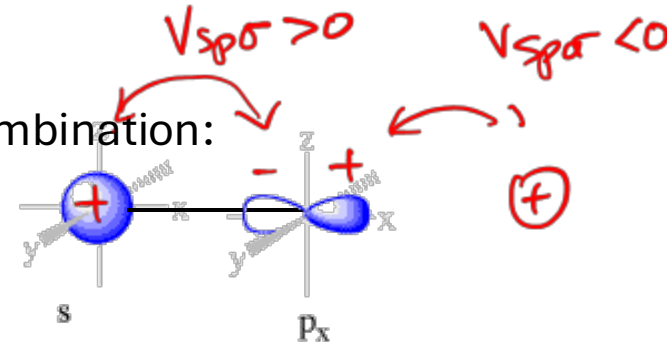
(d) Antibonding probability density

# Diatomic molecules with p orbitals

The p valence electrons (e.g., C, N, O) also participate in bonding

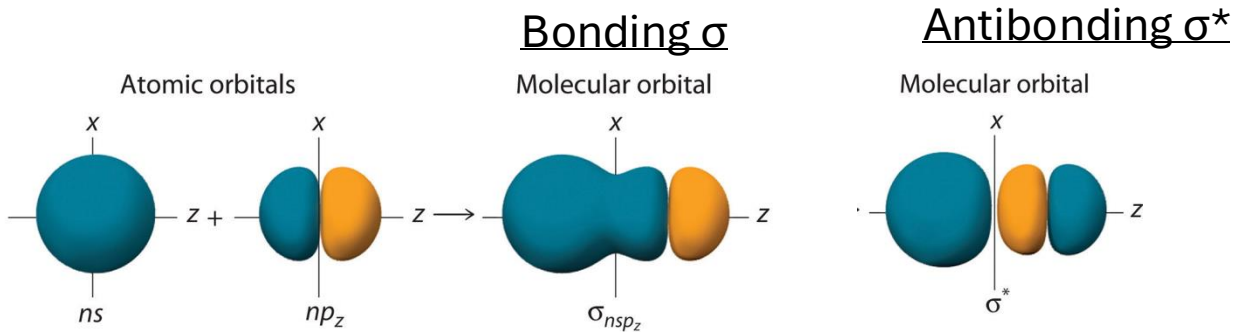


For the  $p_x$ -s combination:



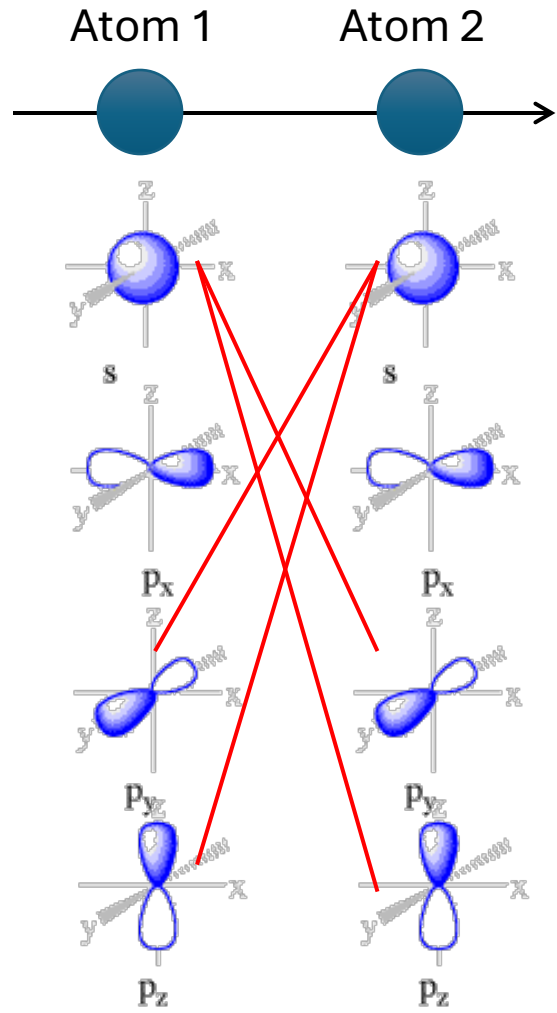
$$\int \psi_1^*(\vec{r}) H \psi_2(\vec{r}) d\vec{r} = V_{sp\sigma}$$

One s and one p orbitals combine to form  $\sigma$  and  $\sigma^*$  orbitals:

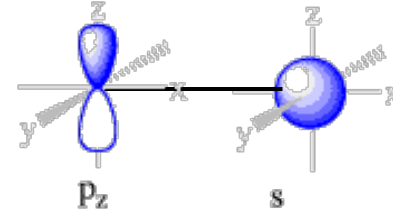


# Diatomic molecules with p orbitals

The p valence electrons (e.g., C, N, O) also participate in bonding



For the p<sub>y</sub>-s or p<sub>z</sub>-s combinations:

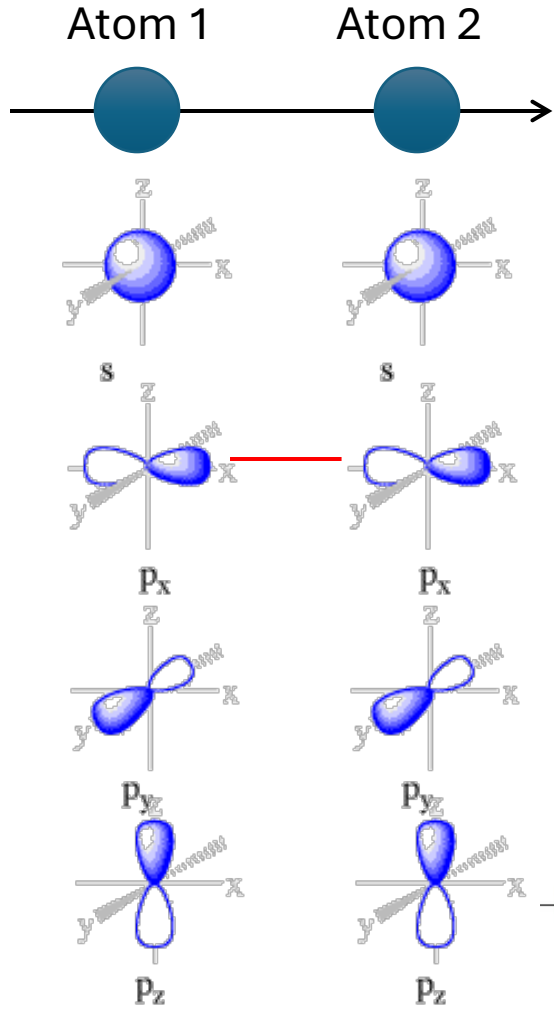


$$\int \psi_1^*(\vec{r}) H \psi_2(\vec{r}) d\vec{r} = 0$$

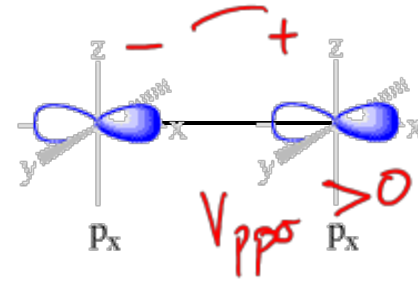
The product of the two orbitals is an odd function with respect to z axis, integral is zero

# Diatomic molecules with p orbitals

The p valence electrons (e.g., C, N, O) also participate in bonding

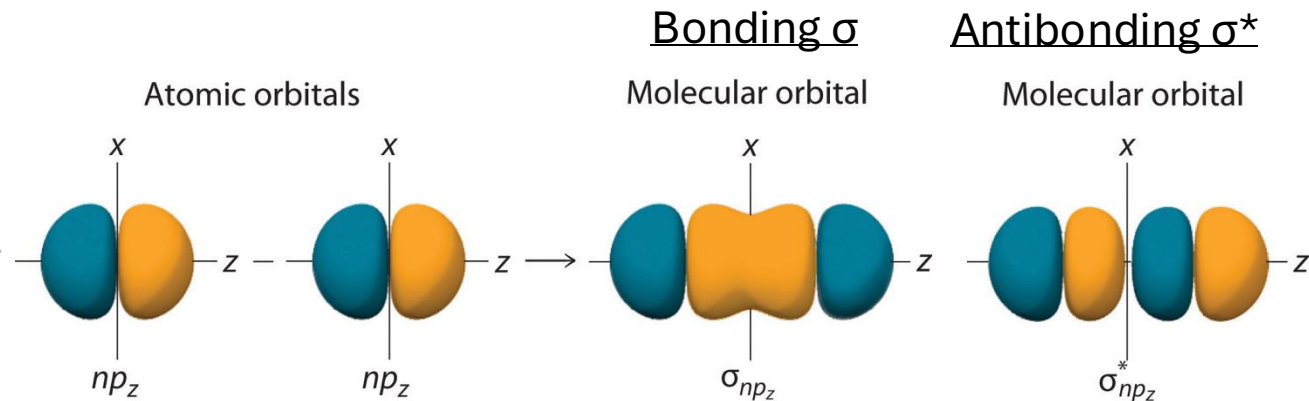


For a pair of p orbitals that interact tip-to-tip :



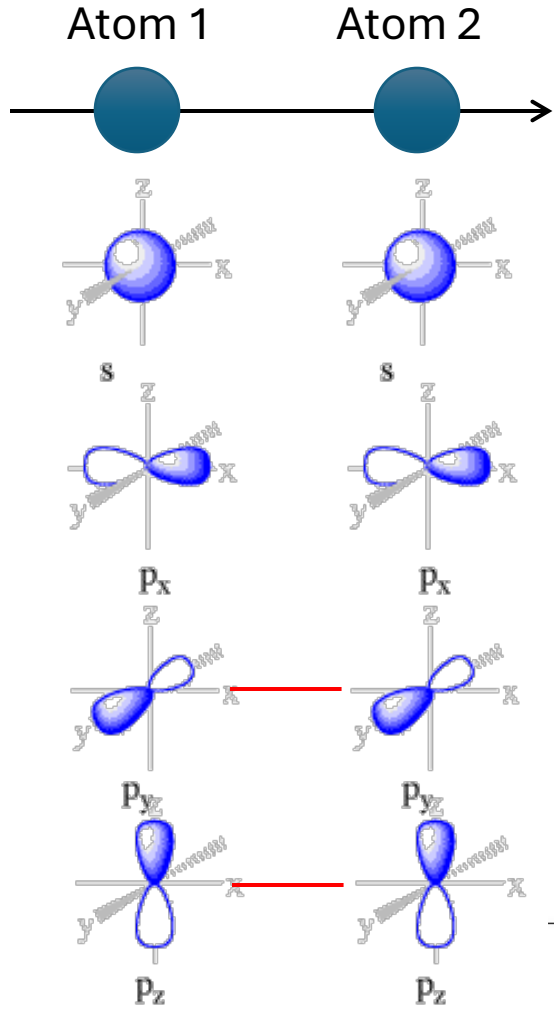
$$\int \psi_1^*(\vec{r}) H \psi_2(\vec{r}) d\vec{r} = \underline{\underline{V_{pp\sigma}}}$$

Two p orbitals combine to form  $\sigma$  and  $\sigma^*$  orbitals

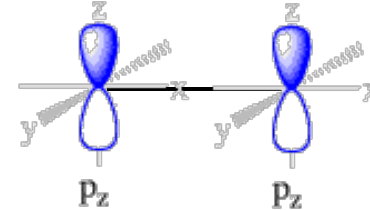


# Diatomic molecules with p orbitals

The p valence electrons (e.g., C, N, O) also participate in bonding

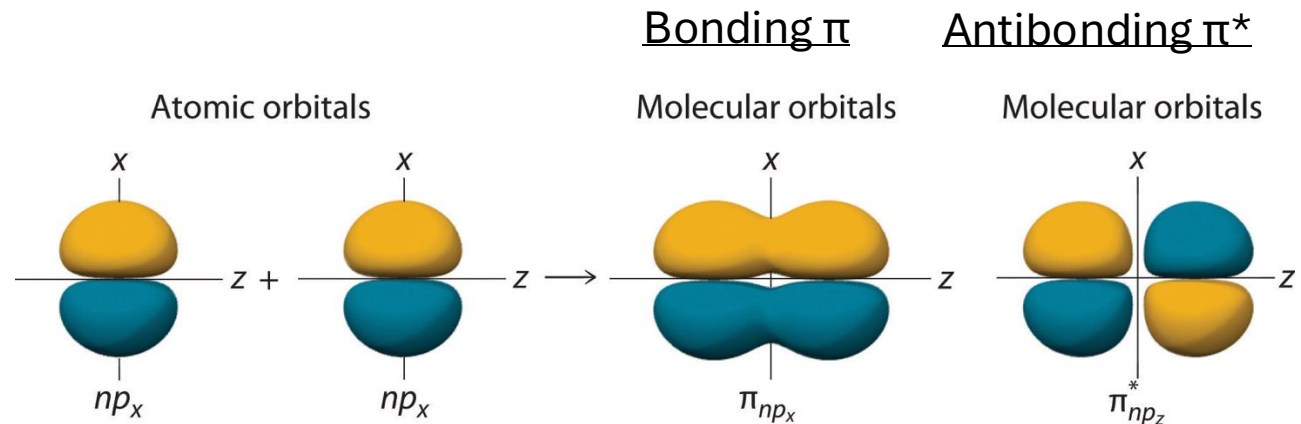


For a pair of p orbitals that interact side-to-side :



$$\int \psi_1^*(\vec{r}) H \psi_2(\vec{r}) d\vec{r} = V_{pp\pi}$$

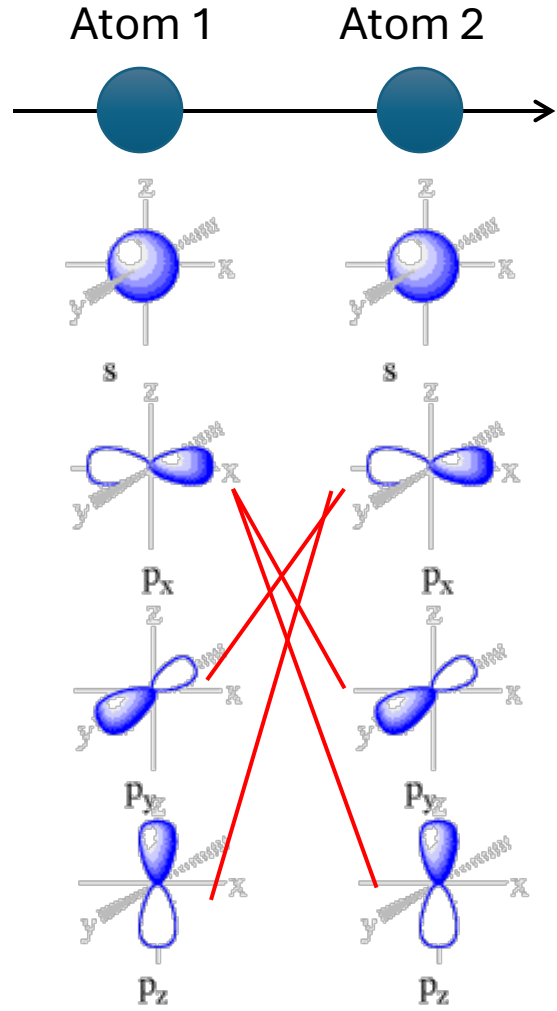
Two p orbitals combine to form  $\pi$  and  $\pi^*$  molecular orbitals:



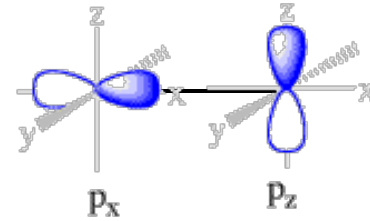


# Diatomic molecules with p orbitals

The p valence electrons (e.g., C, N, O) also participate in bonding



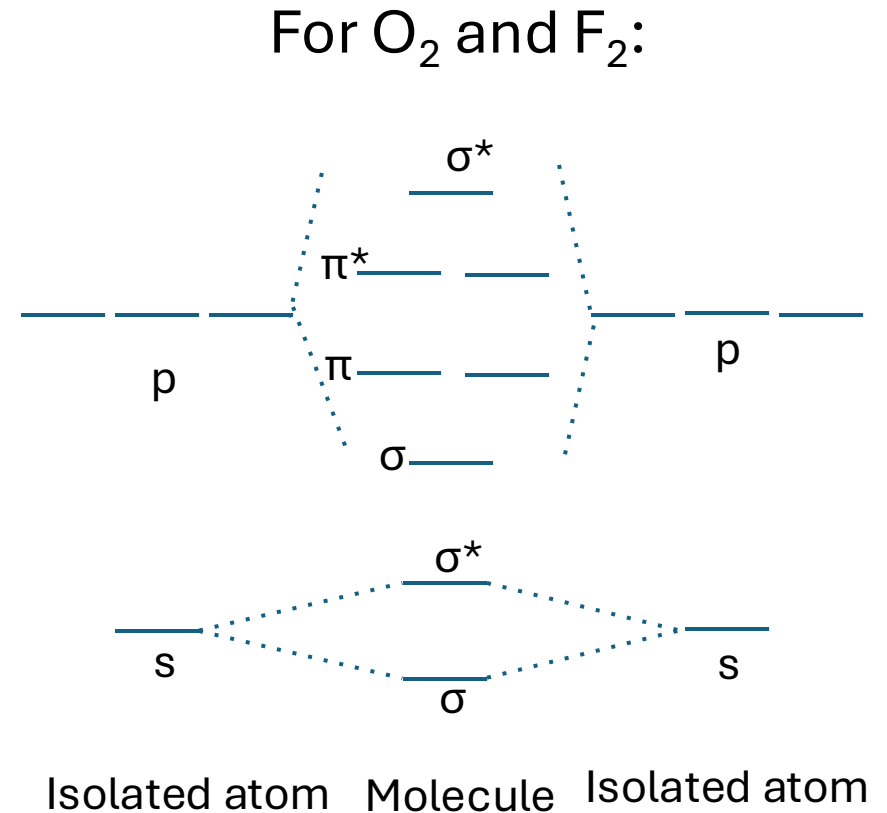
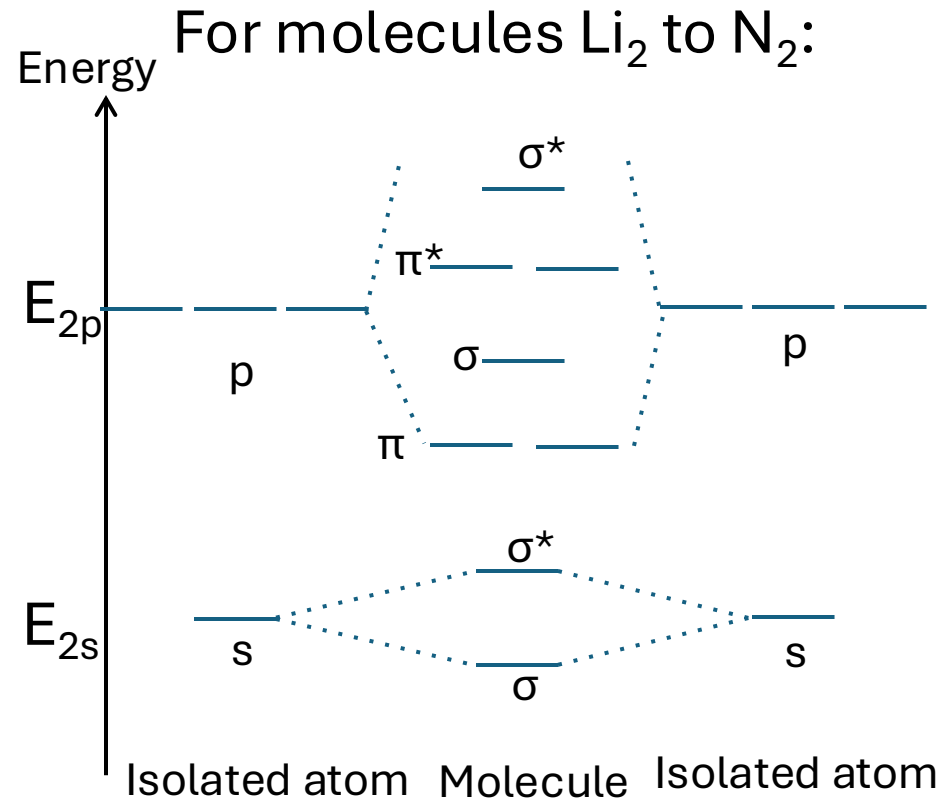
For the tip-to-side p-orbital combination:



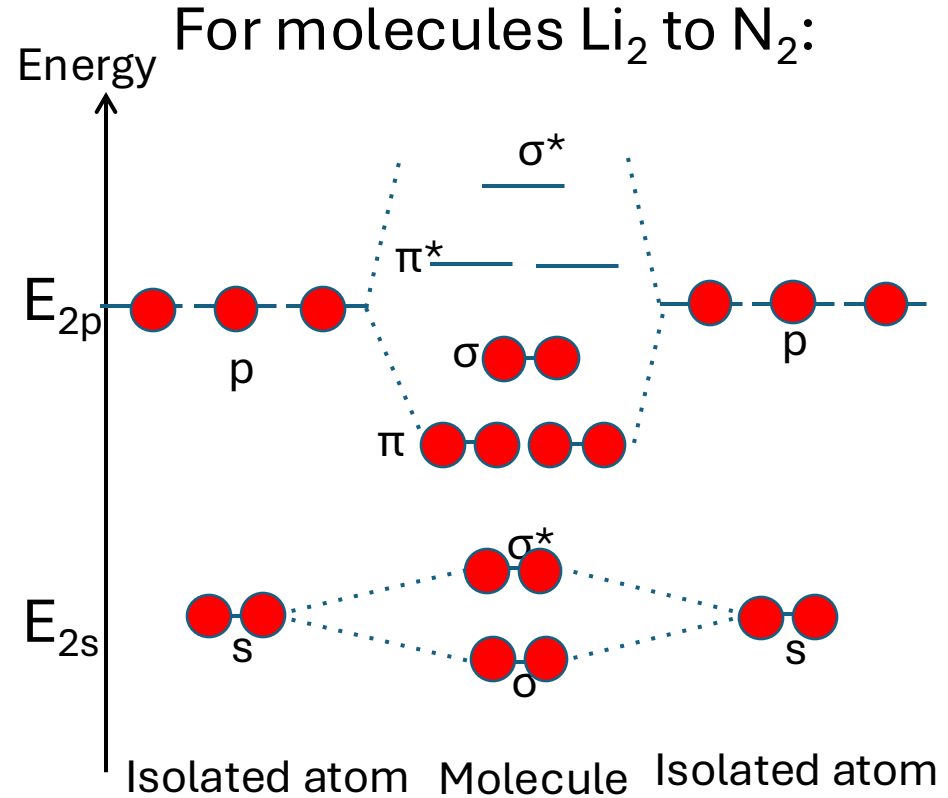
$$\int \psi_1^*(\vec{r}) H \psi_2(\vec{r}) d\vec{r} = 0$$

The product of the two orbitals is an odd function with respect to z, integral is zero

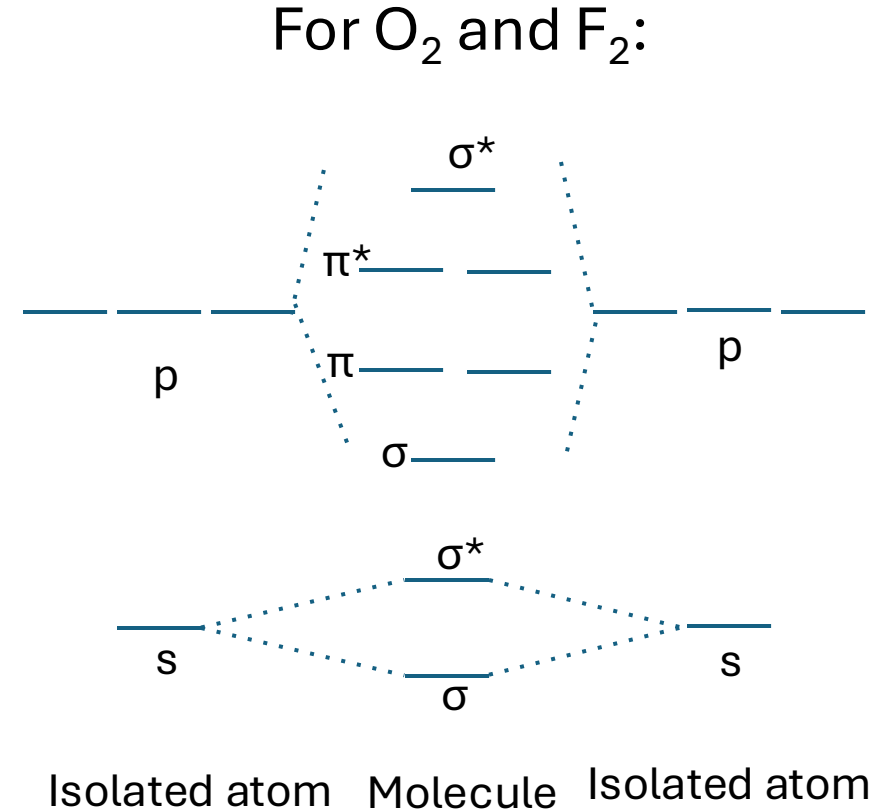
# Energies of molecular orbitals



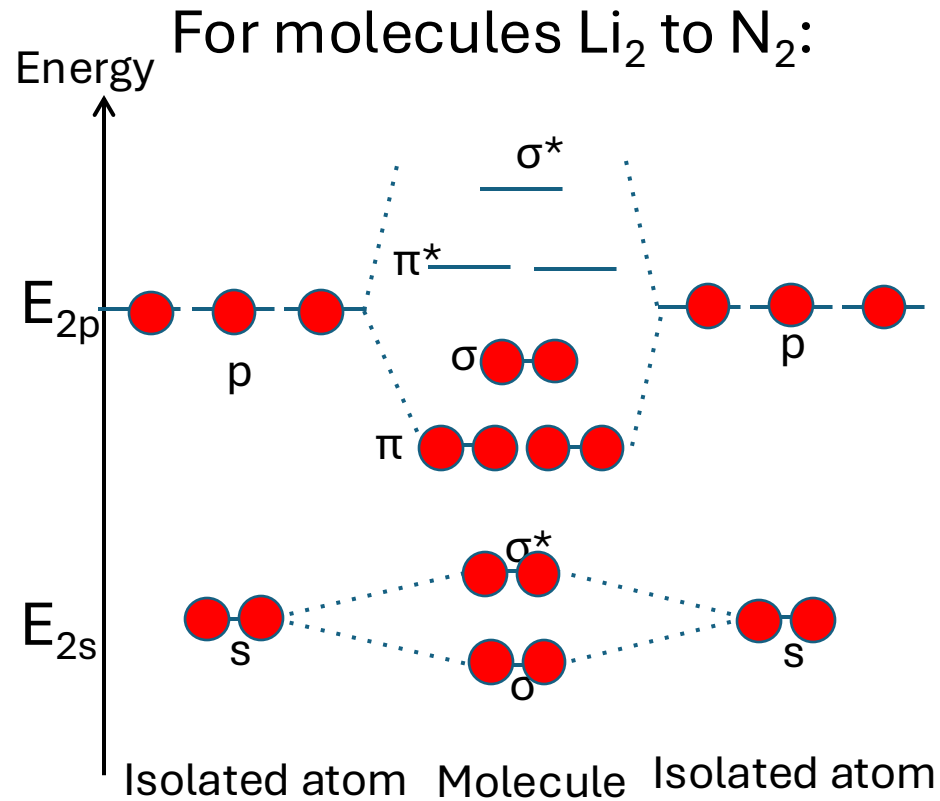
# Exercise: occupations of the molecular orbitals of N<sub>2</sub>



Answer



# Question

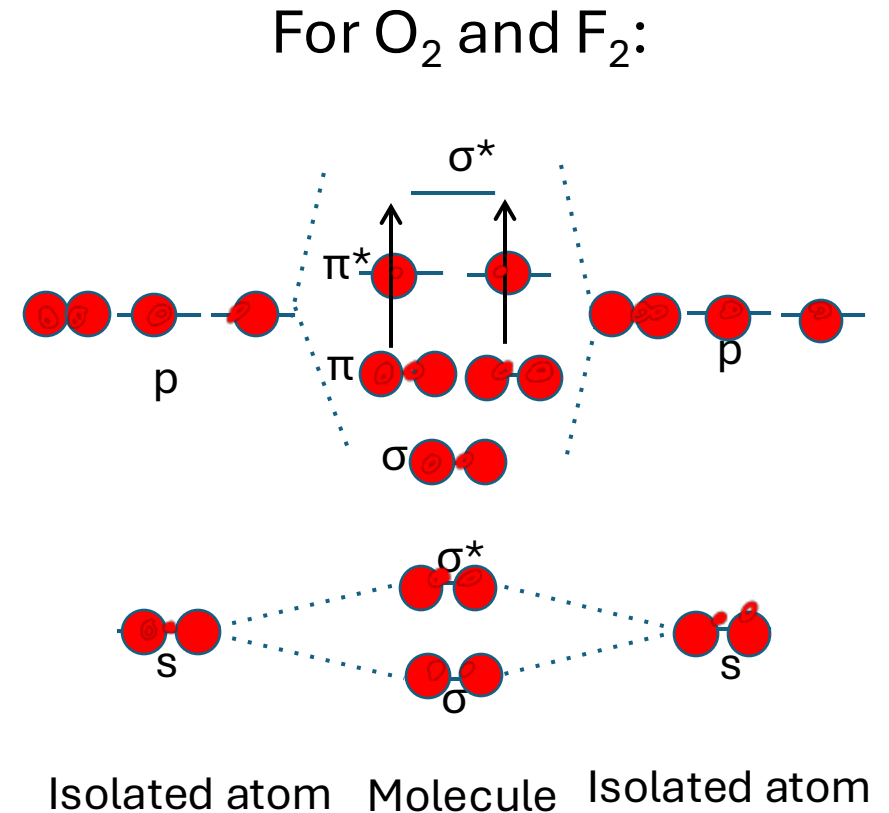
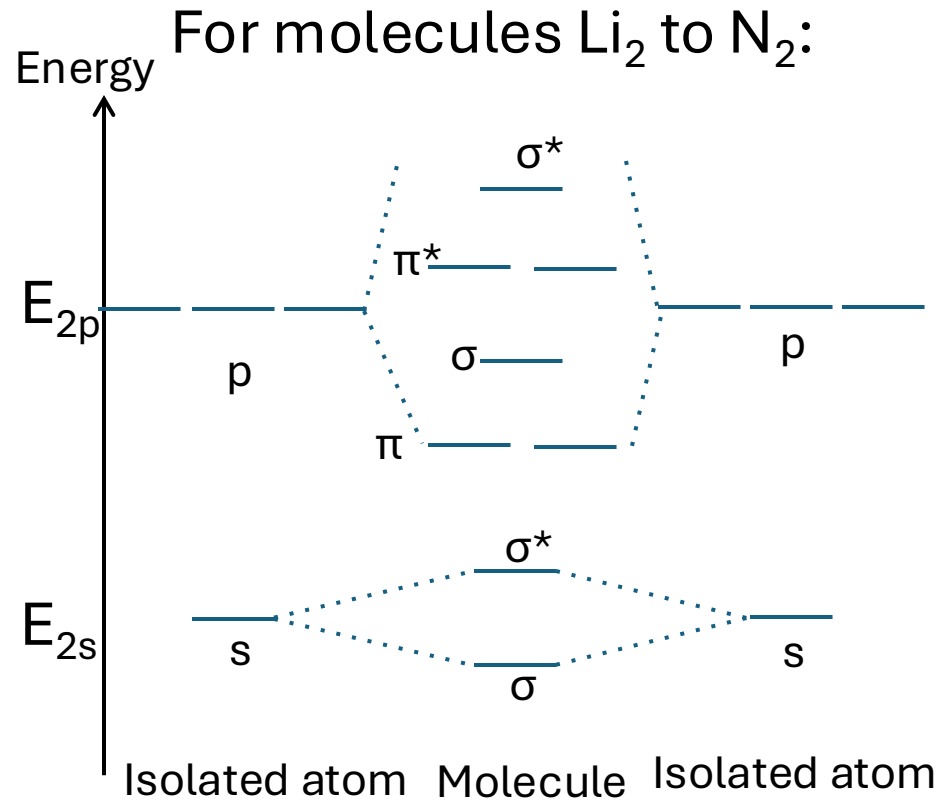


What is the nature of the bond of  $\text{N}_2$ ?

- A. Single
- B. Double
- C. Triple
- D. Quadruple

Answer: C. Three pairs of electrons lower their energy by occupying molecular orbitals

# In-class exercise: find the occupations of the molecular orbitals of $O_2$



Answer

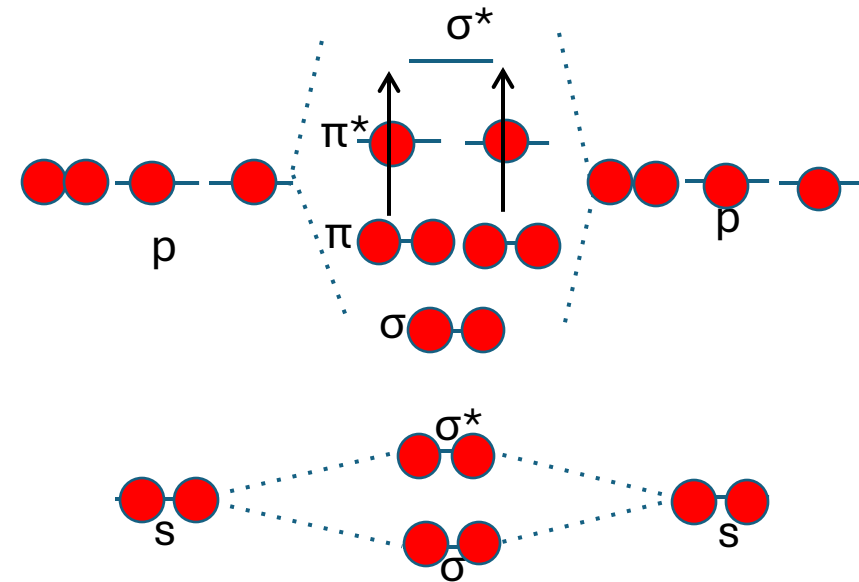
# Question:

What is the nature of the bond of  $O_2$ ?

- A. Single
- B. Double
- C. Triple
- D. Quadruple

Answer: B. The pair of electrons occupying  $\pi^*$  cancel out the energy gain by one of the  $\pi$  pairs

For  $O_2$  and  $F_2$ :



Isolated atom    Molecule    Isolated atom

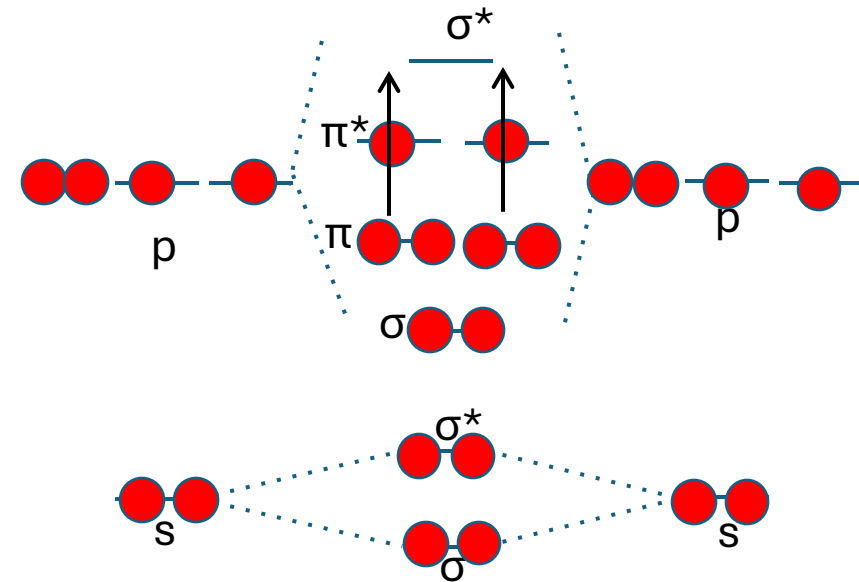
# Question:

Is  $O_2$  really magnetic?

Answer:

<https://www.youtube.com/watch?v=Lt4P6ctf06Q>

For  $O_2$  and  $F_2$ :



Isolated atom    Molecule    Isolated atom



# Hybridization

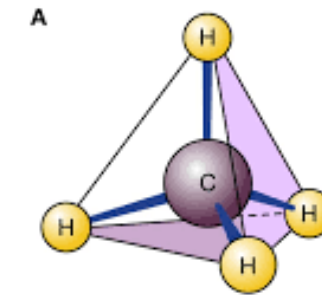
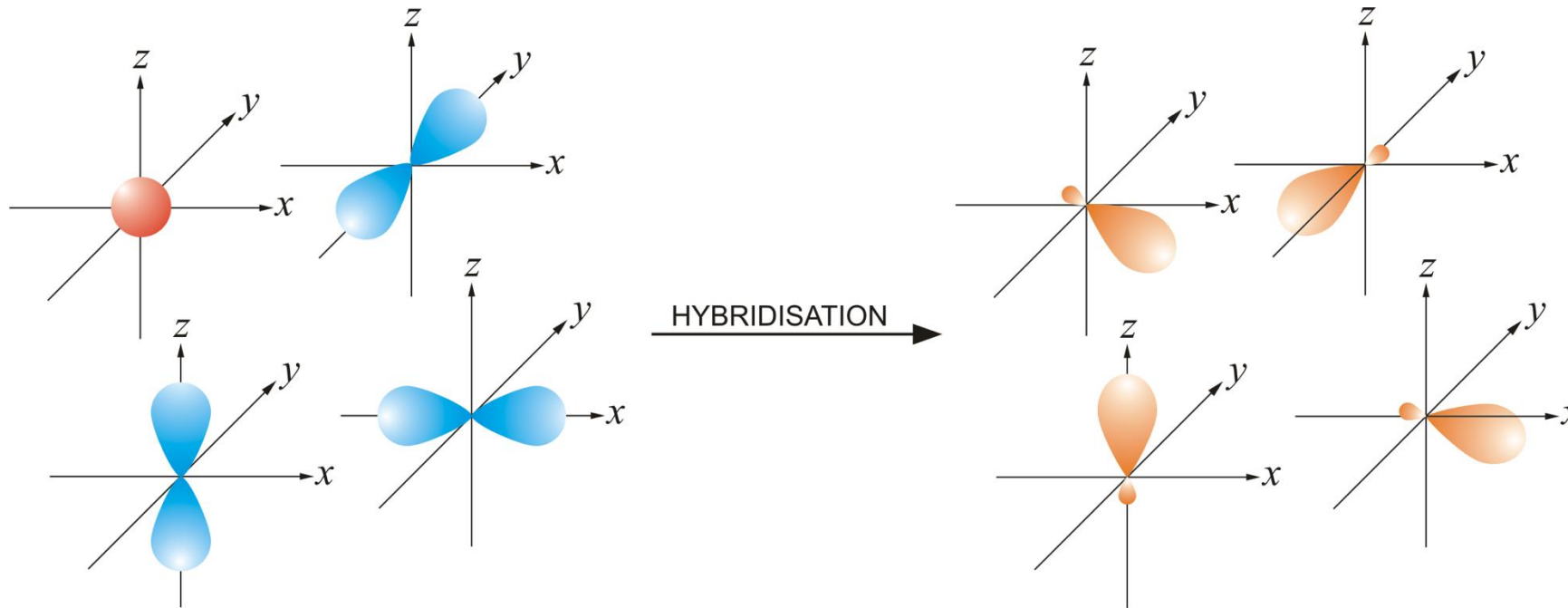
Mix s and p orbitals of the same atom to create new types of hybrid orbitals. More directional, stronger bonding.

<http://glossary.periodni.com/glossary.php?en=sp3+hybrid+orbital>

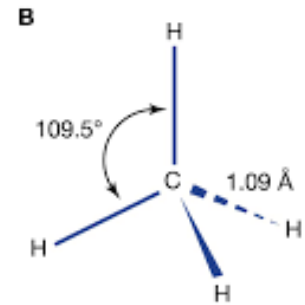
**sp<sup>3</sup> hybridization:** mix all 3 p orbitals with s.

Directional orbitals: more overlap, larger interatomic matrix element, stronger bonds.

Example: CH<sub>4</sub>



© Encyclopædia Britannica, Inc.

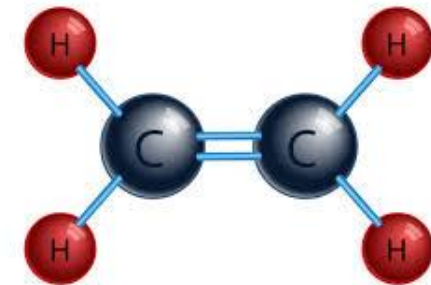
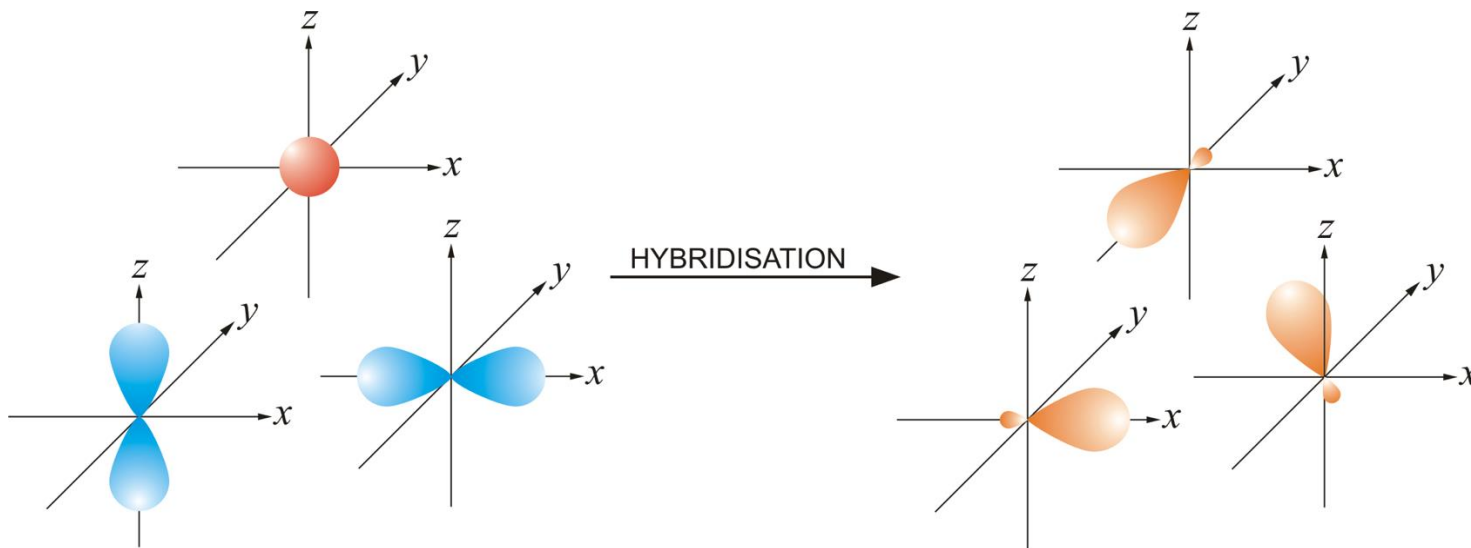


# sp<sup>2</sup> hybridization

Mix s with only two of the three p orbitals, create planar directional molecules.

<http://glossary.periodni.com/glossary.php?en=sp2+hybrid+orbital>

Example: ethene C<sub>2</sub>H<sub>4</sub>

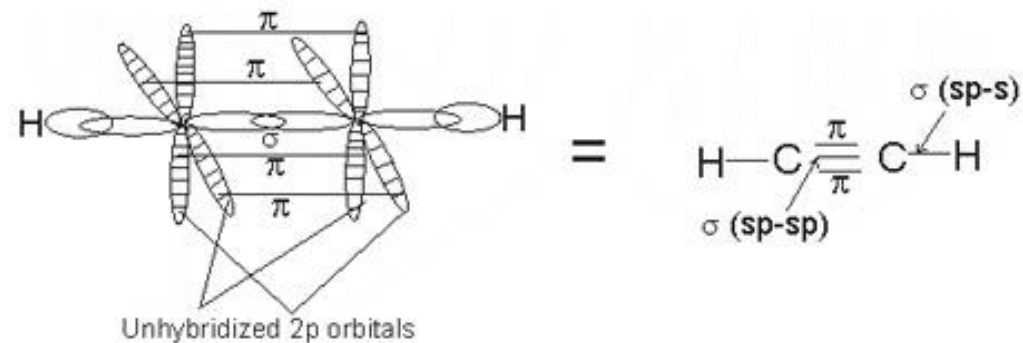
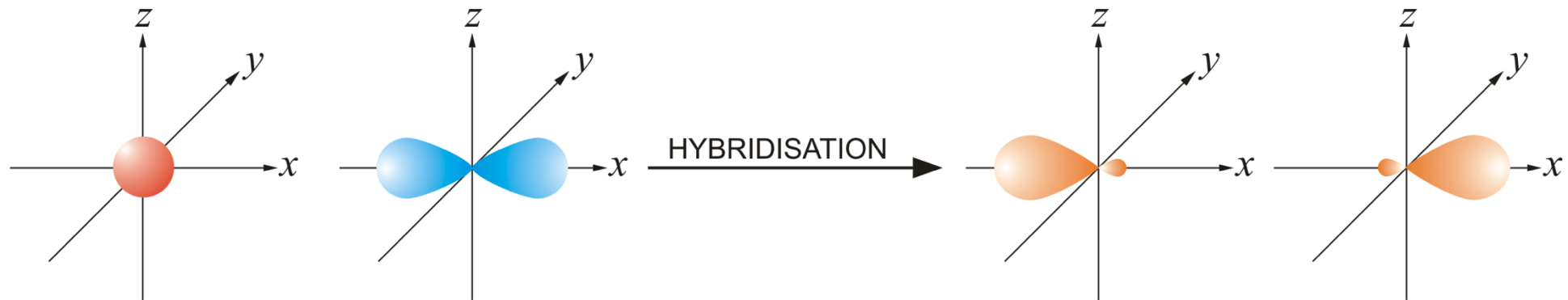


# sp hybridization

Mix s with only one p orbital, create linear bonds.

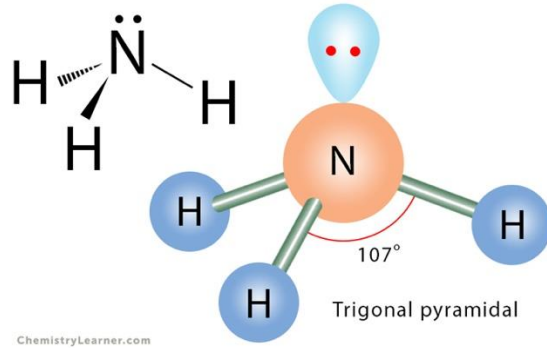
<http://glossary.periodni.com/glossary.php?en=sp+hybrid+orbital>

Example: acetyline (ethyne)  $\text{C}_2\text{H}_2$

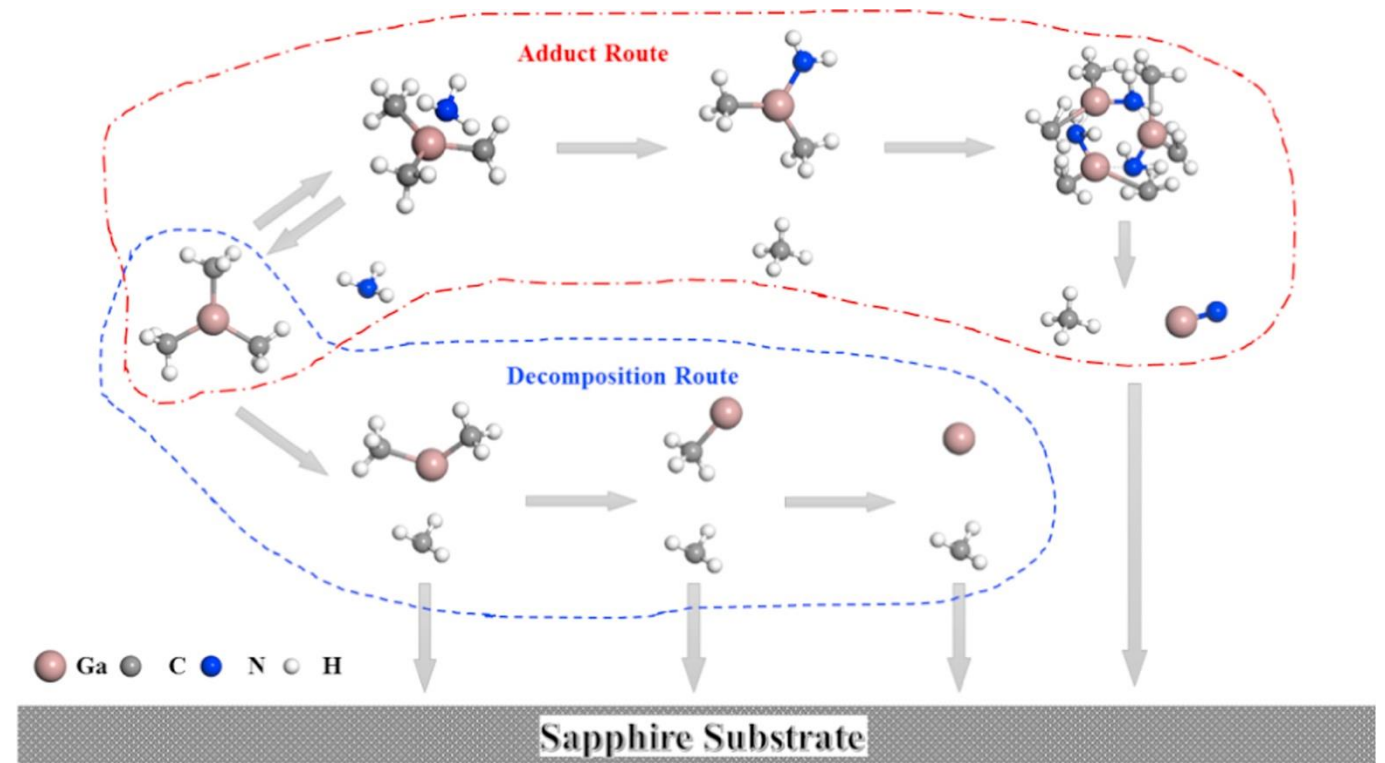
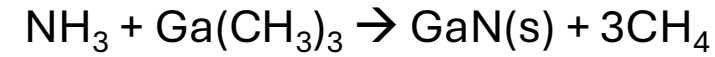
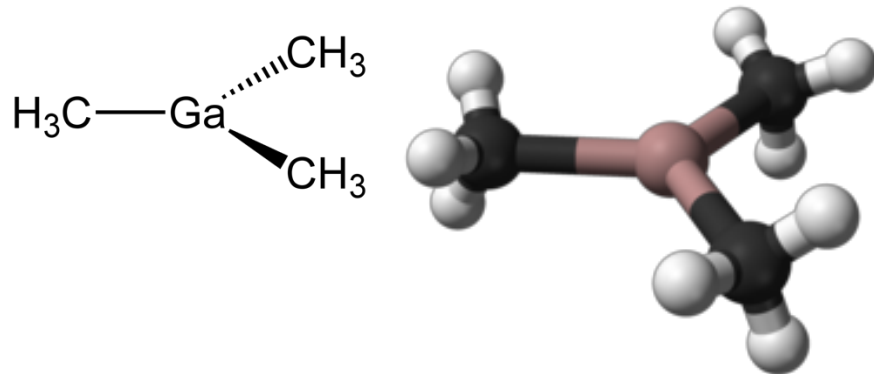


# Precursors for the MOCVD synthesis of nitrides

Ammonia:  $\text{NH}_3$



Trimethylgallium:  $\text{Ga}(\text{CH}_3)_3$



<https://doi.org/10.1016/j.jcrysgro.2016.09.010>

# Conjugation

Alternate single-double bonds

Example:  $\text{CH}_2=\text{CH}-\text{CH}=\text{CH}_2$

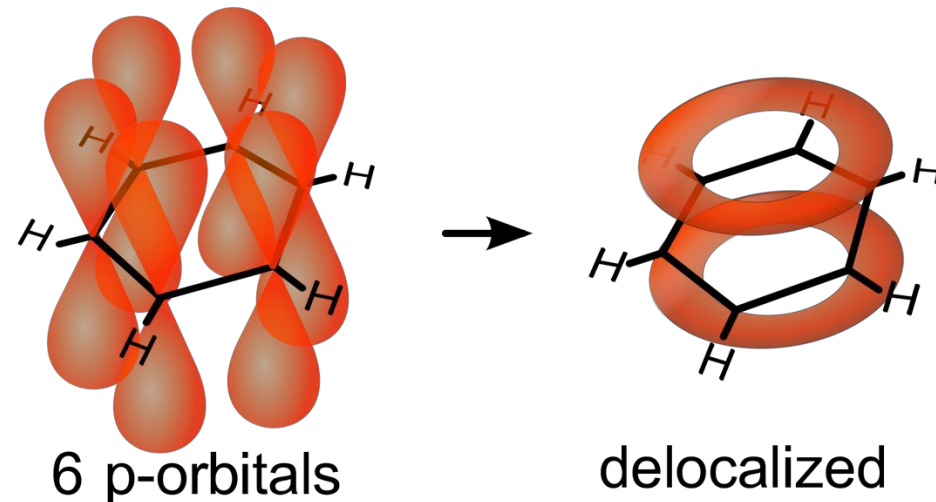
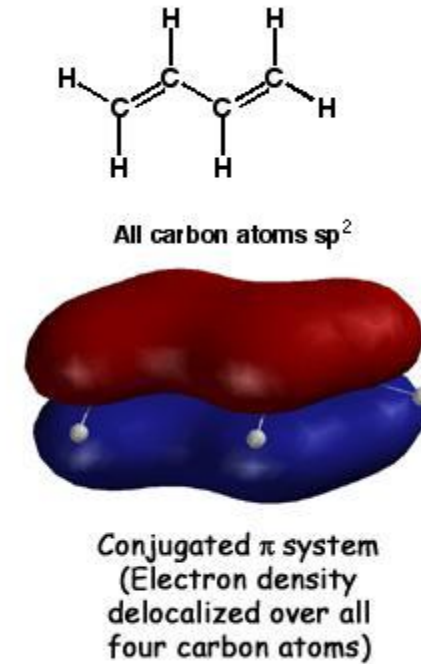
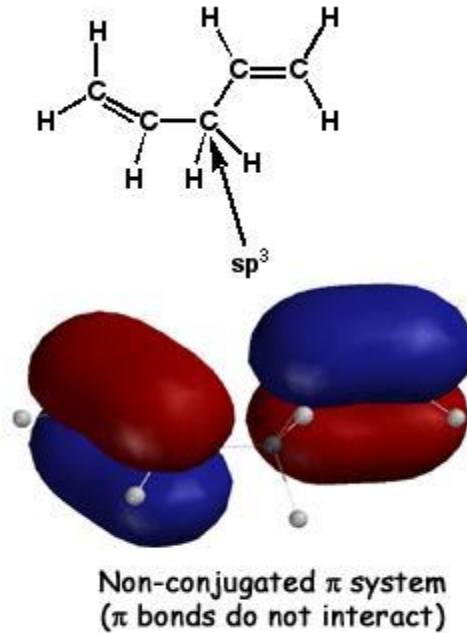
Example: benzene

Conjugated wave functions extend over the entire volume of the molecule.

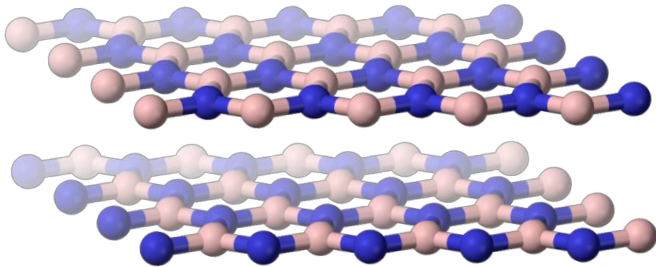
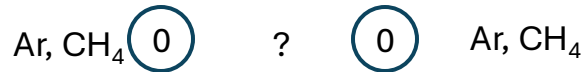
Applications: organic electronics.

Wave functions of extended systems (e.g., nitrides) are also extended over the entire volume of the material.

<https://instruct.uwo.ca/chemistry/373f/Nifty%20Stuff/ethylene.htm>



# Van der Waals interaction



For neutral, non-polar objects (Ar,  $\text{CH}_4$ ):

Expectation value of electric dipole moment:

$$\langle \vec{d}_1 \rangle = \langle \vec{d}_2 \rangle = 0$$

→ No static dipole-dipole interaction

But: from quantum mechanics:  $\langle \vec{d}_1 \cdot \vec{d}_2 \rangle \neq 0$

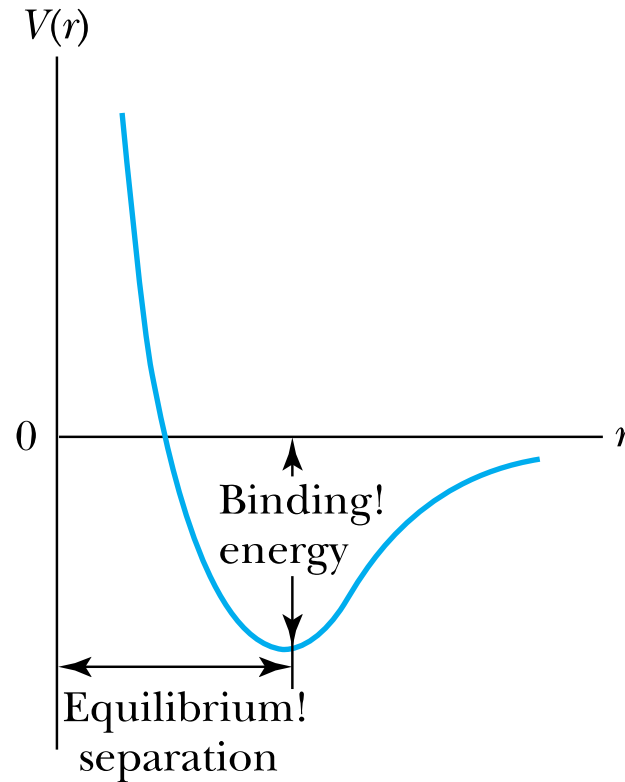
→ Dynamic dipole-dipole interaction

The van der Waals interaction is a purely quantum mechanical interaction. Always attractive.

$$V(r) = -\frac{A}{r^6}$$

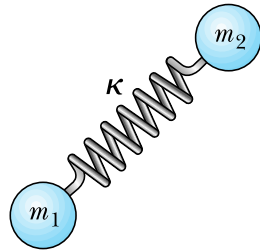
E.g., van der Waals bonding in 2D materials (graphite, hexagonal BN, etc.)

# Binding energy and equilibrium separation



Typical values for the equilibrium distance and for the binding energy:  $R_0 \approx 10^{-10} \text{ m}$ ,  $E_0 \approx 1 \text{ eV}$

# Vibrational spectra of molecules

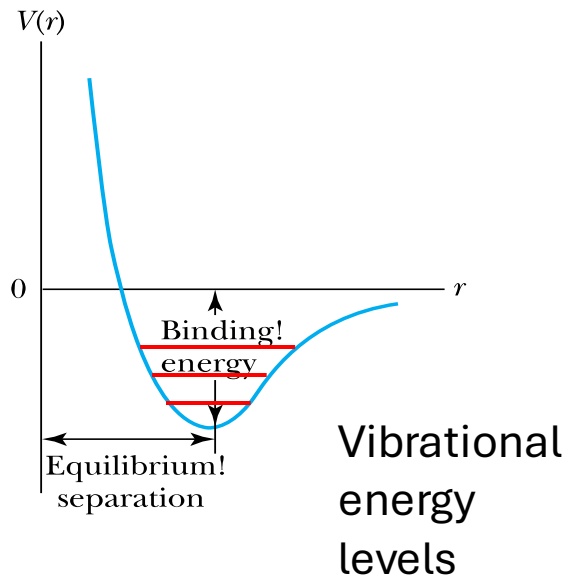


**Figure 10.5** A model of a diatomic molecule with the two masses  $m_1$  and  $m_2$  connected by a massless spring with force constant  $\kappa$ .

$$E_{\text{vibr}} = \left(n + \frac{1}{2}\right) \hbar \omega$$

$$\omega = \sqrt{\frac{\kappa}{\mu}}$$

$$\mu = m_1 m_2 / (m_1 + m_2)$$



Exercise: what is the vibrational frequency of HCl if  $k = 480 \text{ N/m}$ ?

$$m_{\text{H}} = 1 \text{ u}, m_{\text{Cl}} = 35.5 \text{ u}, \text{u} = 1.66054 \times 10^{-27} \text{ kgr}$$

**Answer:**

$$\mu = m_{\text{H}} m_{\text{Cl}} / (m_{\text{H}} + m_{\text{Cl}}) = 1.63 \times 10^{-27} \text{ kgr}$$

$$\omega = 5.43 \times 10^{14} \text{ rad/s}, \hbar \omega = 0.36 \text{ eV}$$

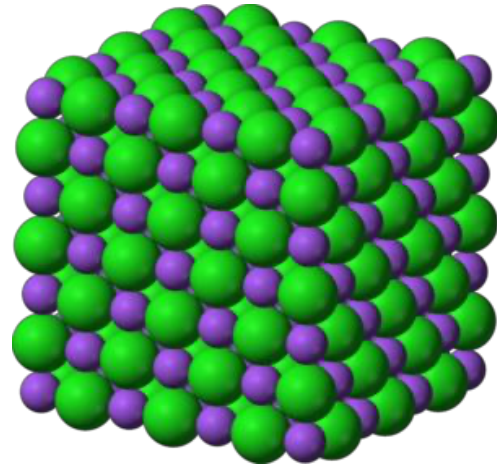
$$f = \omega / 2\pi = 8.64 \times 10^{13} \text{ Hz}$$



# Crystals

**Crystal** = regular periodic arrangement of atoms in a material

As opposed to **amorphous** materials: no long-range order



Question: give examples of crystalline materials:

Answer: salt, sugar, metals, silicon, nitrides,  
all materials at low T

Question: what evidence do we have that materials are crystalline? Answer: X-ray and electron diffraction

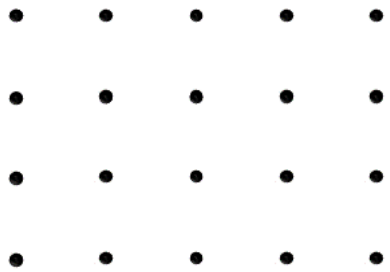
Question: Why are materials crystalline?

Answer: a crystal arrangement minimizes the energy

Question: Why?

Answer: we do not know

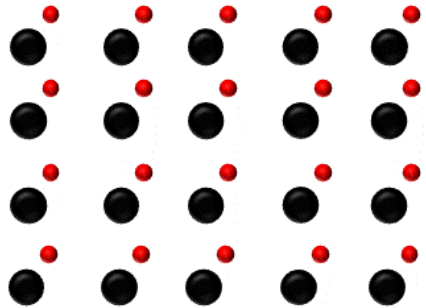
# Crystal structure = space lattice + atomic basis



Space Lattice

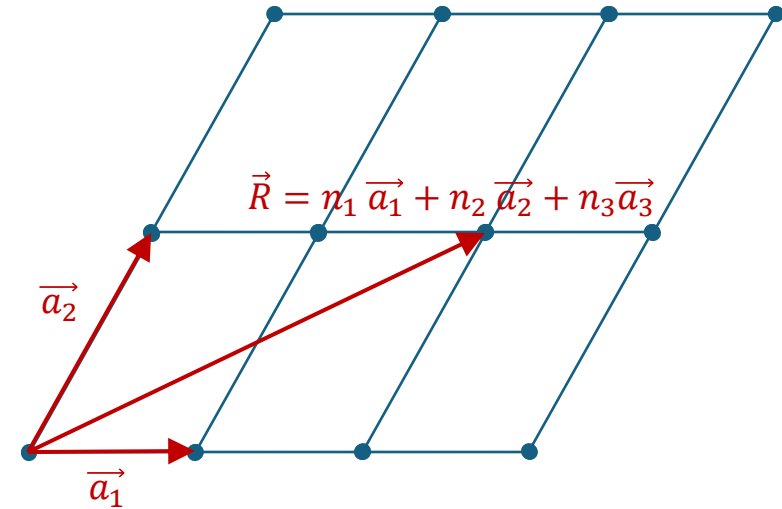


Basis (atoms)



Crystal structure

The crystal structure is formed by adding basis (atoms) to every lattice points of the lattice. The number of atoms in the basis may be one or more than one.



Space lattice = mathematical translation of the crystal that leaves the material unchanged.

A crystal has translational symmetry: if we perform a displacement by a vector  $\vec{R}$ , the potential is identical

$$V(\vec{r} + \vec{R}) = V(\vec{r})$$

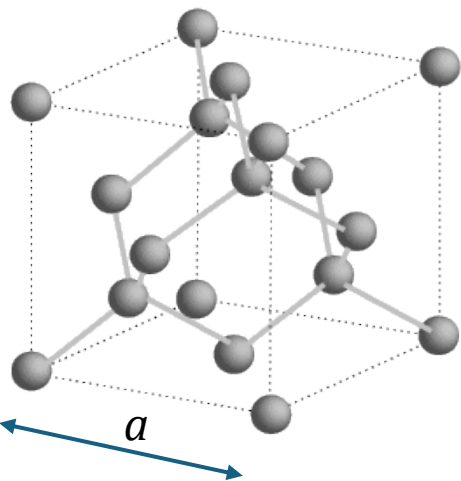
Lattice vectors  $\vec{R} = n_1\vec{a}_1 + n_2\vec{a}_2 + n_3\vec{a}_3$

$n_1, n_2, n_3$  are integers

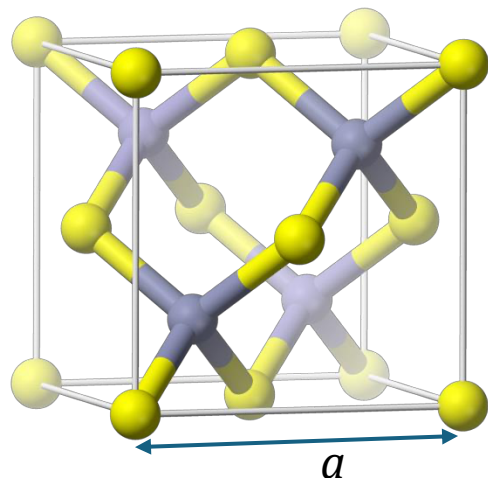
$\vec{a}_1, \vec{a}_2, \vec{a}_3$  are the basis vectors

Each point of the lattice is given by a set of integers, and each set of integers gives a point on the lattice.

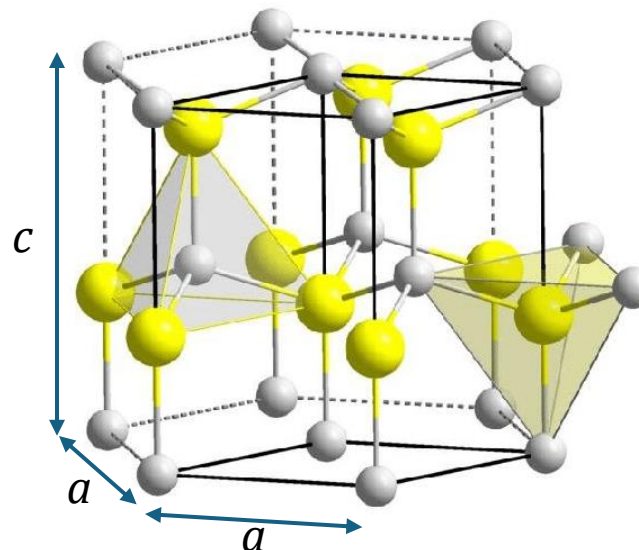
# Crystal structures of semiconductors



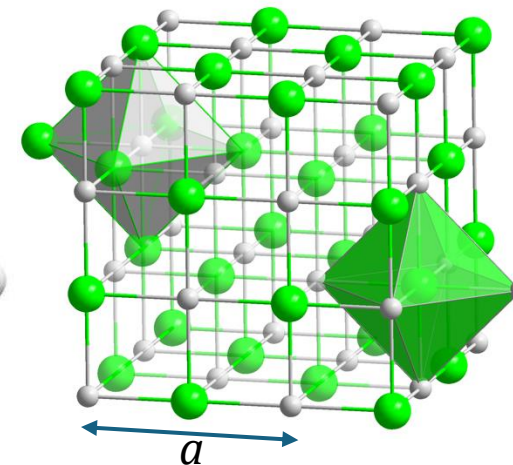
Silicon, germanium:  
diamond structure



GaAs: zincblende (or  
sphalerite) structure



GaN, AlN, InN:  
wurtzite structure



ScN, YN, LaN:  
rocksalt structure

Based on *fcc* lattice:

Lattice vectors:

$$\begin{aligned}\vec{a}_1 &= \frac{a}{2}(\hat{x} + \hat{y}) \\ \vec{a}_2 &= \frac{a}{2}(\hat{y} + \hat{z}) \\ \vec{a}_3 &= \frac{a}{2}(\hat{x} + \hat{z})\end{aligned}$$

Basis:

$$\begin{aligned}\vec{\tau}_1 &= 0 \\ \vec{\tau}_2 &= \frac{a}{4}(\hat{x} + \hat{y} + \hat{z})\end{aligned}$$

Based on *hcp* lattice:

$$\begin{aligned}\vec{a}_1 &= a\hat{x} \\ \vec{a}_2 &= \left(\frac{a}{2}\hat{x} + \frac{\sqrt{3}a}{2}\hat{y}\right) \\ \vec{a}_3 &= c\hat{z}\end{aligned}$$

$$\text{Ga: } \vec{\tau}_1 = 0$$

$$\text{Ga: } \vec{\tau}_2 = \frac{1}{3}(\vec{a}_1 + \vec{a}_2) + \frac{1}{2}\vec{a}_3$$

$$\text{N: } \vec{\tau}_3 = u\hat{z}$$

$$\text{N: } \vec{\tau}_4 = \frac{1}{3}(\vec{a}_1 + \vec{a}_2) + (u + \frac{1}{2})\vec{a}_3$$

Based on *fcc* lattice:

$$\begin{aligned}\vec{a}_1 &= \frac{a}{2}(\hat{x} + \hat{y}) \\ \vec{a}_2 &= \frac{a}{2}(\hat{y} + \hat{z}) \\ \vec{a}_3 &= \frac{a}{2}(\hat{x} + \hat{z})\end{aligned}$$

$$\vec{\tau}_1 = 0$$

$$\vec{\tau}_2 = \frac{a}{2}\hat{x} \text{ [or } \vec{\tau}_2 = \frac{a}{2}(\hat{x} + \hat{y} + \hat{z})]$$