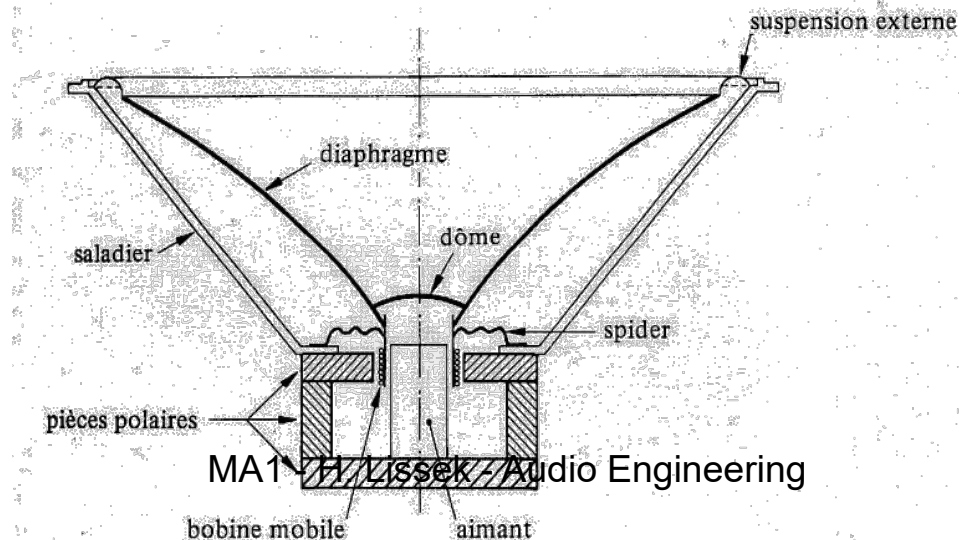


## 5.2 Electrodynamic loudspeaker

# Definitions

Electrodynamic loudspeaker, with moving coil + radiating cone:

- elastically suspended *diaphragm* with *spider* and *external suspension*
- *moving coil*, attached to the diaphragm
- *magnetic circuit* of the driver
- *basket*= open frame (made of metal)



# Radiation

If radiation through a horn: *indirect radiation*

Else: *direct radiation*

In a general sense, loudspeaker on infinite baffle  
(piston behavior):

- after  $f$  and dimensions (w. resp.  $/\lambda$ ),  
loudspeaker= small pulsating source in  $2\pi$  sr
- identical velocity at each point of diaphragm:

$$\underline{q}_d = S_d \underline{v}_d$$

# Bandwidth

zone of  $f$  where the sound pressure generated by the loudspeaker is constant

LF limit:

- omnidirectional radiation  $\rightarrow$  constant  $P_a$   
since  $P_a = R_{ar} q_d^2 \rightarrow q_d$  shall vary in  $1/f$
- mobile system=resonator, force independant of  $f$   
 $\rightarrow$  flow velocity varies in  $1/f$  if controlled by mass  
 $\rightarrow$  low limit of  $B$  = resonance frequency of the mobile system

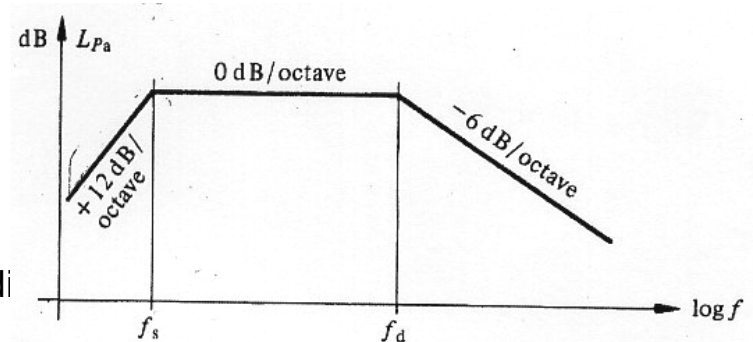
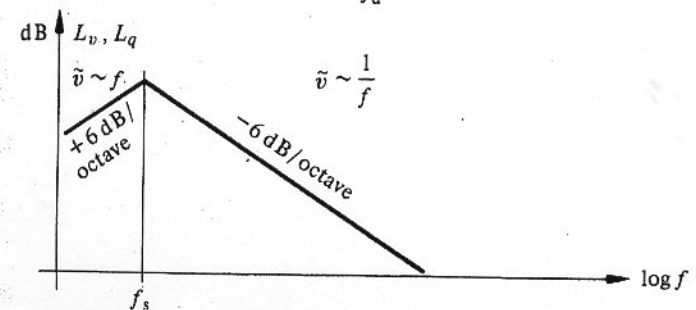
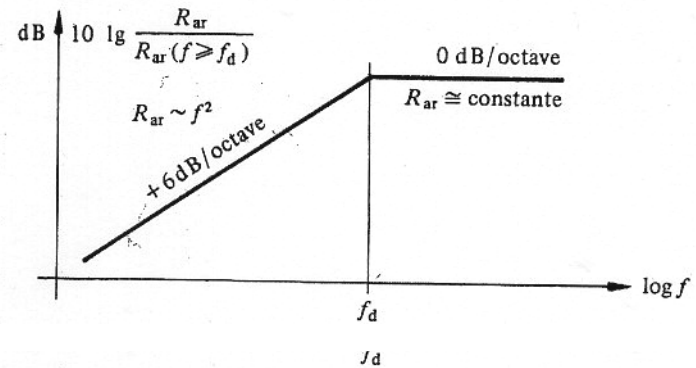
# Bandwidth

HF limit:

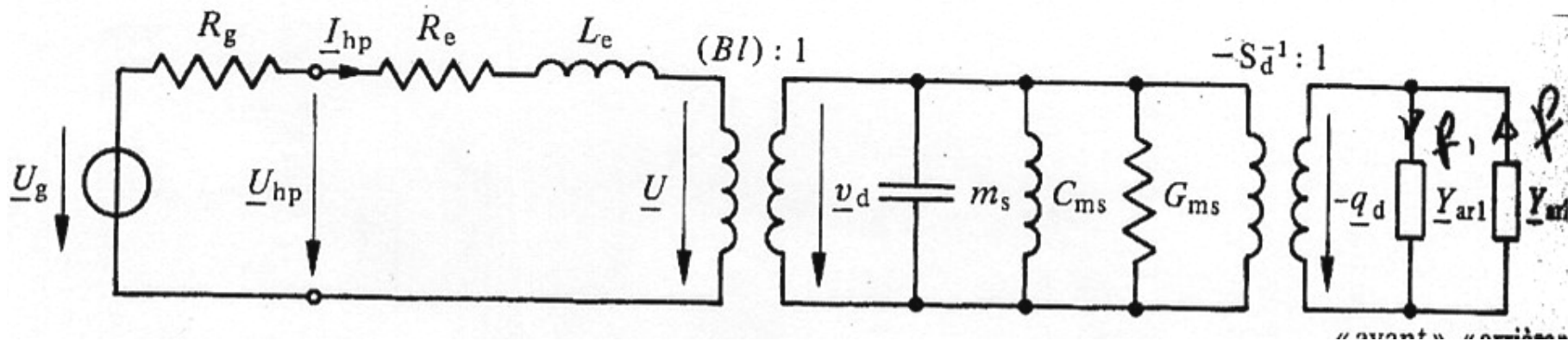
- $ka$  high enough: constant radiation resistance
- since  $q_d$  varies in  $1/f$

➔ HF limit:  $ka = \sqrt{2}$

$$f_d = \frac{c}{\sqrt{2}\pi a}$$



# Scheme of a loudspeaker in infinite baffle



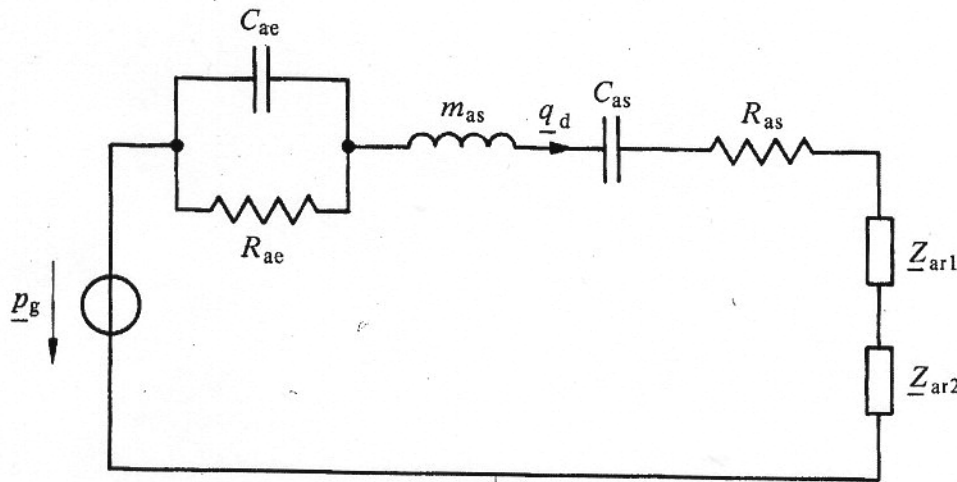
# Equivalent **acoustic** scheme

suppression of 2-ports by denoting:

$$\underline{p}_g = \frac{Bl}{S_d [R_g + R_e + j\omega L_e]} \underline{U}_g$$

$$R_{ae} = \frac{(Bl)^2}{S_d^2 (R_g + R_e)}$$

$$C_{ae} = \frac{L_e S_d^2}{(Bl)^2}$$



$$R_{as} = \frac{R_{ms}}{S_d^2}$$

$$m_{as} = \frac{m_s}{S_d^2}$$

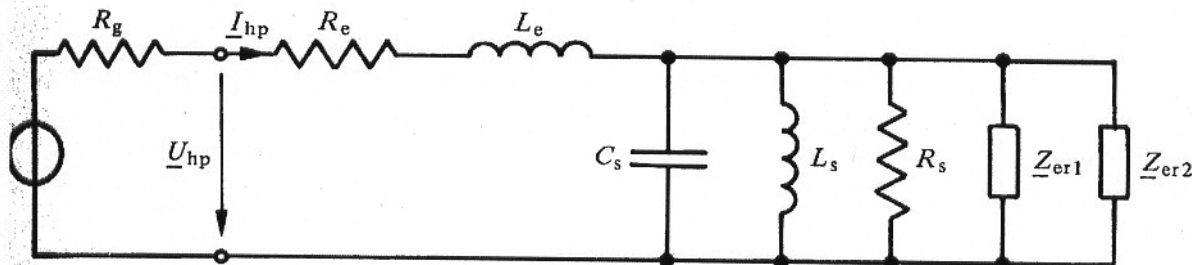
$$C_{as} = S_d^2 C_{ms}$$

# Equivalent **electric** scheme

by denoting:  $L_s = C_{ms} (Bl)^2 = \frac{C_{as}}{S_d^2} (Bl)^2$   $\underline{Z}_{er} = \frac{(Bl)^2}{S_d^2 \underline{Z}_{ar}}$

$$C_s = \frac{m_s}{(Bl)^2} = \frac{m_{as}}{(Bl)^2} S_d^2$$

$$R_s = \frac{(Bl)^2}{R_{ms}} = \frac{(Bl)^2}{S_d^2 R_{as}}$$



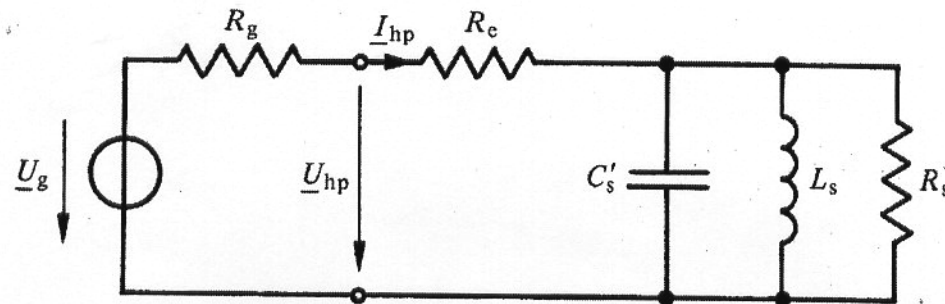


# Low frequency approximation

below resonance frequency  $f_s$ :

- $(\omega L_e) \ll R_e$
- $2R_{ar} < R_{as}$
- $X_{ar}$  = mass  $m_r$ , added to the acoustic masses:

$$m'_{as} = \frac{m_s + 2m_r}{S_d^2} = m_{as} + 2m_{ar}$$



# Small signals parameters

In the LF range,  
small signals (or Thiele/Small) parameters :

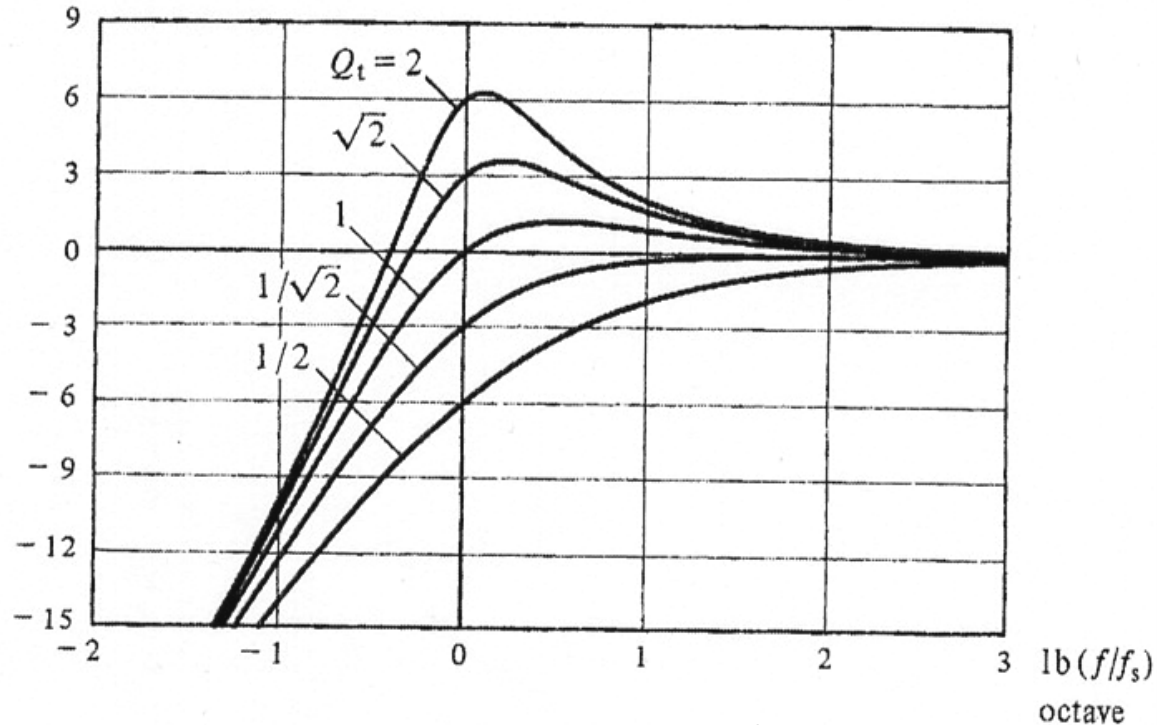
- resonance frequency  $f_s$  
$$f_s = \frac{1}{2\pi\sqrt{m'_s C_{ms}}} = \frac{1}{2\pi\sqrt{m'_{as} C_{as}}} = \frac{1}{2\pi\sqrt{C'_s L_s}}$$
- air volume equivalent to  $C_{as}$  
$$V_{as} = \rho c^2 C_{as}$$
- electrical quality factor  $Q_{es}$  at  $f_s$  
$$Q_{es} = \omega_s C'_s R_e = \frac{1}{\omega_s C_{as} \underbrace{R'_{ae}}_{R_{ae}|_{R_g=0}}}$$
- mechanical quality factor  $Q_{ms}$  at  $f_s$  
$$Q_{ms} = \omega_s C'_s R_s = \frac{R_s}{\omega_s L_s} = \frac{1}{\omega_s C_{as} R_{as}}$$
- total quality factor  $Q_{ts}$  at  $f_s$  
$$Q_{ts} = \frac{Q_{es} \cdot Q_{ms}}{Q_{es} + Q_{ms}}$$
- $S_d, R_e, L_e, Bl$

# Diaphragm volume flow

cf acoustic scheme

leads to:

$20 \lg |G_s|, \text{ dB}$



$$\underline{q}_d = \left[ \underbrace{\frac{S_d}{Q_e Bl} U_g}_{\underline{q}_s} \right] \frac{1}{j\omega / \omega_s} \underbrace{\frac{(j\omega / \omega_s)^2}{(j\omega / \omega_s)^2 + Q_t^{-1}(j\omega / \omega_s) + 1}}_{G_s(j\omega / \omega_s)}$$

# Acoustic power and pressure

Radiated power is  
then

$$P_a = R_{ar} \tilde{q}_d^2$$

$$P_a = \underbrace{\frac{2\pi\rho}{c} \cdot f_s^2 \cdot \tilde{q}_s^2}_{P_{as}} \cdot |G_s(j\omega / \omega_s)|^2$$

Sound pressure (omnidirectional radiation in  $2\pi$  sr)

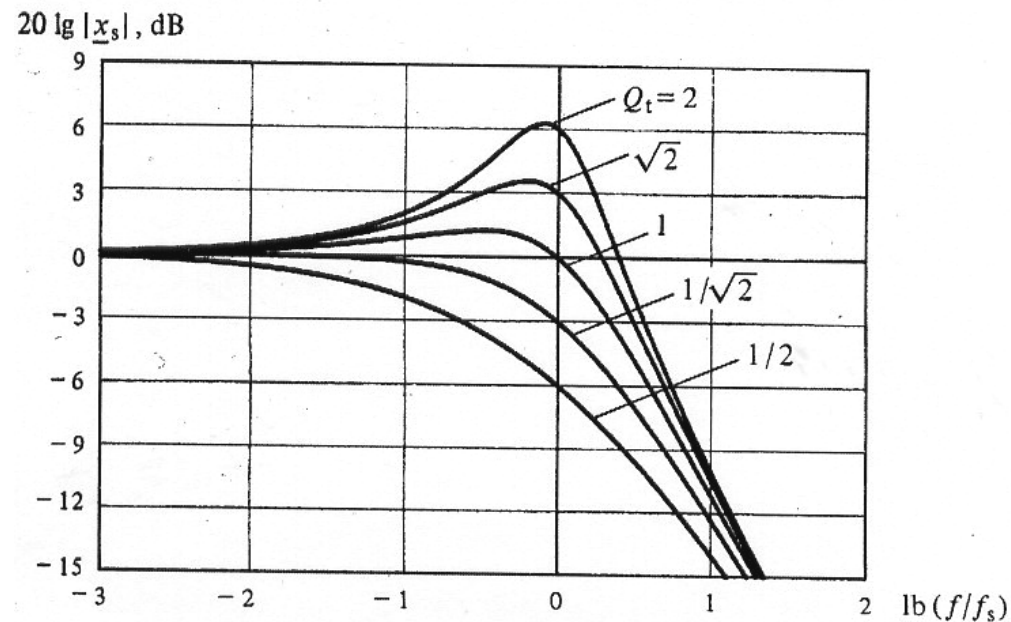
$$\tilde{p}(r) = \sqrt{Z_c \underbrace{\frac{P_a}{2\pi r^2}}_{I(r)}} = \tilde{p}_s |G_s(j\omega / \omega_s)|$$

where  $p_s$  is independant of  $f$

# Diaphragm elongation

can be computed by integrating  $q_d$   
(multiply with  $1/j\omega$  in harmonic state)

$$\underline{\xi}_d = \underbrace{\frac{\underline{U}_g}{\omega_s Q_e(Bl)}}_{\underline{\xi}_s} \cdot \underbrace{\frac{G_s(j\omega/\omega_s)}{(j\omega/\omega_s)^2}}_{\underline{x}_s}$$

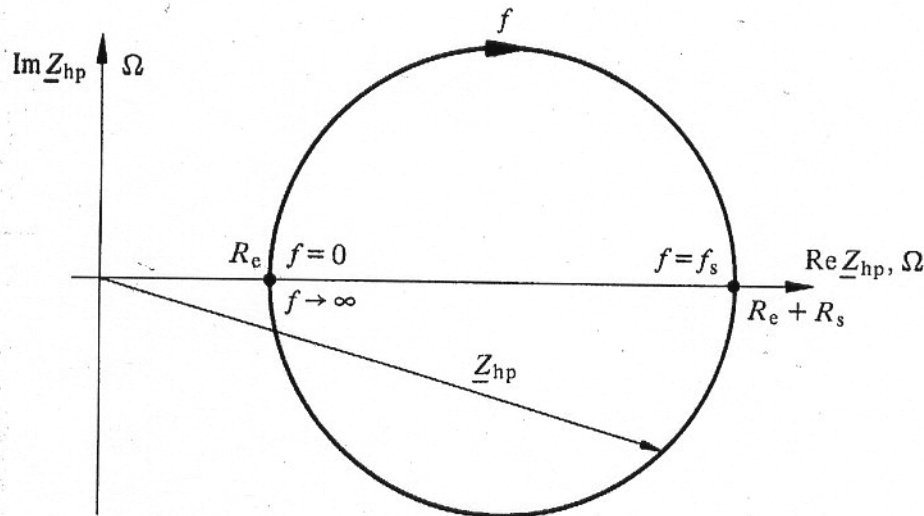


# Input impedance

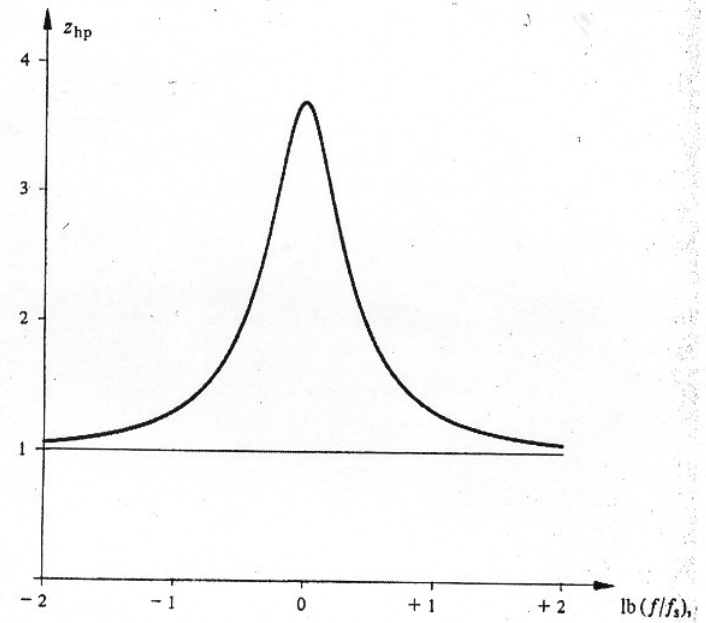
$$\underline{Z}_{HP} = \underline{U}_{HP} / \underline{I}_{HP}$$

At LF:

$$\underline{Z}_{HP} = \frac{\underline{Z}_{HP}}{R_e} = 1 + \underbrace{\left( \frac{Q_{ms}}{Q_{es}} \right) \frac{(j\omega / \omega_s) Q_{ms}^{-1}}{(j\omega / \omega_s)^2 + (j\omega / \omega_s) Q_{ms}^{-1} + 1}}_{\gamma / R}$$



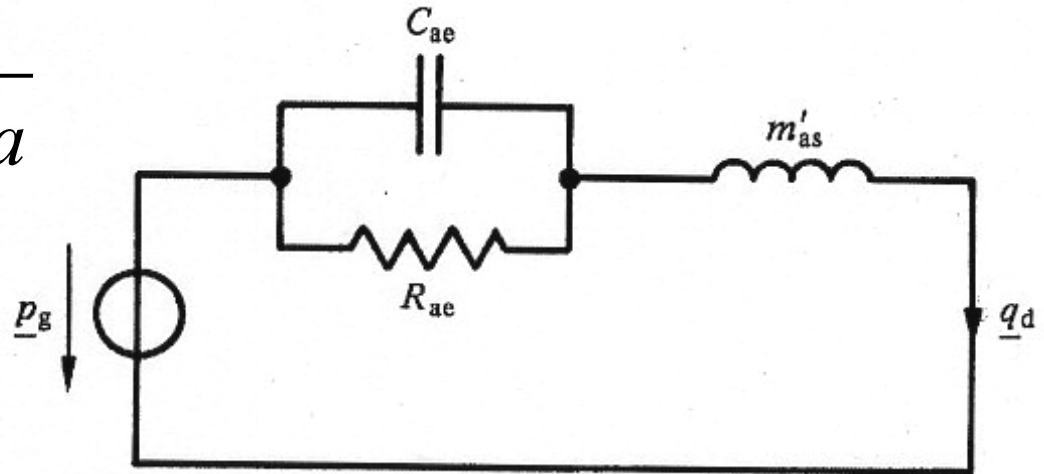
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# Behavior at high frequencies

HF=around  $f_d = \frac{c}{\sqrt{2\pi}a}$   
 $(\omega L_e) \gg R_e$

$\omega m'_{as} \gg R_{as} + 2R_{ar}$



→ resonance for:

$$f_s' = \frac{1}{2\pi} \sqrt{\underbrace{\frac{(Bl)^2}{L_e m'_{as} S_d^2} - \frac{(R_g + R_e)^2}{L_e^2}}_{\text{if positive}}}$$

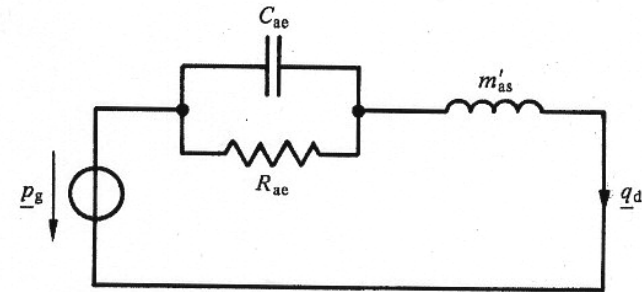
$$\underline{p}_g = \frac{Bl}{j\omega L_e S_d} \bar{U}_g$$

$$m'_{as} = m_{as} + 2 \frac{X_{ar}}{\omega}$$

# Behavior at high frequencies

below  $f''_s$ ,

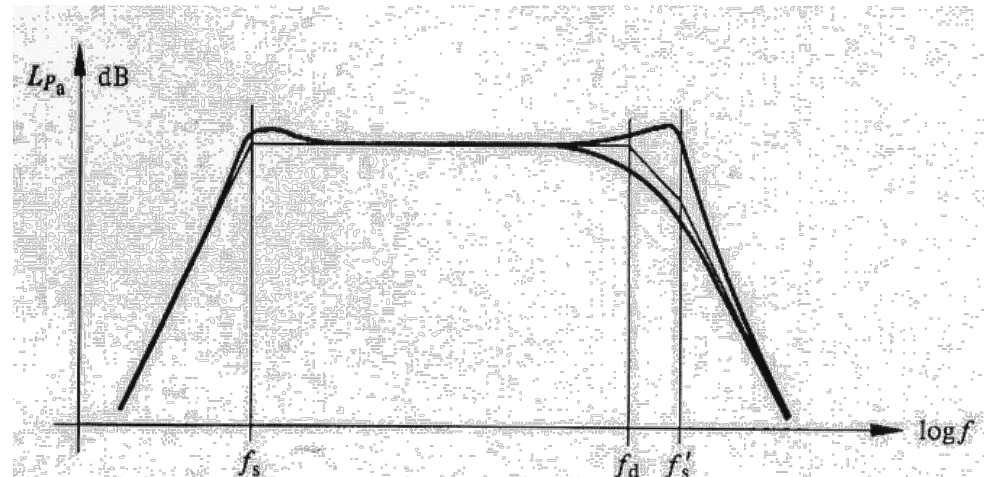
$$\underline{q}_d \cong \frac{\underline{p}_g}{j\omega m'_{as}} = \left( \frac{Bl}{j\omega L_e \cdot j\omega m'_{as} \cdot S_d} \right) \cdot \underline{U}_g$$



decreases in  $1/f^2$

$R_{ar} \sim \text{const.}$  since  $f > f_d$ ,

$P_a$  in  $1/f^4$  (-12 dB/octave)

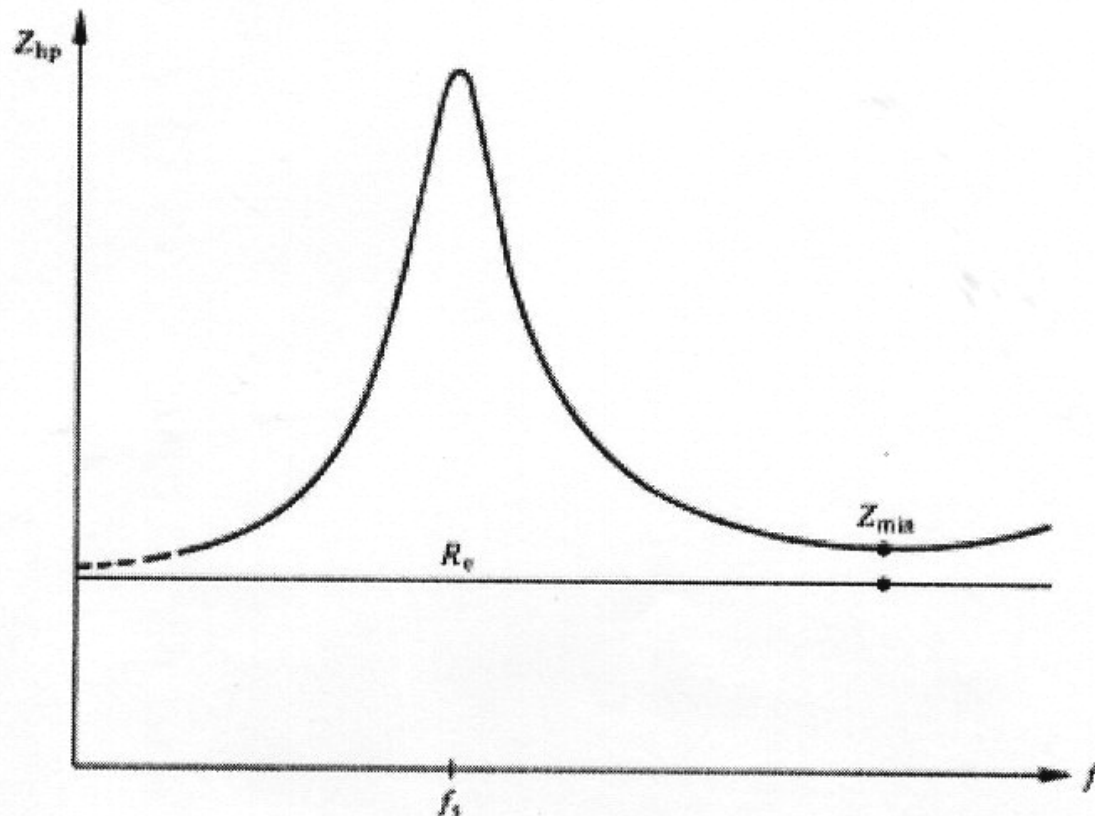




# Input impedance in HF

$L_e$  not negligible anymore

$f > 4.f_s$ ,  $Z_{em}$  is negligible



# Efficiency

efficiency =  $\frac{\text{radiated acoustic power}}{\text{provided electrical power}}$

$$P_{er} = \left( \frac{U_g}{R_g + R_e} \right)^2 \cdot R_e$$

$$\begin{aligned} \eta_s &= \frac{P_{as}}{P_{er}} = \left\{ \frac{\rho}{2\pi c} \right\} \left\{ \frac{(Bl)^2}{S_d^2 R_e m_{as}'^2} \right\} \\ &= 100 \cdot \left\{ \frac{4\pi^2}{c^3} \right\} \left\{ \frac{f_s^3 V_{as}}{Q_{es}} \right\} \end{aligned}$$

# Large signal parameters

Usage limitations:

- $P_{th}$  nominal thermal power handling capacity of the driver : dissipation by heat in the driver
- $\hat{\xi}_{\max}$  : maximum linear peak excursion of the cone.
- $\hat{\xi}_h$  : maximum excursion of the driver before distortion, defining a volume  $\hat{V}_d = S_d \hat{\xi}_h$
- $U_g \leq \frac{\hat{V}_d Q_e \omega_s (Bl)}{\sqrt{2} S_d \underbrace{x_{\max}}_{\max(|x_s|)}}$  voltage limited by distortion

# Specifications of a loudspeaker

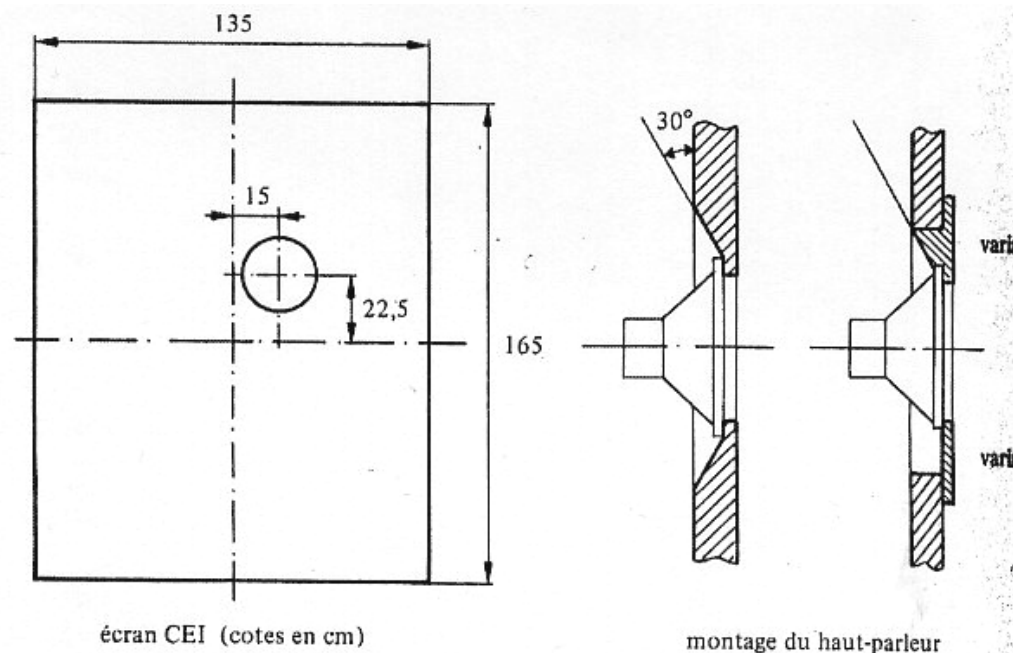
- Physical characteristics
- Polarities of electrical terminals
- Small and large signals parameters
- Performances of the loudspeaker, nominal and limit conditions of use

# Nominal quantities

- Nominal power  $P_n$
- Nominal bandwidth  $B_n$
- Nominal impedance  $Z_n$

# Response curve

- Free field
- Mounted on a normalized screen (IEC or AES)
- Distance of measurement
- Input power = 1/10 of nominal power



# Equivalent Thévenin source

- Loudspeaker = equivalent Thévenin source
- Force source  $\underline{F}_g = S_d \underline{p}_g$
- Source impedance  $\underline{Z}_{mg}$  (mechanical + mechanical equivalent of  $R_e$  and  $L_e$ )
- Load impedance  $\underline{Z}_{mc}$  = front and rear radiation impedances

