

Problem Set 11

CIVIL-425: Continuum Mechanics and Applications

22 May 2025

Exercise. Bingham fluid flow between two parallel plates.

Question 1: Momentum Balance Under steady-state assumption momentum balance is

$$\nabla \cdot \boldsymbol{\sigma} = 0$$

or

$$-\nabla p + \nabla \cdot \mathbf{s} = 0$$

The expanded form in two dimensions is

$$\begin{aligned} -\frac{\partial p}{\partial x} + \frac{\partial s_{xx}}{\partial x} + \frac{\partial s_{xy}}{\partial y} &= 0 \\ -\frac{\partial p}{\partial y} + \frac{\partial s_{xy}}{\partial x} + \frac{\partial s_{yy}}{\partial y} &= 0 \end{aligned}$$

where $s_{xx} = 0$ and $s_{yy} = 0$ because $\dot{\mathbf{u}} = (\dot{u}_x(y), 0)$. Assuming that $\frac{\partial p}{\partial y}$ is very small in a narrow channel $w \ll L$, we obtain

$$\begin{aligned} -\frac{\partial p}{\partial x} + \frac{\partial s_{xy}}{\partial y} &= 0 \\ \frac{\partial s_{xy}}{\partial x} &= 0 \end{aligned}$$

Question 2: Shear Stress Profile Integrating momentum balance we obtain

$$s_{xy} = \frac{dp}{dx}y$$

This linear profile indicates a maximum and minimum at the plates, depending on the sign of $\frac{dp}{dx}$.

Question 3: Velocity Profile The fluid starts to flow if $\sqrt{J_2(\mathbf{s})} > \sigma_Y$. Deviatoric part of stress tensor is such that $tr(\mathbf{s}) = 0$, therefore $\sqrt{J_2(\mathbf{s})} = |s_{xy}|$. The flow condition becomes

$$|s_{xy}| > \sigma_Y$$

In pure shear, plastic strain rate can be written as

$$\boldsymbol{\epsilon}^p = \begin{bmatrix} 0 & \dot{\gamma} \\ \dot{\gamma} & 0 \end{bmatrix}$$

where $\dot{\gamma} = \frac{du_x}{dy}$.

For $|y| \leq y_c$ where y_c is where $|s_{xy}(y_c)| = \sigma_Y$ ($\Rightarrow y_c = \pm \frac{\sigma_Y}{|dp/dx|}$), the fluid behaves as rigid and $\dot{\gamma} = \frac{du_x}{dy} = 0$, which leads to $\dot{u}_x = A = \text{const.}$

On the other hand, when $|y| > y_c$, we have

$$\dot{\gamma} = \frac{d\dot{u}_x}{dy} = \frac{|s_{xy}| - \sigma_Y}{\eta} \frac{s_{xy}}{|s_{xy}|}$$

recalling that $s_{xy} = \frac{dp}{dx}y$, we obtain

$$\dot{\gamma} = \frac{d\dot{u}_x}{dy} = \frac{\left| \frac{dp}{dx} \right| |y| - \sigma_Y}{\eta} \text{sign}(y)$$

First, assume, that $y > y_c$, then

$$\dot{\gamma} = \frac{d\dot{u}_x}{dy} = \frac{\left| \frac{dp}{dx} \right| y - \sigma_Y}{\eta}$$

integrating and taking into account $\dot{u}_x(y = +h) = 0$ we obtain

$$\dot{u}_x(y) = \frac{1}{\eta} \left(\frac{1}{2} \left| \frac{dp}{dx} \right| (y^2 - h^2) + \sigma_Y(h - y) \right)$$

Analogously for $y < -y_c$

$$\dot{\gamma} = \frac{d\dot{u}_x}{dy} = \frac{\left| \frac{dp}{dx} \right| y + \sigma_Y}{\eta}$$

integrating and taking into account $\dot{u}_x(y = -h) = 0$ we obtain

$$\dot{u}_x(y) = \frac{1}{\eta} \left(\frac{1}{2} \left| \frac{dp}{dx} \right| (y^2 - h^2) + \sigma_Y(y + h) \right)$$

Now we can compute constant velocity at $y = \pm y_c$

$$\dot{u}_x(y) = \frac{1}{\eta} \left(\frac{1}{2} \left| \frac{dp}{dx} \right| (y_c^2 - h^2) + \sigma_Y(h - y_c) \right)$$

We can summarise our solution

$$\begin{aligned} \dot{u}_x(y) &= \frac{1}{\eta} \left(\frac{1}{2} \left| \frac{dp}{dx} \right| (y^2 - h^2) + \sigma_Y(h - y) \right) \text{ for } y > y_c \\ \dot{u}_x(y) &= \frac{1}{\eta} \left(\frac{1}{2} \left| \frac{dp}{dx} \right| (y_c^2 - h^2) + \sigma_Y(h - y_c) \right) \text{ for } -y_c \leq y \leq y_c \\ \dot{u}_x(y) &= \frac{1}{\eta} \left(\frac{1}{2} \left| \frac{dp}{dx} \right| (y^2 - h^2) + \sigma_Y(y + h) \right) \text{ for } y < -y_c \end{aligned}$$

