

Molecular quantum dynamics: Solutions 8

a)

The coefficients of the exact solution for the TDSE are

$$\dot{c}_k(t) = -\frac{i}{\hbar} \sum_m e^{-i\omega_{km}t} V_{km}(t) c_m(t), \quad (1)$$

where $\omega_{km} = (E_m - E_k)/\hbar$. Therefore, we get

$$\dot{c}_a(t) = -\frac{i}{2\hbar} U e^{i(\omega - \omega_{ab})t} c_b(t) \quad (2)$$

$$\dot{c}_b(t) = -\frac{i}{2\hbar} U e^{-i(\omega - \omega_{ab})t} c_a(t). \quad (3)$$

Differentiating $\dot{c}_a(t)$ and substituting $\dot{c}_b(t)$ and $c_b(t)$, we obtain

$$\ddot{c}_a(t) = -\frac{i}{2\hbar} U e^{i(\omega - \omega_{ab})t} \dot{c}_b(t) + \frac{\omega - \omega_{ab}}{2\hbar} U e^{i(\omega - \omega_{ab})t} c_b(t) \quad (4)$$

$$= -\frac{U^2}{4\hbar^2} c_a(t) + i(\omega - \omega_{ab}) \dot{c}_a(t) \quad (5)$$

This is a second order differential equation of the following type:

$$\ddot{x} + b\dot{x} + kx = 0 \quad (6)$$

and if $b^2 < 4k$, the solution is

$$x(t) = e^{\alpha t} [C_1 \cos(\beta t) + C_2 \sin(\beta t)], \quad (7)$$

where $\alpha = -b/2$ and $\beta = (1/2)\sqrt{4k - b^2}$. In our case, $b = -i(\omega - \omega_{ab})$ and $k = U^2/(4\hbar^2)$ so $b^2 < 4k$. The solution for $c_a(t)$ is given as

$$c_a(t) = e^{i(\omega - \omega_{ab})t/2} [C_1 \cos(\omega_r t) + C_2 \sin(\omega_r t)] \quad (8)$$

Similarly for $c_b(t)$,

$$\ddot{c}_b(t) + i(\omega - \omega_{ab})\dot{c}_b(t) + \frac{U^2}{4\hbar^2} c_b(t) = 0 \quad (9)$$

with the solution

$$c_b(t) = e^{-i(\omega - \omega_{ab})t/2} [D_1 \cos(\omega_r t) + D_2 \sin(\omega_r t)]. \quad (10)$$

From $c_a(t=0) = 1$ we obtain $C_1 = 1$ and from $c_b(t=0) = 0$, $D_1 = 0$.

We need to know $\dot{c}_a(t=0)$ and $\dot{c}_b(t=0)$ to figure out C_2 and D_2 .

$$\dot{c}_a(0) = -\frac{iU}{2\hbar} c_b(0) = 0 \quad (11)$$

$$\dot{c}_b(0) = -\frac{iU}{2\hbar} c_a(0) = -\frac{iU}{2\hbar}. \quad (12)$$

On the other hand, differentiating the solutions of ODEs give

$$\dot{c}_a(t) = \frac{i(\omega - \omega_{ab})}{2} c_a(t) + e^{i(\omega - \omega_{ab})t/2} [\omega_r C_2 \cos(\omega_r t) - \omega_r \sin(\omega_r t)] \quad (13)$$

$$\dot{c}_b(t) = -\frac{i(\omega - \omega_{ab})}{2} c_b(t) + e^{-i(\omega - \omega_{ab})t/2} \omega_r D_2 \cos(\omega_r t) \quad (14)$$

Substituting $t = 0$ to Eq. (13) and Eq. (14) gives

$$\dot{c}_a(0) = \frac{i(\omega - \omega_{ab})}{2} + \omega_r C_2 \quad (15)$$

$$\dot{c}_b(0) = \omega_r D_2 \quad (16)$$

Finally, we obtain $C_2 = -i(\omega - \omega_{ab})/(2\omega_r)$ and $D_2 = (-iU)/(2\hbar\omega_r)$ to get the final solution as

$$c_a(t) = e^{i(\omega - \omega_{ab})t/2} [\cos(\omega_r t) - \frac{i(\omega - \omega_{ab})}{2\omega_r} \sin(\omega_r t)] \quad (17)$$

$$c_b(t) = e^{-i(\omega - \omega_{ab})t/2} [\frac{-iU}{2\hbar\omega_r} \sin(\omega_r t)] \quad (18)$$

b)

The transition probability $P_{a \rightarrow b}(t)$ is given by

$$P_{a \rightarrow b}(t) = |c_b(t)|^2 = \frac{U^2}{4\hbar^2} \frac{\sin^2(\omega_r t)}{\omega_r^2} = \frac{(U/\hbar)^2}{(\omega - \omega_{ab})^2 + (U/\hbar)^2} \sin^2(\omega_r t). \quad (19)$$

Both the fraction and $\sin^2(\omega_r t)$ do not exceed 1 so the transition probability also doesn't.

$$|c_a(t)|^2 = \cos^2(\omega_r t) + \frac{(\omega - \omega_{ab})^2}{4\omega_r^2} \sin^2(\omega_r t) \quad (20)$$

$$= \cos^2(\omega_r t) + \left(1 - \frac{U^2/\hbar^2}{4\omega_r^2}\right) \sin^2(\omega_r t) \quad (21)$$

Finally,

$$|c_a(t)|^2 + |c_b(t)|^2 = \cos^2(\omega_r t) + \sin^2(\omega_r t) = 1 \quad (22)$$

c)

When

$$\left(\frac{U}{\hbar}\right)^2 \ll (\omega - \omega_{ab})^2, \quad (23)$$

i.e, if $U \ll \hbar(\omega - \omega_{ab})$ then $\omega_r \approx (\omega - \omega_{ab})/2$. This gives

$$P_{a \rightarrow b}(t) = \frac{U^2}{4\hbar^2} \frac{\sin^2[(\omega - \omega_{ab})t/2]}{[(\omega - \omega_{ab})/2]^2} \quad (24)$$

$$= \left(\frac{U}{2}\right)^2 \frac{\sin^2[(E_a - E_b + E)t/2\hbar]}{[(E_a - E_b + E)/2]^2}, \quad (25)$$

which agrees with the TDPT result we derived in the lecture.

d)

For a normalized quantum state ψ , both ψ and $\lambda\psi$ represent the same physical state, where λ is a non-zero complex number satisfying $|\lambda| = 1$. [Such a λ can be written as $\lambda = \exp(i\phi)$.] Therefore, as long as $|c_b|^2 = 0$, the system returns to the original state. The system first returns to its original state at

$$t^* = \frac{\pi}{\omega_r} \quad (26)$$