

# Structural Analysis

## Part III - X-ray tools

### Session 3

#### X-ray scattering and diffraction

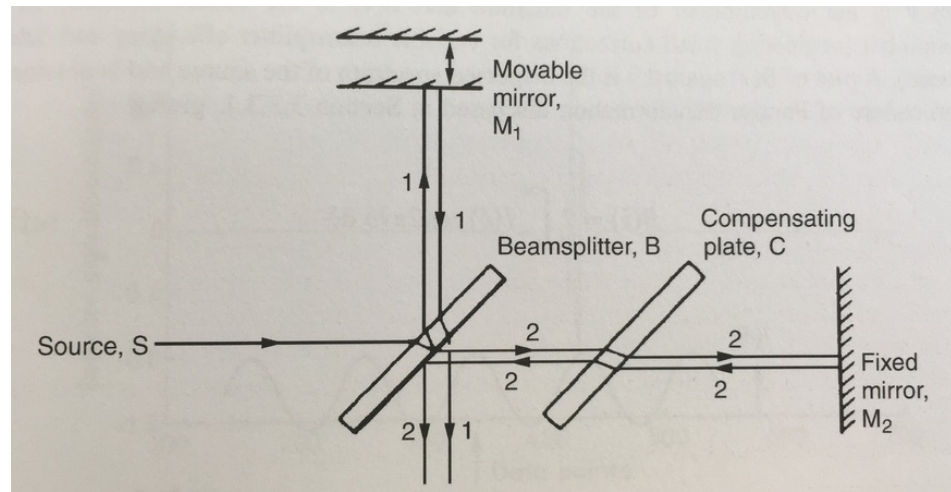
## Intermezzo: Fourier Transform

## Alternative approach: Fourier Transform (Infrared) Spectrometer

- Not a dispersive measurement but based on interferometry
- Use all wavelength at the same time
- Manipulation in  $x$  - cm (real space) to get information in wave numbers  $x^{-1} - \text{cm}^{-1}$  (frequencies)
- Measure interferogram, Fourier transform into spectrum

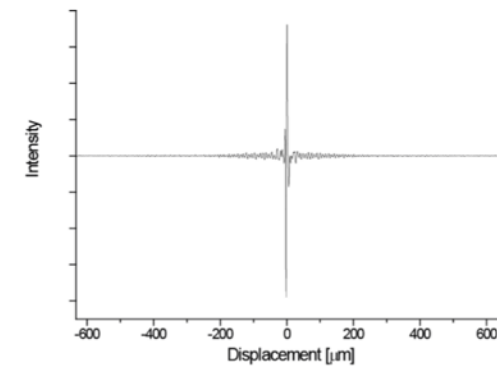


# Measuring an interferogram

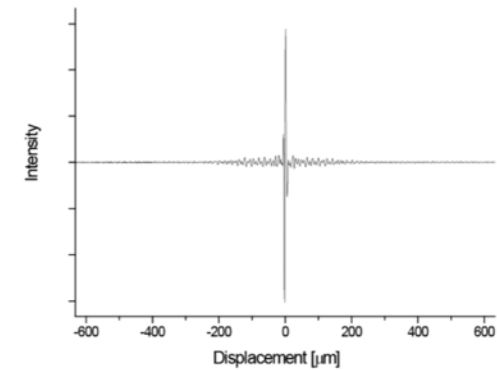
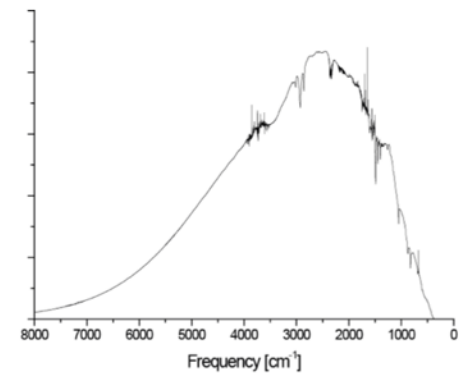


Sample

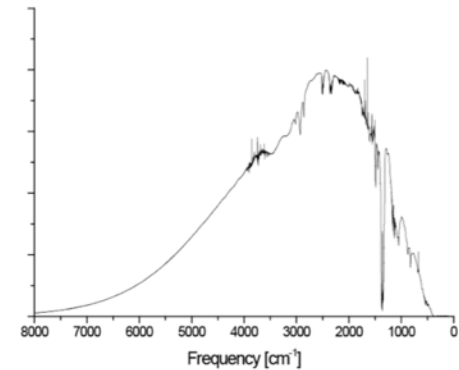
Detector



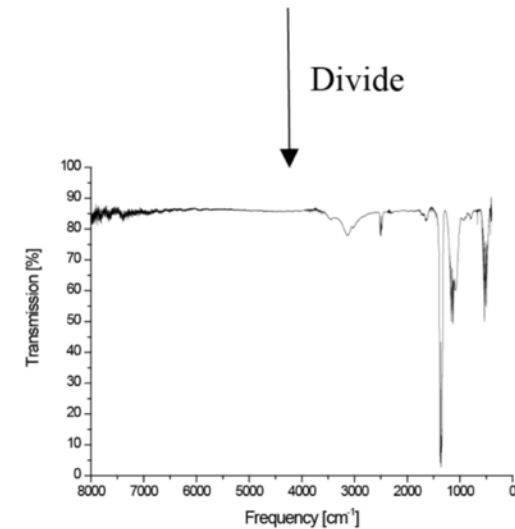
Empty Cell  
FT



Filled Cell  
FT



Divide



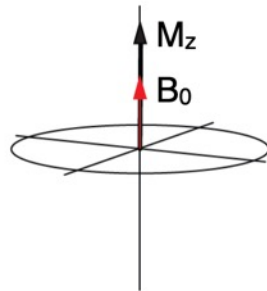
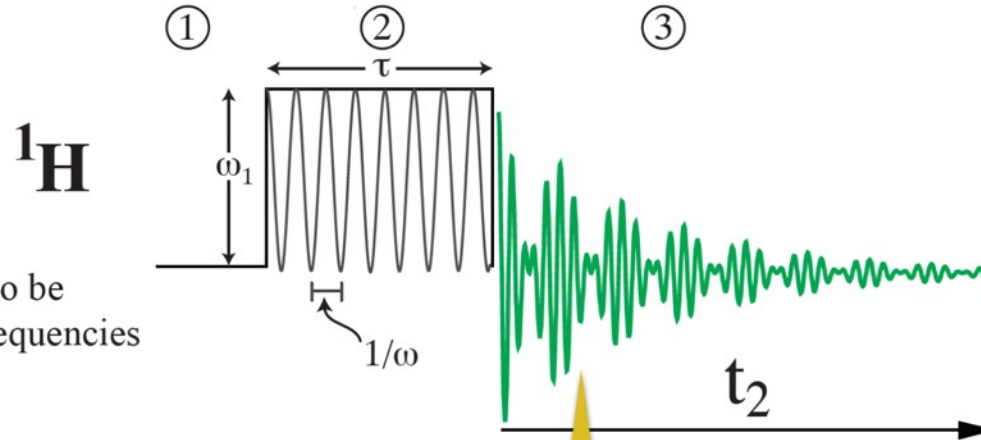
Prof. Emsley part of this class, 5 weeks ago!

## Pulsed FTNMR Spectroscopy

$$\omega_1 = -\gamma B_1$$

$$\omega_1 \tau = \pi/2$$

$\omega$  = carrier frequency,  
chosen by the operator to be  
near to the resonance frequencies



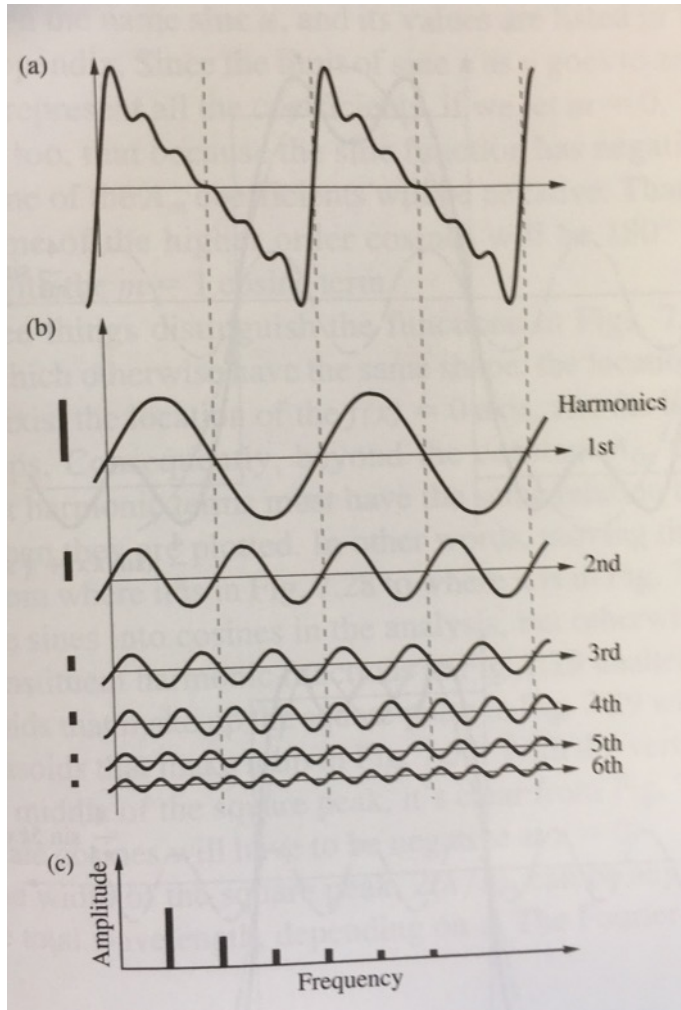
what does the signal actually look like?

1. Equilibrium. The net magnetization is aligned along the direction of the main field (z-axis).

2. A field is applied in the transverse plane. The magnetization of the ensemble precesses around the field.

3. The field is removed leaving a net transverse component of the ensemble magnetization. This *coherence* then starts to precess around the main field.

# Principle idea



## Mathematical description

- The principle idea of the Fourier analysis / transformation is that any function can be represented by an (infinite) series of harmonic functions.
- The Fourier transform decomposes a function into its constituent frequencies.

Thus we can write

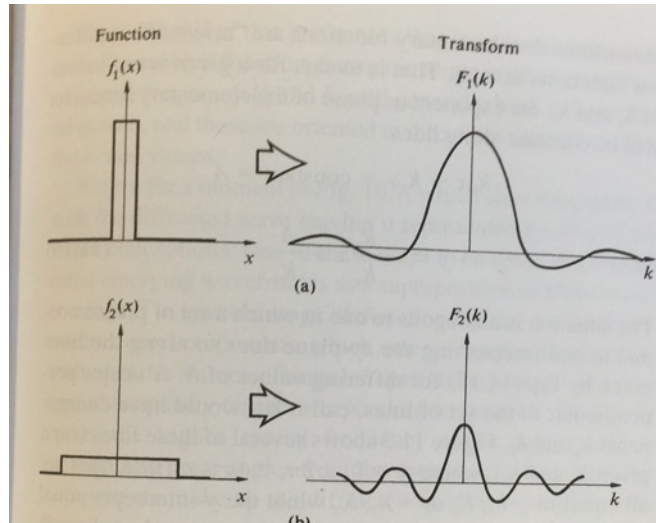
$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} F(k) e^{-ikx} dk$$

provided that

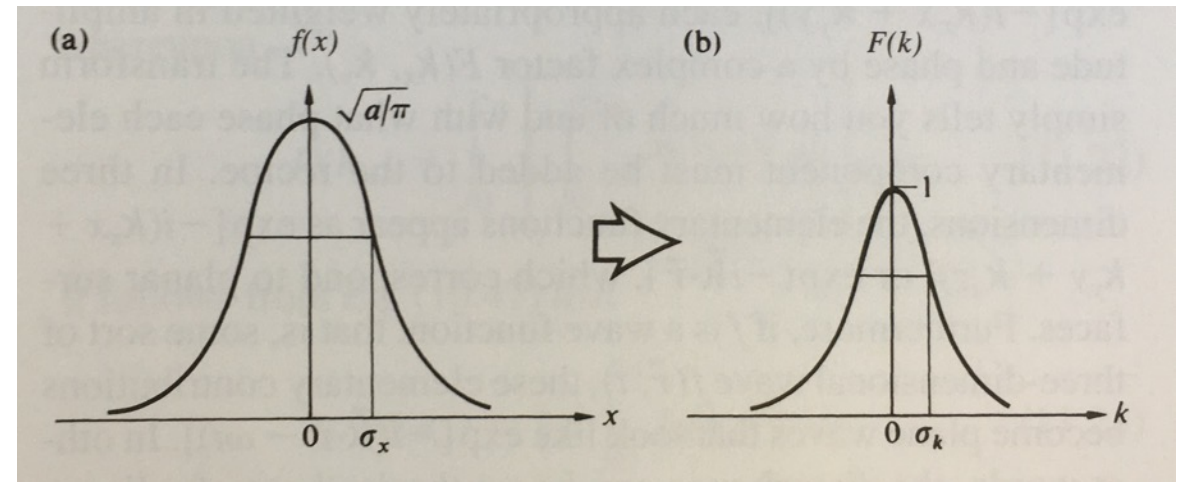
$$F(k) = \int_{-\infty}^{+\infty} f(x) e^{ikx} dx$$

# Examples

## Square functions



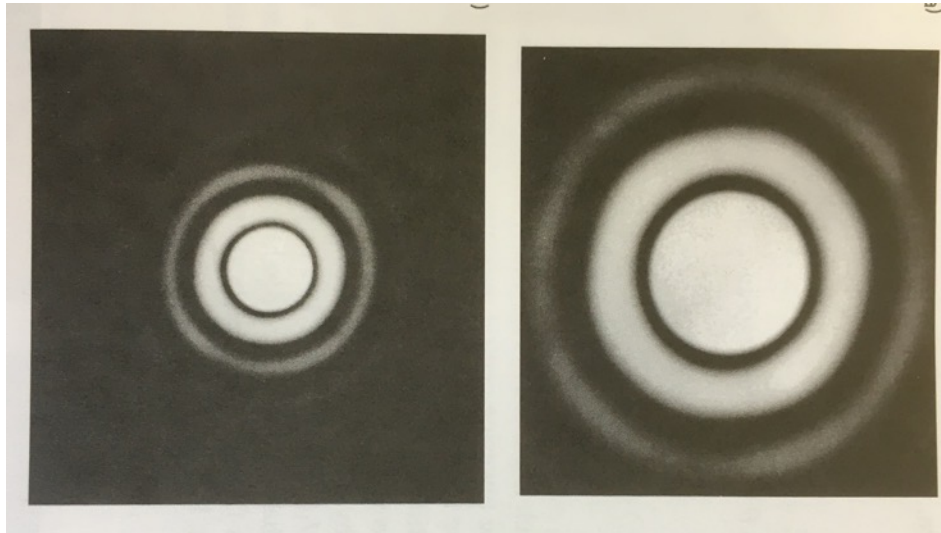
## Gauss function



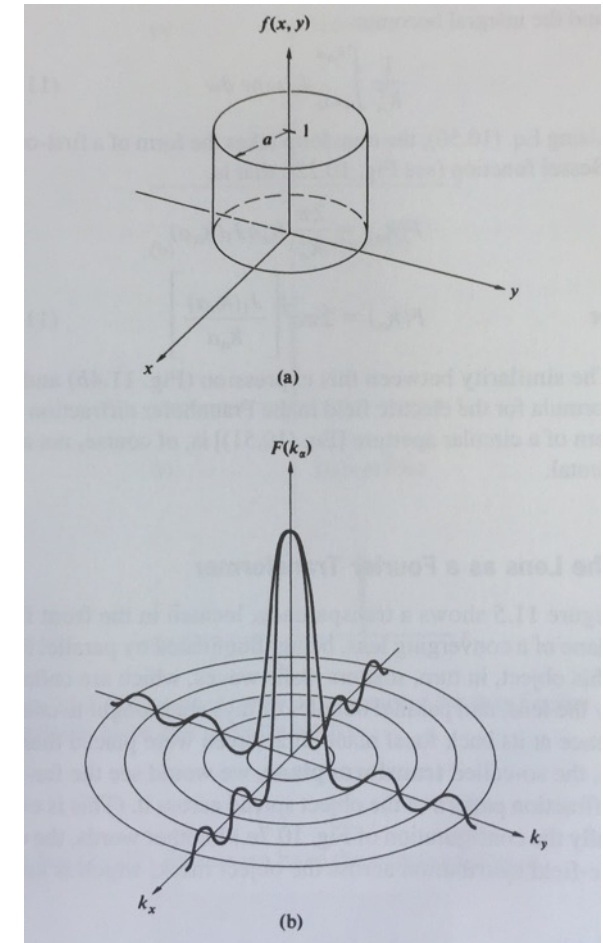


# Fourier transform of round aperture and airy pattern

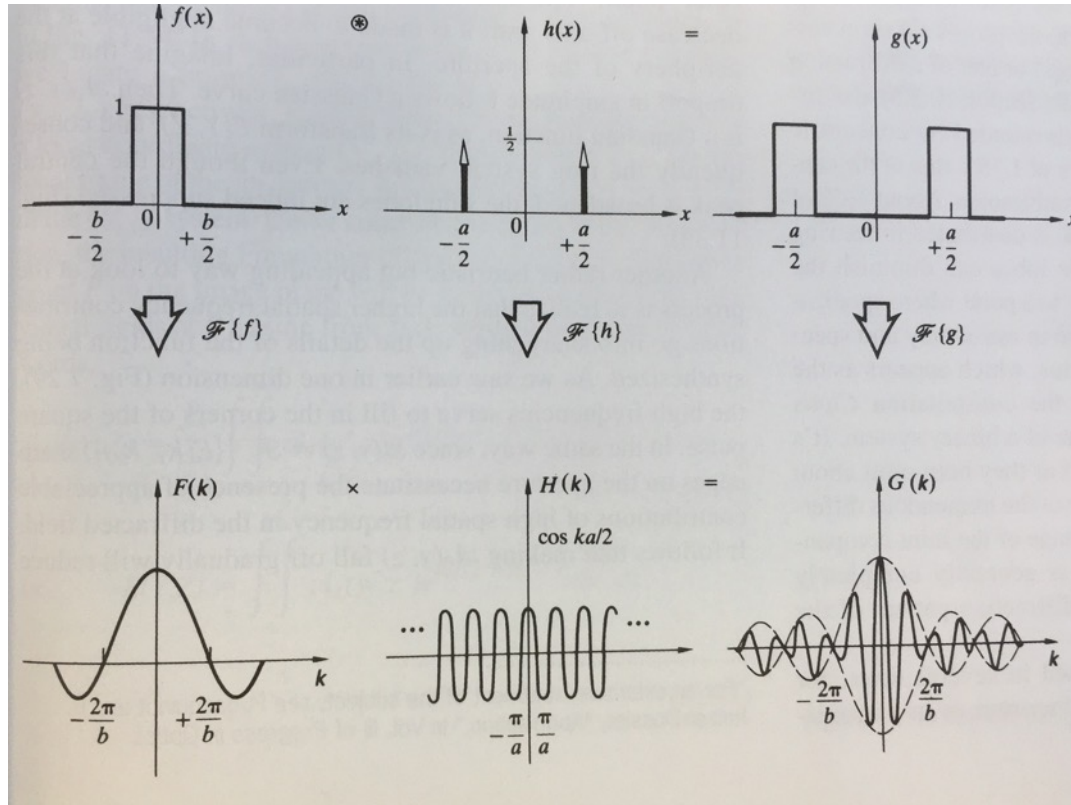
Experiment: Diffraction pattern of circular aperture, Airy pattern



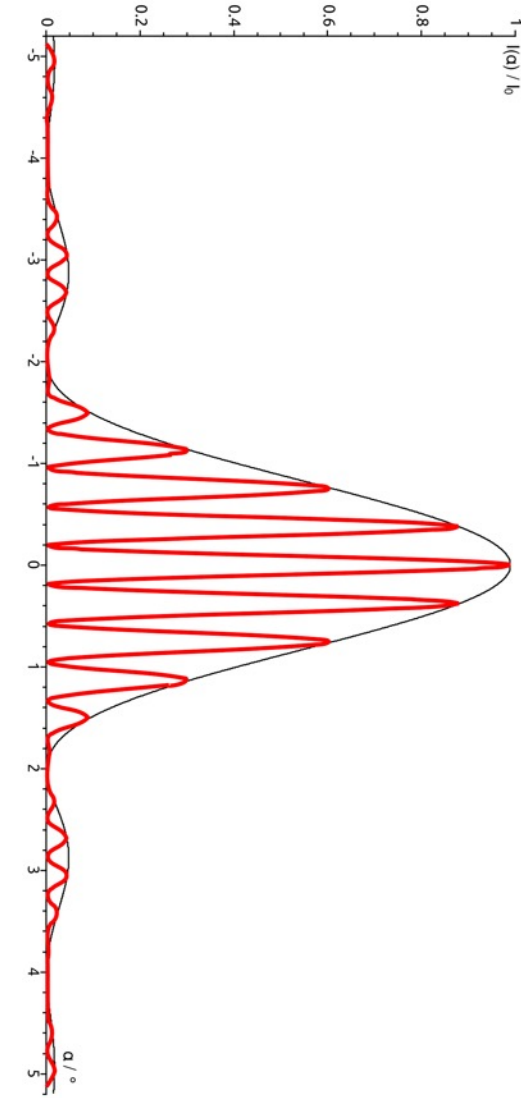
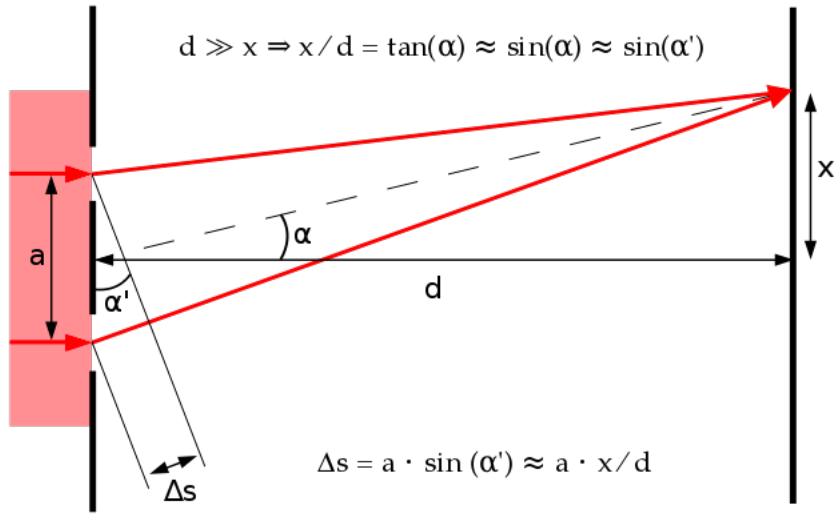
Theory: Fourier transform of cylinder or “top-hat” function



# Diffraction as Fourier transform:



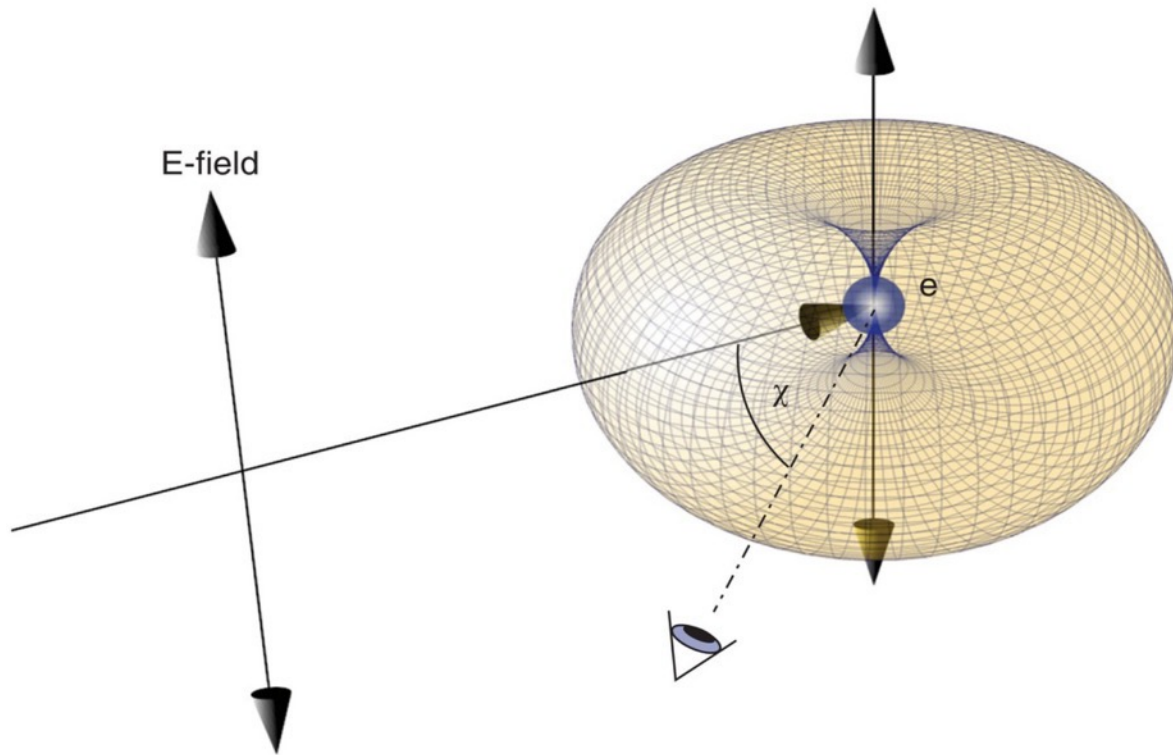
# The double slit



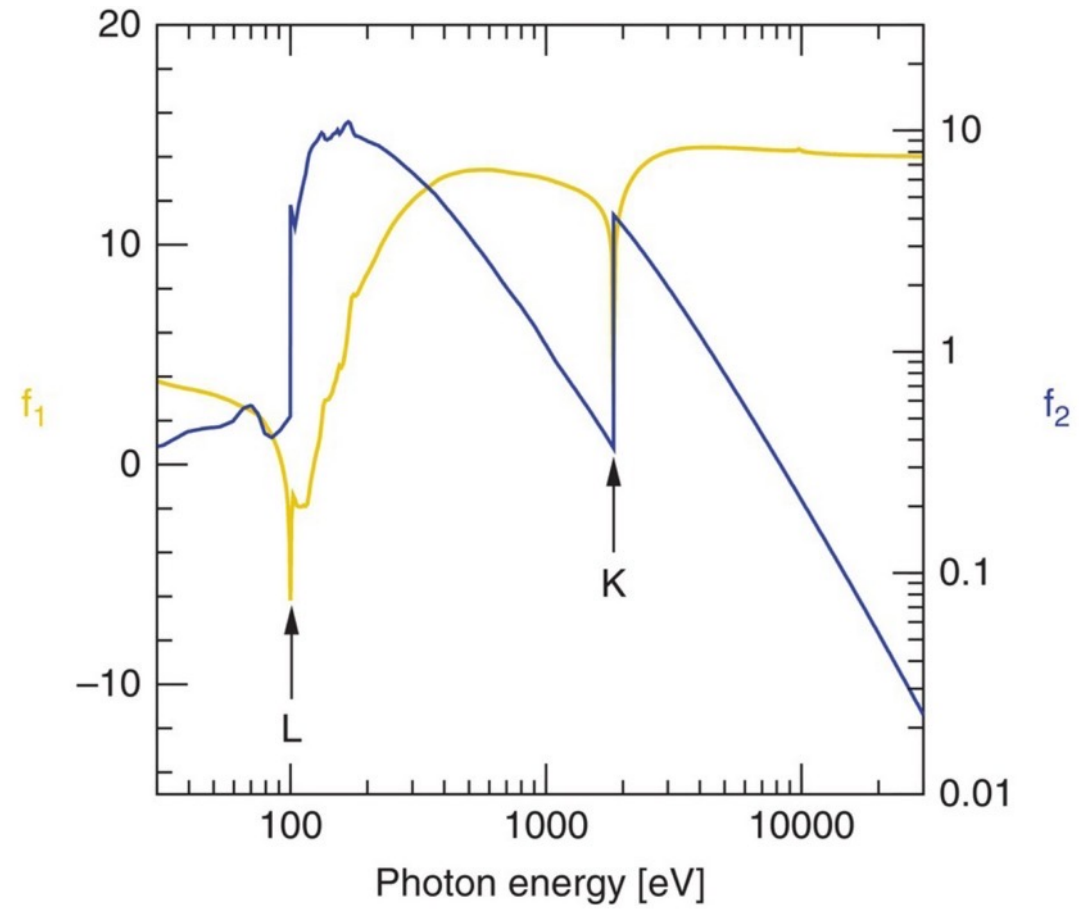
## Fourier relationships – some examples

From Session 1

# Thomson scattering from a single electron

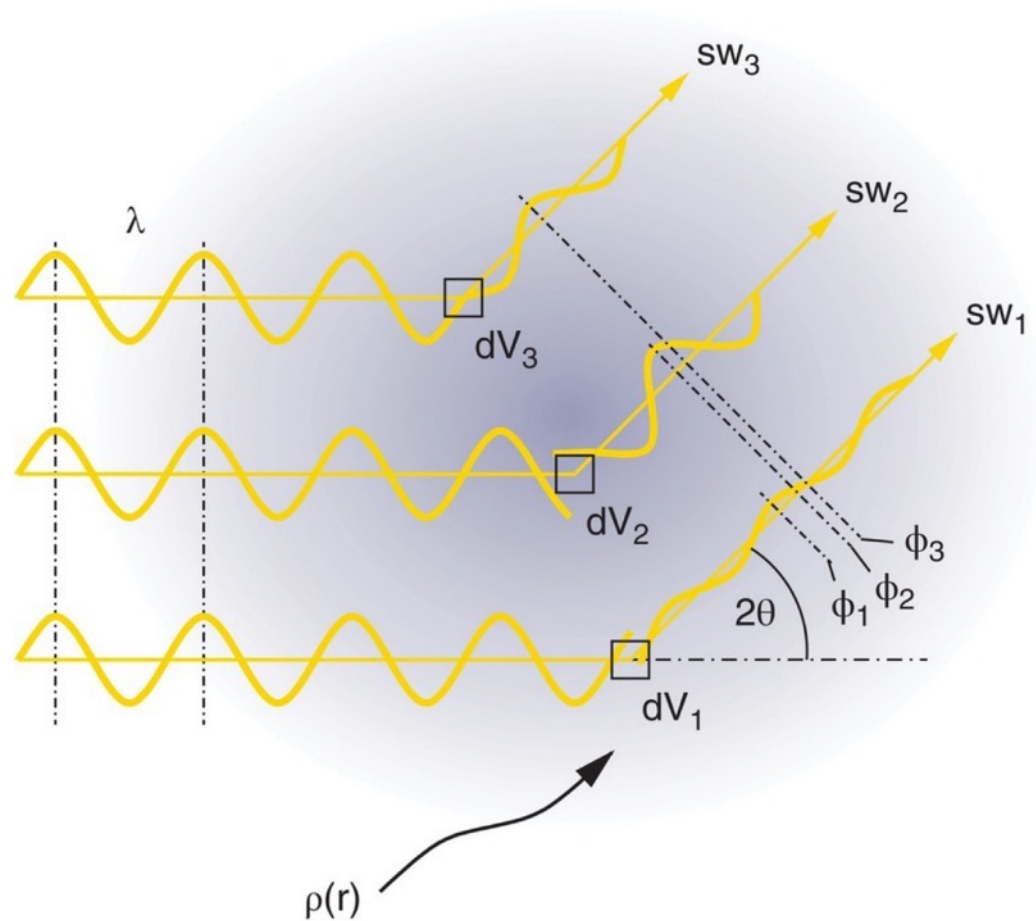


# Atomic scattering factors and refractive index

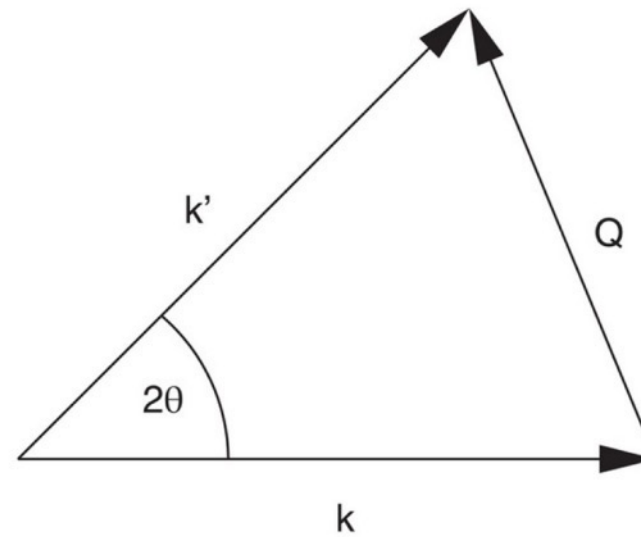


# Definition of the scattering vector $Q$

(a)



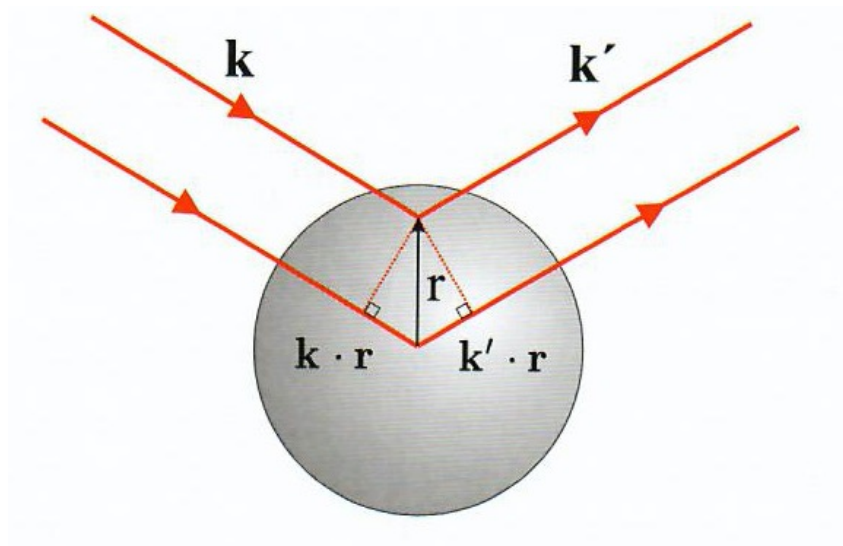
(b)



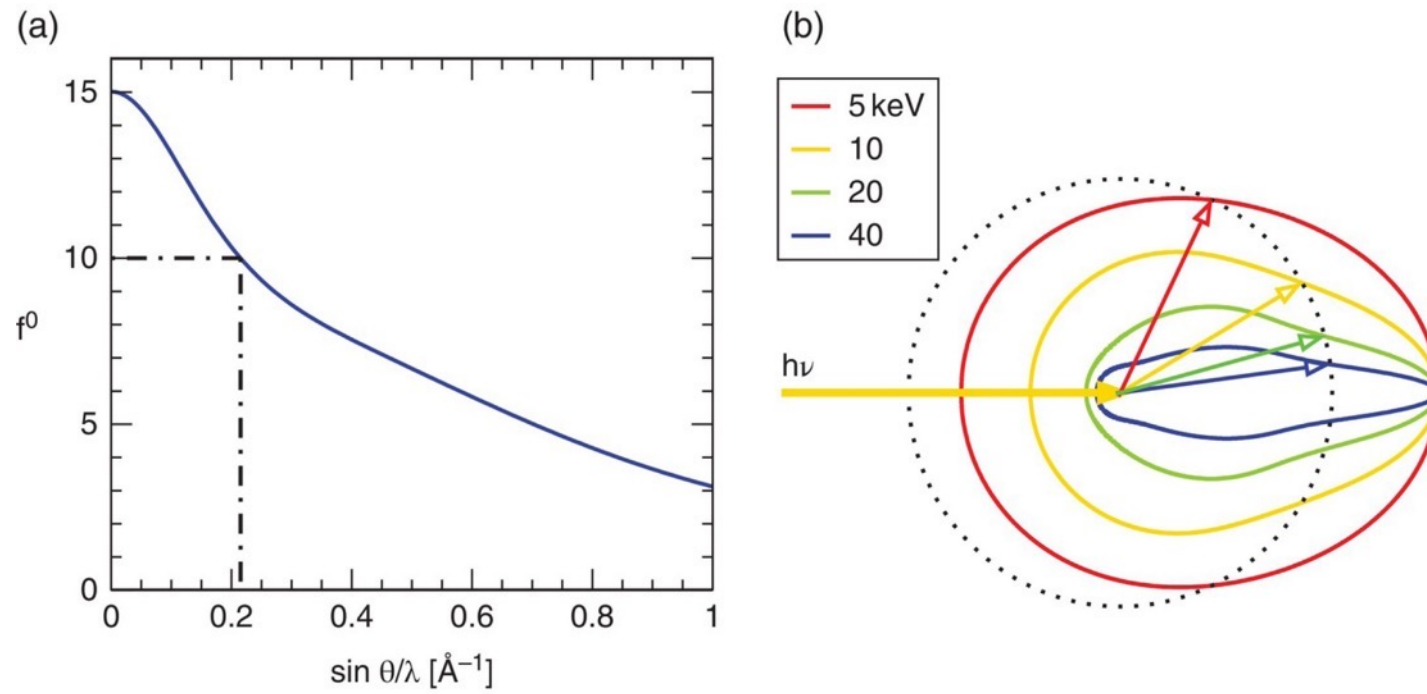
$$|k| = |k'| = 2\pi/\lambda$$



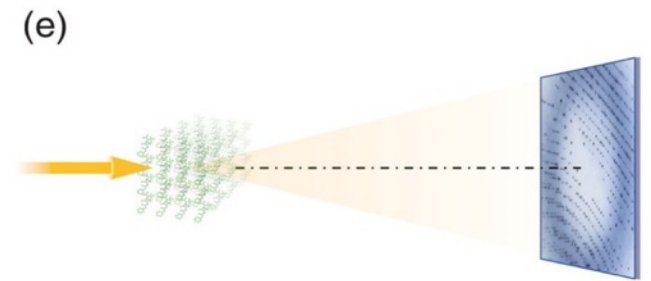
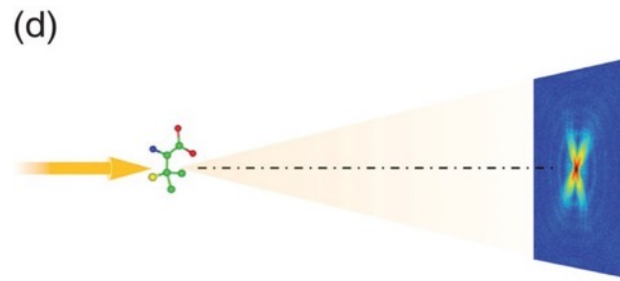
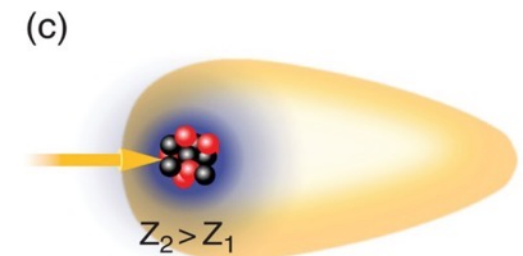
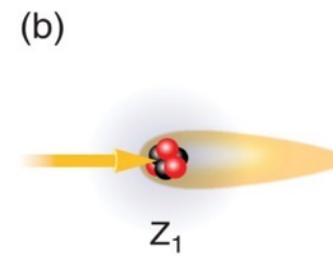
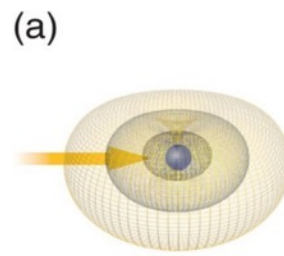
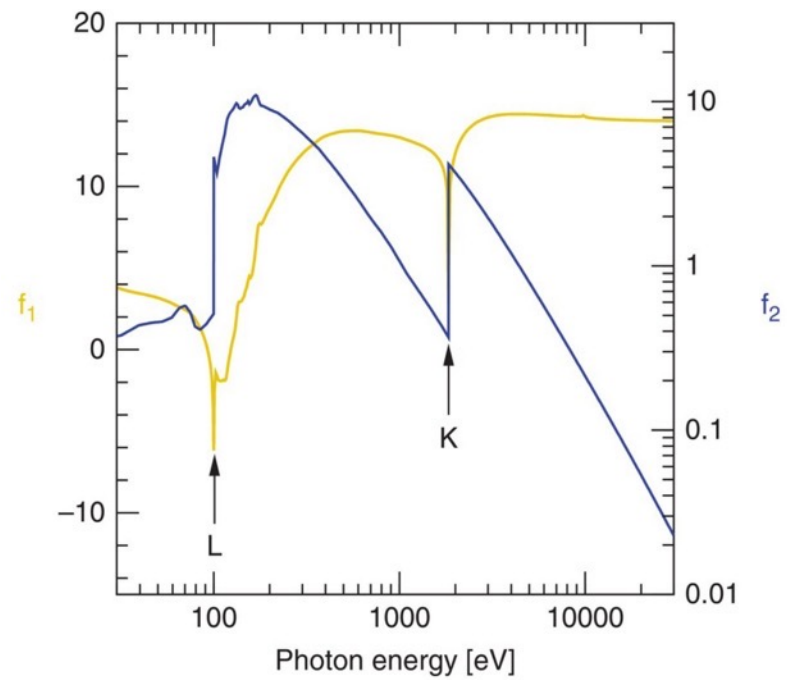
## Scattering from an electron cloud



## Now full cross section / atomic scattering factors



# Atomic scattering factors again



From Session 2

# Bragg scattering

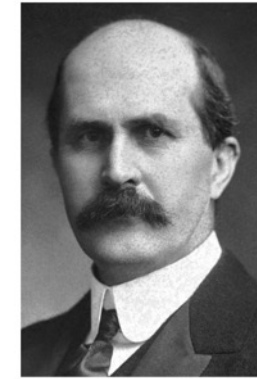
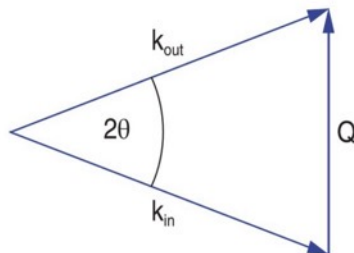
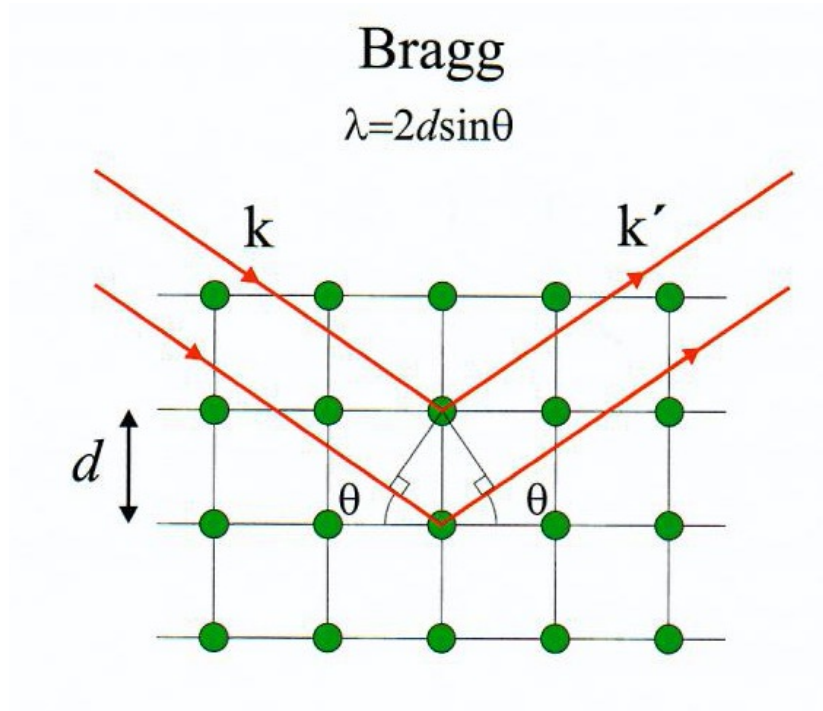


Photo from the Nobel Foundation archive.

**Sir William Henry Bragg**

Prize share: 1/2

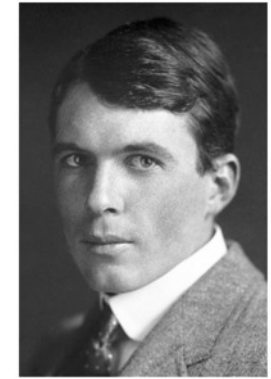
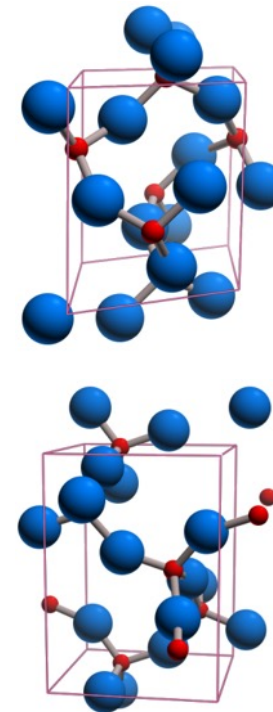
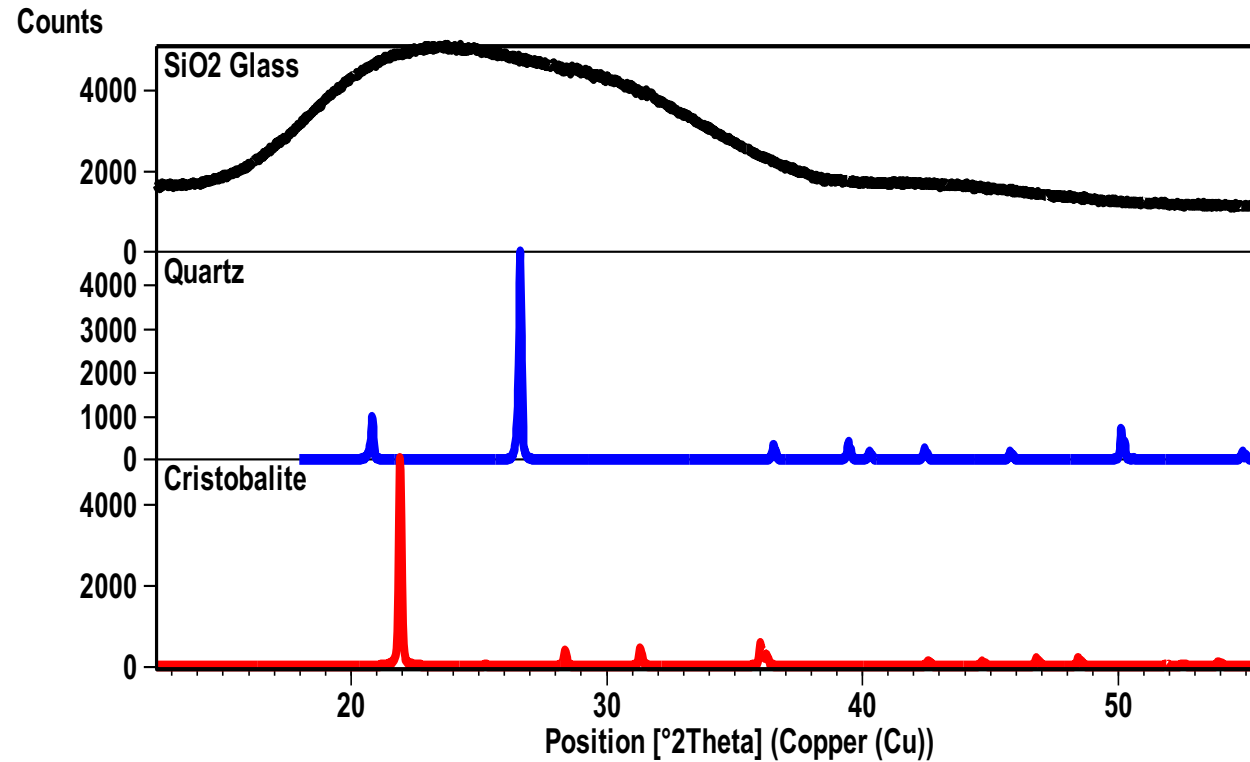


Photo from the Nobel Foundation archive.

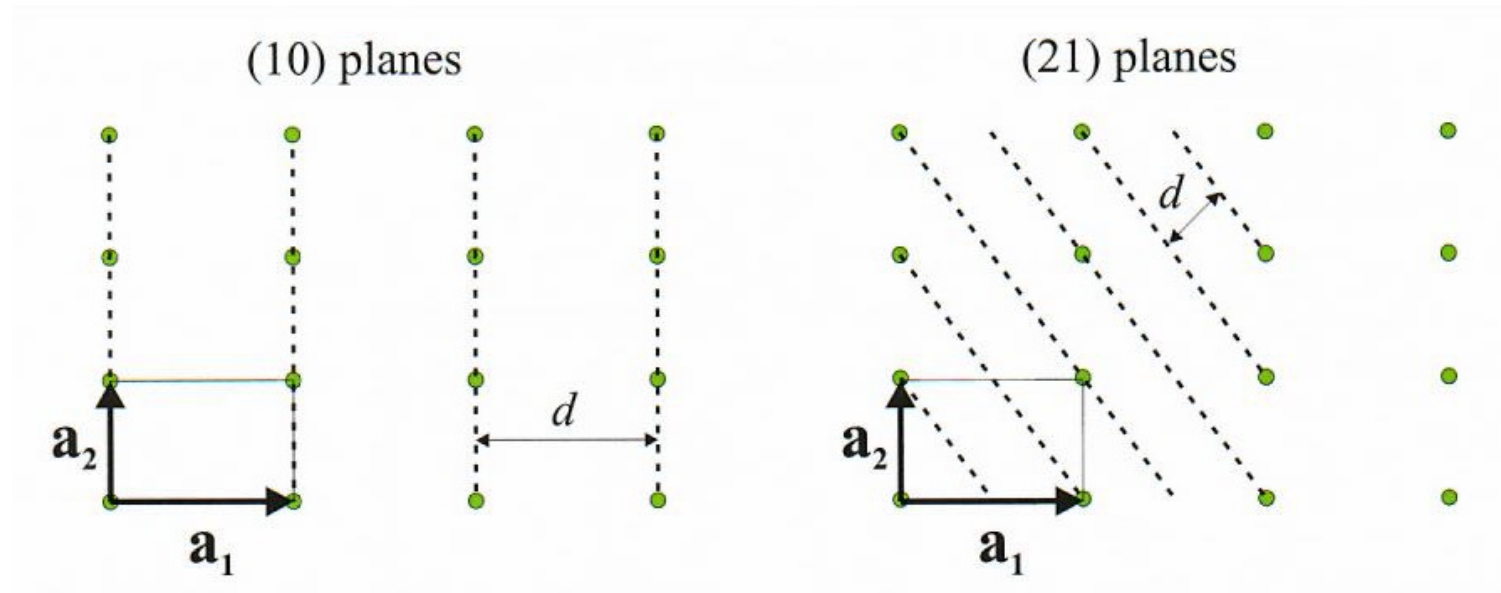
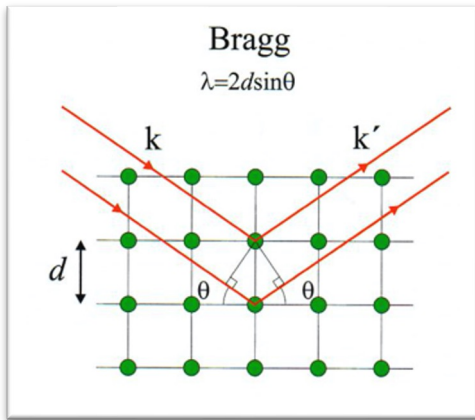
**William Lawrence Bragg**

Prize share: 1/2

# Characteristic signals from different – chemically identical – samples

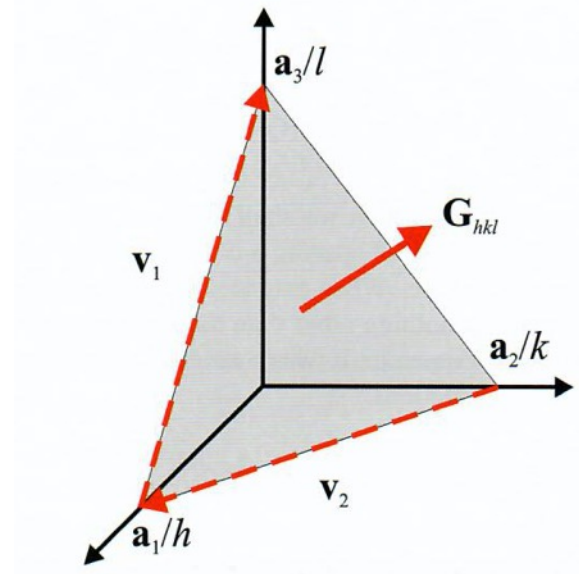
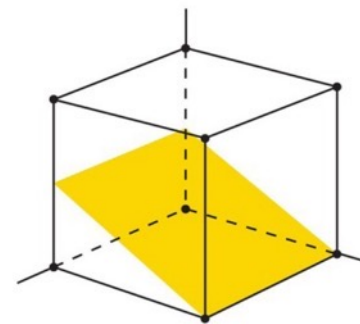
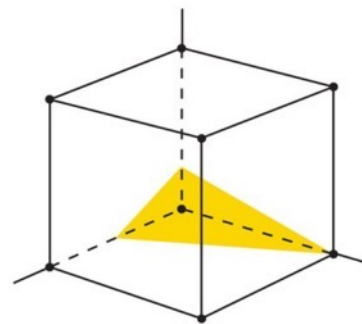
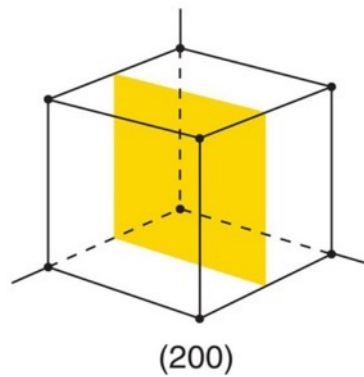
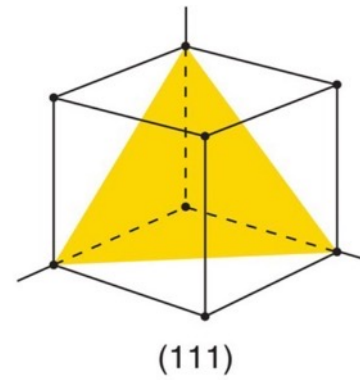
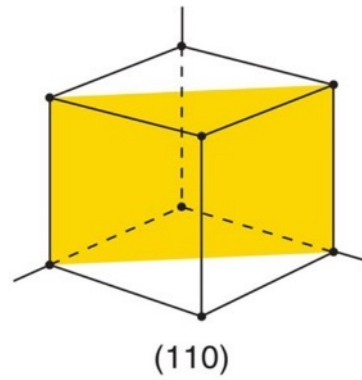
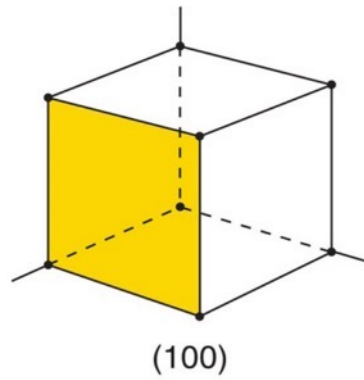


# Lattice planes



- Find the intercepts of the plane with the respective crystal axis
- Take the reciprocal of these numbers, reduce to smallest integer

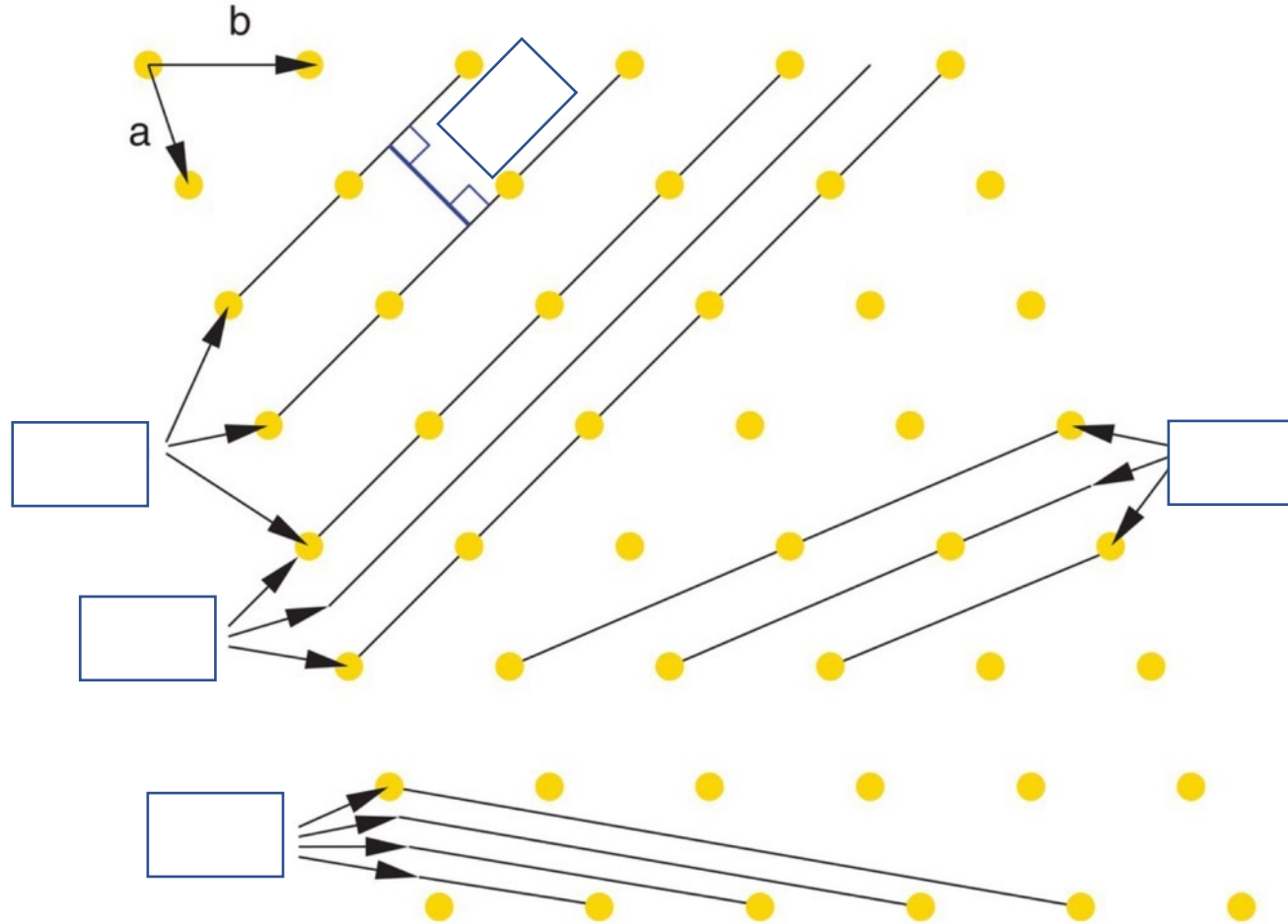
# Miller indices



- Find the intercepts of the plane with the respective crystal axis
- Take the reciprocal of these numbers, reduce to smallest integer

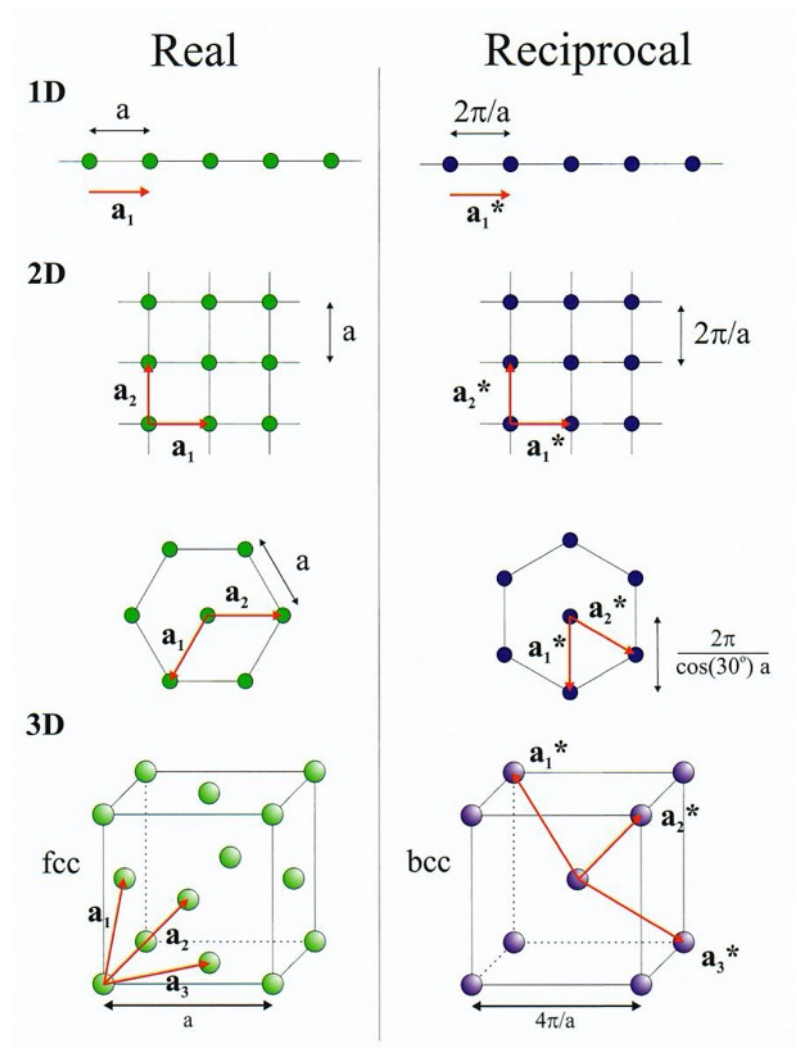


## Exercise: Identify lattice planes



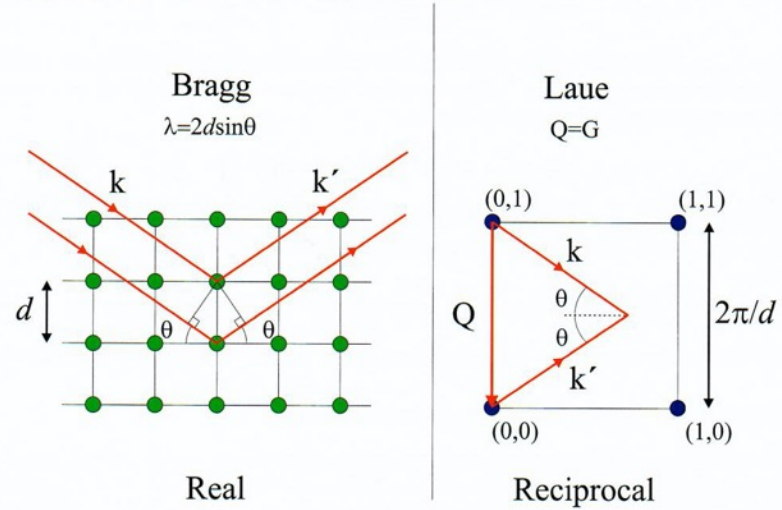
## More on diffraction of single crystals

For each lattice a reciprocal lattice can be defined



# Bragg and Laue conditions

(a) Equivalence of Bragg and Laue



(b) Miller indices and reciprocal lattice vectors

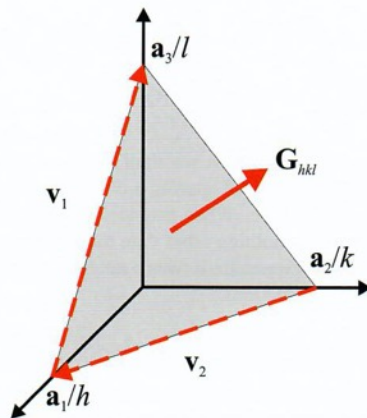
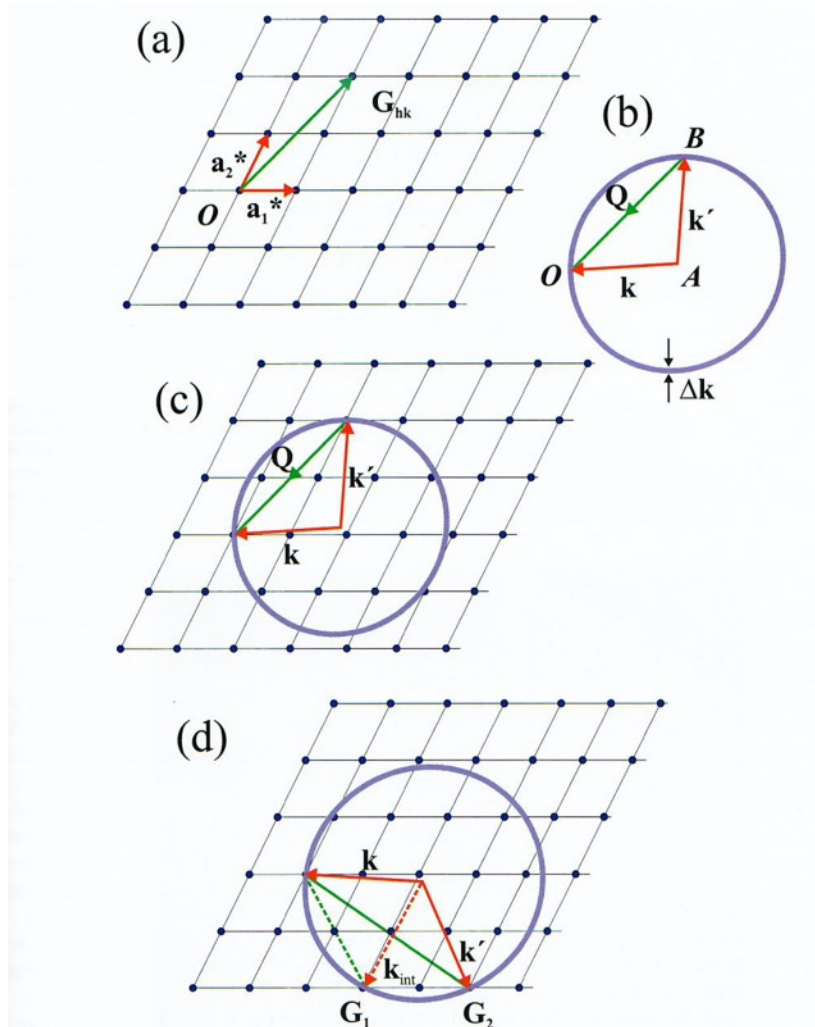


Photo from the Nobel Foundation archive.

**Max von Laue**

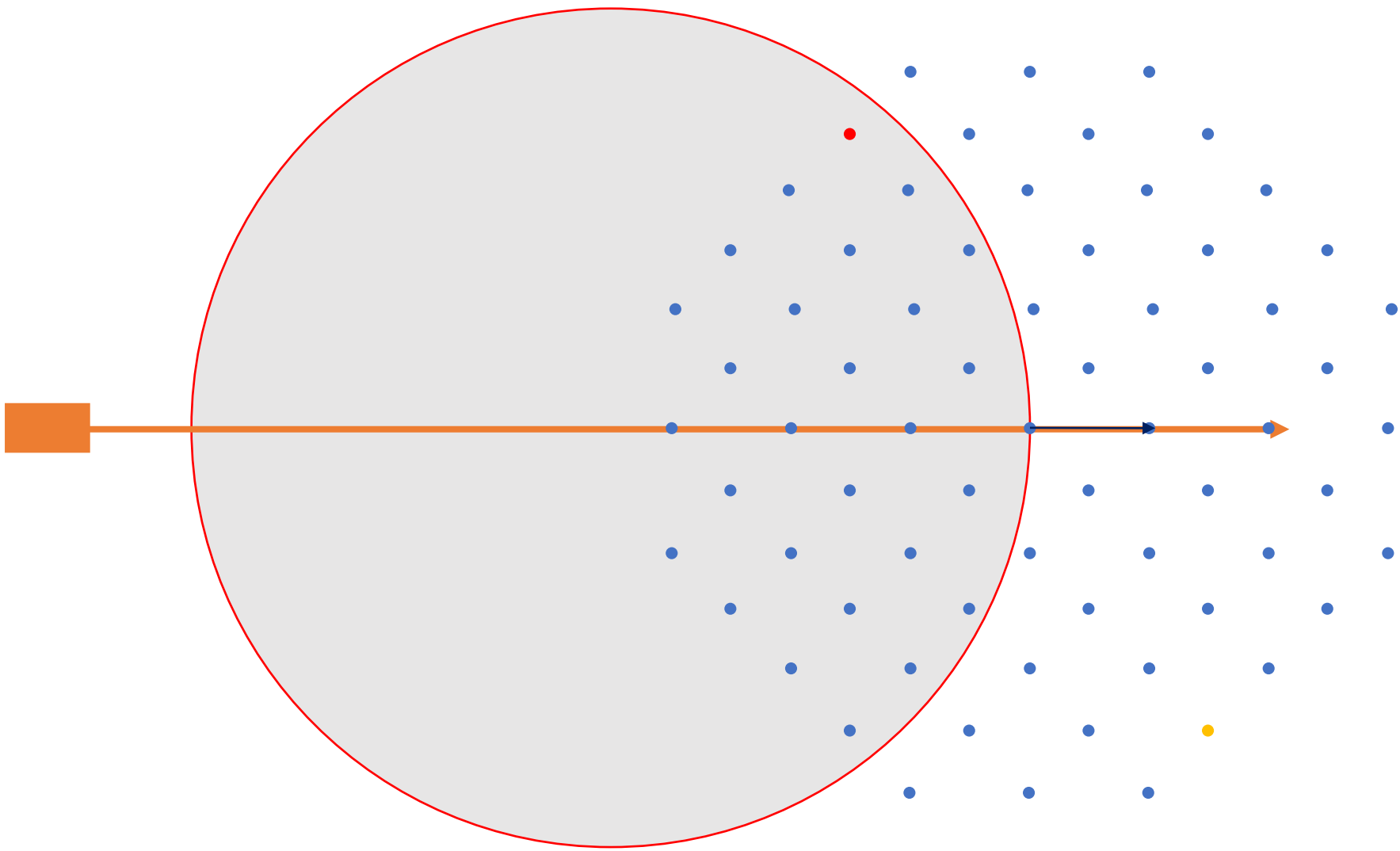
Prize share: 1/1

# The Ewald Sphere

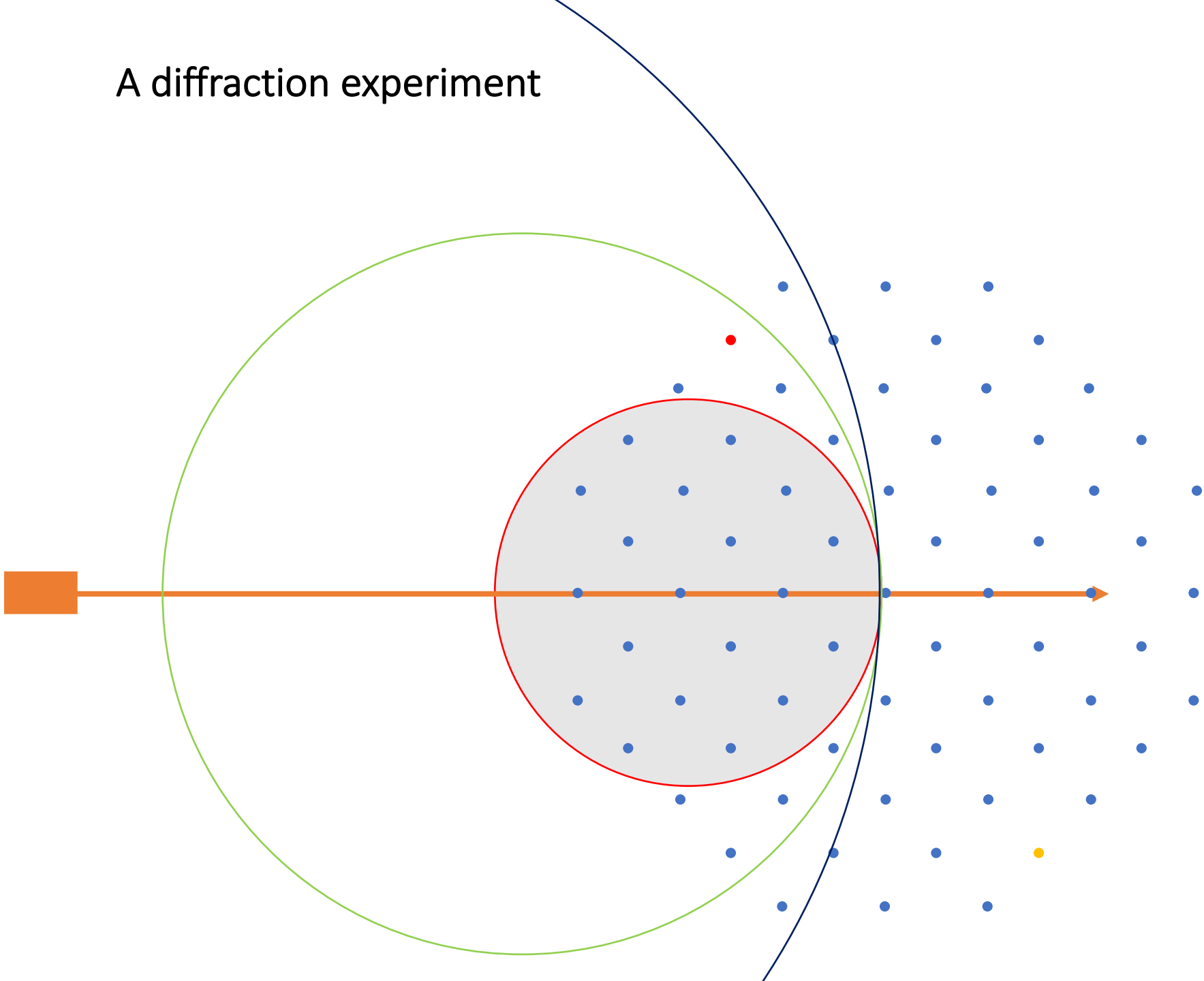


## A Laue diffraction experiment

# A diffraction experiment

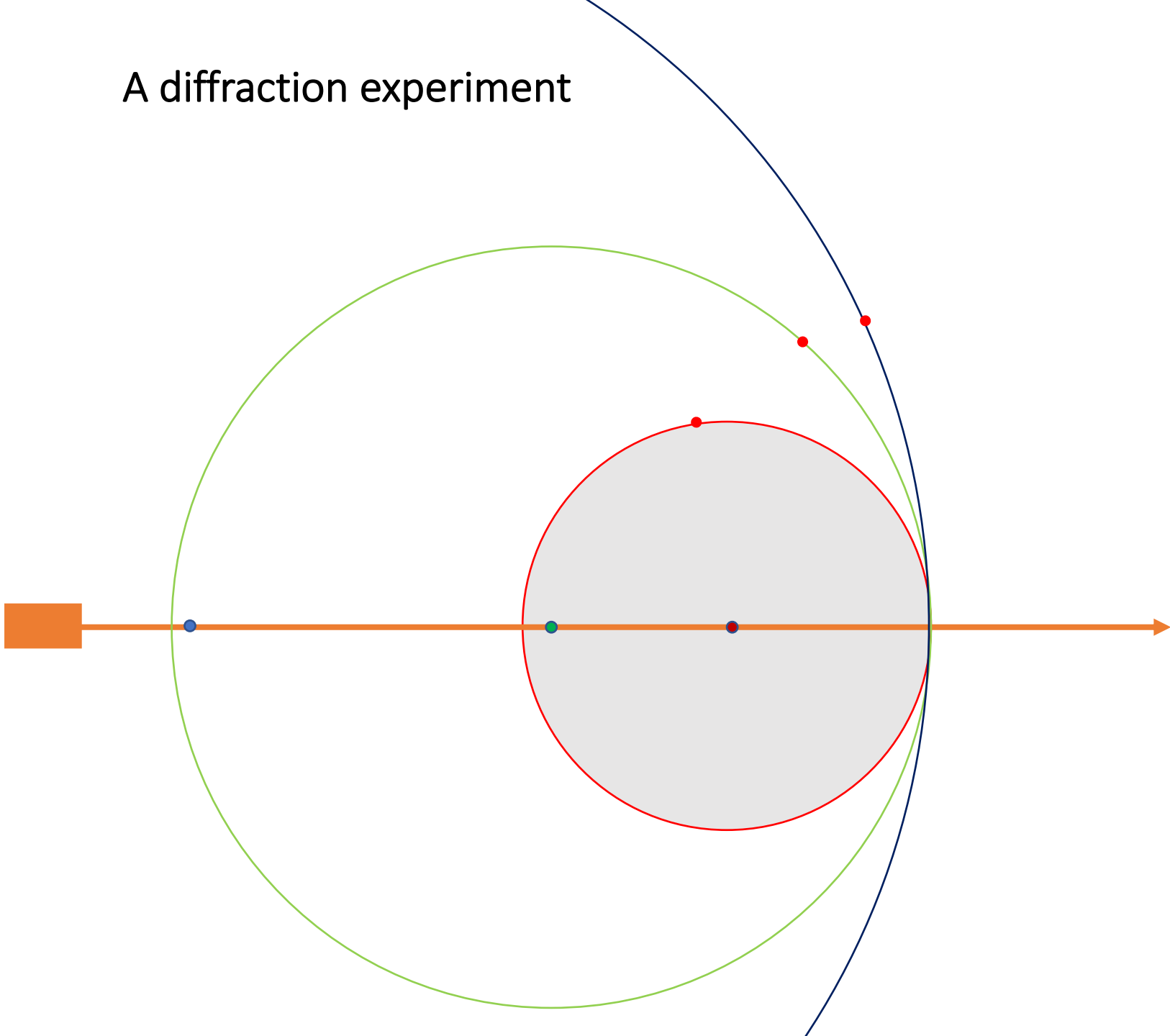


A diffraction experiment

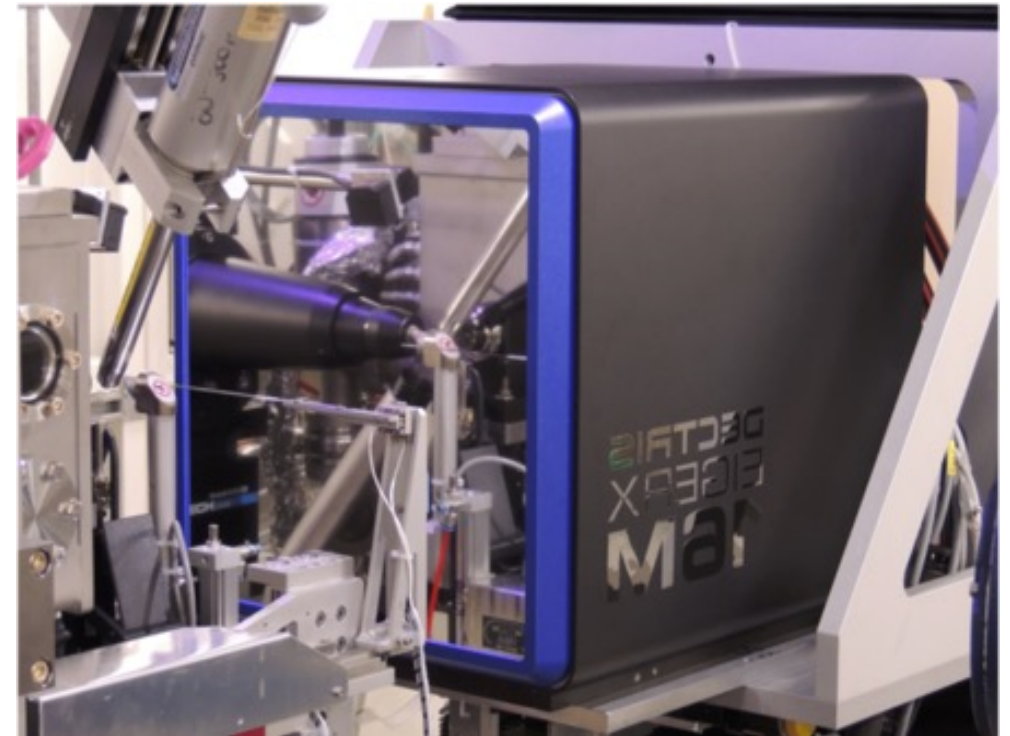
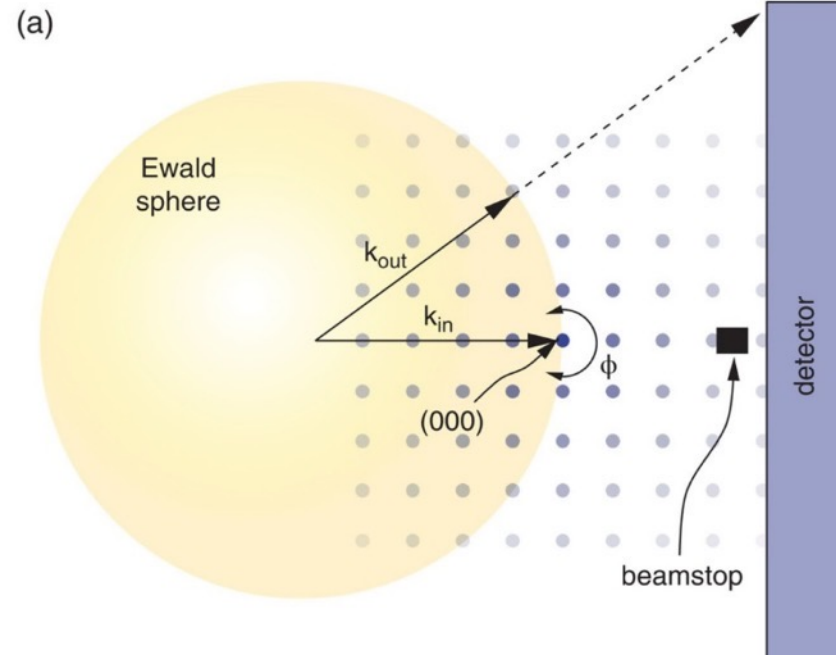




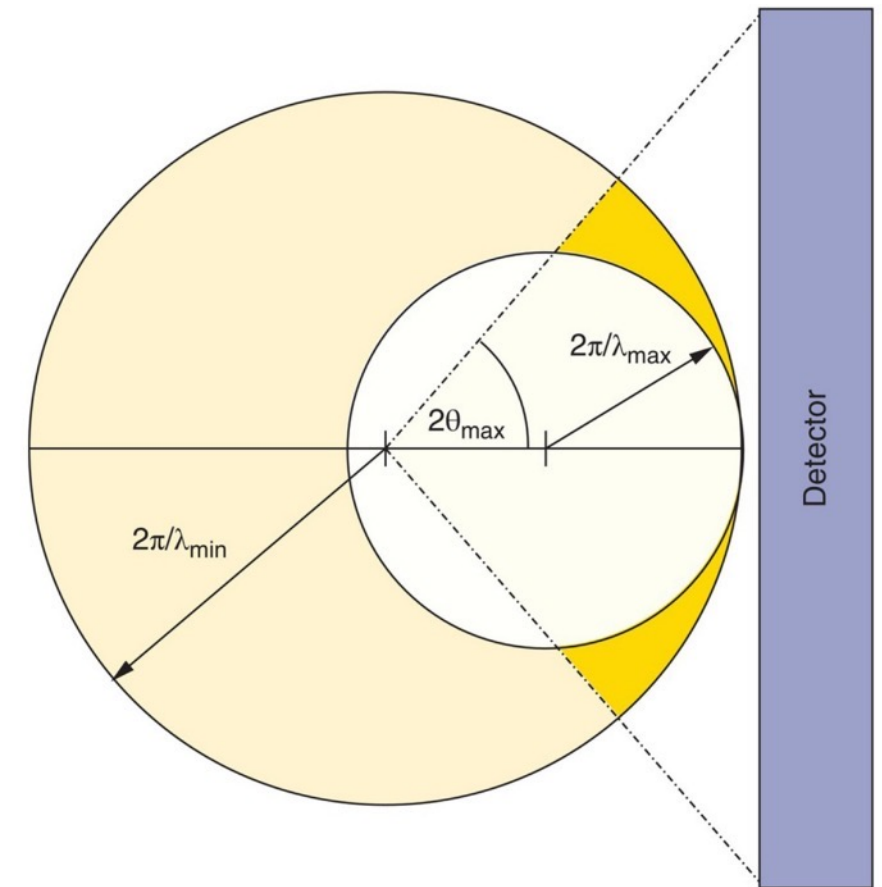
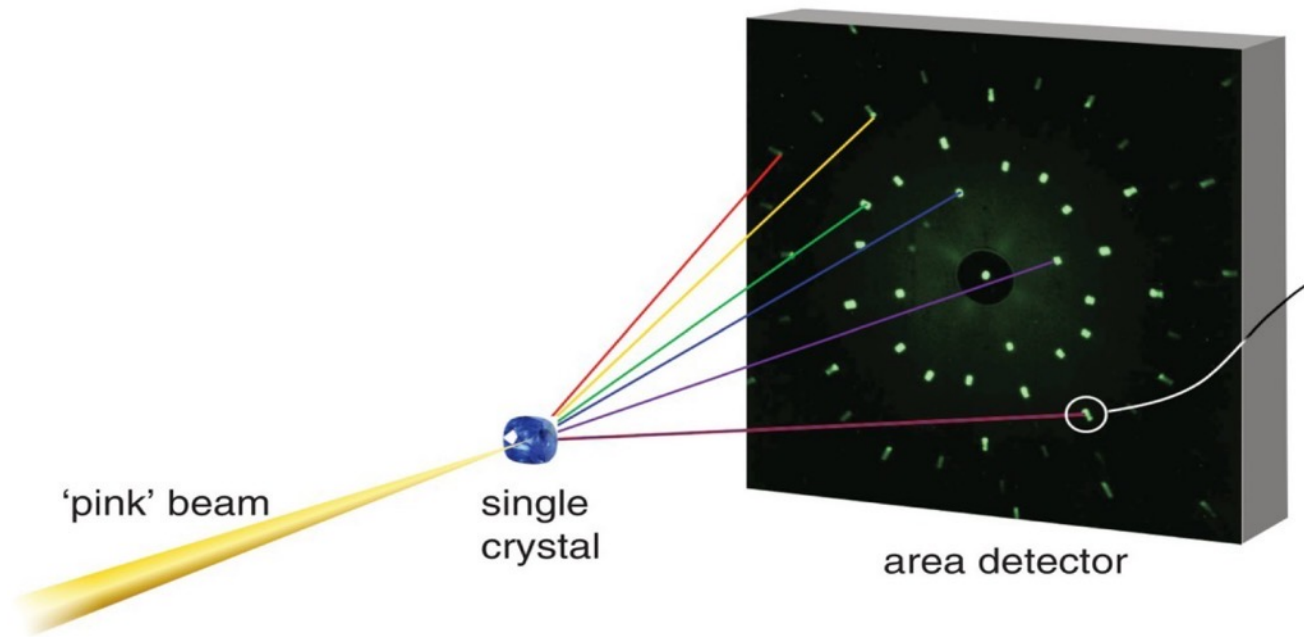
A diffraction experiment



# A diffraction experiment

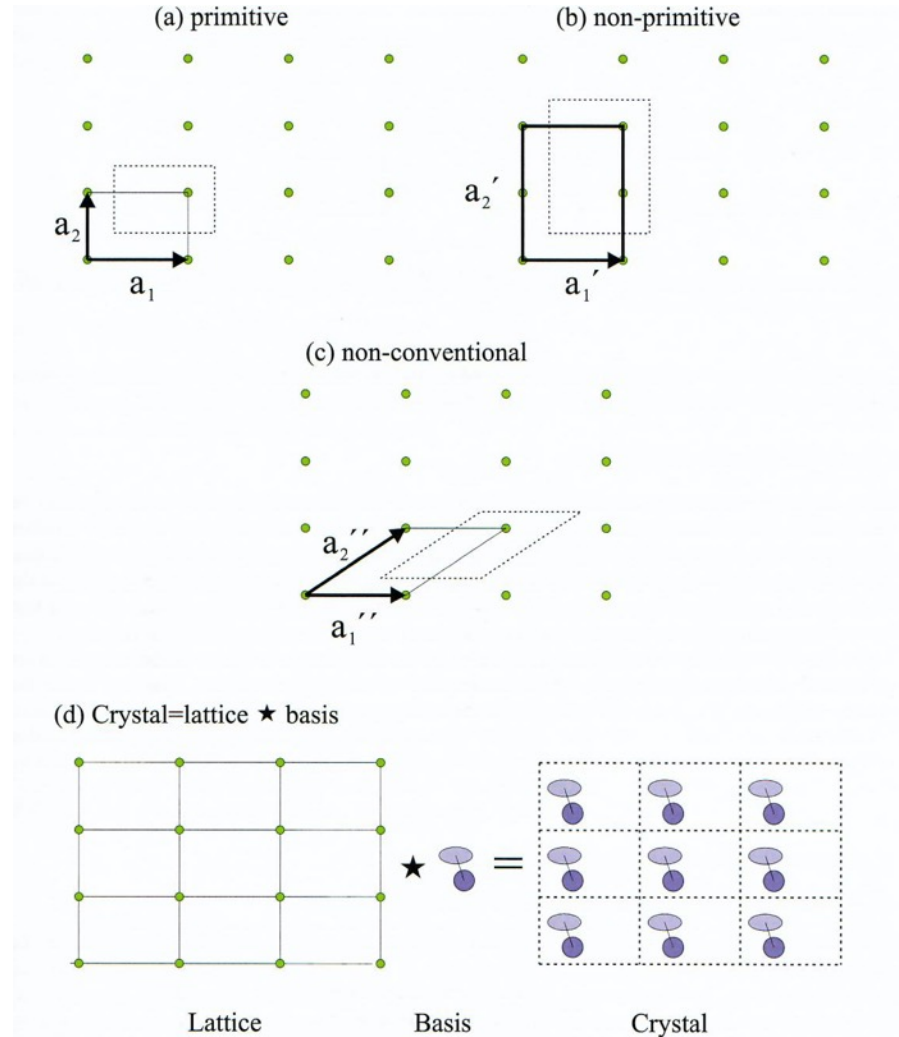


# Laue method

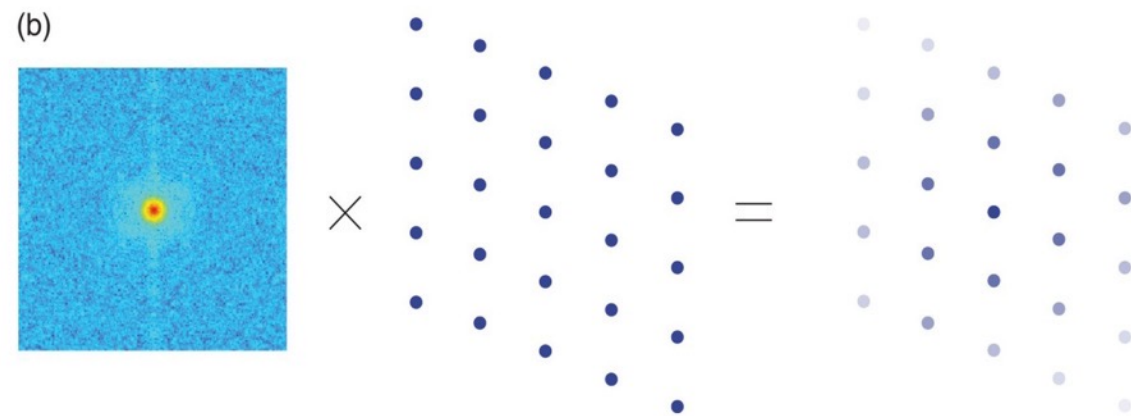
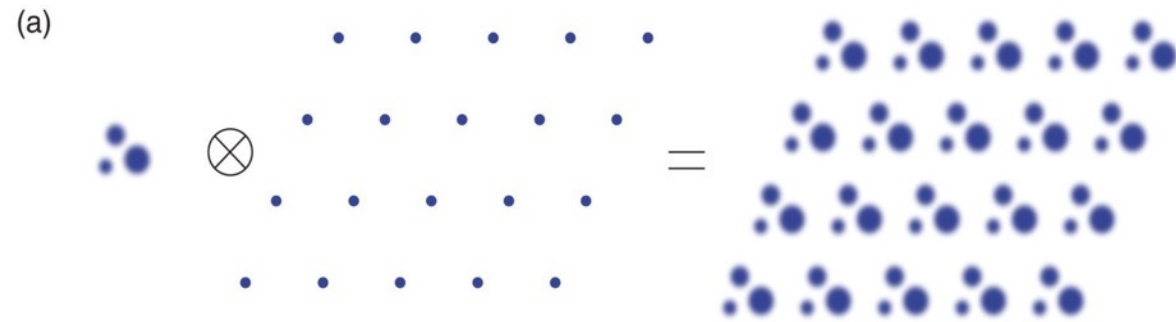


## More on Bragg peaks

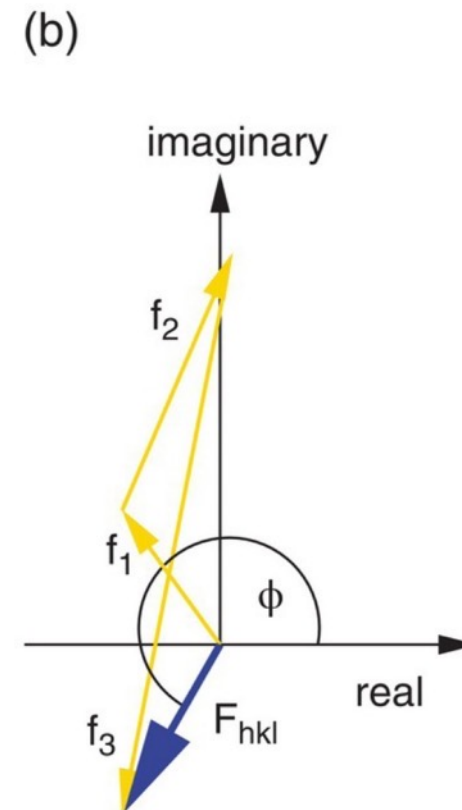
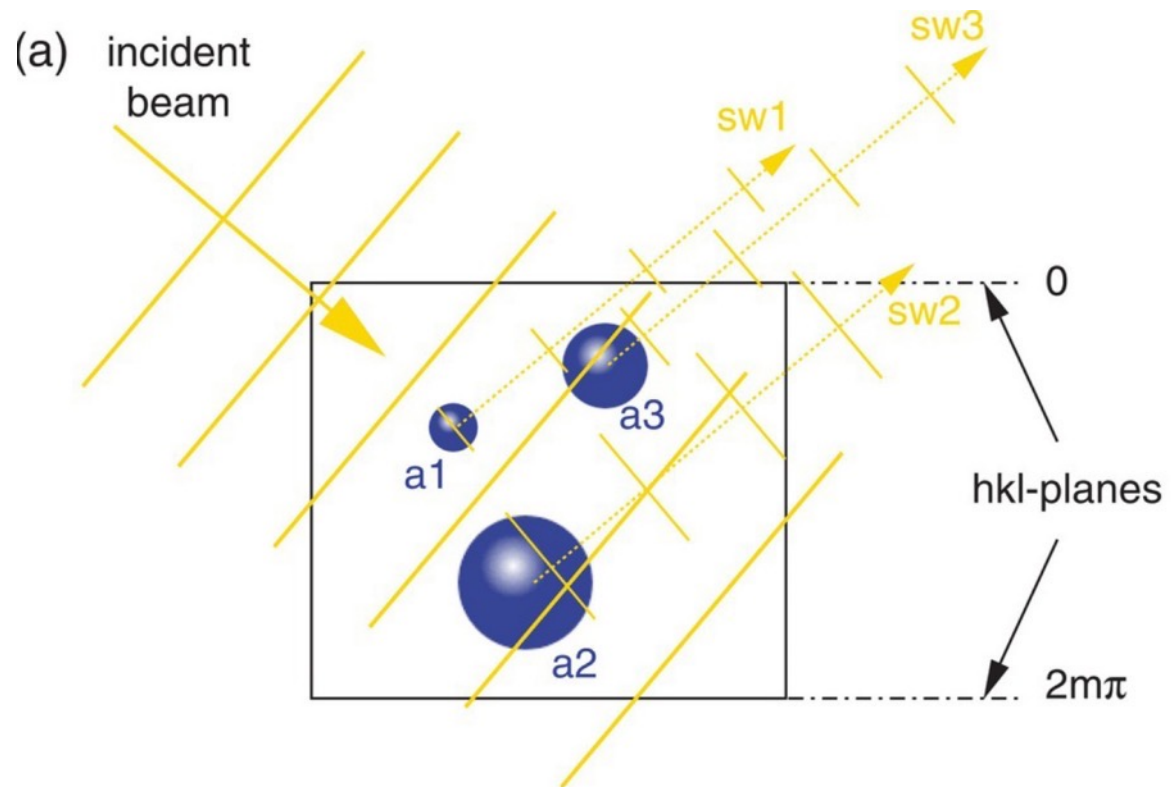
# A crystal is defined by its lattice and basis



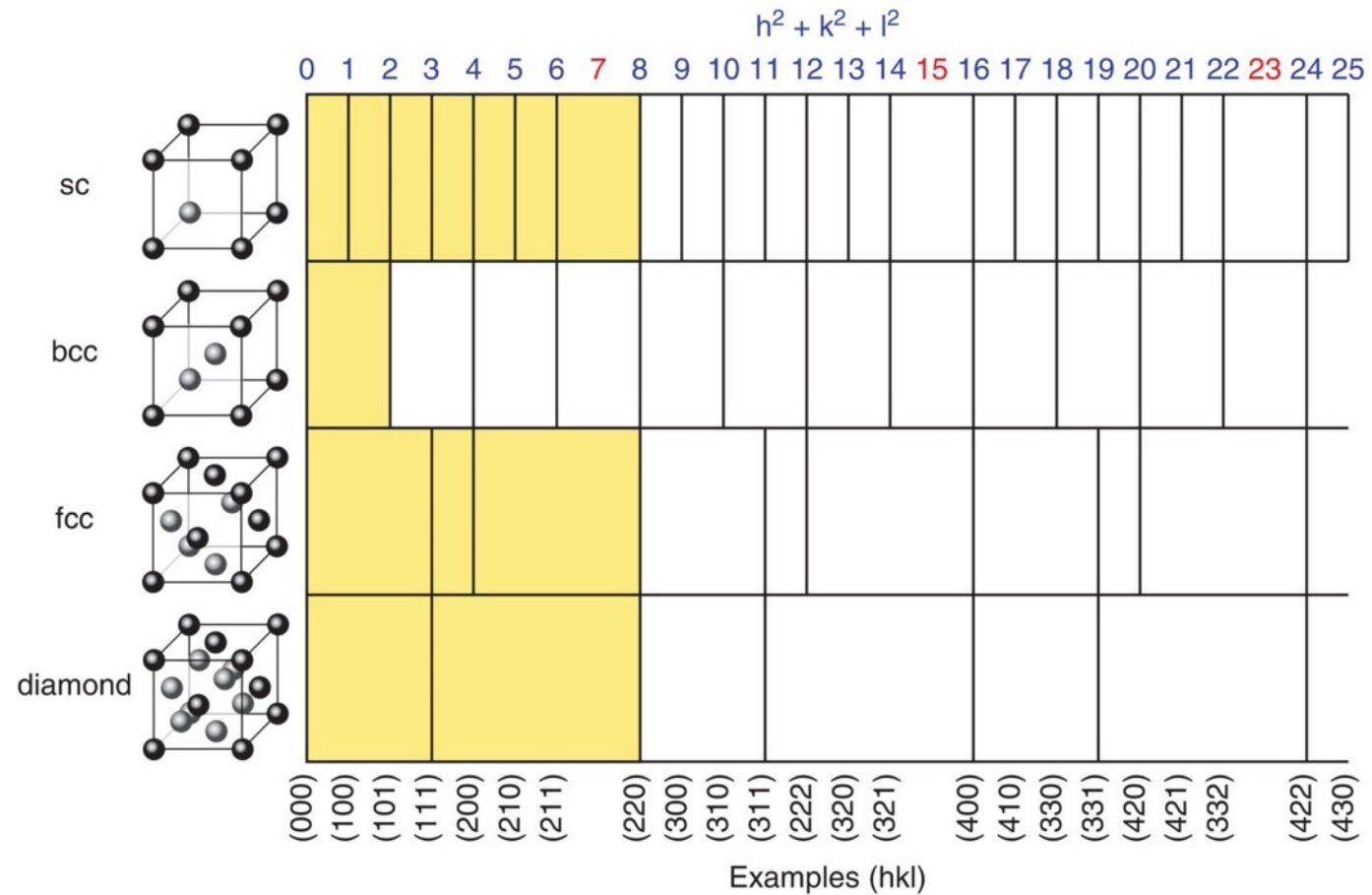
# The structure factor



# The structure factor

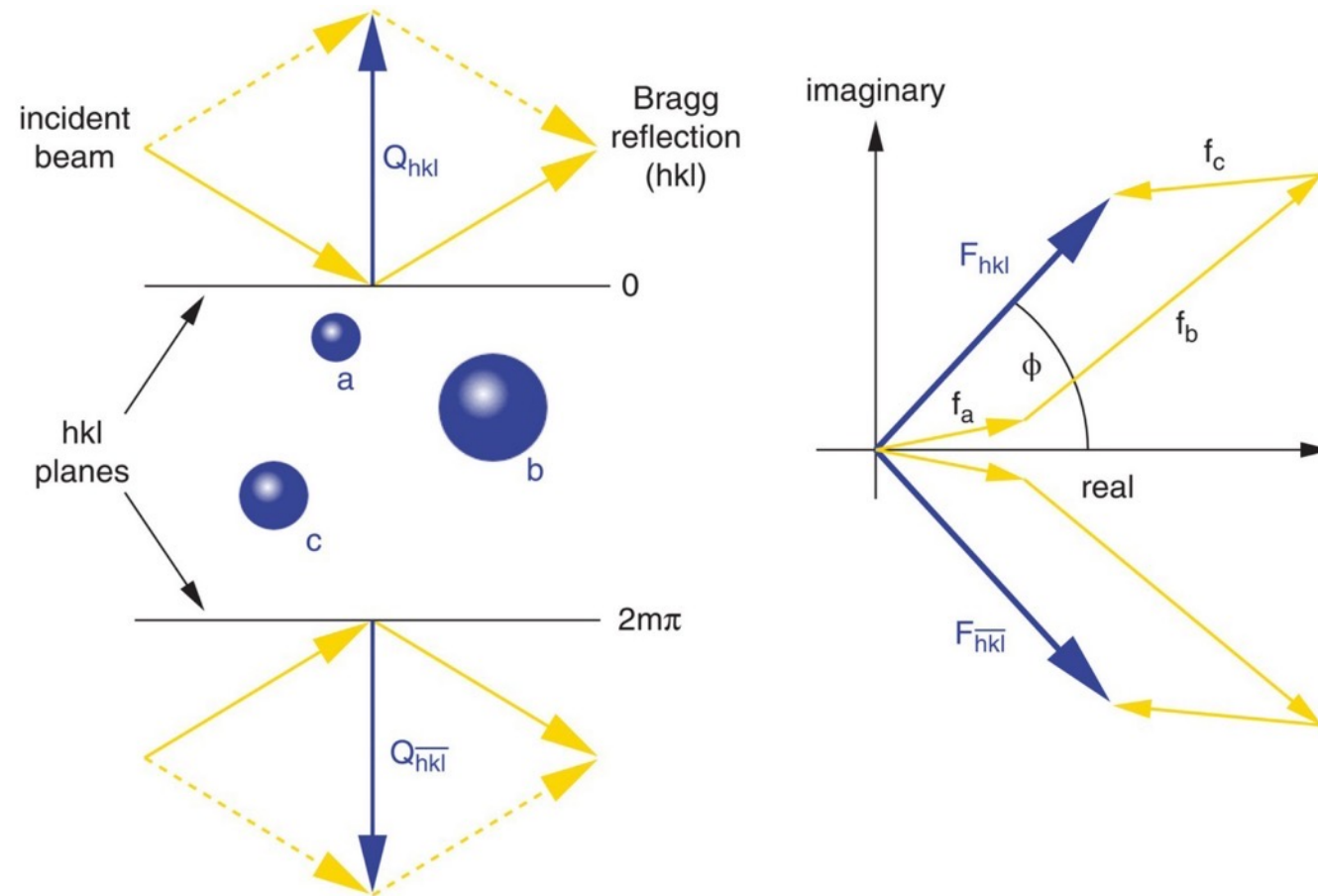


# Allowed and forbidden reflections

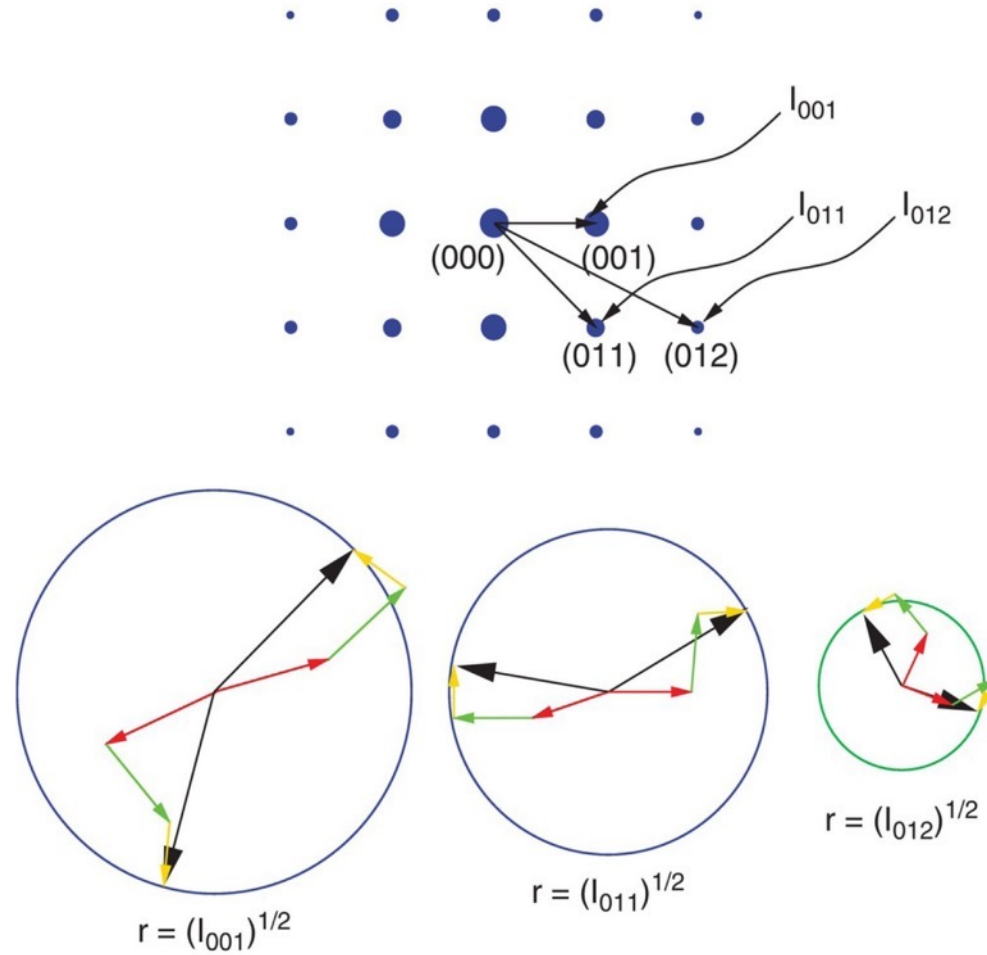




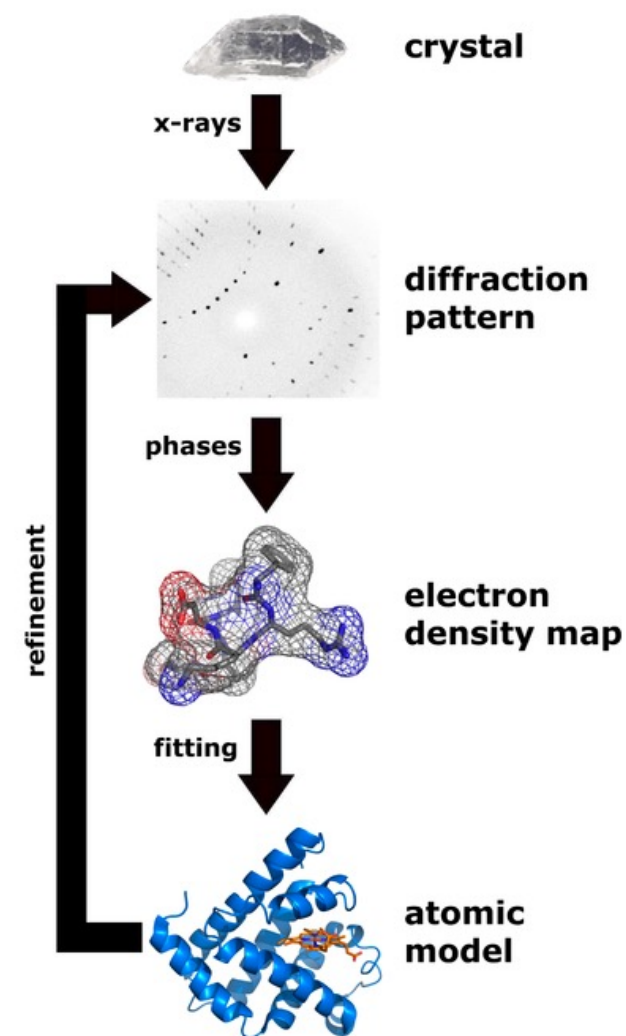
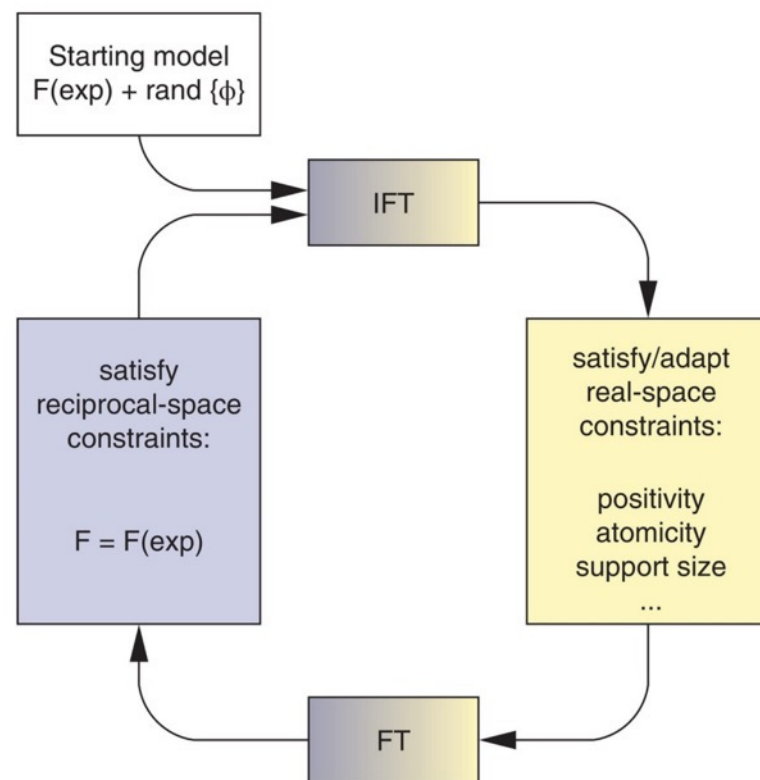
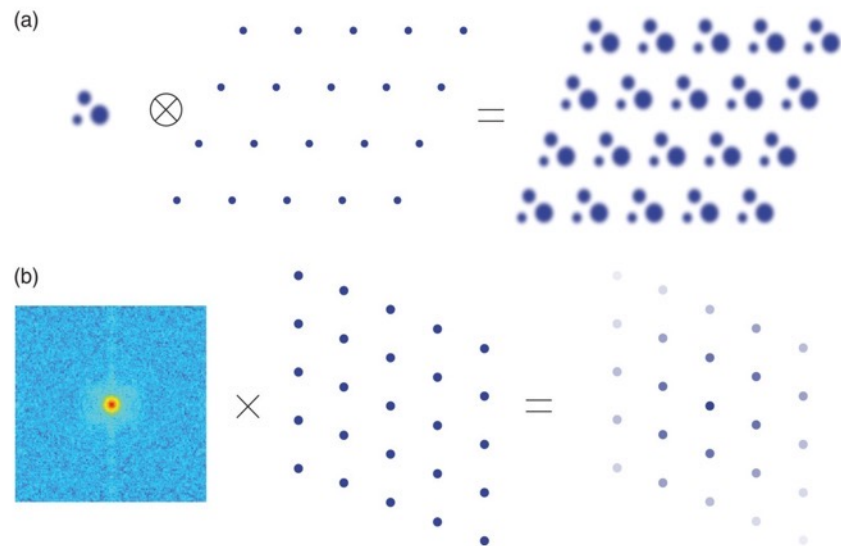
# Friedels law



# The phase problem



# Solving a crystal structure

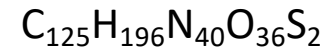


A few examples and case studies

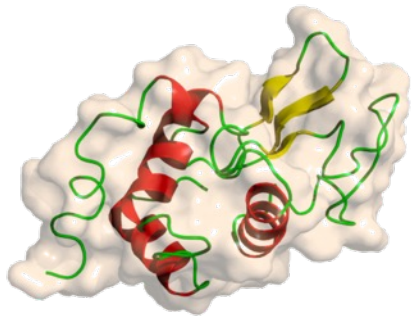
# Remember last week example for diffraction of lysozyme crystal

Lysozym (Protein, Immunsystem)

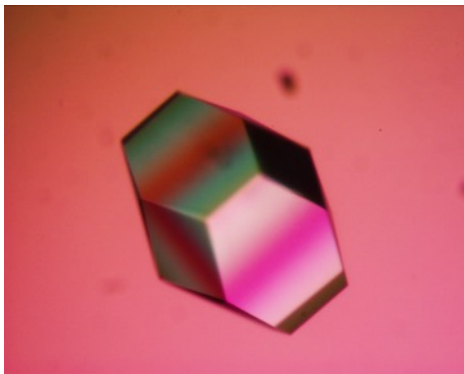
Stoichiometric formula



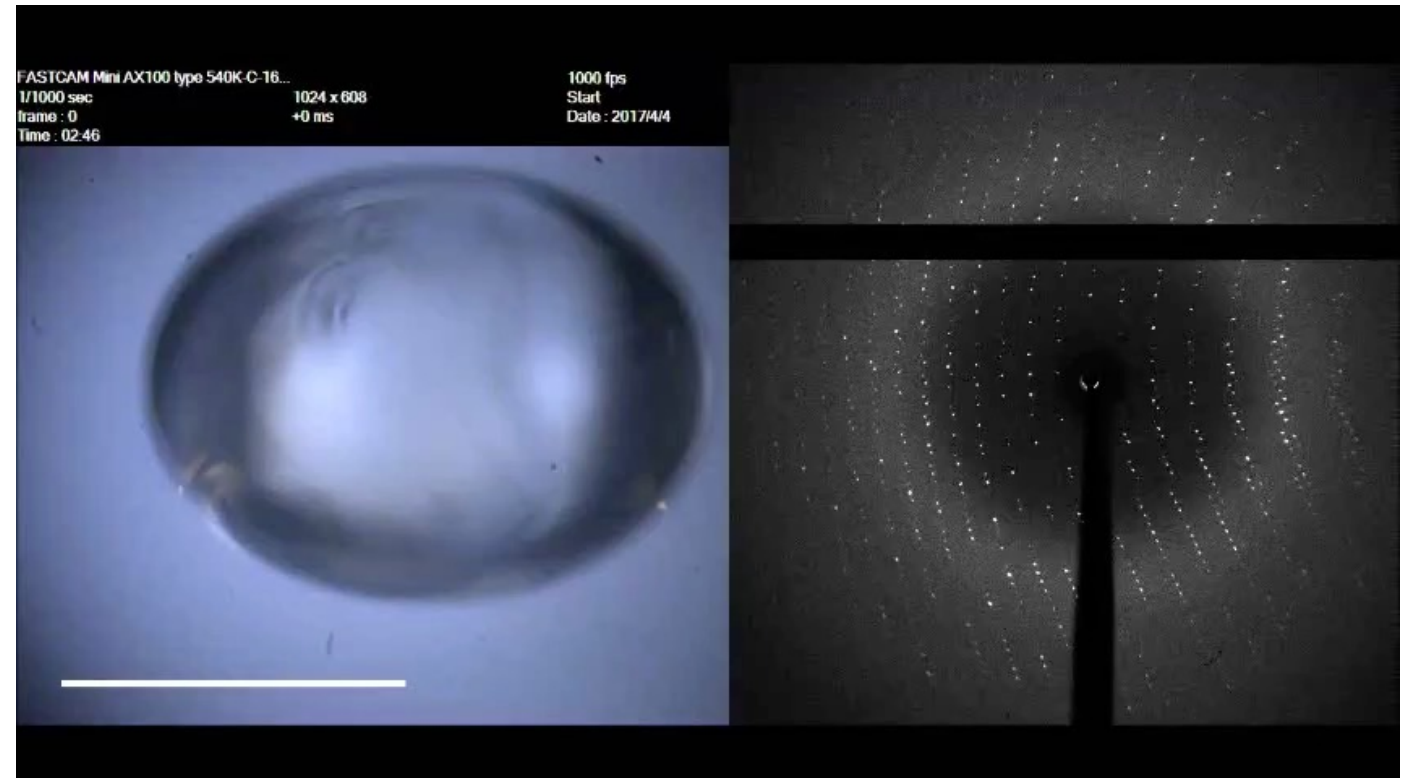
Structure



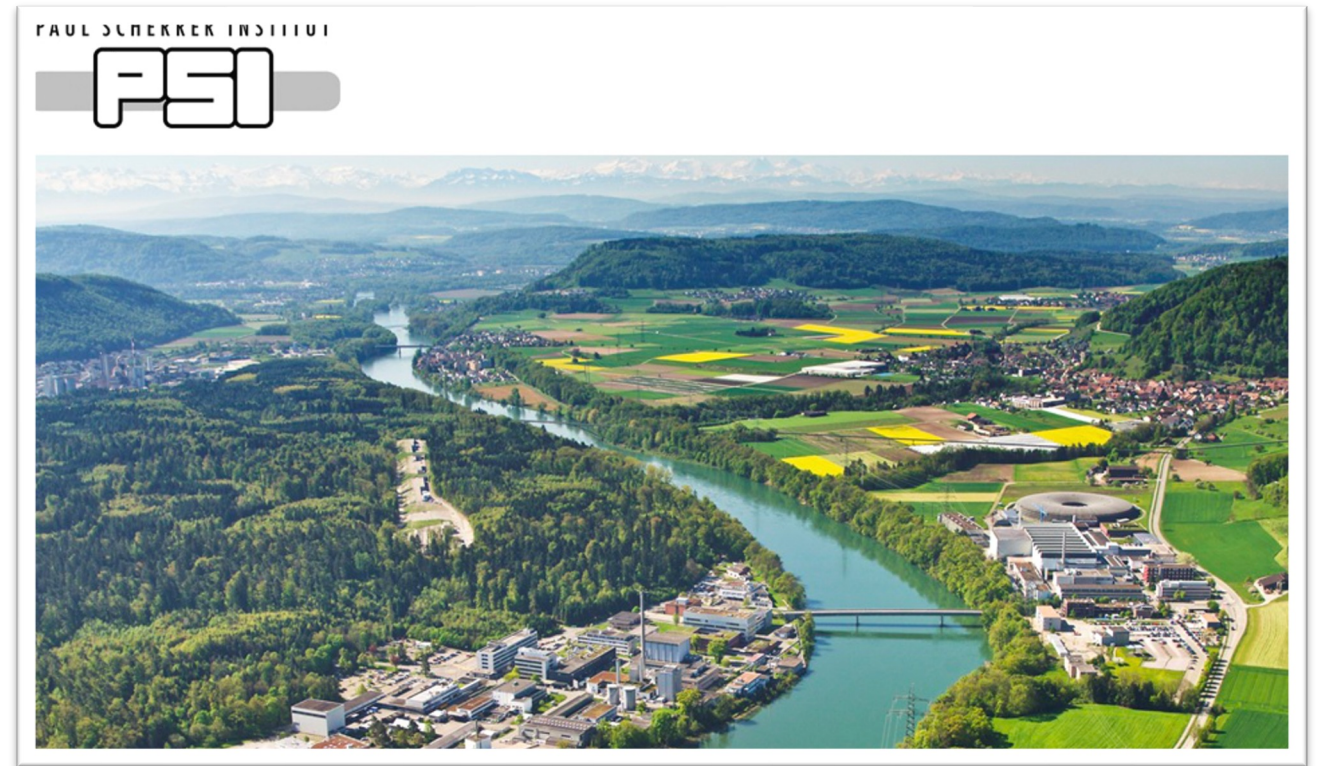
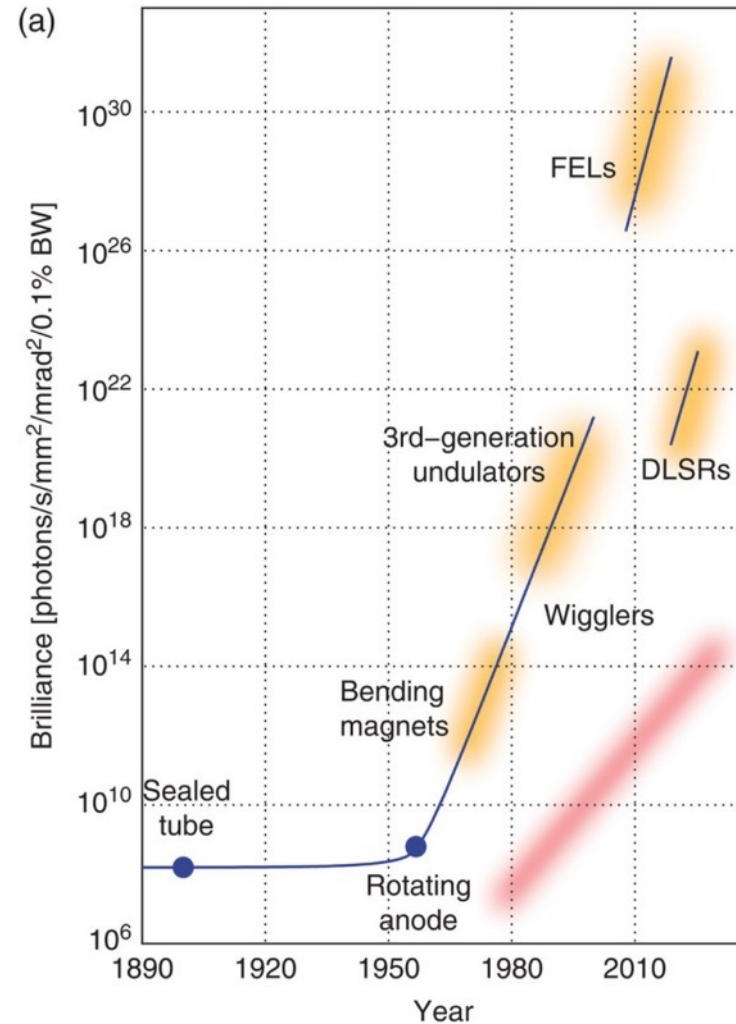
Crystal



Experiment at the Swiss Light Source



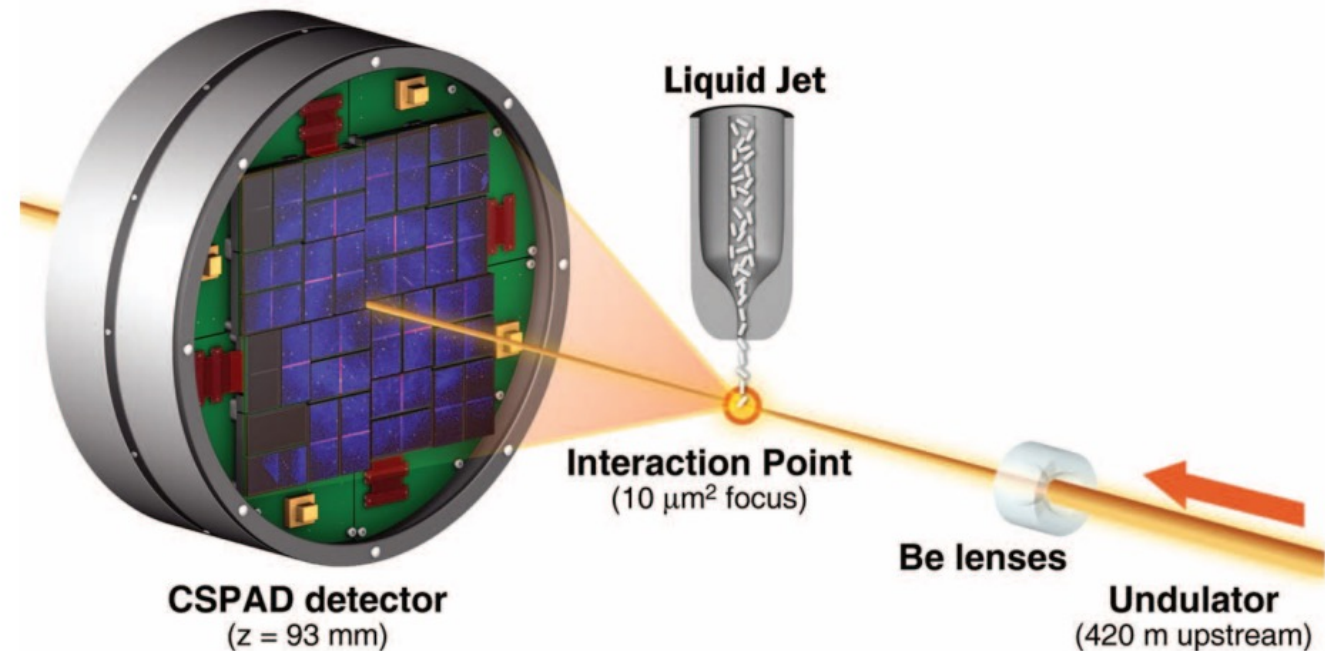
Each X-ray source allows characteristic experiments. „Output“ varies by >20 orders of magnitude



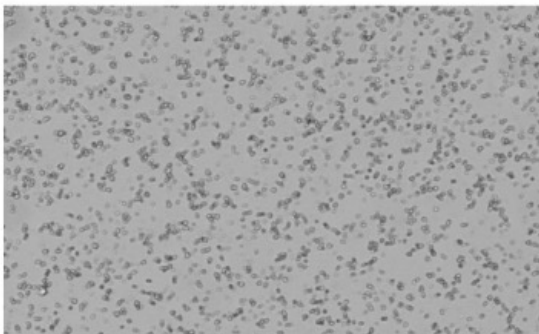


## High-Resolution Protein Structure Determination by Serial Femtosecond Crystallography

Sébastien Boutet,<sup>1\*</sup> Lukas Lomb,<sup>2,3</sup> Garth J. Williams,<sup>1</sup> Thomas R. M. Barends,<sup>2,3</sup> Andrew Aquila,<sup>4</sup> R. Bruce Doak,<sup>5</sup> Uwe Weierstall,<sup>5</sup> Daniel P. DePonte,<sup>4</sup> Jan Steinbrener,<sup>2,3</sup> Robert L. Shoeman,<sup>2,3</sup> Marc Messerschmidt,<sup>1</sup> Anton Barty,<sup>4</sup> Thomas A. White,<sup>4</sup> Stephan Kassemeyer,<sup>2,3</sup> Richard A. Kirian,<sup>5</sup> M. Marvin Seibert,<sup>1</sup> Paul A. Montanez,<sup>1</sup> Chris Kenney,<sup>6</sup> Ryan Herbst,<sup>6</sup> Philip Hart,<sup>6</sup> Jack Pines,<sup>6</sup> Gunther Haller,<sup>6</sup> Sol M. Gruner,<sup>7,8</sup> Hugh T. Philipp,<sup>7</sup> Mark W. Tate,<sup>7</sup> Marianne Hromalik,<sup>9</sup> Lucas J. Koerner,<sup>10</sup> Niels van Bakel,<sup>11</sup> John Morse,<sup>12</sup> Wilfred Ghonsalves,<sup>1</sup> David Arnlund,<sup>13</sup> Michael J. Bogan,<sup>14</sup> Carl Caleman,<sup>4</sup> Raimund Fromme,<sup>15</sup> Christina Y. Hampton,<sup>14</sup> Mark S. Hunter,<sup>15</sup> Linda C. Johansson,<sup>13</sup> Gergely Katona,<sup>13</sup> Christopher Kupitz,<sup>15</sup> Mengning Liang,<sup>4</sup> Andrew V. Martin,<sup>4</sup> Karol Nass,<sup>16</sup> Lars Redecke,<sup>17,18</sup> Francesco Stellato,<sup>4</sup> Nicusor Timneanu,<sup>19</sup> Dingjie Wang,<sup>5</sup> Nadia A. Zatsepin,<sup>5</sup> Donald Schafer,<sup>1</sup> James DeFeaver,<sup>1</sup> Richard Neutze,<sup>13</sup> Petra Fromme,<sup>15</sup> John C. H. Spence,<sup>5</sup> Henry N. Chapman,<sup>4,16</sup> Ilme Schlichting<sup>2,3</sup>



B



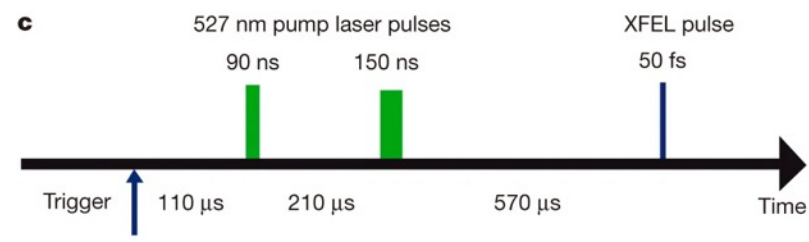
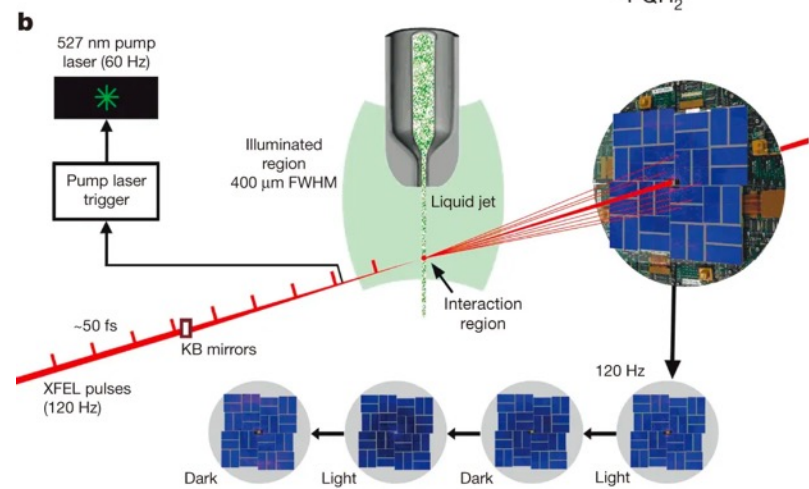
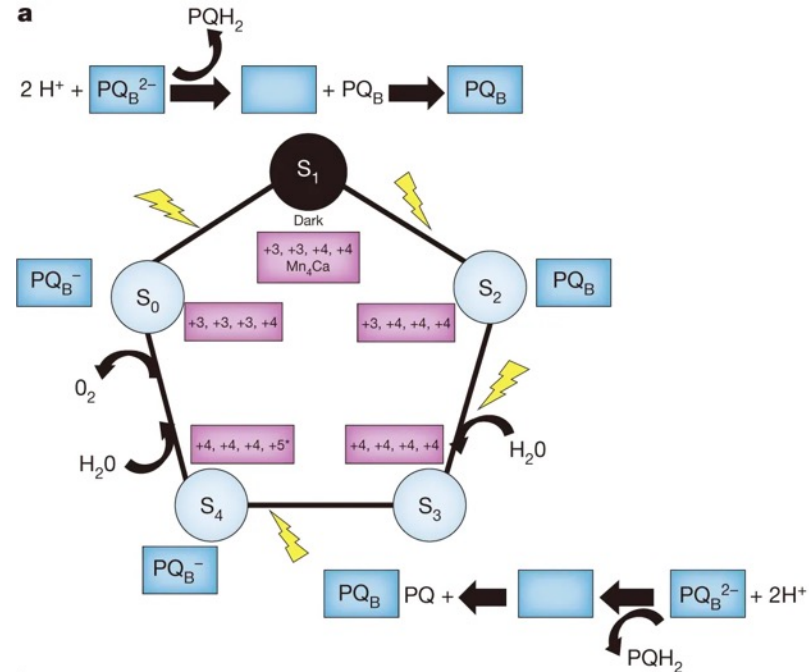
C



## Letter | Published: 09 July 2014

## Christopher Kupitz, Shibom Basu, [...] Petra From

*Nature* **513**, 261–265(2014) | [Cite this article](#)





# Recent results from SwissFEL

## Article

## Femtosecond-to-millisecond structural changes in a light-driven sodium pump

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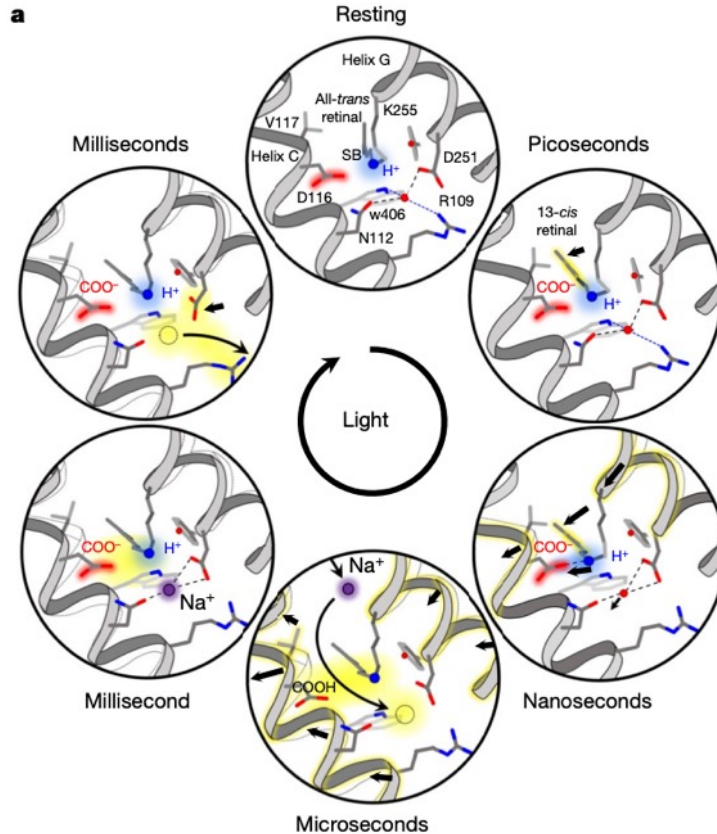
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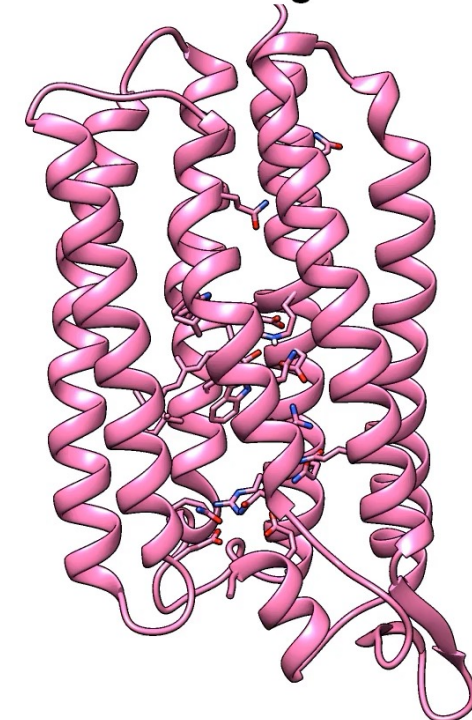
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Light-driven sodium pumps actively transport small cations across cellular membranes<sup>1</sup>. These pumps are used by microorganisms to convert light into membrane potential and have become useful optogenetic tools with applications in neuroscience. Although the resting state structures of the prototypical sodium pump *Krokinobacter eikastus* rhodopsin 2 (KR2) have been solved<sup>2,3</sup>, it is unclear how structural alterations over time allow sodium to be translocated against a concentration gradient. Here, using the Swiss X-ray Free Electron Laser<sup>4</sup>, we have collected serial crystallographic data at ten pump-probe delays from femtoseconds to milliseconds. High-resolution structural snapshots throughout the KR2 photocycle show how retinal isomerization is completed on the femtosecond timescale and changes the local structure of the binding pocket in the early nanoseconds. Subsequent rearrangements and deprotonation of the retinal Schiff base open an electrostatic gate in microseconds. Structural and spectroscopic data, in combination with quantum chemical calculations, indicate that a sodium ion binds transiently close to the retinal within one millisecond. In the last structural intermediate, at 20 milliseconds after activation, we identified a potential second sodium-binding site close to the extracellular exit. These results provide direct molecular insight into the dynamics of active cation transport across biological membranes.



## KR2 resting state



The end