



8

Electron spectroscopy 1



Physics of Surfaces and Interfaces

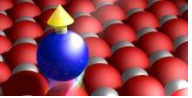
H. Ibach

Springer

Surface Science - An Introduction

K.Oura, V. G. Lifshits, A. A. Saranin, A.V. Zotov, M. Katayama

Springer



Electronic devices (FET, FGT, ...), created by successive processes (lithography, etching, growth, implantation,), are composed of many “layers”

Detection of electron confinement in quantum wells (thin films, heterostructures)

Magnetic tunnel junction (ferromagnet/insulating layers): increased performances by improving the sharpness of interfaces

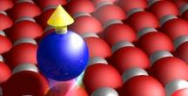
How to verify composition and cleanliness?

How to verify mixing / alloying at interfaces?

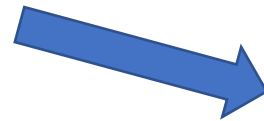
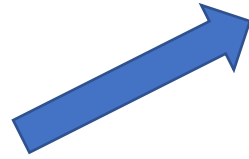
How to determine the density of states?

How to determine the band structure?

How to detect hybridization or bond formation?



Electron spectroscopy



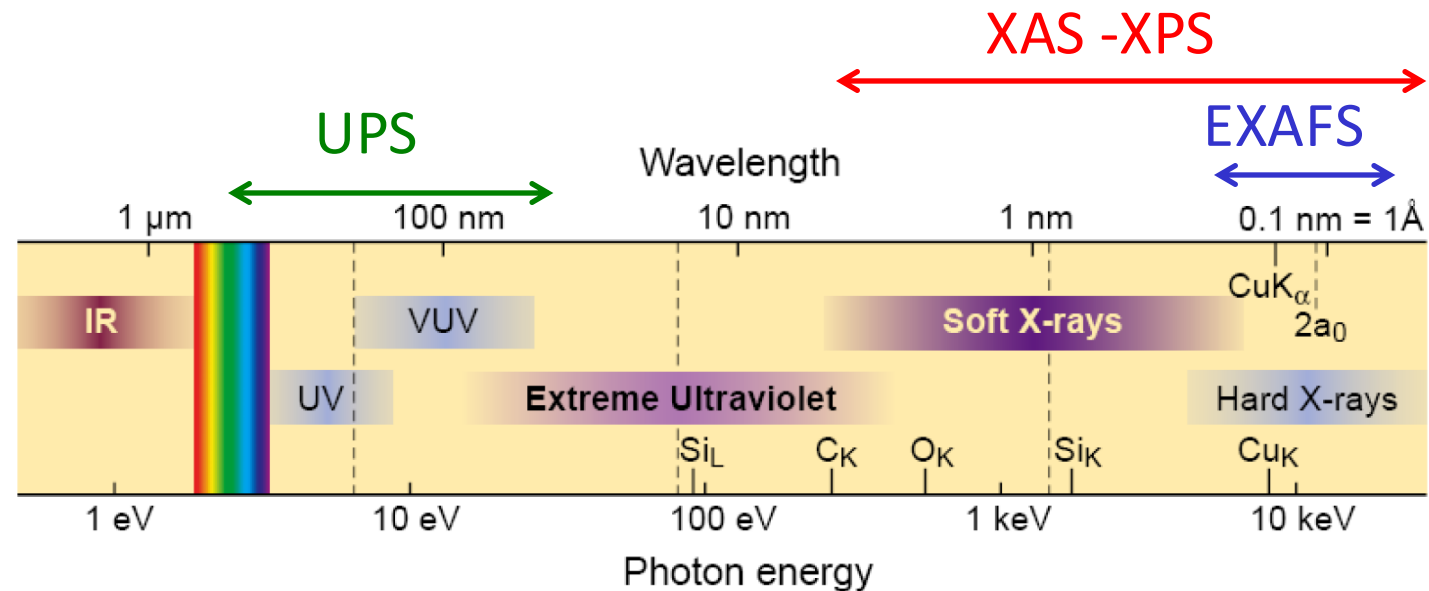
Chemical composition	→ XPS, Auger
Valence DOS	→ UPS
Band structure	→ ARUPS (ARPES)
Magnetism	→ XMCD
Bond orientation	→ XLD
Structure modification	→ EXAFS
...	
...	



Technique	Probe	Measure
Auger	electron	electron
XPS, UPS, ARPES	photon	electron
XAS, EXAFS	photon	photon / electron

The choice of the technique depends on the information we want to acquire

Electromagnetic spectrum

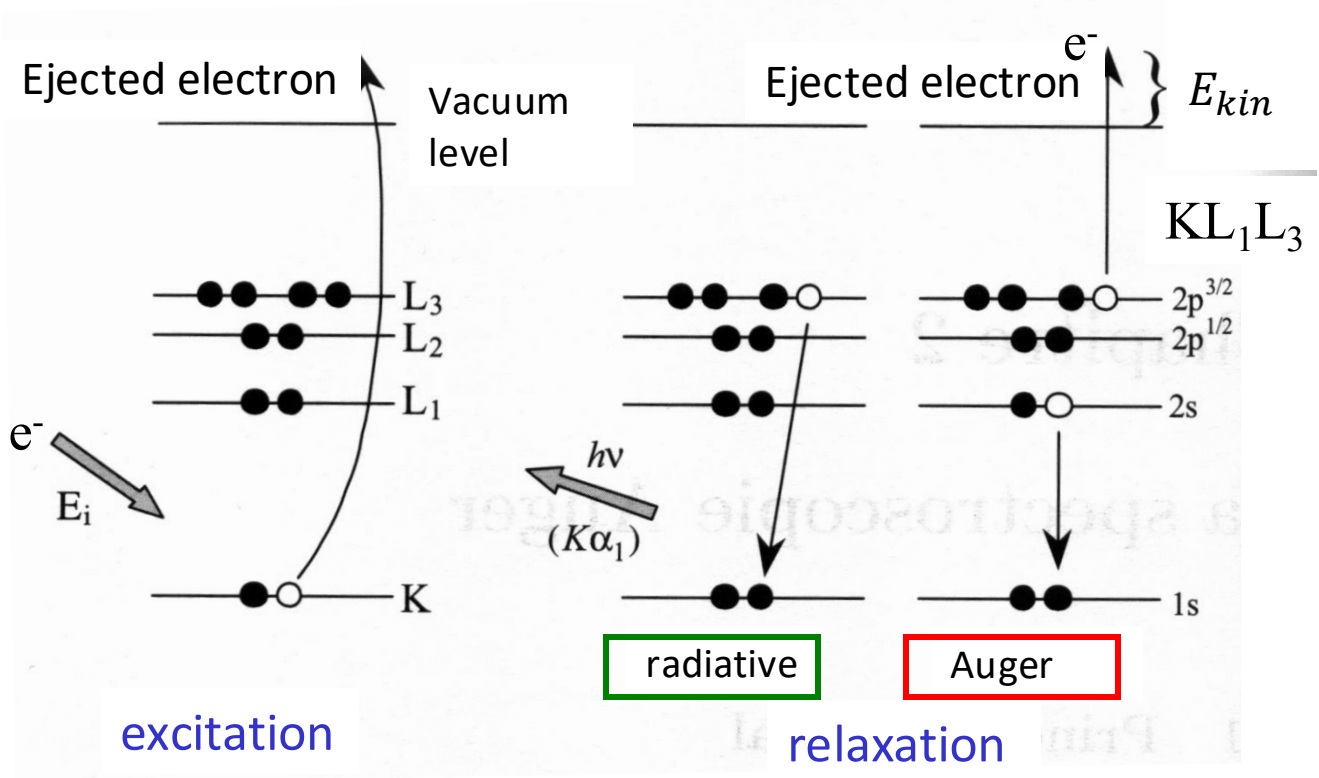


$$E = h\nu = hc/\lambda; c = 3 \cdot 10^8 \text{ m/s}; h = 6.6 \cdot 10^{-34} \text{ Js}$$



Exercise 8.1

Auger spectroscopy is based upon an electron in - electron out process.



E_{bind} : binding energy of the electronic level typically referred to E_F and not to E_{vac} → in principle one needs to consider also the work function Φ (3-5 eV)

$$E_{kin} = E_{bind}(K) - E_{bind}(L_1) - E_{bind}(L_3) - \Phi$$

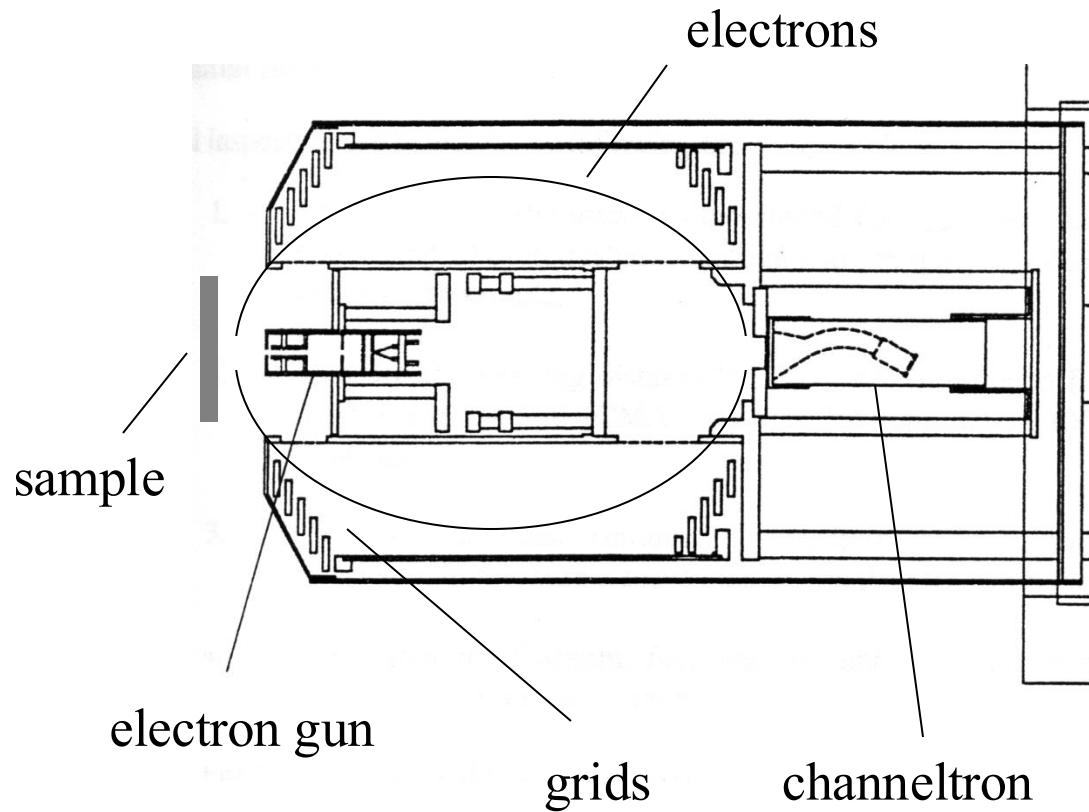
$E_{bind}(L_i)$ and $E_{bind}(K)$ depend on the element

E_{kin} does not depend on $E_{incident}$



chemical sensitivity

simple description neglecting screening and final state effects



Electron energy: in 1- 10 keV
 out 0.01 – 2 keV

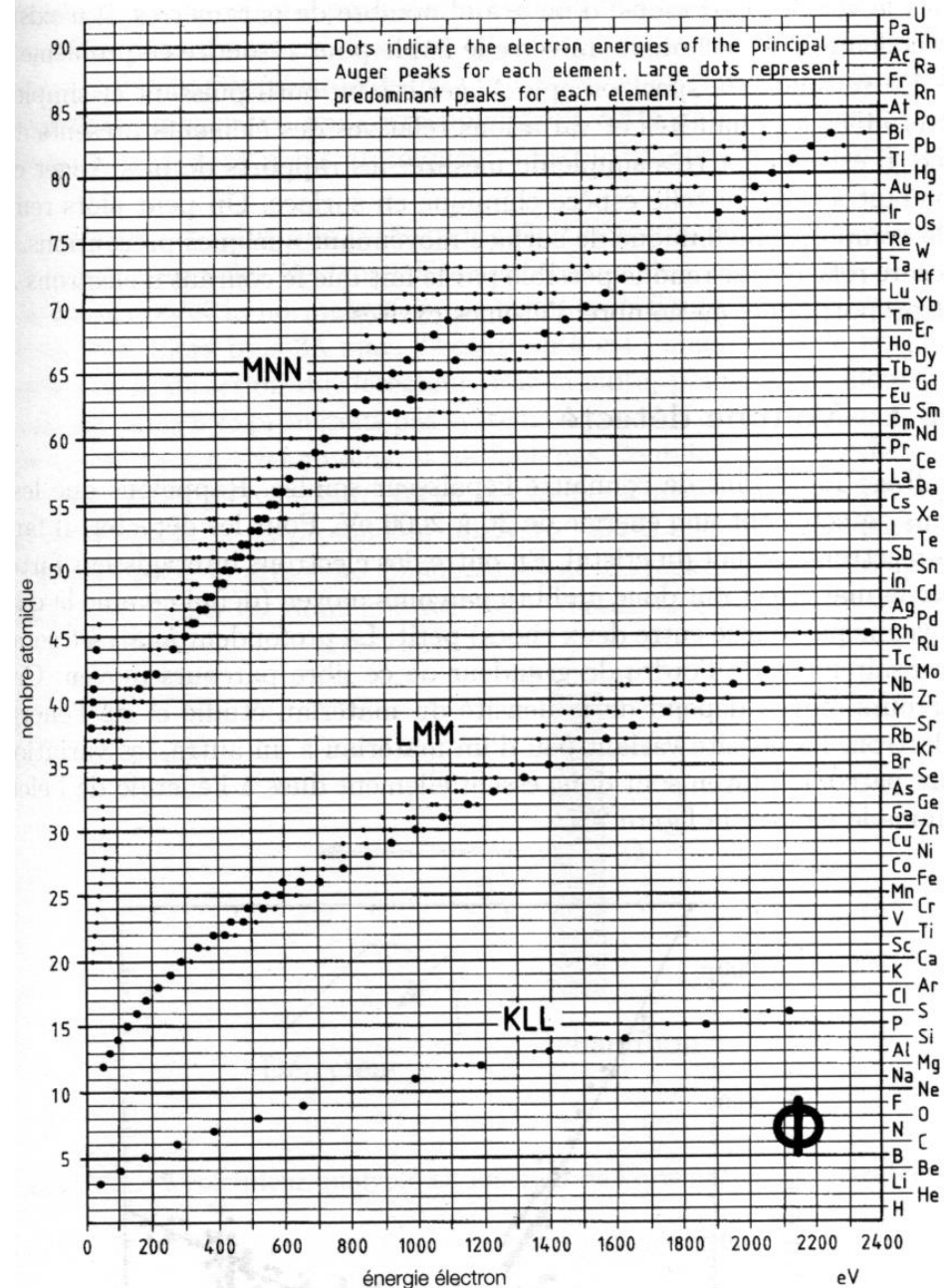
detection limit ~ 0.5% of a monolayer

Auger spectrum: number of emitted electrons as a function of their kinetic energy



Auger transitions for the different elements

The energy transitions are characteristic of the element
→ chemical sensitivity

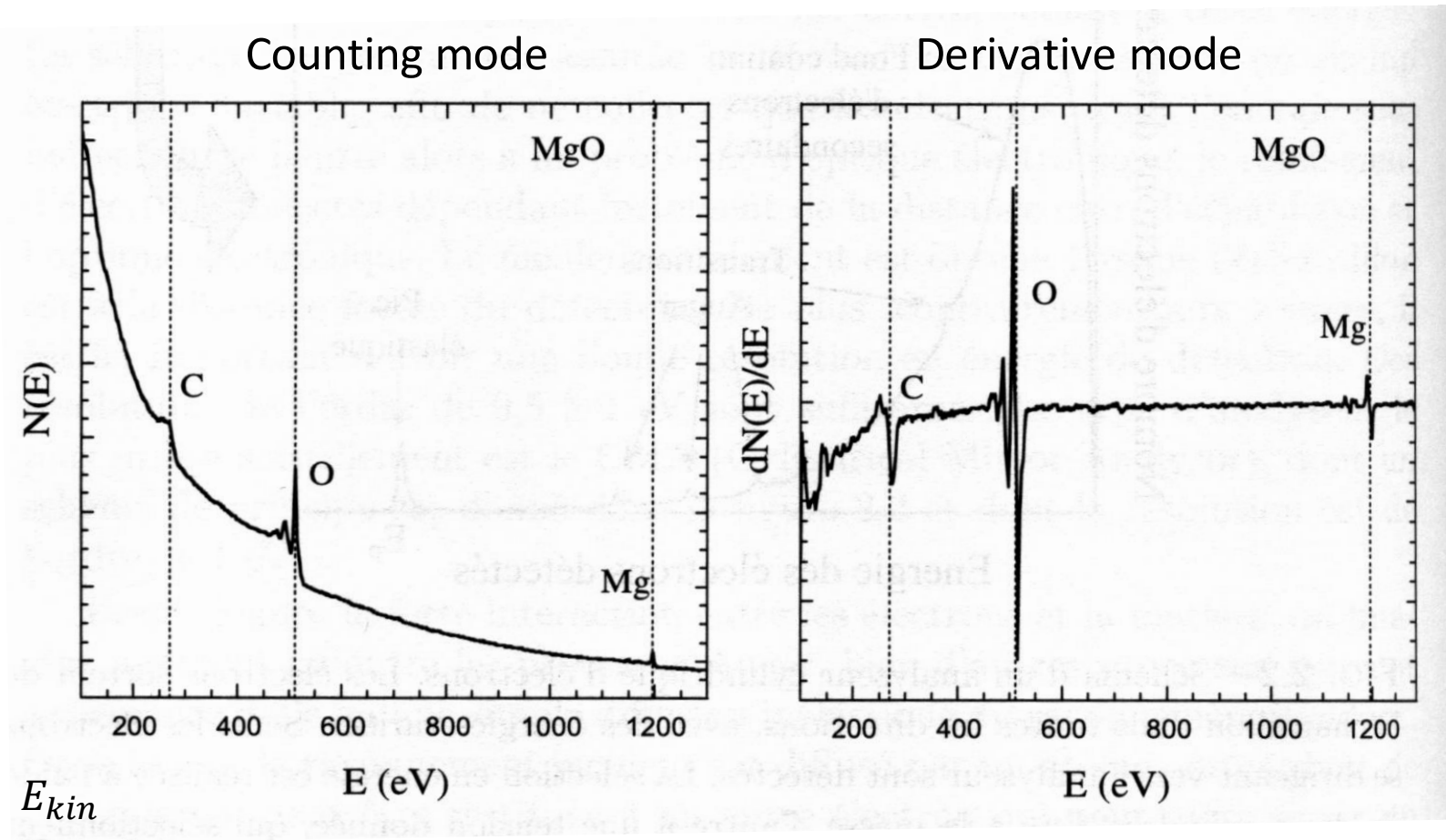
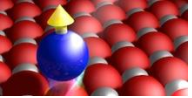


n: principal quantum number

MNN
involved electrons:
n=3 (excitation)
n=4

LMM
involved electrons:
n=2 (excitation)
n=3

KLL
involved electrons:
n=1 (excitation)
n=2



The monotonic background is due to multi-scattered electrons

Relative sensitivity: correction factors must be considered to obtain the real intensity



Auger electrons (detected at surface) coming from depth z :

$$I(z) = I(0) \exp(-z/\lambda)$$

total Auger electron current measured at surface:

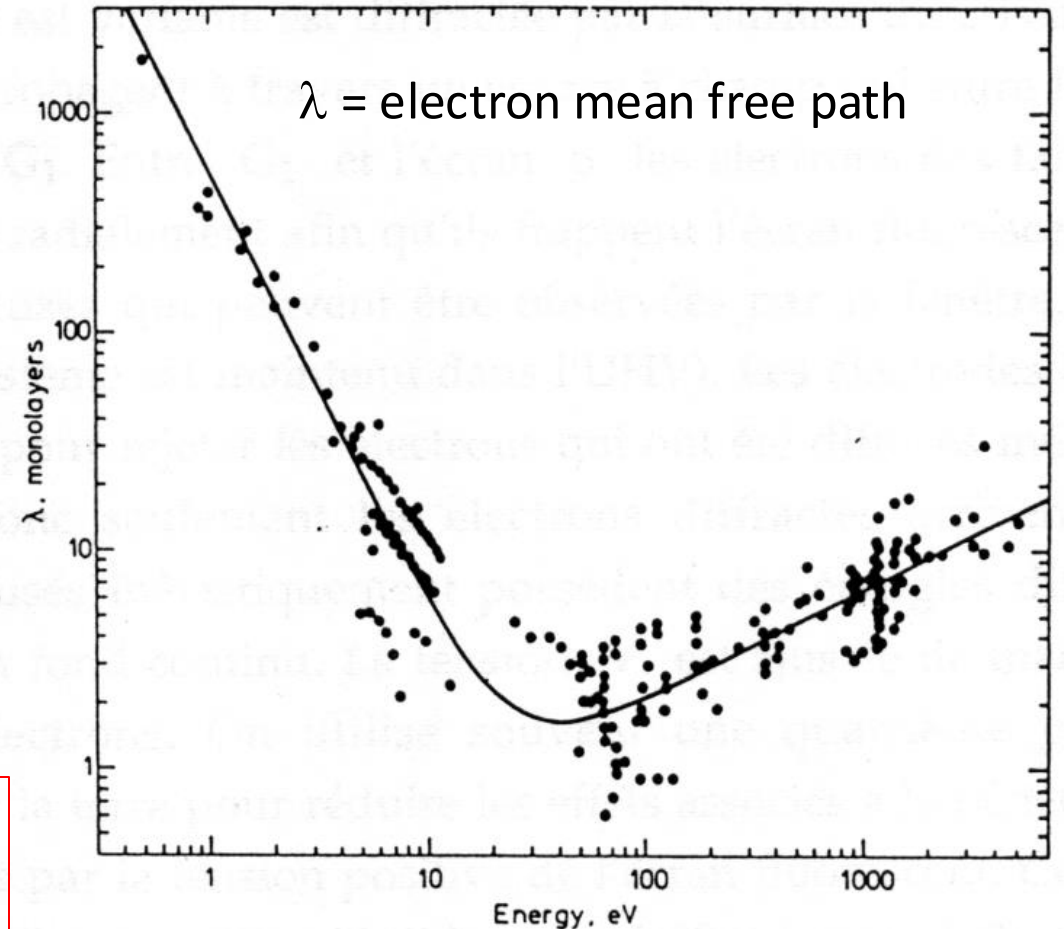
$$I = \int_0^z I(z) dz = I(0) \lambda (1 - \exp(-z/\lambda))$$

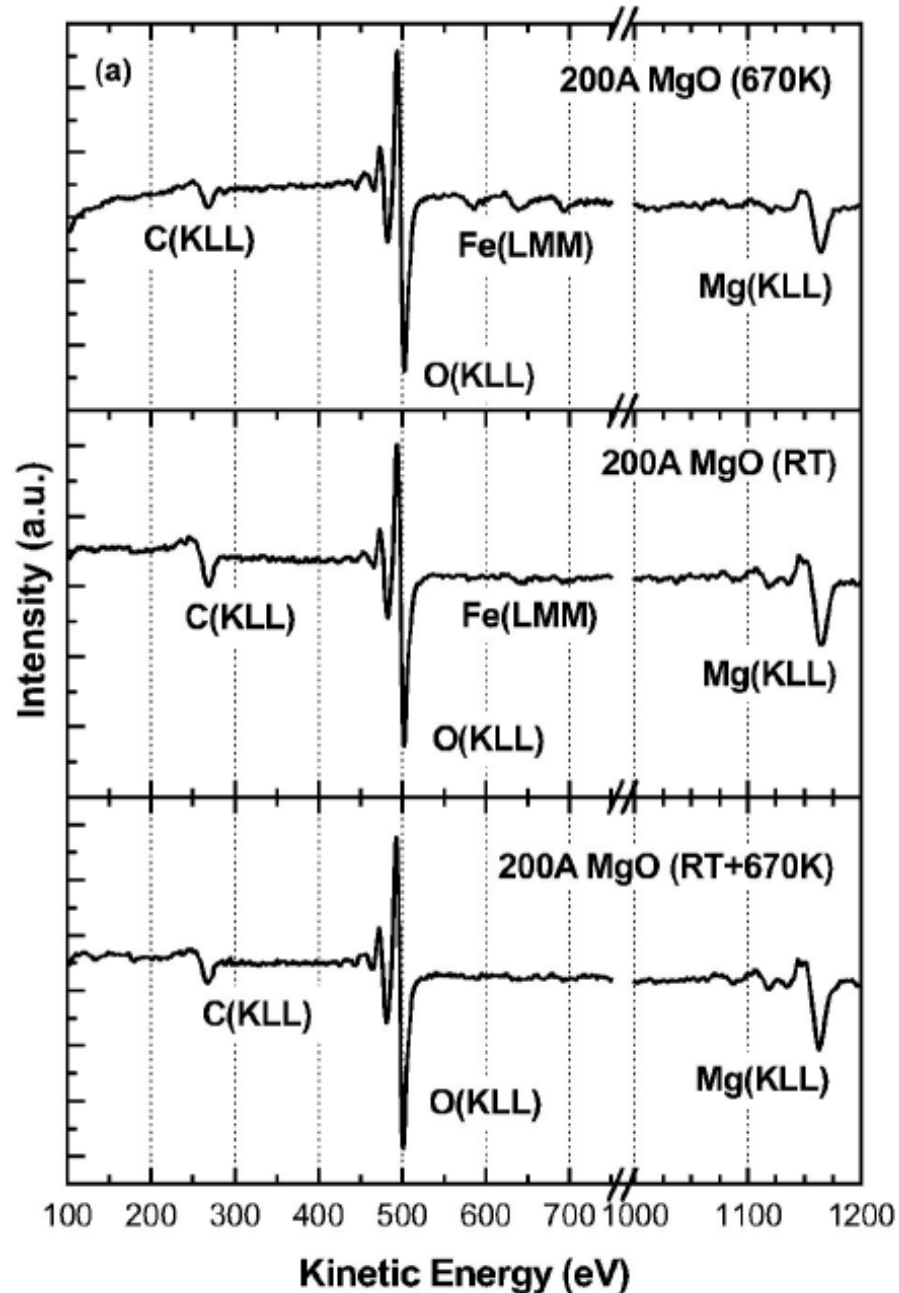
Surface sensitivity results from the limited mean free path of the outgoing electrons

Thickness D from which originates 95% of the signal:

$$I(D) = I(0) \lambda (1 - \exp(-D/\lambda)) = I(0) \lambda 0.95$$

$$D = -\lambda \ln(0.05) \approx 3\lambda$$





Visible Fe absorption peaks:
MgO does not wet completely the Fe bottom electrode
(columnar growth with pits)

Almost no Fe absorption peaks, but rough MgO layer

No Fe signal:
annealing of the RT grown sample gives rise to a
continuous layer (no pits)



Exciting particle: photon

Emitted particle: electron

The energy of a photon is given by $E = h\nu$

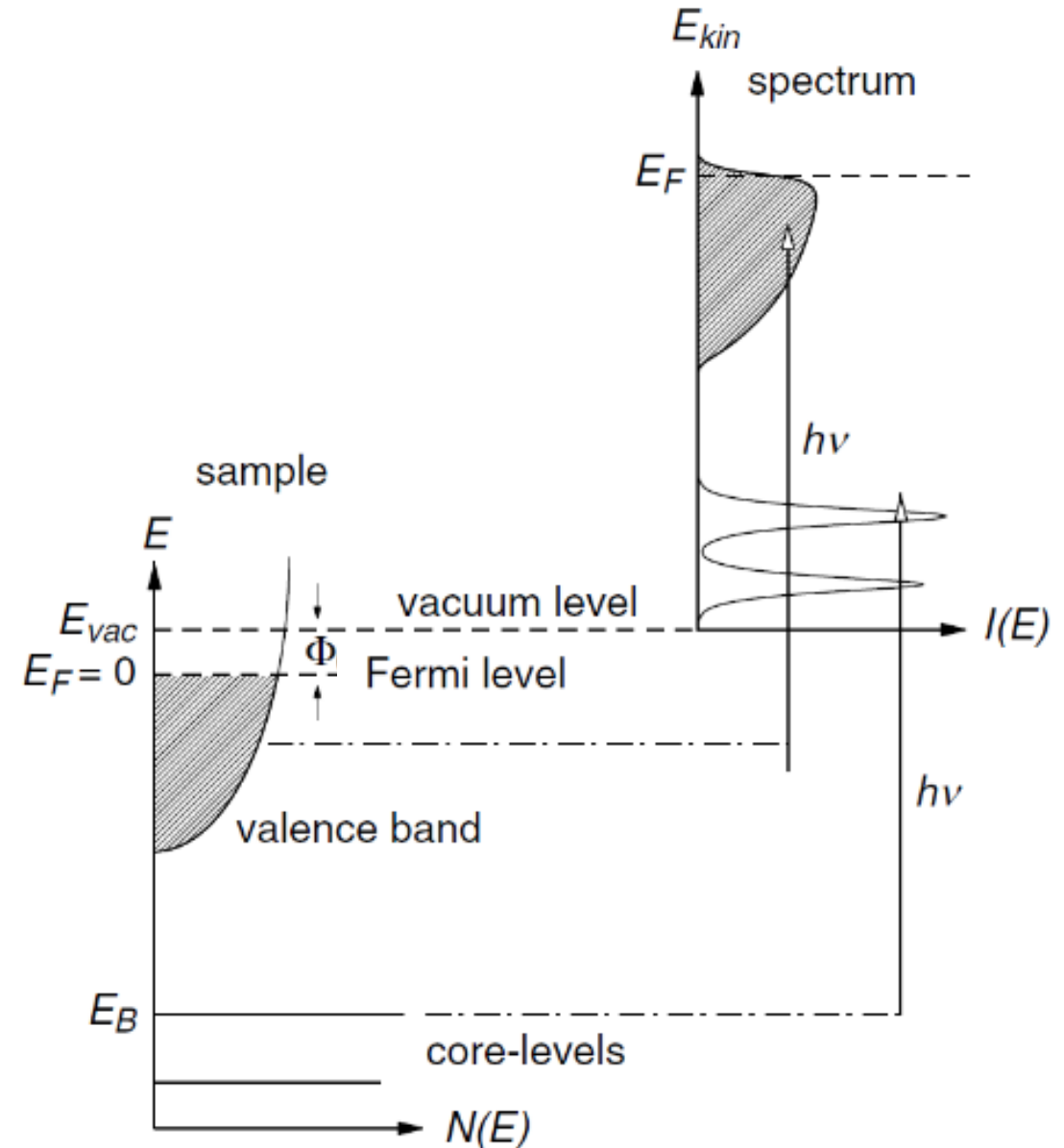
h Planck constant

ν radiation frequency (Hz)

$$E_{kin} = h\nu - E_{bind} - \Phi$$

E_{kin} is the kinetic energy of the emitted electron

it depends on the level that has been excited

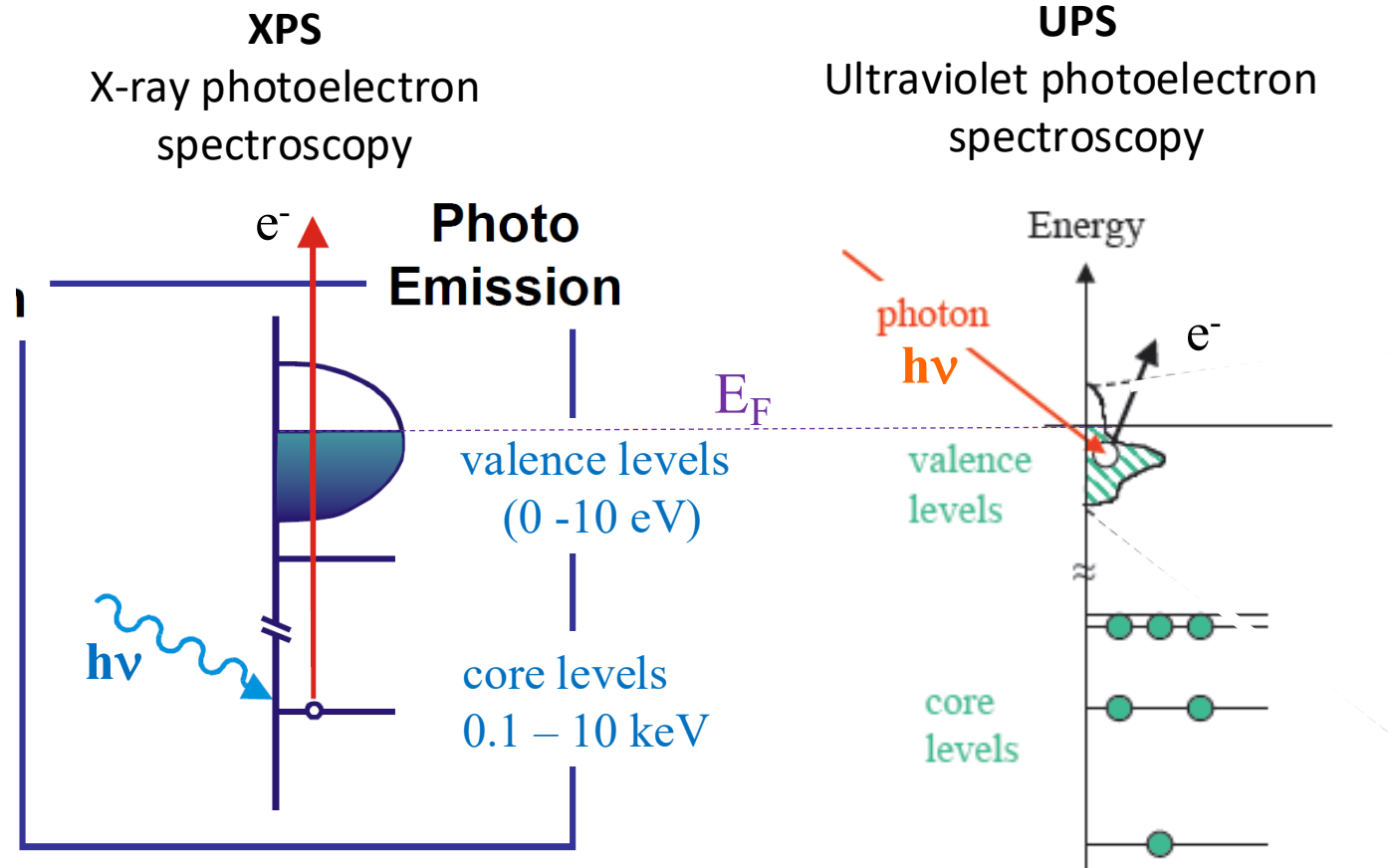


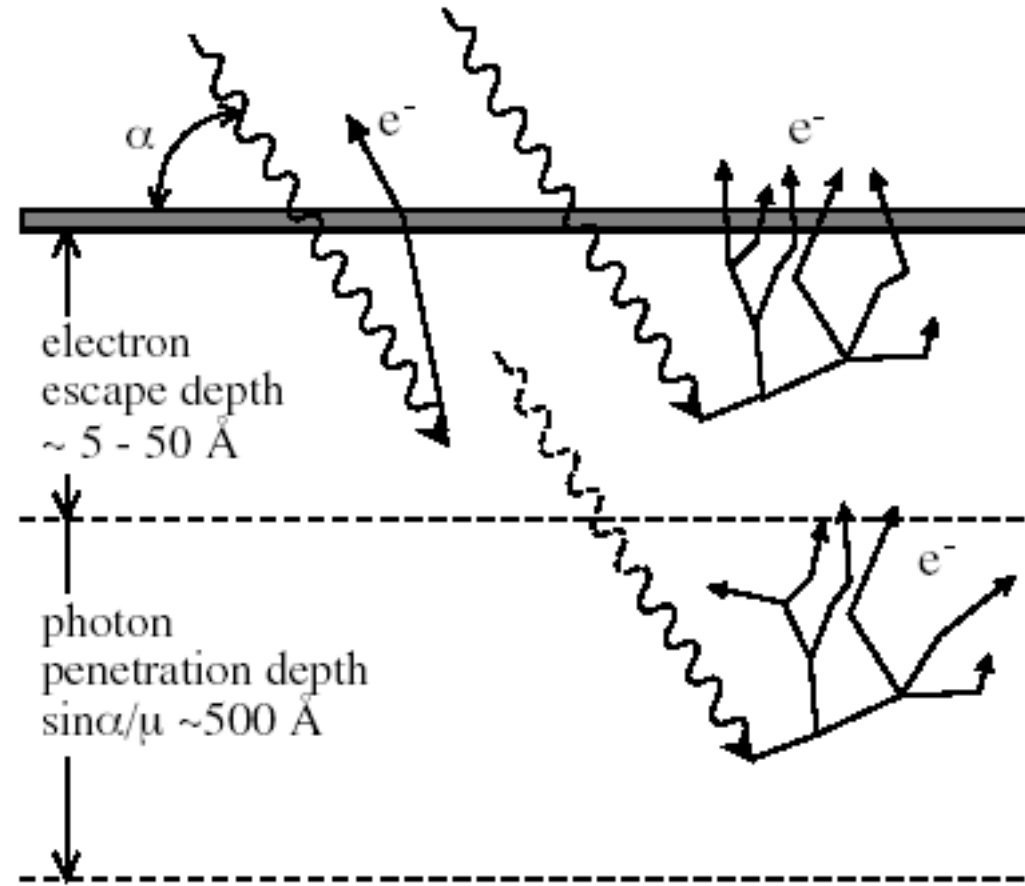
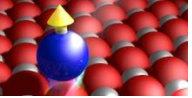


X ray photons (0.2 –10 keV) → to investigate core levels

UV photons (10 - 100 eV) → to investigate valence levels

although energy distinction between core and valence is not strict



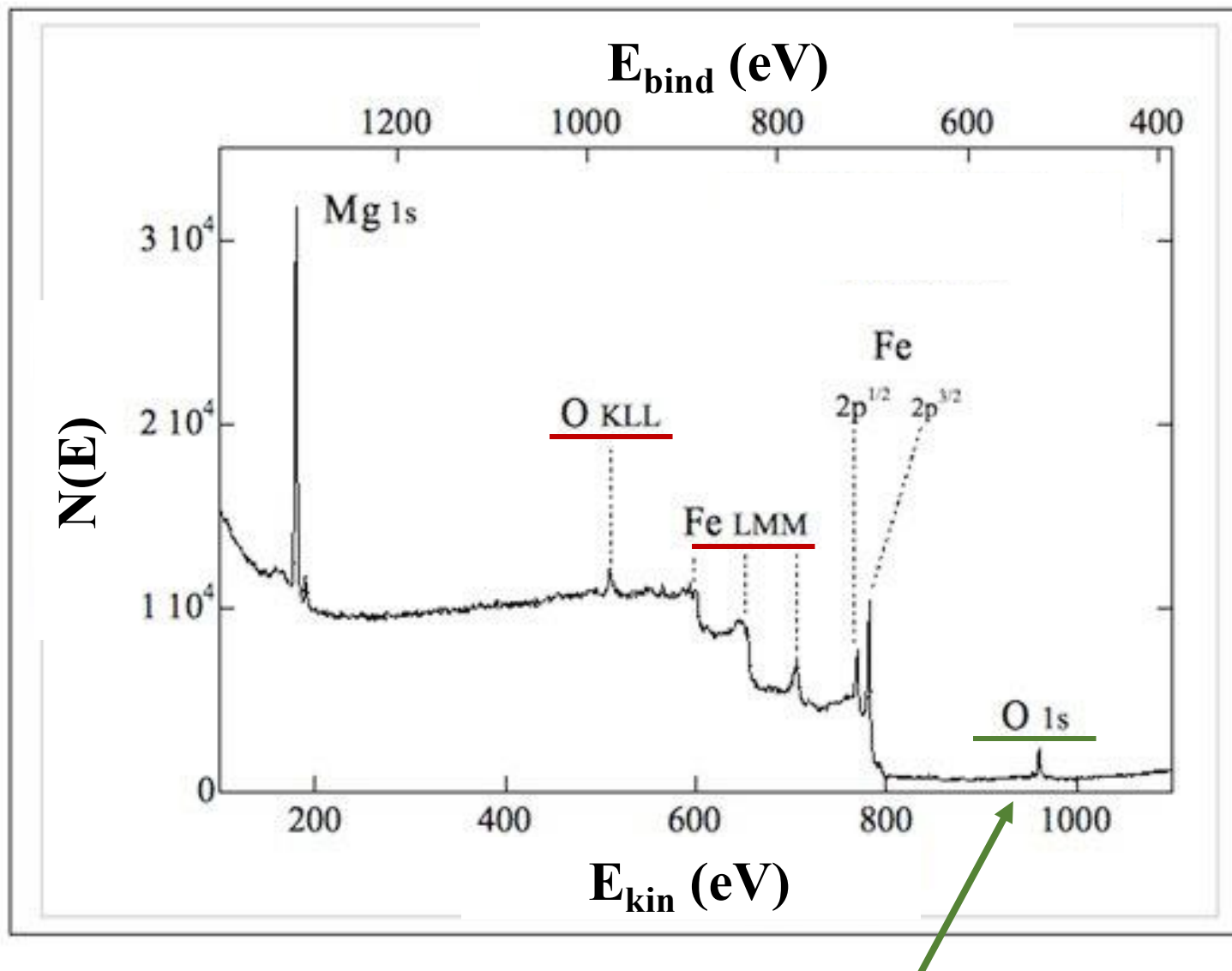


Photon penetration depth $> 1-10 \text{ mm}$

→ surface sensitivity given by the reduced mean free path (escape depth) of the outgoing electrons



Exercise 8.2



$$h\nu = 1486 \text{ eV}$$

Auger electrons are also detected.

Their energy does not depend on the energy of the photons

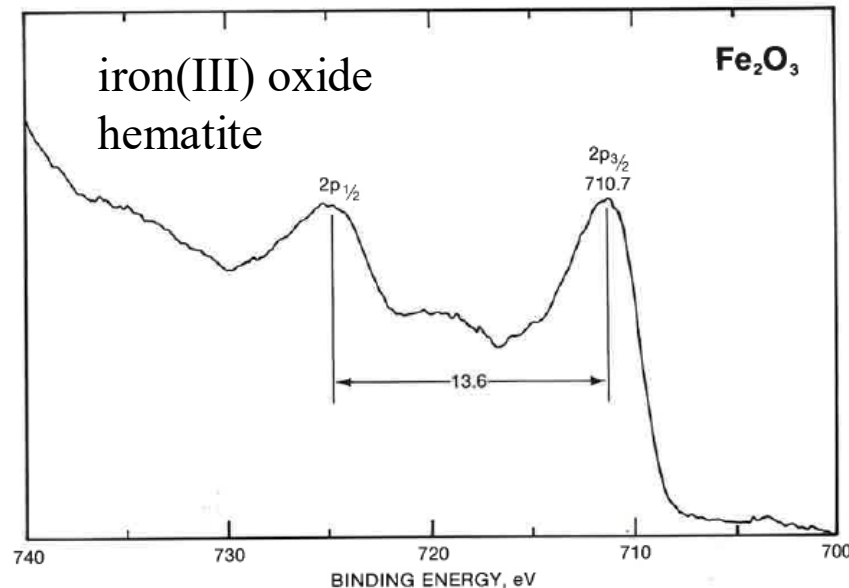
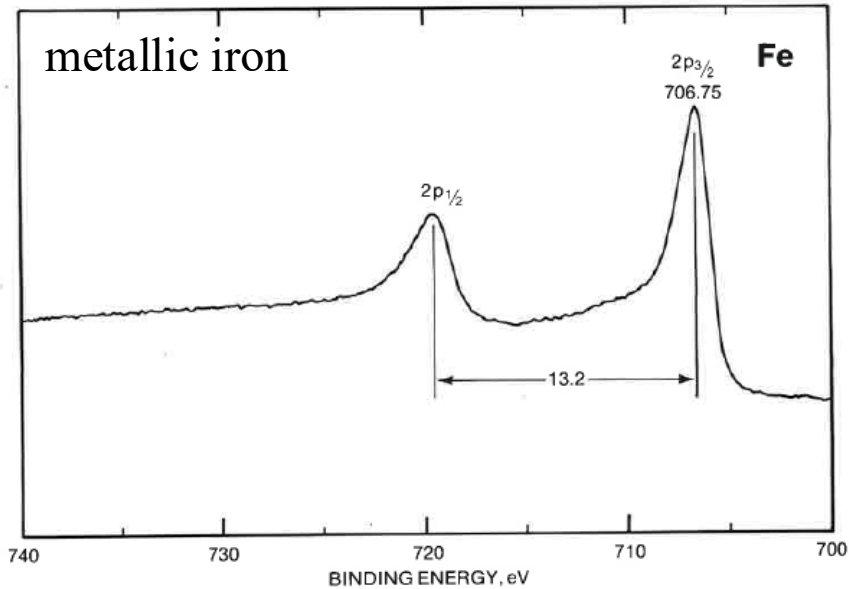
$$\text{energy of O 1s (K edge)} : E_{bind} = h\nu - E_{kin} - \Phi \approx 1486 - 945 \approx 540 \text{ eV}$$



XPS spectrum: Fe 2p (L_2 , L_3), chemical shift

atomic:
[Ar] 3d⁶ 4s²

HANDBOOK OF X-RAY PHOTOELECTRON SPECTROSCOPY



2p levels are split by the spin-orbit coupling

$$j = l + s, l - s$$

$$l = 1$$

$$s = 1/2$$

$$j = 3/2, 1/2$$

based on degeneracy ($2j + 1$), intensity

$$\text{ratio: } 2p_{1/2} : 2p_{3/2} = 1 : 2$$

The anti-parallel alignment is more favorable, has a lower energy and therefore appears at higher binding energy

+3 oxidation state:

Fe “loses” the two 4s electrons and one 3d electron

chemical shift correlated with overall charge on atom:

reduced charge → reduced coulomb screening

→ increased binding energy

(mainly an initial state effect)

peak energy → oxidation state



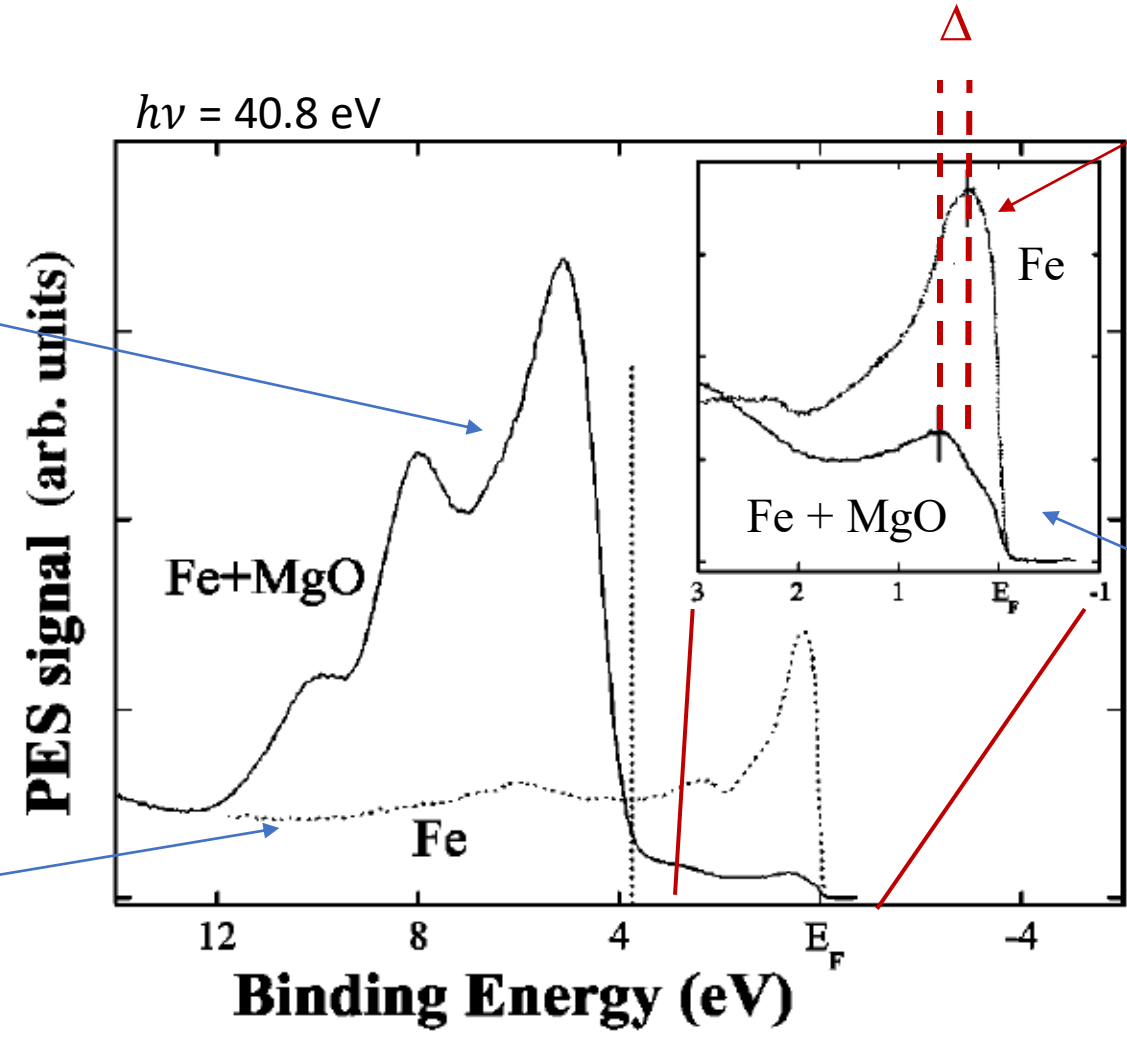
Example UPS spectrum: MgO/Fe (interface in MTJ)

Exercises 8.3 – 8.5

MgO is an insulator with a gap of 8 eV
Fe electronic structure: [Ar] 3d⁶ 4s²

Spectral intensity dominated by the MgO contribution

Spectrum of Fe before MgO growth



Weak modification of the Fe DOS

Spectral intensity dominated by the Fe contribution (no MgO state in the energy gap)



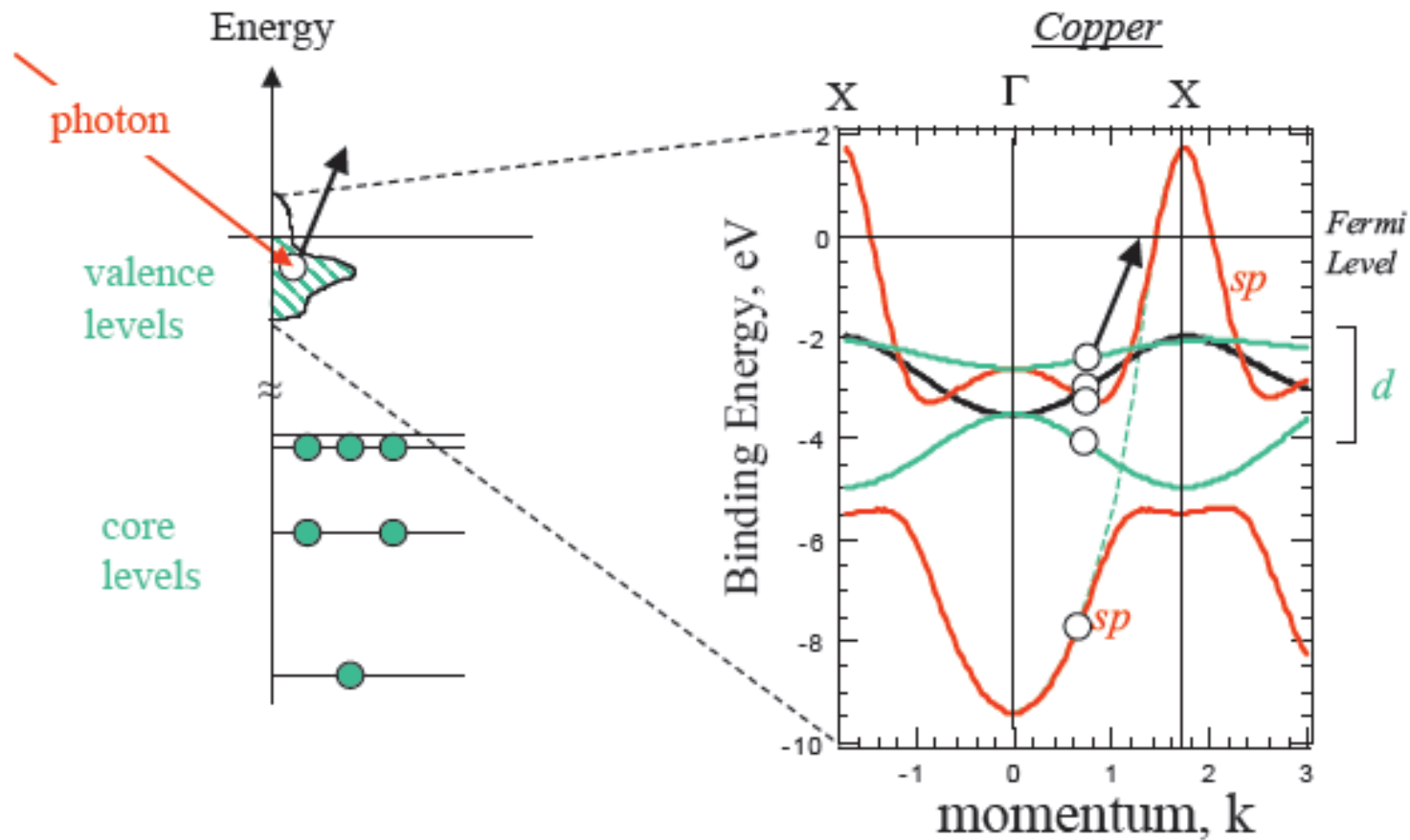
UV photons (10 - 100 eV) → to investigate valence levels

ARUPS: Angle Resolved UPS

ARPES: Angle Resolved PhotoElectron Spectroscopy

UPS is k -integrated (DOS)

ARPES can measure the wave vector k and the energy at the same time (band structure)



STM gives access to the DOS as a function of energy, but poor information on k (with exceptions, as for example the standing wave of 2D gas at metal surfaces)

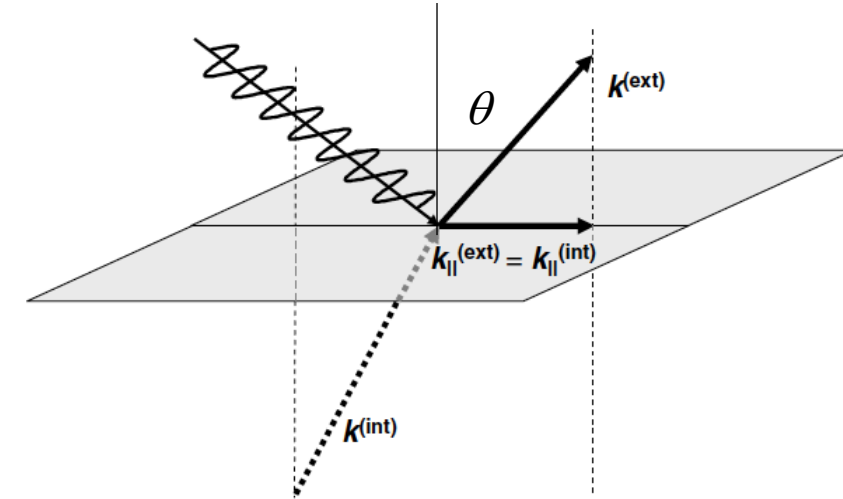


We have already found the condition for energy conservation: $E_{kin} = h\nu - E_{bind} - \Phi$
What about the momentum?

Three-step model (valid also for XPS and UPS):

- 1) The electron is excited from an initial to a final state within the crystal;
- 2) The electron travels through the solid towards the surface without scattering;
- 3) The electron crosses the surface and is emitted into the vacuum with a certain kinetic energy

Momentum conservation: due to symmetry translation breaking at the surface, only the electron momentum component parallel to the surface is conserved, modulo a reciprocal surface vector (photon momentum is negligible)



$$\hbar \mathbf{k}_{\parallel,i} = \hbar \mathbf{k}_{\parallel,f} + \mathbf{G}$$

$$\hbar k_{\parallel,f} = \sqrt{2mE_{kin}} \sin \theta$$

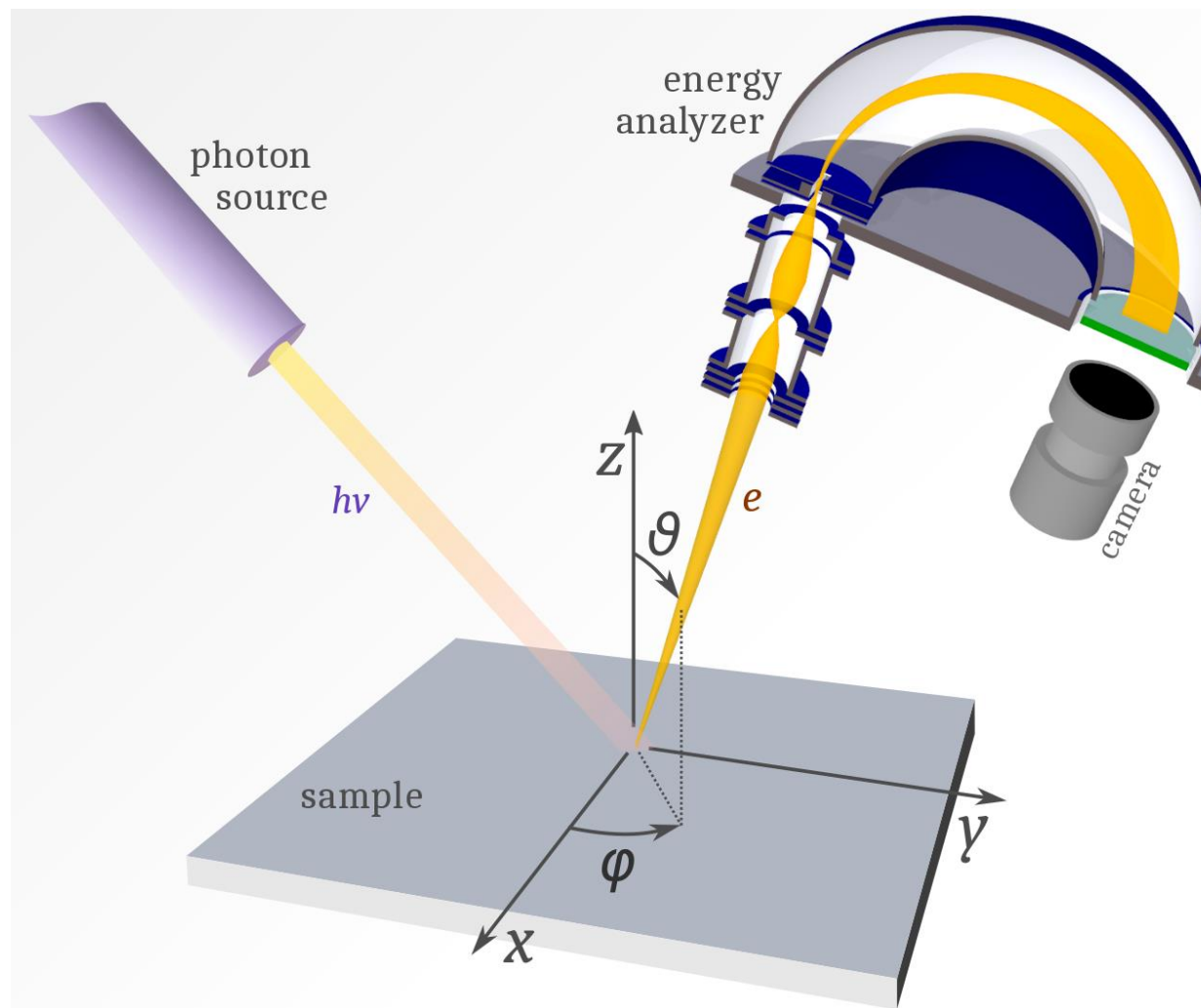
$$\hbar k_{\parallel,i} = \sqrt{2m(h\nu - E_{bind} - \Phi)} \sin \theta$$

Because of their 2D nature, the k -vector of surface states is fully determined by the k_{\parallel} -vector of the photoemitted electron.

The perpendicular component of the k -vector of bulk states remains unknown, but can be established by special techniques



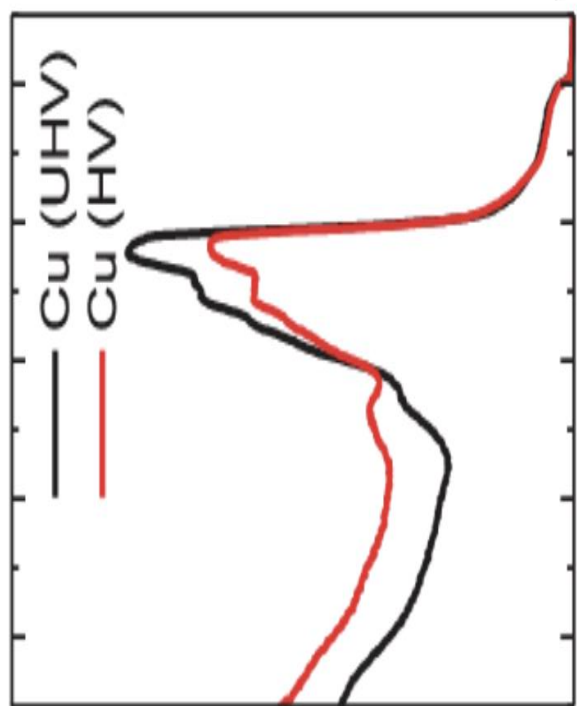
Measurement of E vs k
Band structure





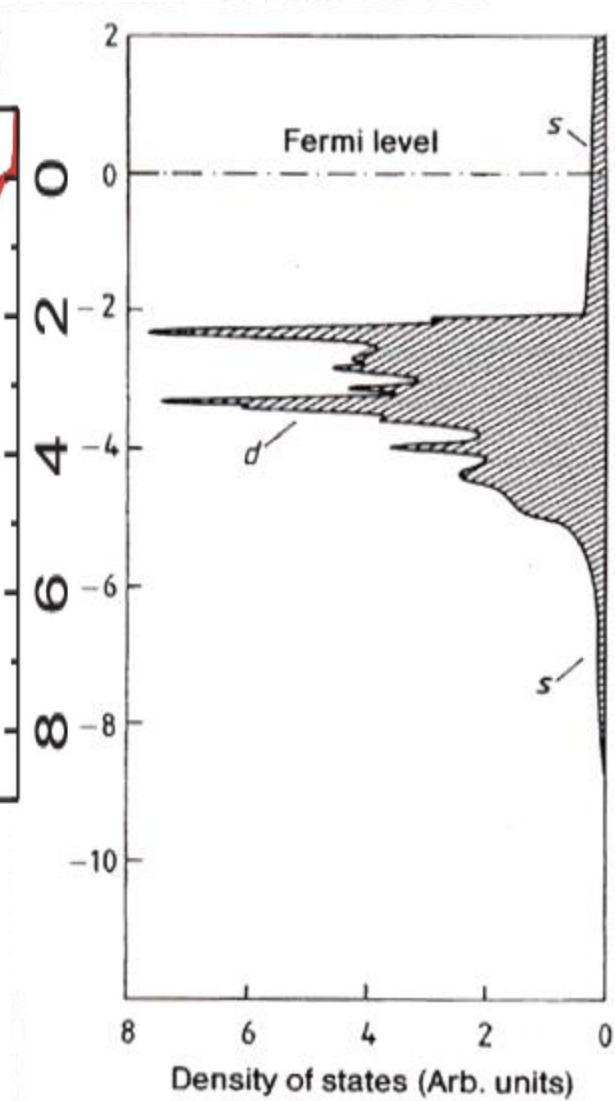
UPS vs ARPES

UPS

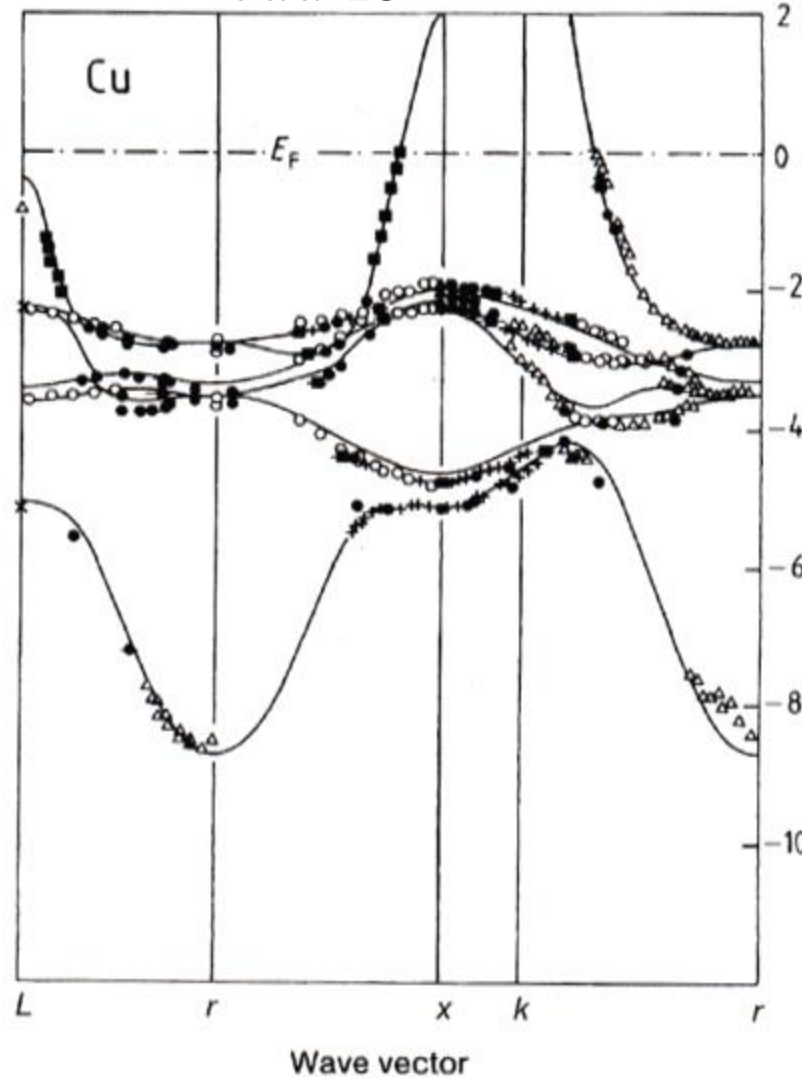


Intensity (counts)

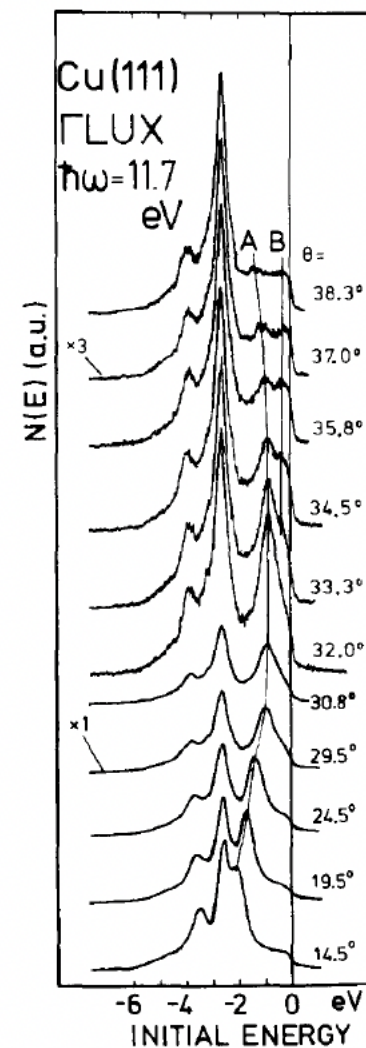
Integrated in k-space

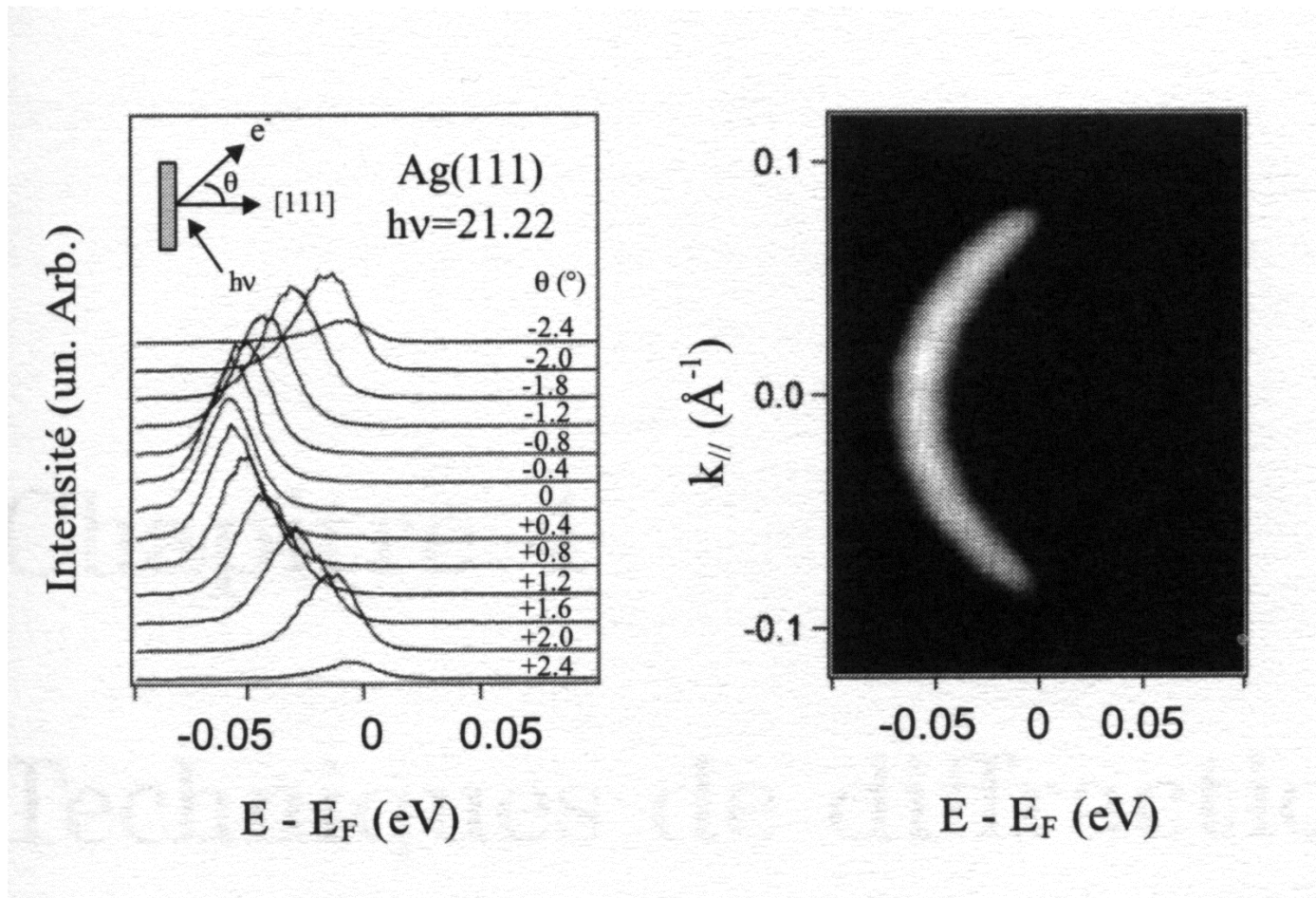


ARPES



Resolved in k-space:
band structure





Surface states behave like a 2D gas of free electrons

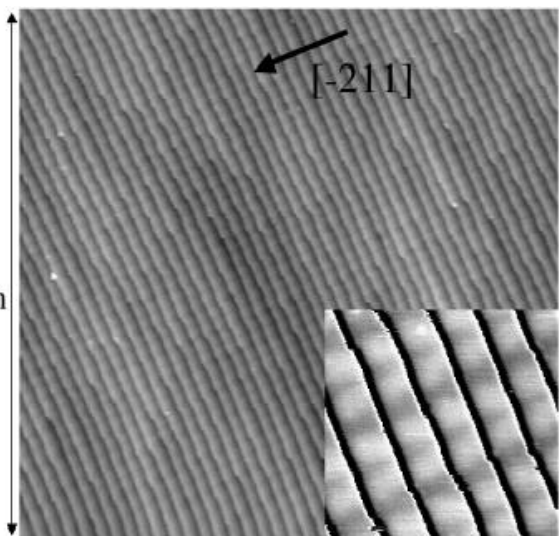
(see lecture on local spectroscopy)

$$E = E_F - E_0 + \frac{\hbar^2 k^2}{2m^*}$$

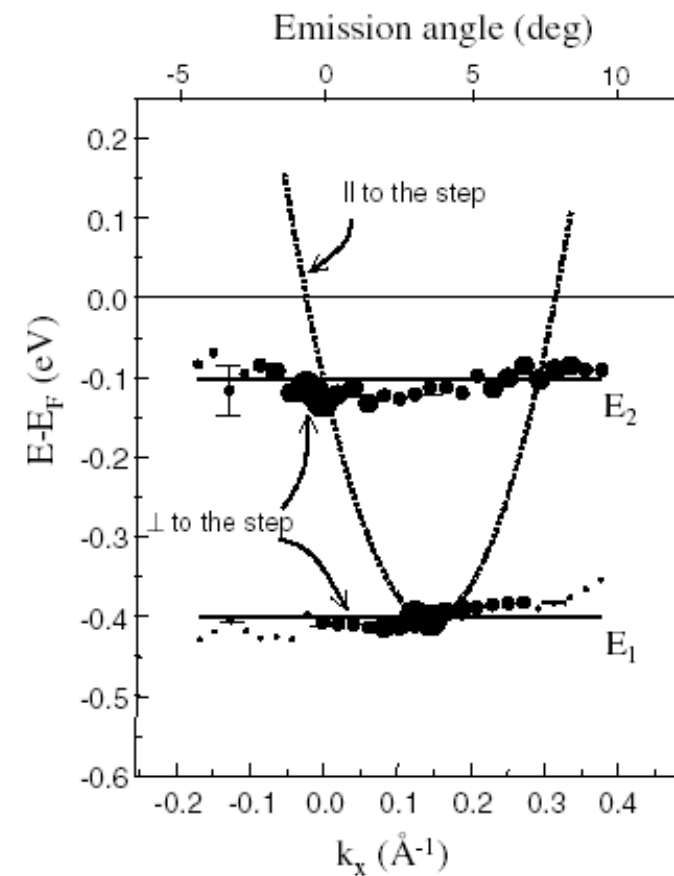
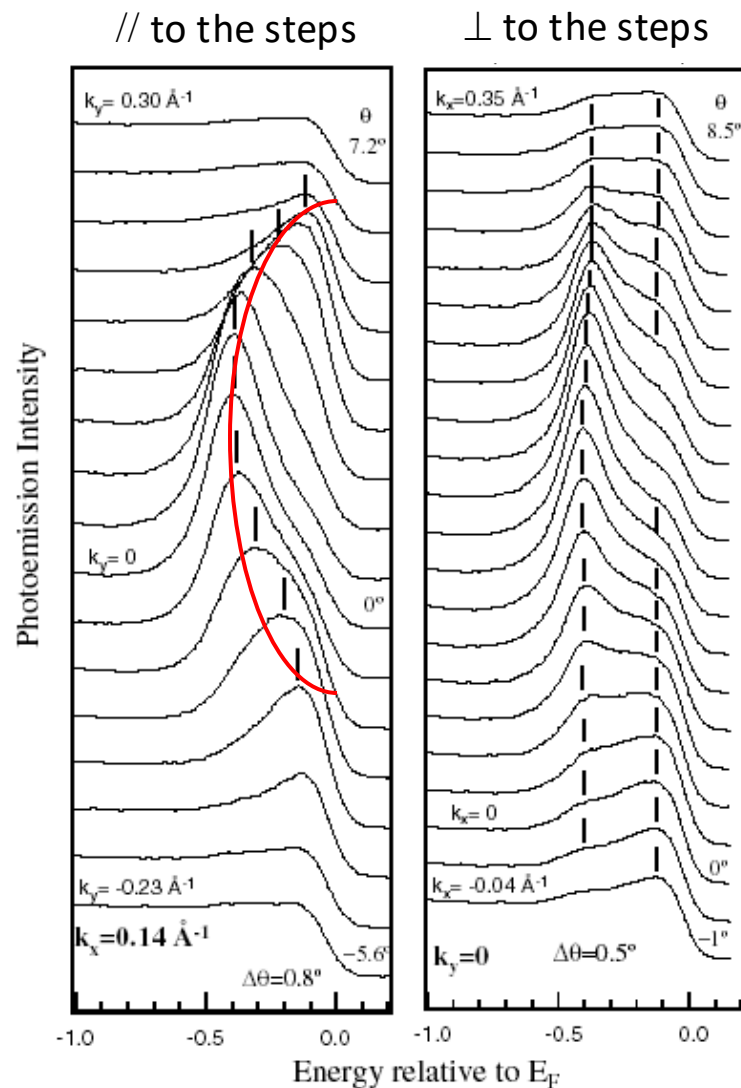
m^* : effective mass of the surface state electron



Exercise 8.6



$L = \text{terrace width} = 3.8 \text{ nm}$



$$\Psi_{\mathbf{k}}(\mathbf{r}) \propto \exp(ik_{\parallel}y) \sin(k_{\perp}x)$$

$$k_{\perp} = \frac{n\pi}{L}$$

$$E_n = E_F - E_0 + \frac{\hbar^2}{2m^*} \frac{\pi^2}{L^2} n^2$$

One dimensional quantum well of width L
perpendicular to the steps



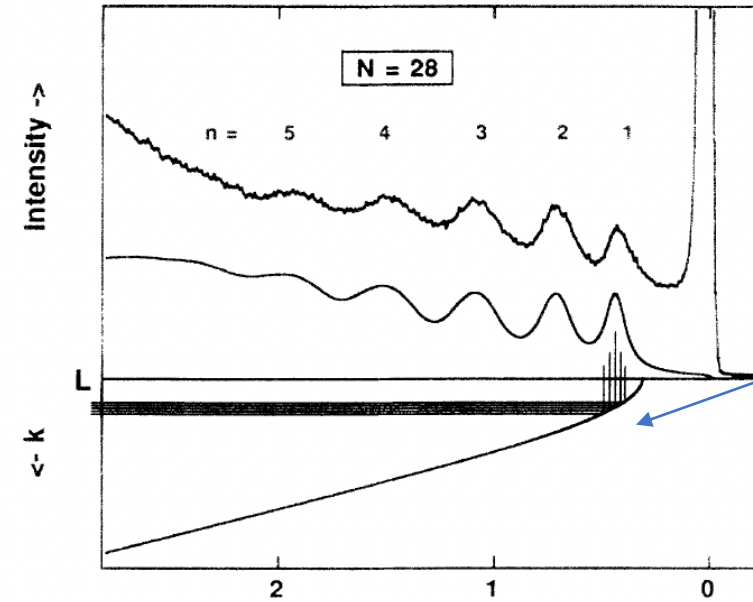
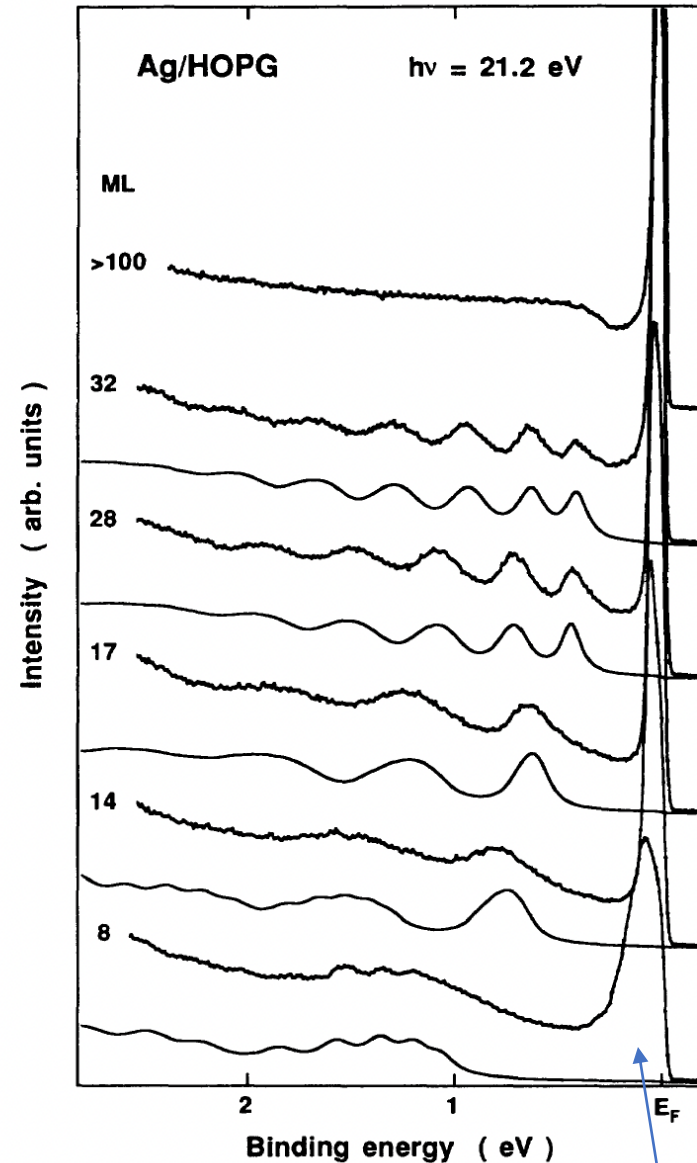
Example: 1D quantum well states

Ag thin films on highly-oriented pyrolytic graphite

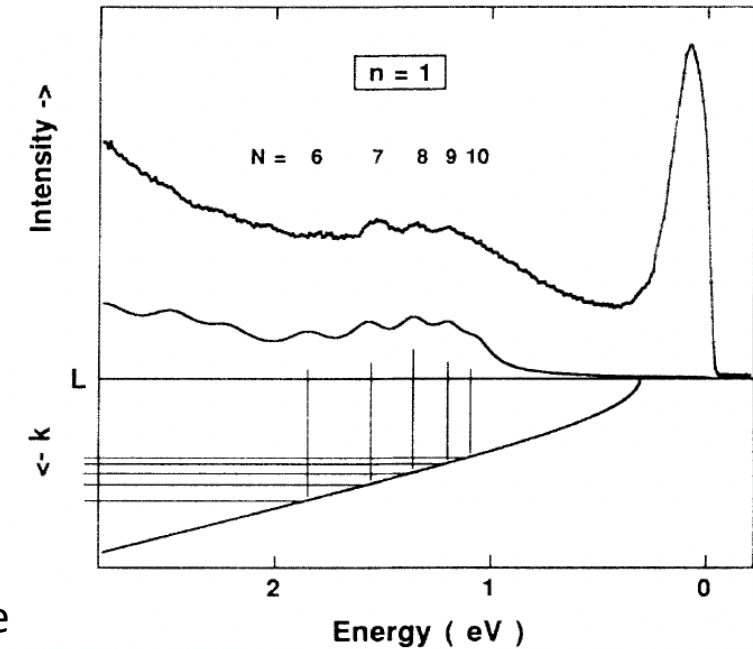
Confinement in the direction perpendicular to the surface

N number of layers

n quantum number

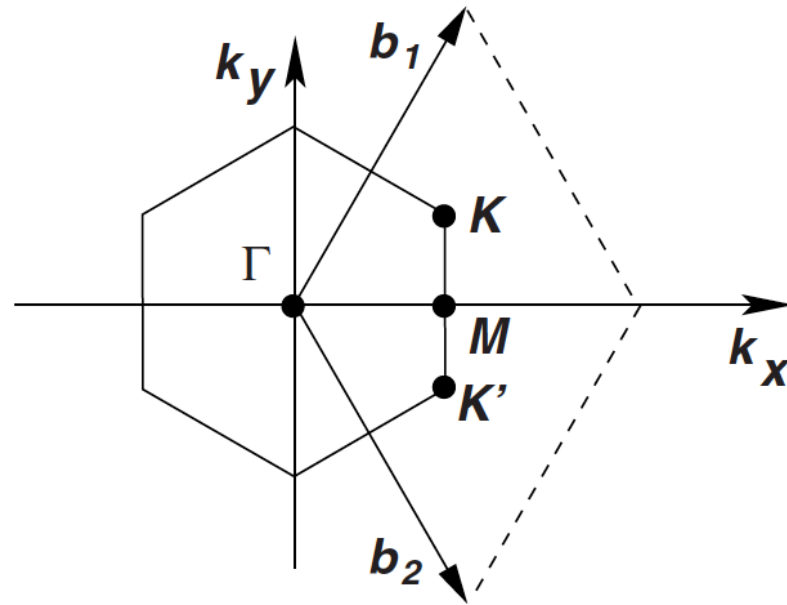
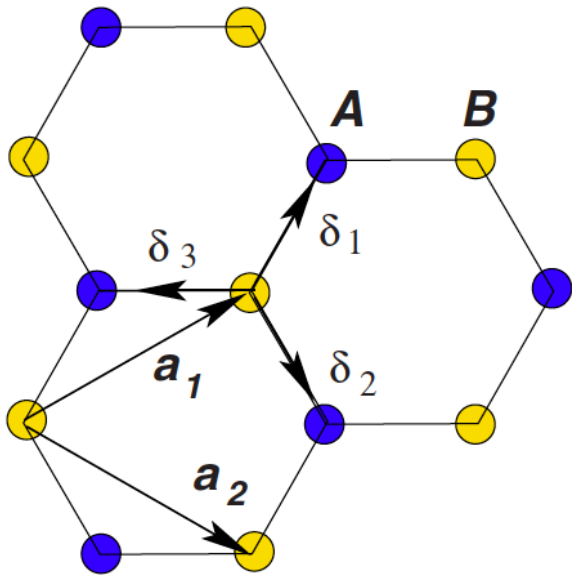


n = 1 for thicknesses around N=28

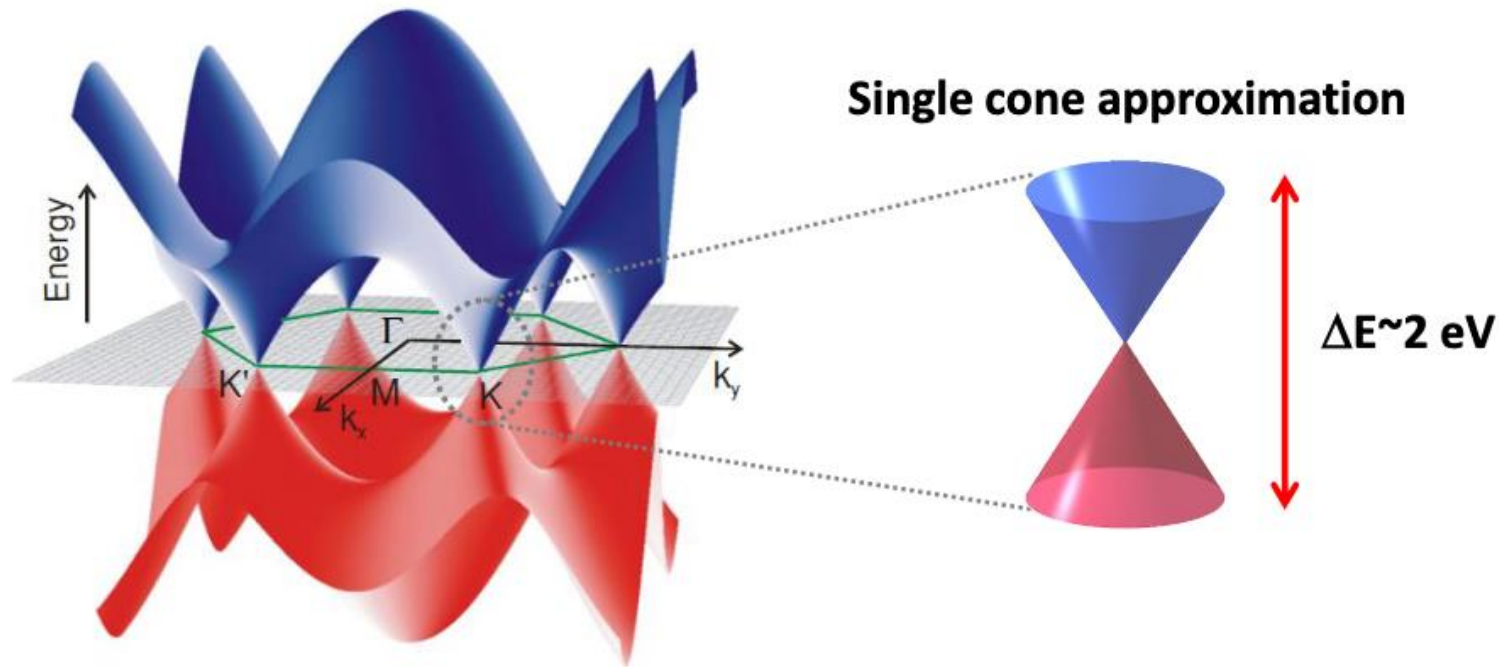




Tight-binding description of the π bands (originating from p_z)



- triangular Bravais lattice with a base formed of two atoms (A and B)
- A and B sublattices
- for each atom, the 3 nearest neighbors belong to the other sublattice
- non-zero hopping parameter only for nn
- starting from Bloch wavefunctions localized on A and B, we build the wavefunctions describing the π bands
- find the expression for their energy vs k in the FBZ



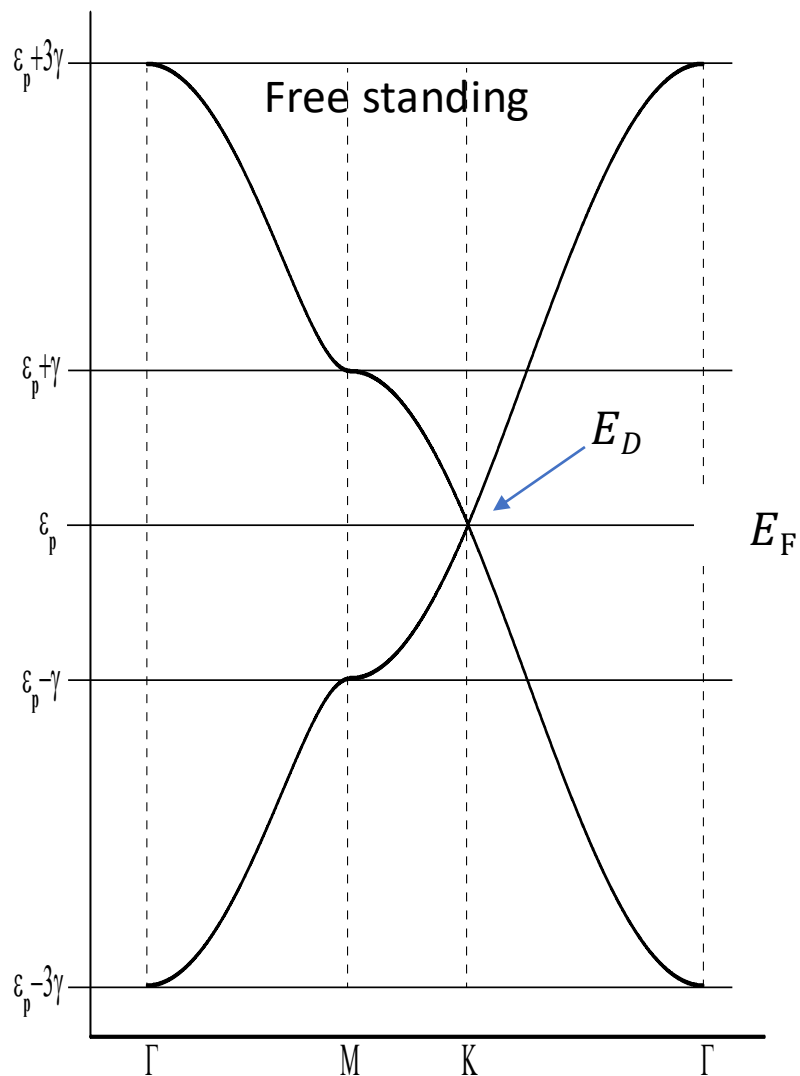
Around K and K' points, the energy dispersion is linear, very different from the usual free-electron like quadratic dependency (Dirac points)

It can be shown that (wavevector k centered on K or K'):

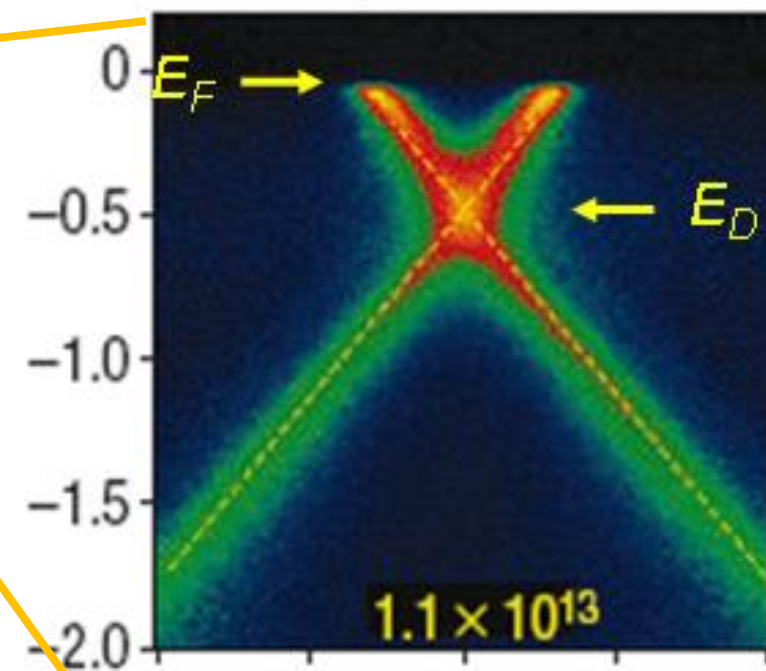
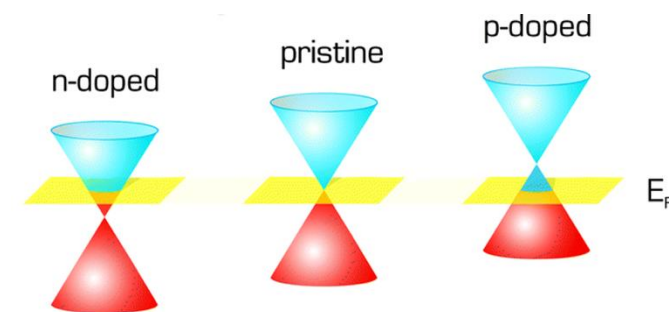
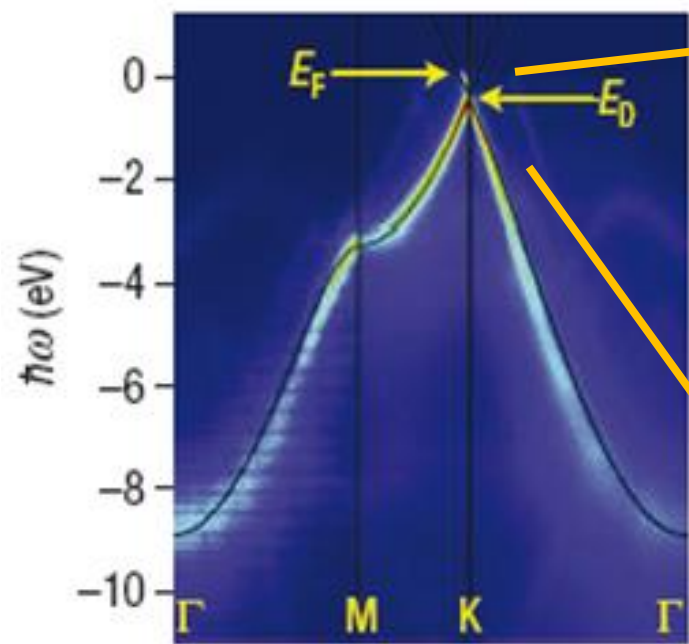
$$E = \pm \hbar v_F k$$

$$v_F = \frac{3|\gamma|a}{2\hbar}$$

for ideal, free-standing graphene E_F coincides with the Dirac points

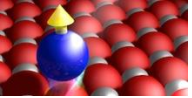


Graphene on SiC



The Dirac point E_D coincides with the Fermi level in free standing graphene

On SiC, $E_D = -0.5$ eV (below Fermi level) due to graphene-substrate hybridization



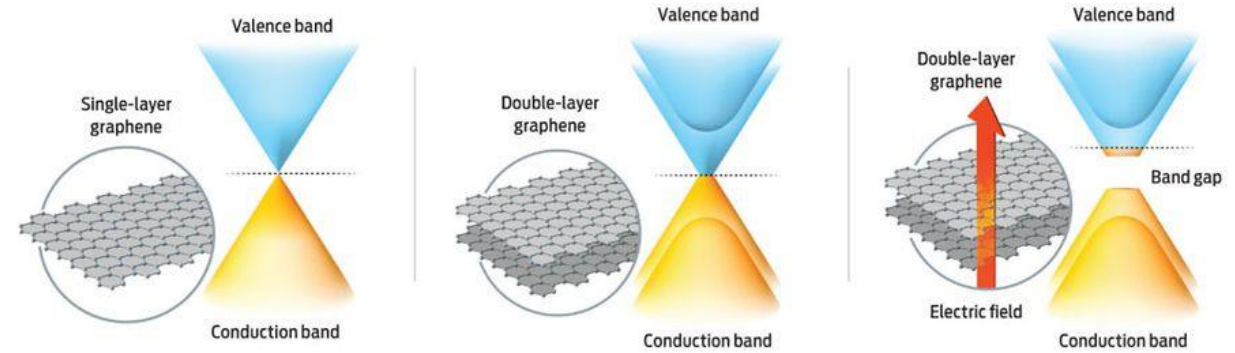
graphene: semimetal or zero-gap semiconductor

technological advantages for graphene to be a semiconductor instead of a semimetal, specifically in electronic applications (transistor)

possible paths:

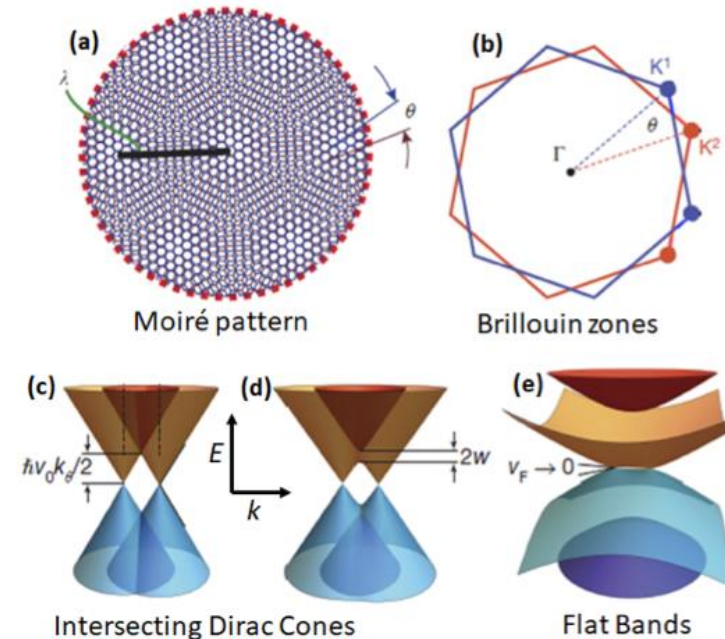
- geometrically confine graphene into nanoribbons or quantum dots
- grow graphene on substrates that induce lattice potentials that open a gap
- bilayer graphene + field effect
- twisted bilayer graphene
- bilayer graphene + strain
- other 2D materials (h-BN, MoS₂, ...)

Widely tunable bandgap in bilayer graphene



doi.org/10.1038/nature08105

Twisted Bilayer Graphene: Moiré-induced Electronic Structure



Images from Y. Cao, et al., *Nature* 556, 80 (2019)