



7

Diffraction



## Chapters on diffraction

Physics of Surfaces and Interfaces

H. Ibach

Springer

Surface Science - An Introduction

K.Oura, V. G. Lifshits, A. A. Saranin, A.V. Zotov, M. Katayama

Springer

Introduction to Surface and Thin Film Processes

J. A. Venables

Cambridge University Press

## More specific

Surface Structure Determination by LEED and X-rays

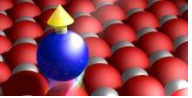
W. Moritz and M. A. Van Hove

Cambridge University Press

Reflection high-energy electron diffraction

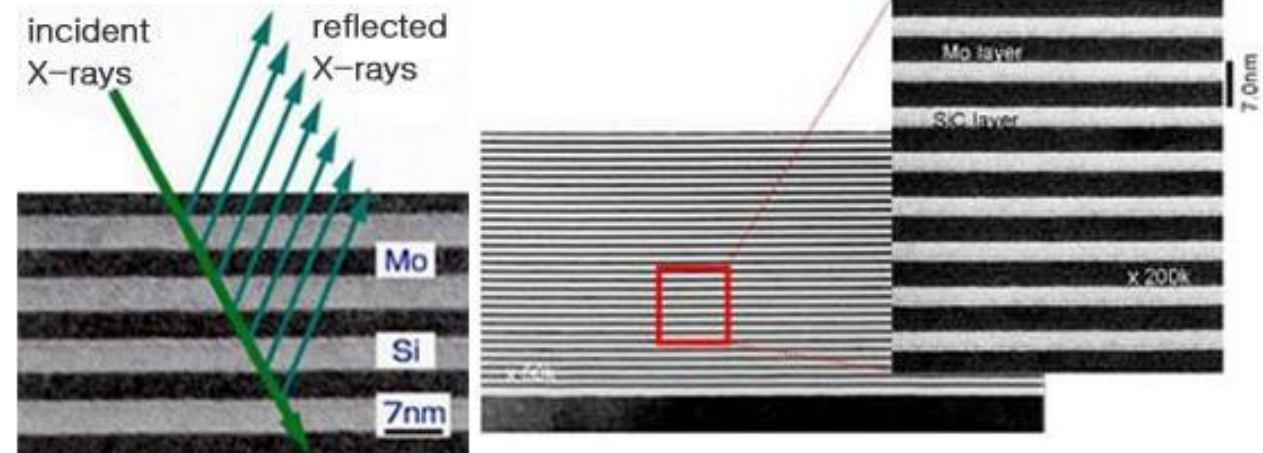
A. Ichimiya and Ph. I. Cohen

Cambridge University Press



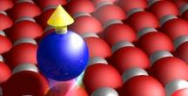
Electronic devices (FET, FGT, ...), created by successive processes (lithography, etching, growth, implantation, ....) are composed of many “layers”

Creation of mirrors and masks for EUV-lithography (thickness, interface roughness...)



Magnetic tunnel junction (spintronics): increased by improving and controlling the crystallographic order at the interfaces

Control of surface and thin film quality (crystallinity), of film thickness, of interfaces is needed  
→  
**Diffraction**



Different probes can be used, given an appropriate the de Broglie wavelength: x-rays, electrons, He-atoms, neutrons ...

Examples:

**X-ray:**  $E = h\nu = \frac{hc}{\lambda} \iff \lambda = \frac{hc}{E}$

$\lambda \approx 1 \text{ \AA} \implies E \approx 12 \text{ k eV}$

**Electrons:**  $p = \hbar k = \frac{h}{\lambda} \iff \lambda = \frac{h}{p} = \frac{h}{\sqrt{2mE}}$

$\lambda \approx 1 \text{ \AA}$       $m_e = 9.1 \cdot 10^{-31} \text{ kg}$

$\implies E \approx 150 \text{ eV}$

**Neutrons:**  $\lambda = \frac{h}{p} = \frac{h}{\sqrt{2mE}}$

$\lambda \approx 1 \text{ \AA}$       $m_n = 1.6749 \cdot 10^{-27} \text{ kg}$

$\implies E \approx 0.08 \text{ eV}$

Here we focus on electrons (surface sensitivity, easy to integrate into other setups)

The interpretation is not as straightforward as for x-rays: electrons interact strongly with matter, dynamic theory should be used, but useful information can be deduced from simple descriptions

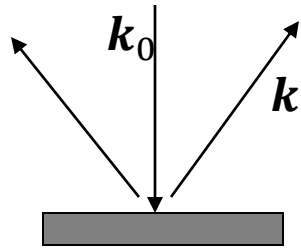


LEED: Low Energy Electron Diffraction

RHEED: Reflection High Energy Electron Diffraction

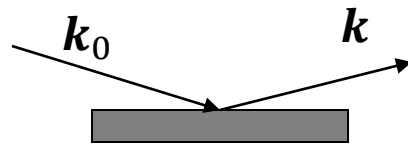
TEM: Transmission Electron Microscopy

LEED



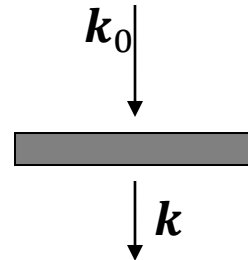
$E = 10 - 300 \text{ eV}$

RHEED

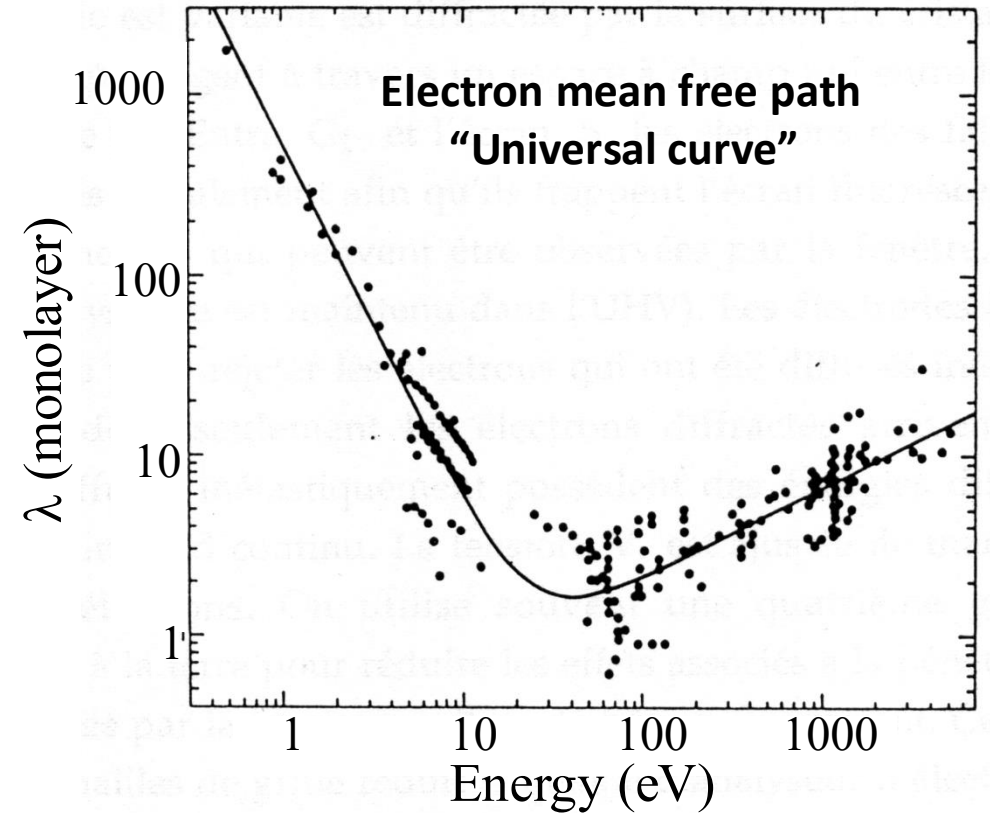


$E = 10 - 50 \text{ keV}$

TEM

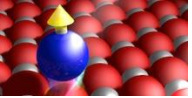


$E = 10 - 200 \text{ keV}$



Electron energy and geometry are chosen considering the electron mean free path and the spatial resolution

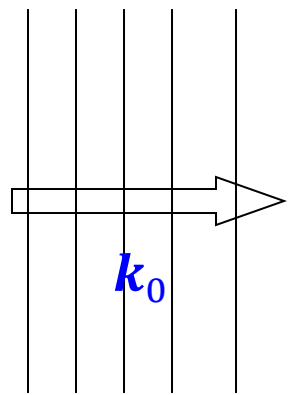
Note: Electron energy = 2 eV  $\rightarrow$  Wavelength = 8.5 Å  
The hypothetical "RLEED" would have low resolution



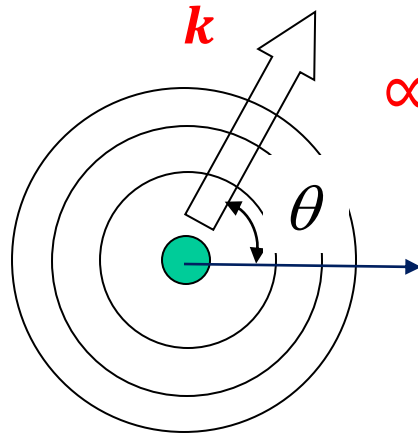
## Single scatterer

$$p_0 = \hbar k_0$$

incoming electron



$$\propto e^{ik_0 \cdot r}$$

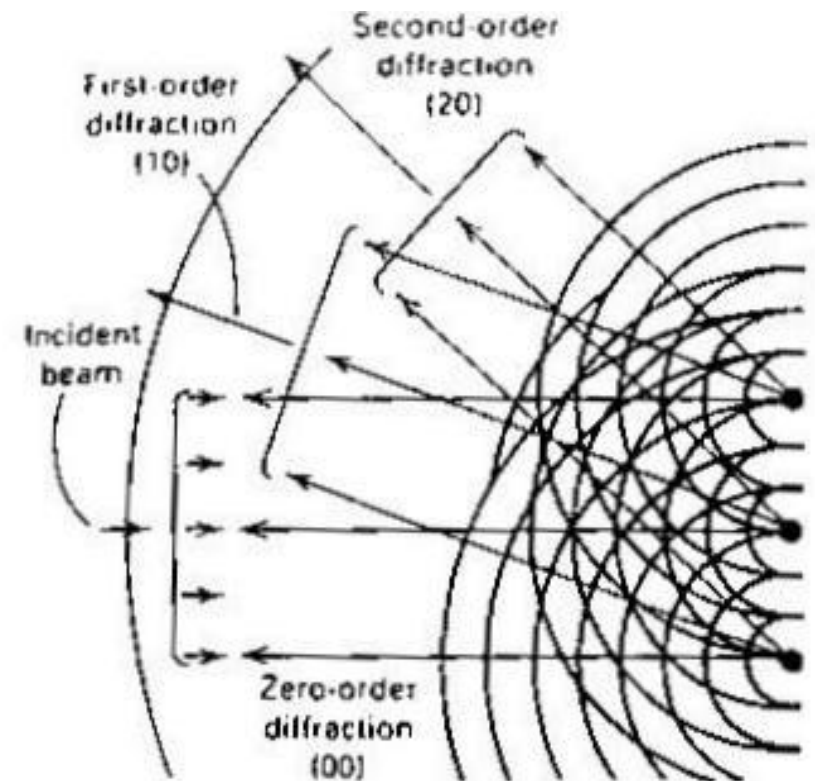


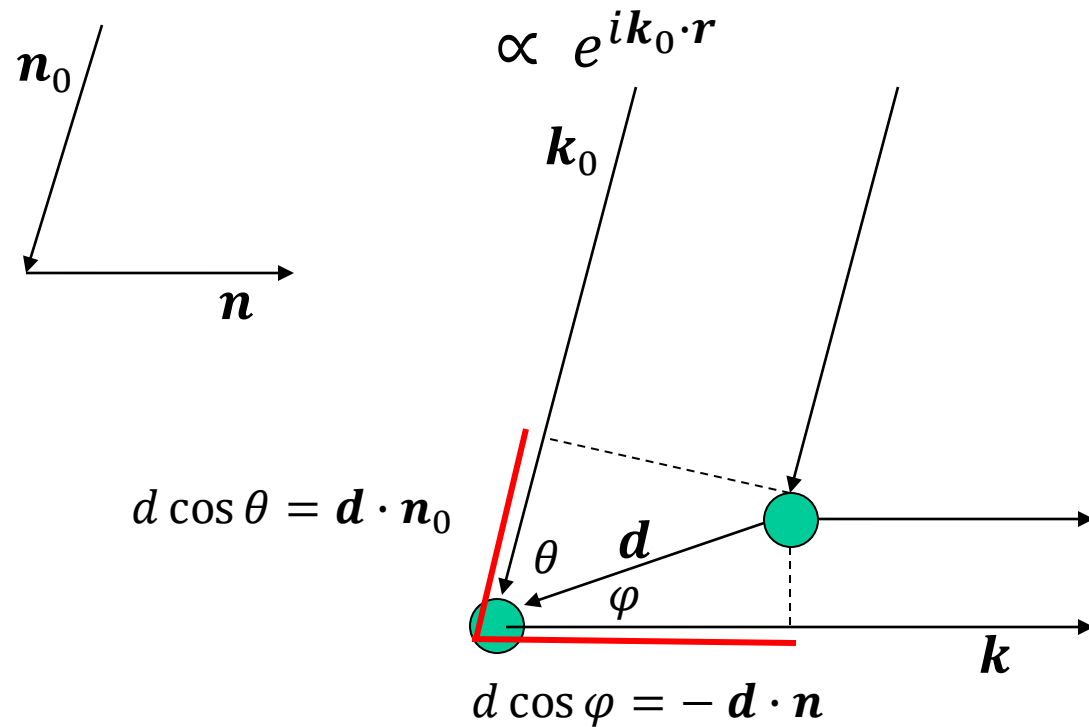
scattered electron

$$\propto f(\theta, r) \frac{e^{ik \cdot r}}{r}$$

$f(\theta, r)$  is the atomic form factor:  
it reflects the atomic electron density

In the case of many scatter points, diffraction spots are the result of the interference pattern generated by the spherical waves scattered by the atoms in the crystal lattice





$$\mathbf{k}_0 = \frac{2\pi}{\lambda} \mathbf{n}_0$$

$$\mathbf{k} = \frac{2\pi}{\lambda} \mathbf{n}$$

## 1) Elastic scattering:

incident and reflected rays have the same wavelength  $\lambda$   
 $|\mathbf{k}| = |\mathbf{k}_0|$

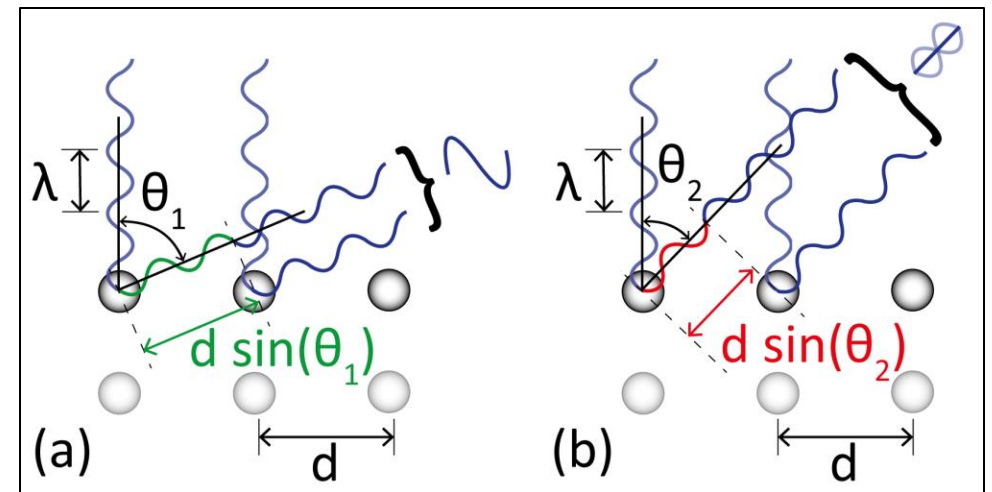
## 2) Constructive interference:

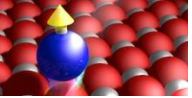
difference in path length = multiple of wavelength

$$d \cos \theta + d \cos \varphi = \mathbf{d} \cdot (\mathbf{n} - \mathbf{n}_0) = m\lambda$$

or

$$\mathbf{d} \cdot (\mathbf{k} - \mathbf{k}_0) = 2\pi m \quad m \in \mathbb{Z}$$





Constructive interference verified by all the lattice points

$$\mathbf{R} \cdot (\mathbf{k} - \mathbf{k}_0) = 2\pi m$$

with  $\mathbf{R}$  Bravais lattice vector,  $m \in \mathbb{Z}$



$$\mathbf{k} - \mathbf{k}_0 = \mathbf{G}$$

$\mathbf{G}$  is a reciprocal lattice vector

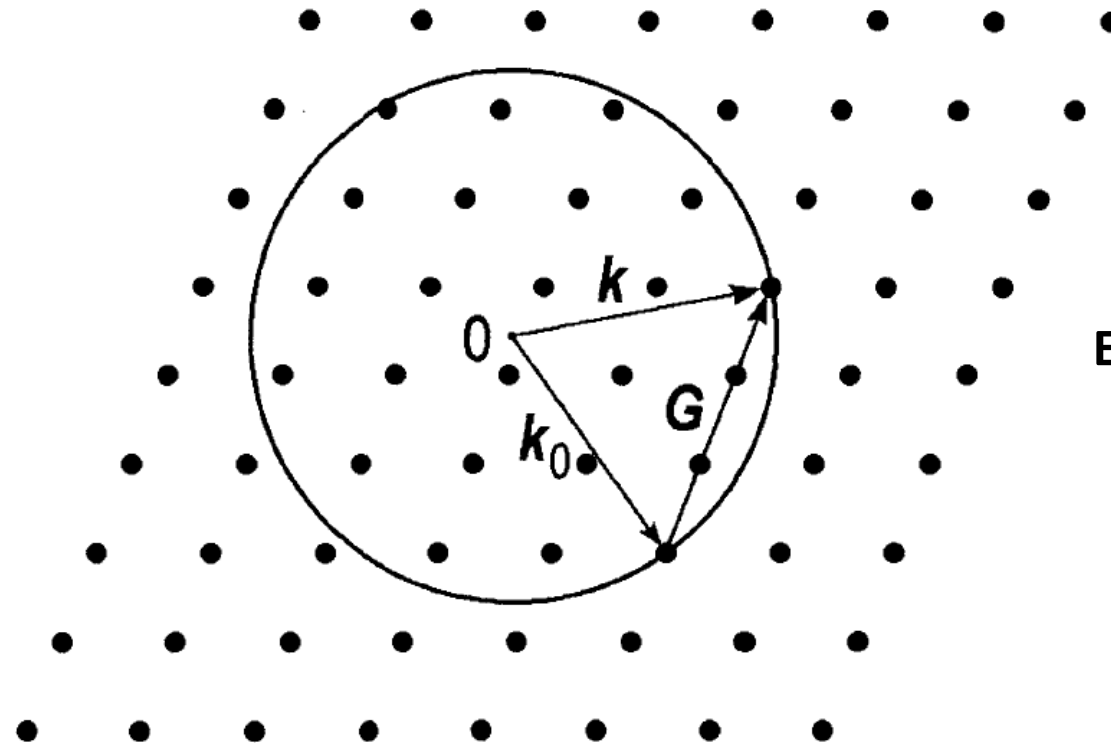
$$\mathbf{G} = \sum_i m_i \mathbf{a}_i^*, \quad \mathbf{a}_1^* = 2\pi \frac{\mathbf{a}_2 \times \mathbf{a}_3}{\mathbf{a}_1 \cdot \mathbf{a}_2 \times \mathbf{a}_3}$$

$$\mathbf{a}_i \cdot \mathbf{a}_j^* = 2\pi \delta_{ij}$$

## Laue diffraction condition

$$|\mathbf{k}| = |\mathbf{k}_0|$$

$$\mathbf{k} - \mathbf{k}_0 = \mathbf{G}$$



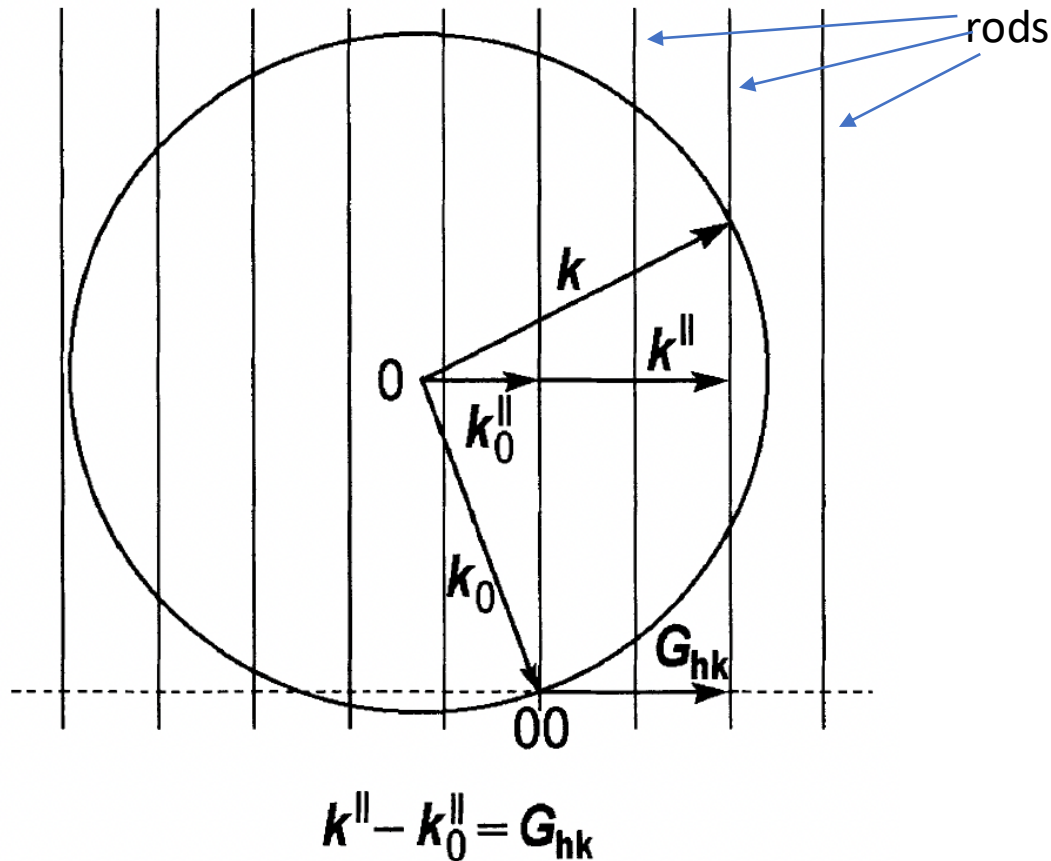
Ewald sphere

2D plane in a  
3D reciprocal  
lattice



At the surface, the bulk periodicity is truncated, and the Laue equation (3 equations, one for each direction) reduces to 2 equations concerning the components of the incident and scattered wave vectors parallel to the surface.

In contrast to the 3D reciprocal lattice points, reciprocal lattice **rods** perpendicular to the surface are attributed to every 2D reciprocal lattice point. (An ideal 2D lattice can be conceived as a 3D lattice with infinite periodicity in the normal direction. This will lead to infinitely dense reciprocal lattice points along the normal direction, thus forming rods.)

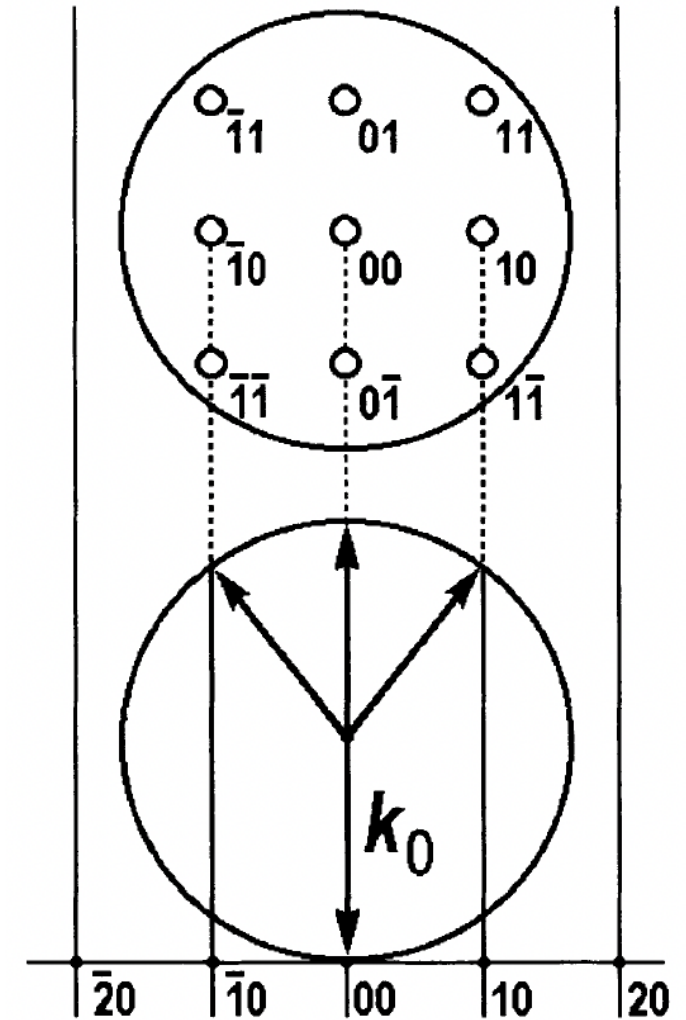
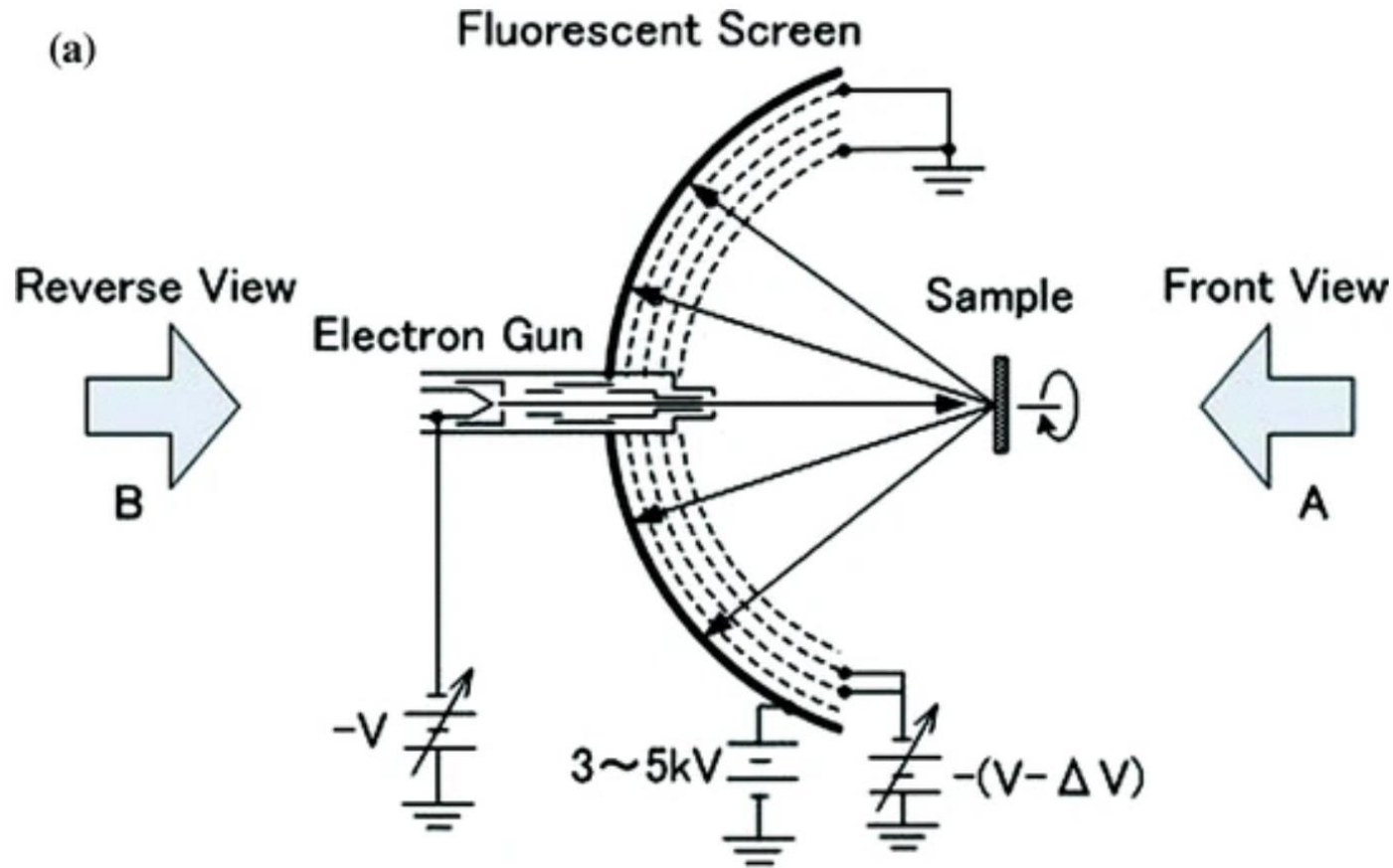


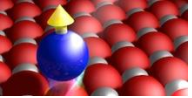
- Scattering from a surface lattice leads to diffracted beams for all incident wave vectors
- Intensity will be higher in correspondence of the points of the bulk reciprocal lattice (actually rods are “modulated”)

$$k_{\parallel} - k_{0\parallel} = G_{hk}$$

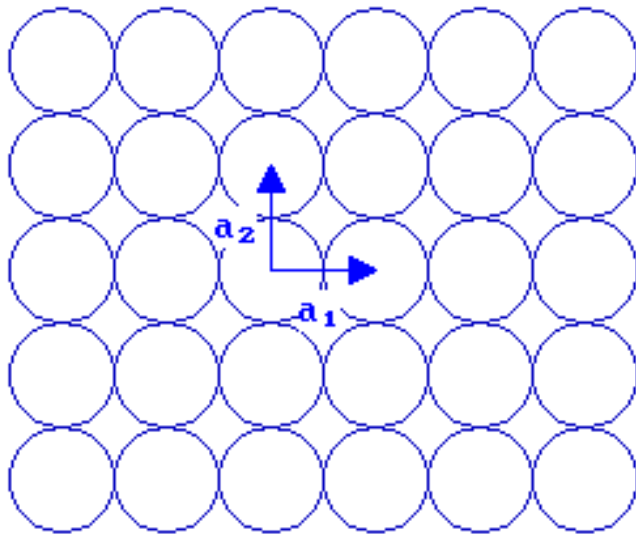
with  $G_{hk}$  a vector of the 2D reciprocal lattice

The diffraction pattern is a map of the 2D reciprocal lattice



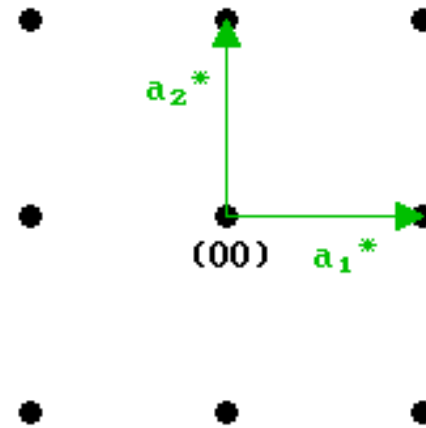


## fcc(100)



$\mathbf{a}_1$  and  $\mathbf{a}_2$  are the primitive vectors in real space

### Diffraction Pattern



$\mathbf{a}_1^*$  and  $\mathbf{a}_2^*$  are the primitive vectors in the reciprocal space

$$\mathbf{a}_i \cdot \mathbf{a}_j^* = 2\pi \delta_{ij}$$

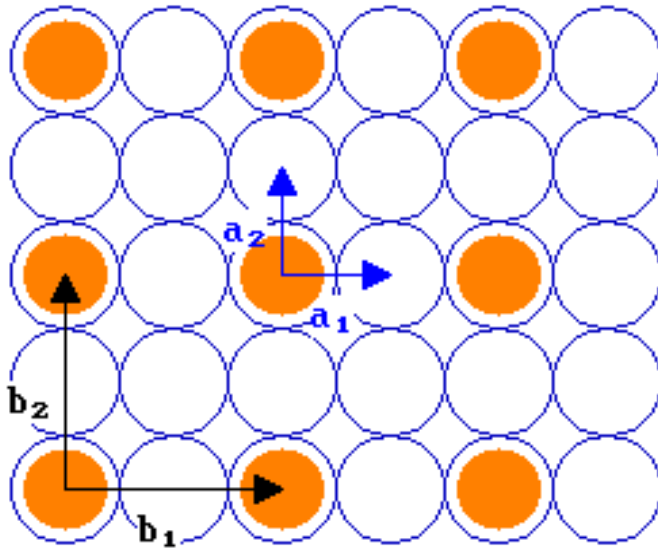
$$\mathbf{a}_1^* = \frac{2\pi}{a_1} \frac{\mathbf{a}_1}{a_1}$$

$$\mathbf{a}_2^* = \frac{2\pi}{a_2} \frac{\mathbf{a}_2}{a_2}$$

(0,0) corresponds the reflected beam along the surface normal



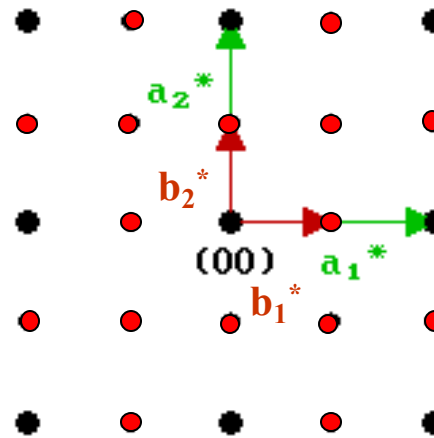
## fcc(100) (2x2)



Orange disks represents adsorbate atoms: their periodicity is responsible of the red spots in the diffraction pattern

2x2 means that the adsorbates have a double periodicity in both directions with respect to the substrate

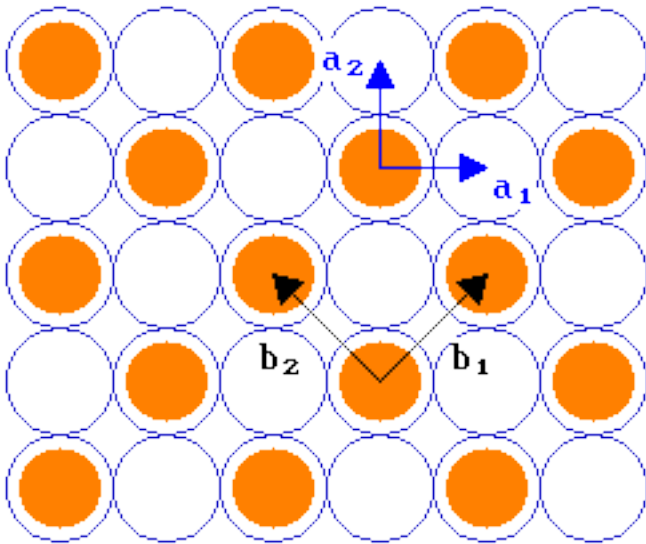
Diffraction Pattern



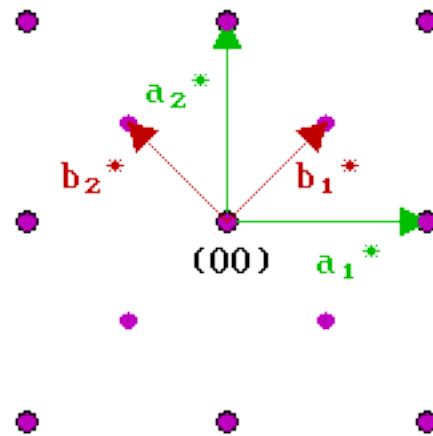
The LEED pattern provides information on the periodicity, but not on the “internal” structure of the superstructure unit cell, or on the position of the orange atoms in the unit cell.

More sophisticated measurements and analysis are needed (Spot-Profile LEED, simulations)

## fcc(100) c(2x2)



## Diffraction Pattern



The surface unit cell is rotated by  $45^\circ$  with respect to the substrate one:  
 $\sqrt{2} \times \sqrt{2} R45^\circ$

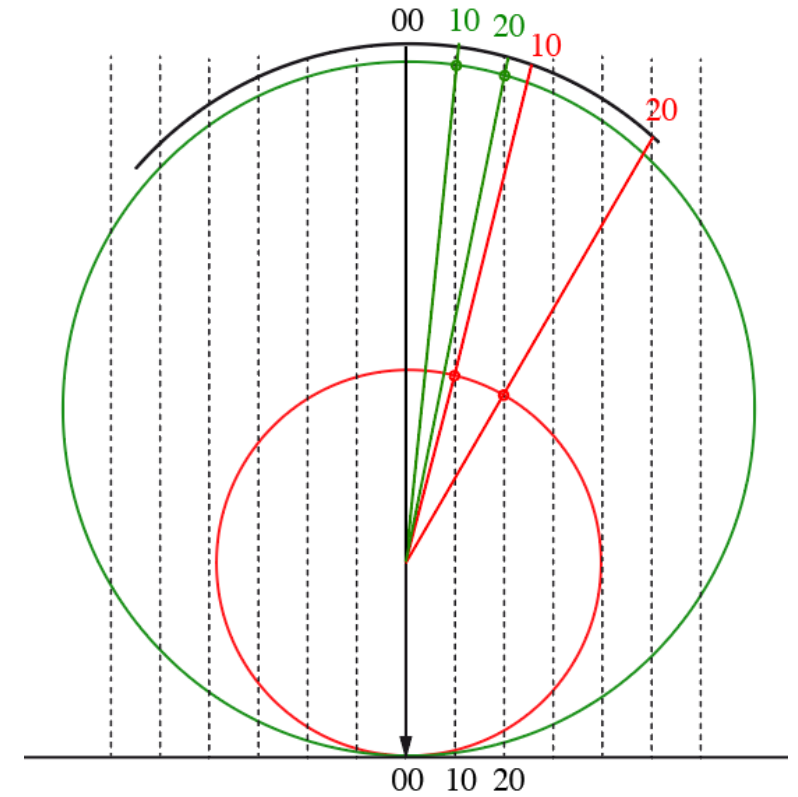
However, it is usual to describe this structure as c(2x2), where “c” stays for “centered” to indicate that there is an additional atom in the center of the 2x2 cell. (This kind of nomenclature is used especially for structures with square or rectangular symmetry).



diffraction pattern recorded as a function of the energy of the incident electrons (increasing energy)



$$E = \frac{\hbar^2 k^2}{2m}$$





Ir(111)



graphene/Ir(111): modulated surface

Electrons scattered by the potential due to :

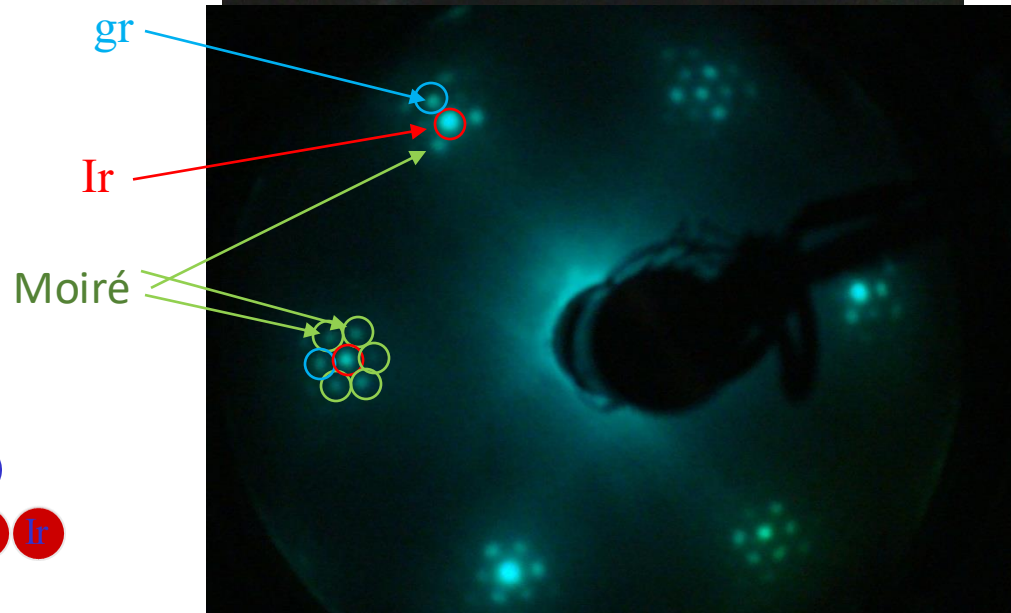
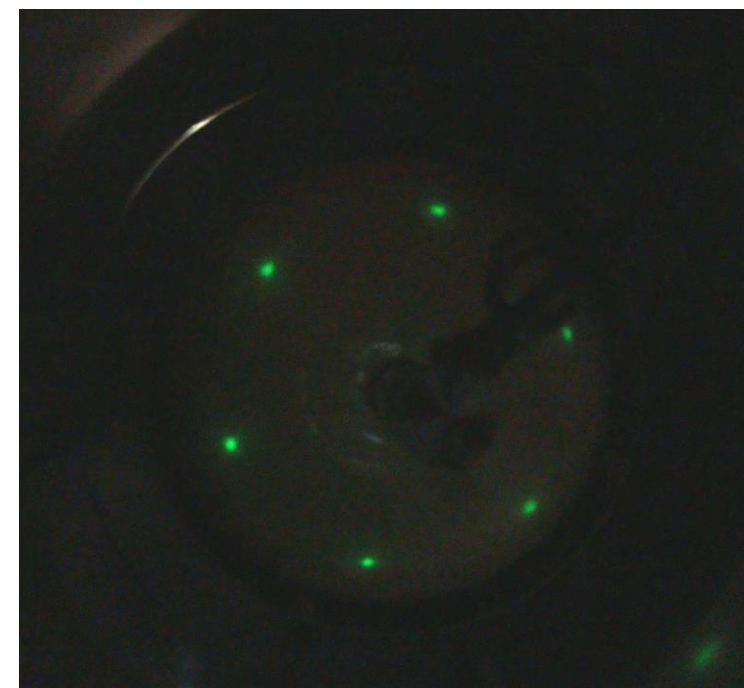
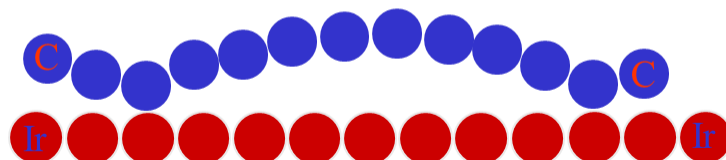
- 1) Ir atoms
- 2) C atoms
- 3) C-Ir bond formed at sites with minimum C-Ir separation (Moiré periodicity)

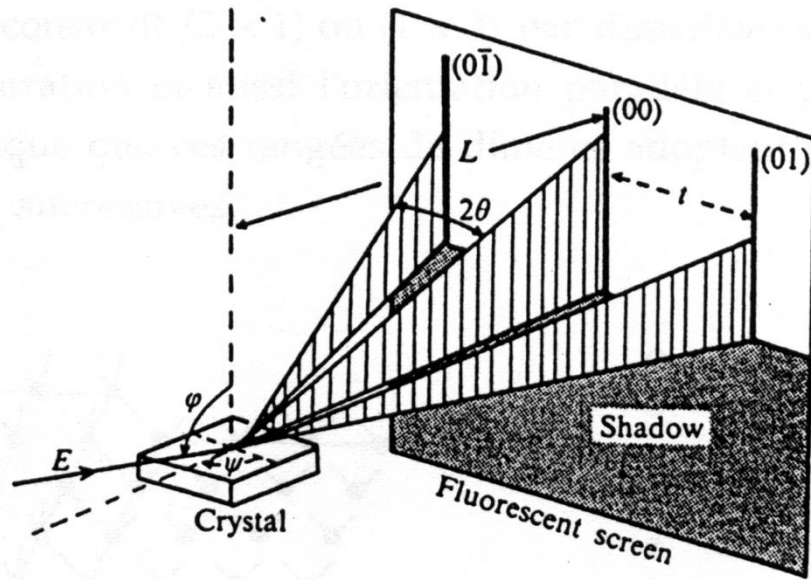
satellite reflections (only the reflections surrounding the main reflections are visible)

$$a_{\text{Ir}} = 0.27 \text{ nm}$$

$$a_{\text{C}} = 0.245 \text{ nm}$$

10 x 10 C atoms over  
9 x 9 Ir atoms

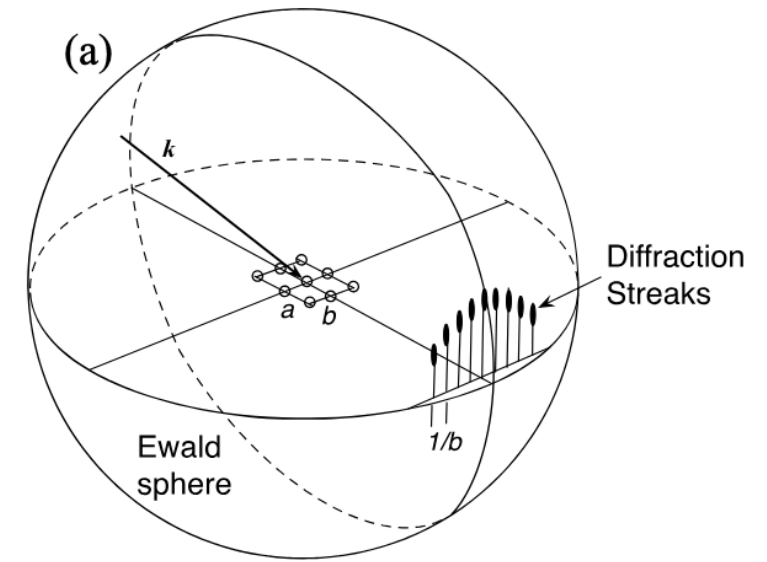




- (0,0) is the specular bar
- Lateral bars are the diffraction ones

Electron penetration depth  $\lambda = 10 - 30$  nm at 40 keV for normal incidence

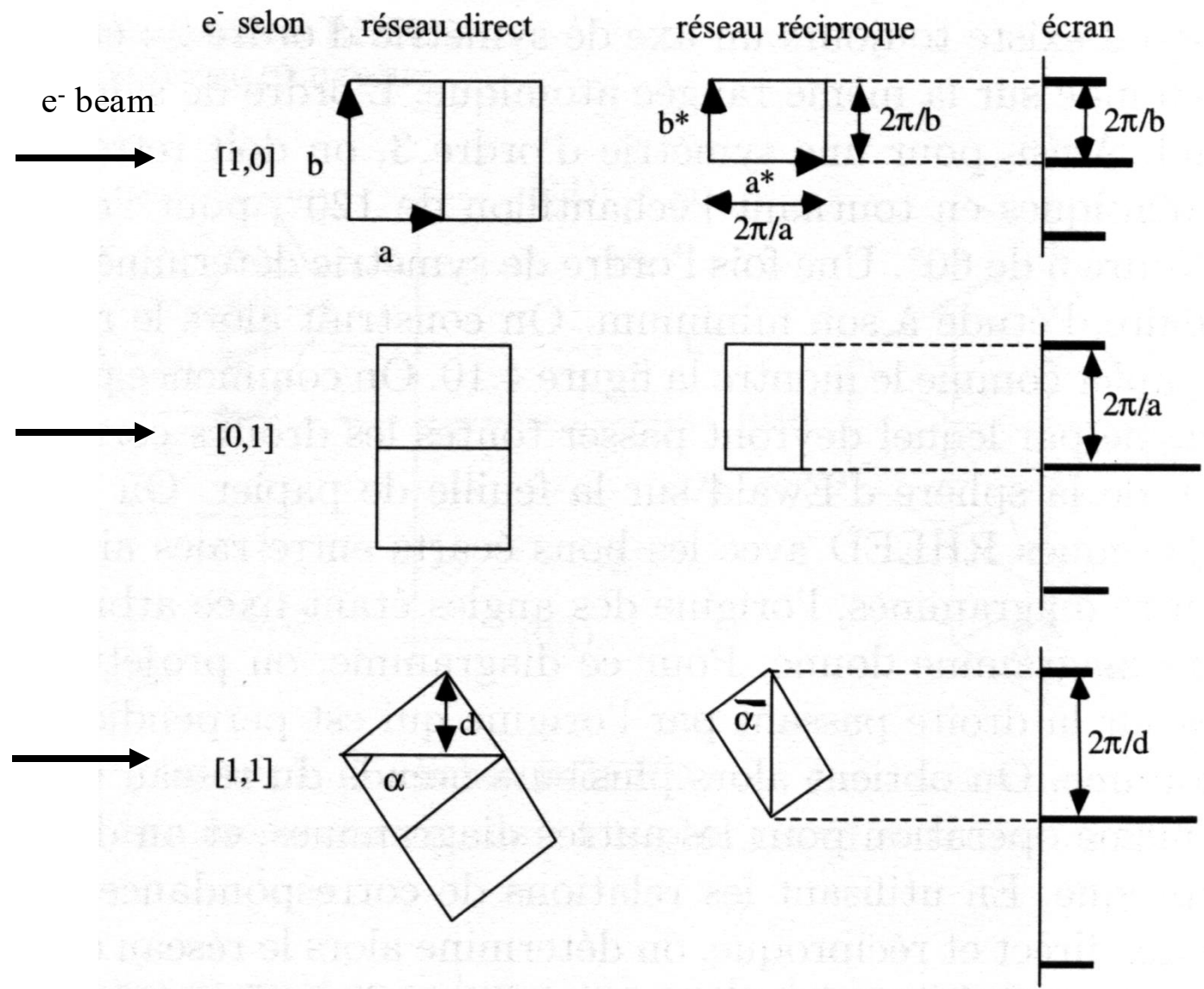
Surface sensitivity if  $\cos \varphi = d/\lambda$  with  $d = 1-2$  atomic layers  
or  
 $\varphi$  about  $89^\circ$



Due to the grazing incidence, the surface is seen as continuous along the incidence direction (no diffraction) and consequently only the periodicity perpendicularly to the incidence direction is detected



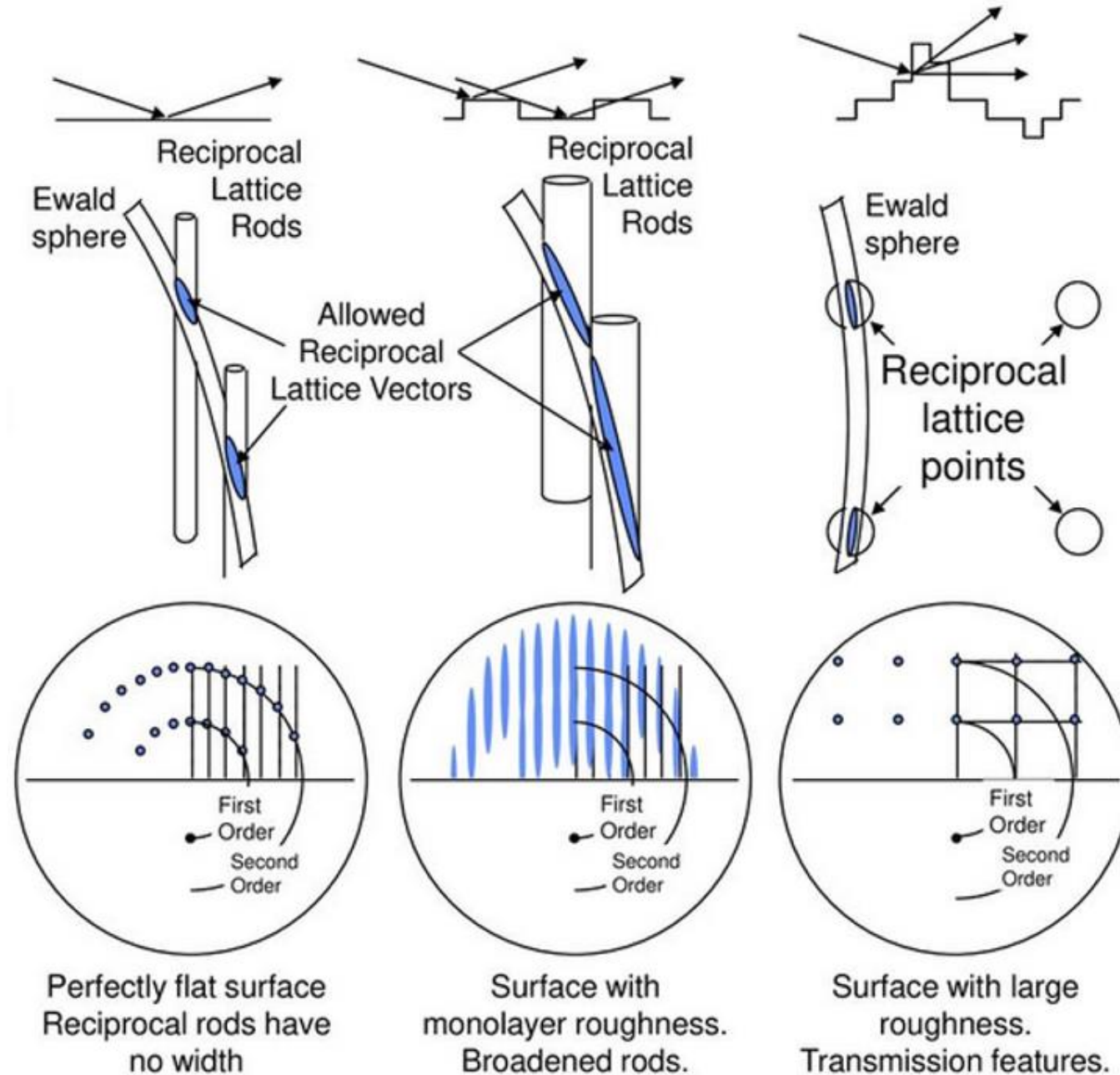
## Exercise 7.5

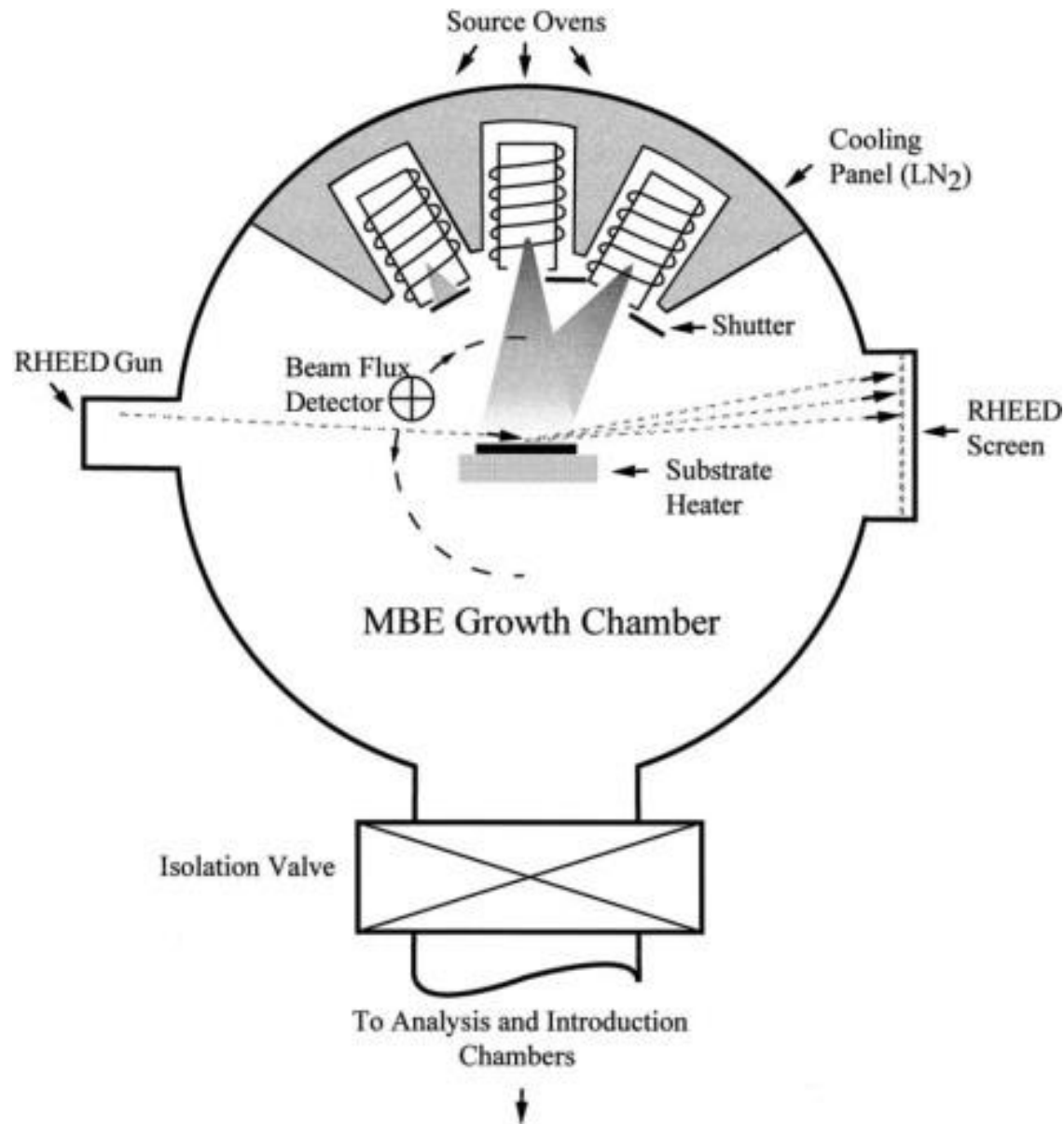
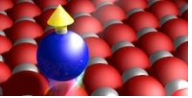


At least measurements at two different angles are needed to get the 2D reciprocal lattice



# RHEED: qualitative pattern interpretation



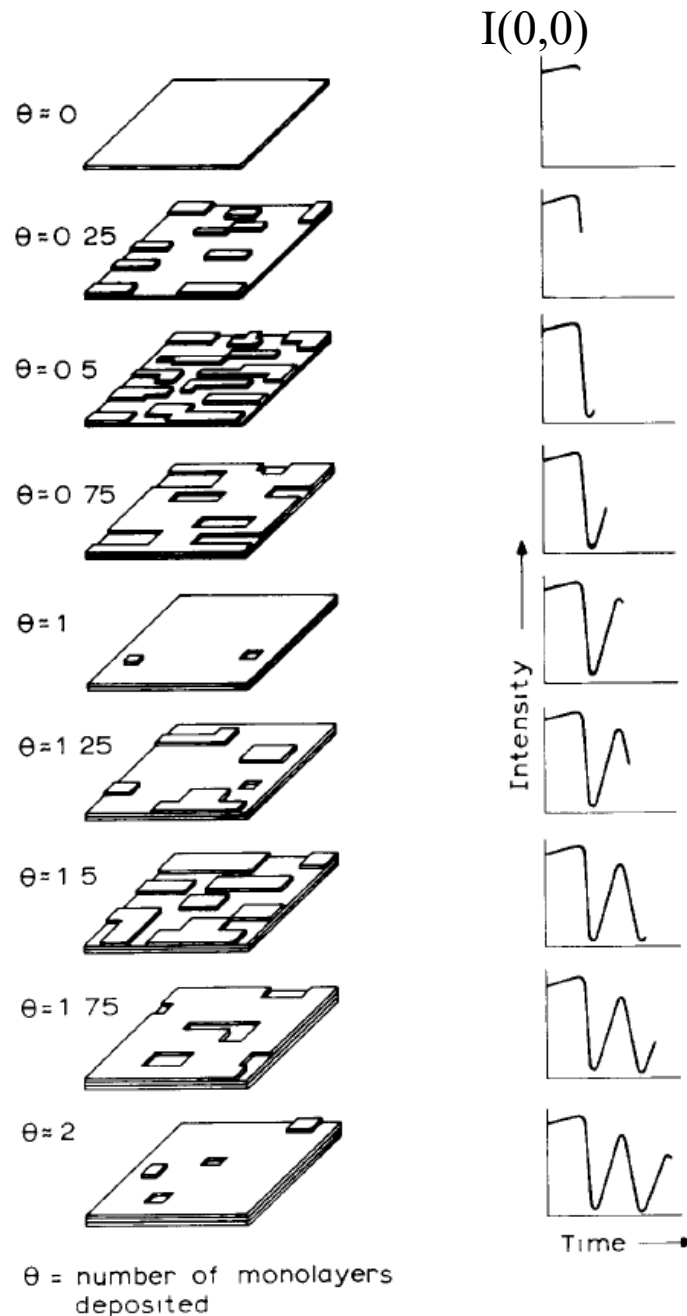


Evolution of the diffraction spots / strikes intensities are used to check the evolution of the crystal structure



Intensity of the specular signal vs time

It depends on the surface roughness (step edge density)



During layer-by-layer growth:  
maximum roughness at 0.5 ML

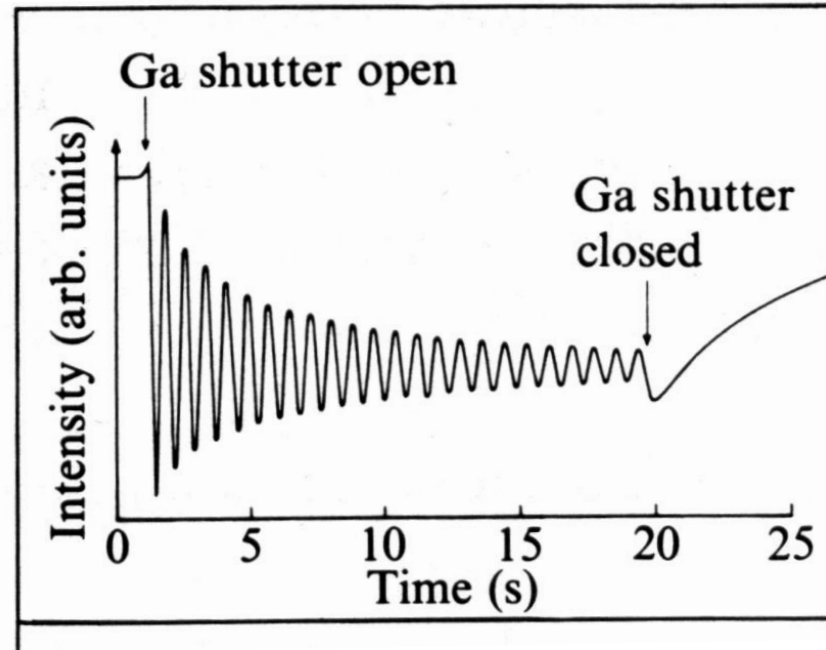
Oscillation period of  $I(0,0)$ : 1 atomic layer



## Exercise 7.6

Growth of GaAs(100) by  
Molecular Beam Epitaxy (MBE)

Oscillation periodicity: 1 atomic layer  
Intensity: decreasing with number of layers



Imperfect layer-by-layer growth

Interpretation of the evolution of the  
intensity minimum is more subtle

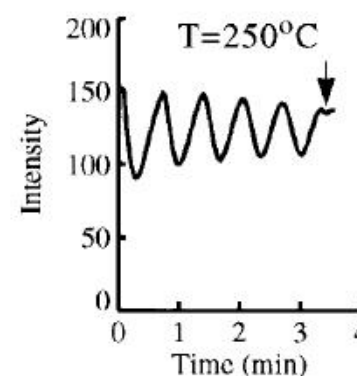
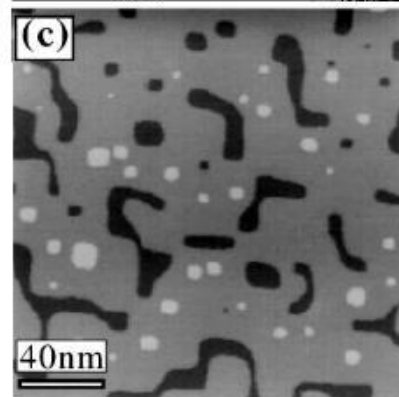
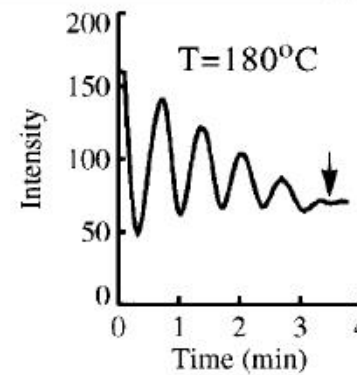
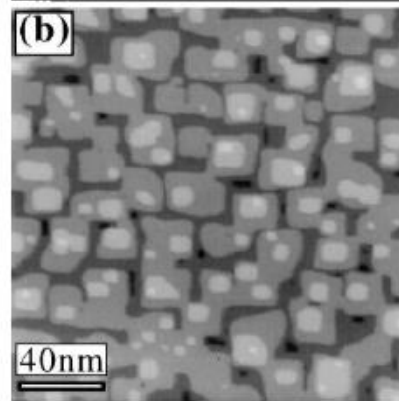
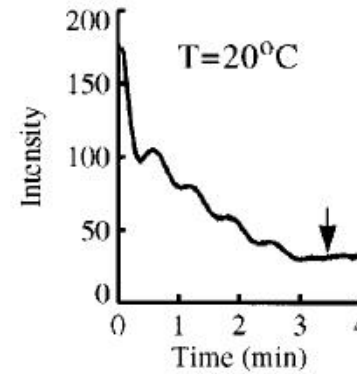
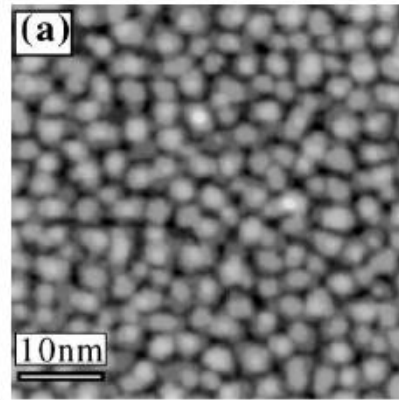
The intensity oscillations of the specular spot (0,0) can be used to  
measure the number of deposited layers and surface roughness



# Example: growth of Fe / Fe(100)

STM

0,0 Intensity



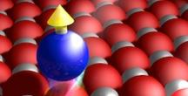
Growth of Fe on Fe(100) as seen by STM and RHEED:

change of growth mode from

three-dimensional island growth at room temperature

to

layer-by-layer growth at higher temperature



Due to the limited penetration depth of electrons in solids, the samples should be very thin: the acceptable thickness is 10-100 nm for conventional microscopes with accelerating voltages of 50-200 keV.

The electron beam loses part of its intensity due to scattering.

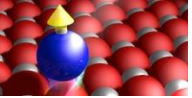
The loss is greater for thicker regions and regions with species of higher atomic number → thicker regions and regions of higher atomic number appear dark.

With high-resolution TEM, images showing the atomic structure can be obtained.

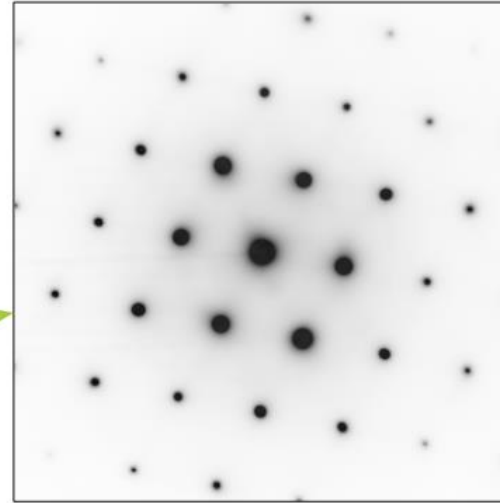
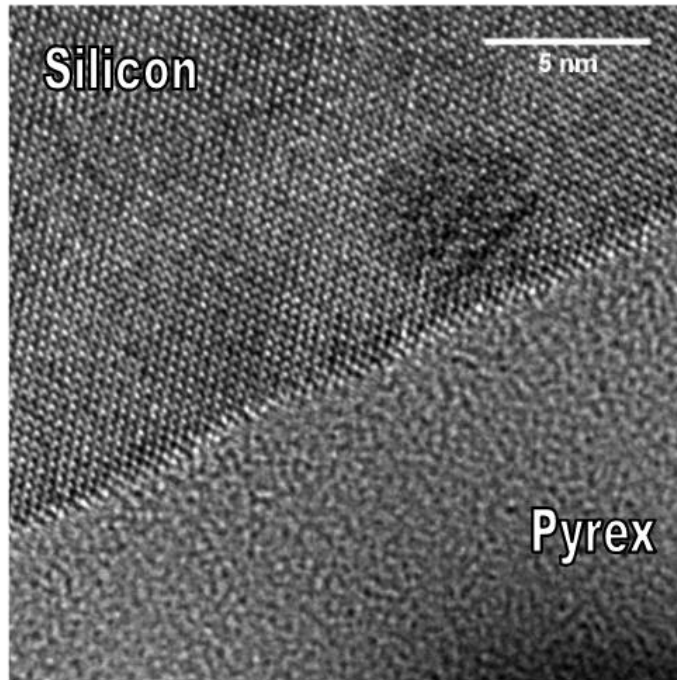
For complex systems, numerical image simulations are required for a reliable interpretation of the image features in terms of atomic structure.

Very sophisticated optics are required

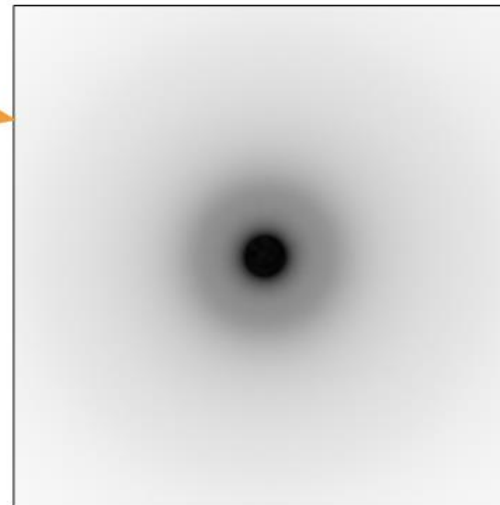
A diffraction mode is also possible



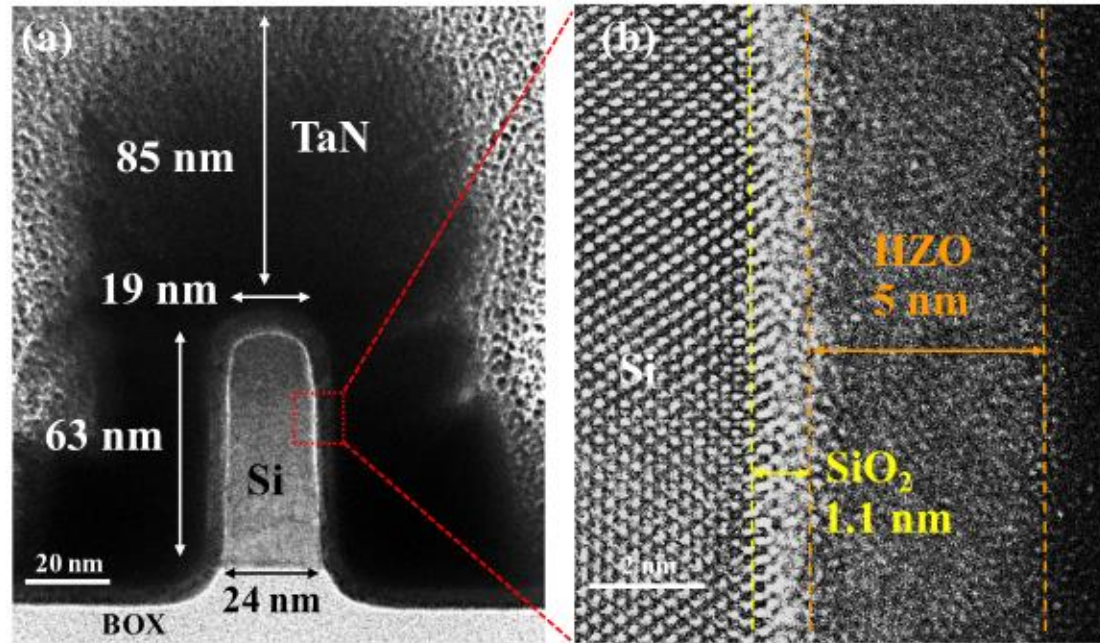
## Silicon – Pyrex (glass) interface



Diffraction pattern of a crystallized material



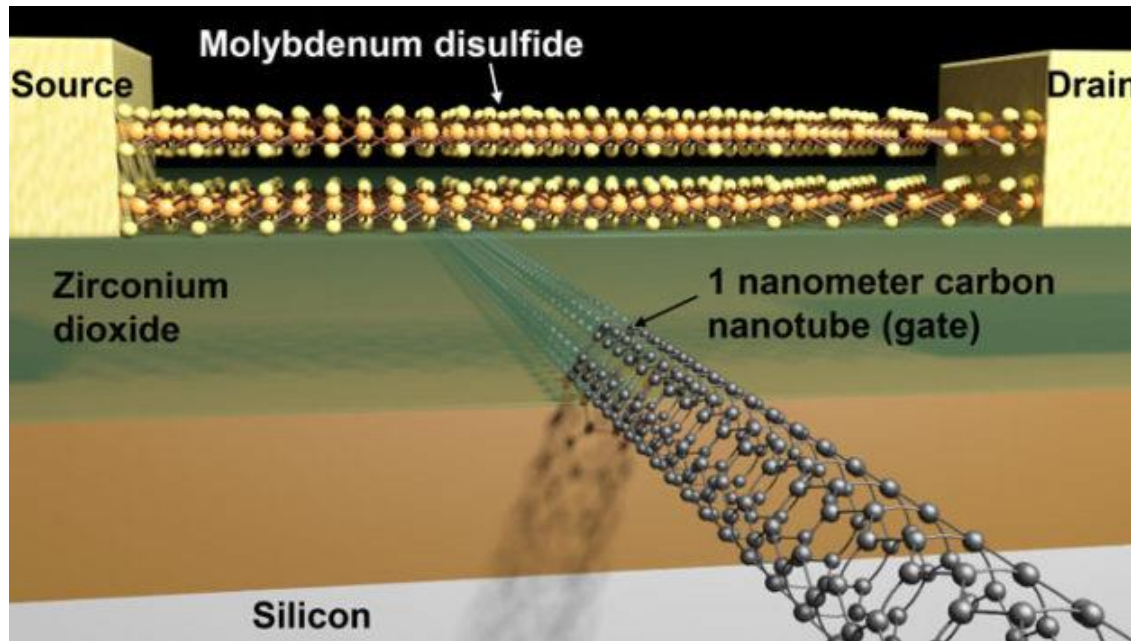
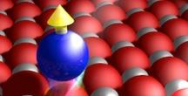
Diffuse rings typical of amorphous material



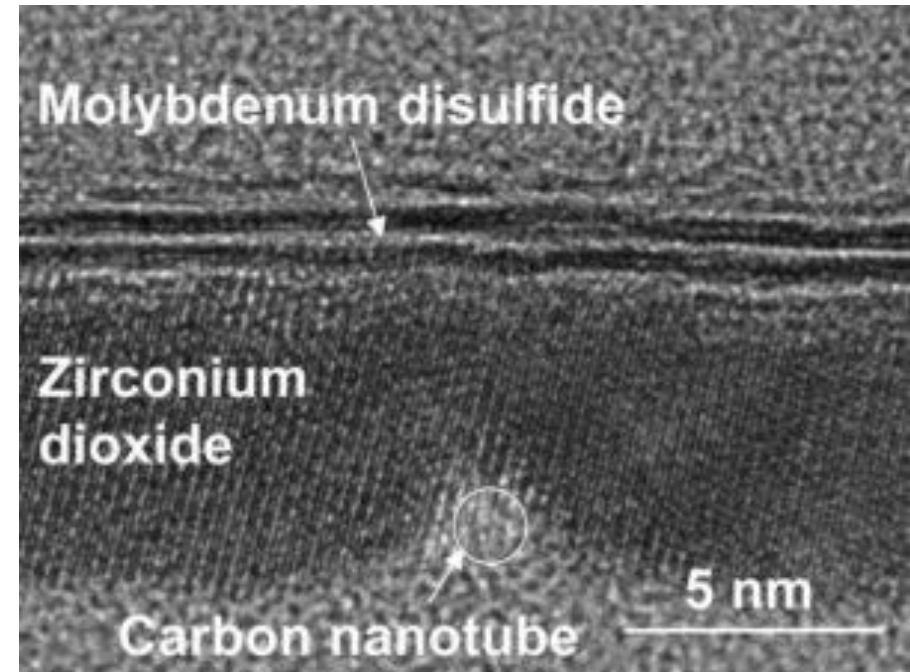
$\text{Hf}_{0.5}\text{Zr}_{0.5}\text{O}_2$  (HZO)      $\kappa \approx 20 - 40$

(a) TEM image of a FinFET

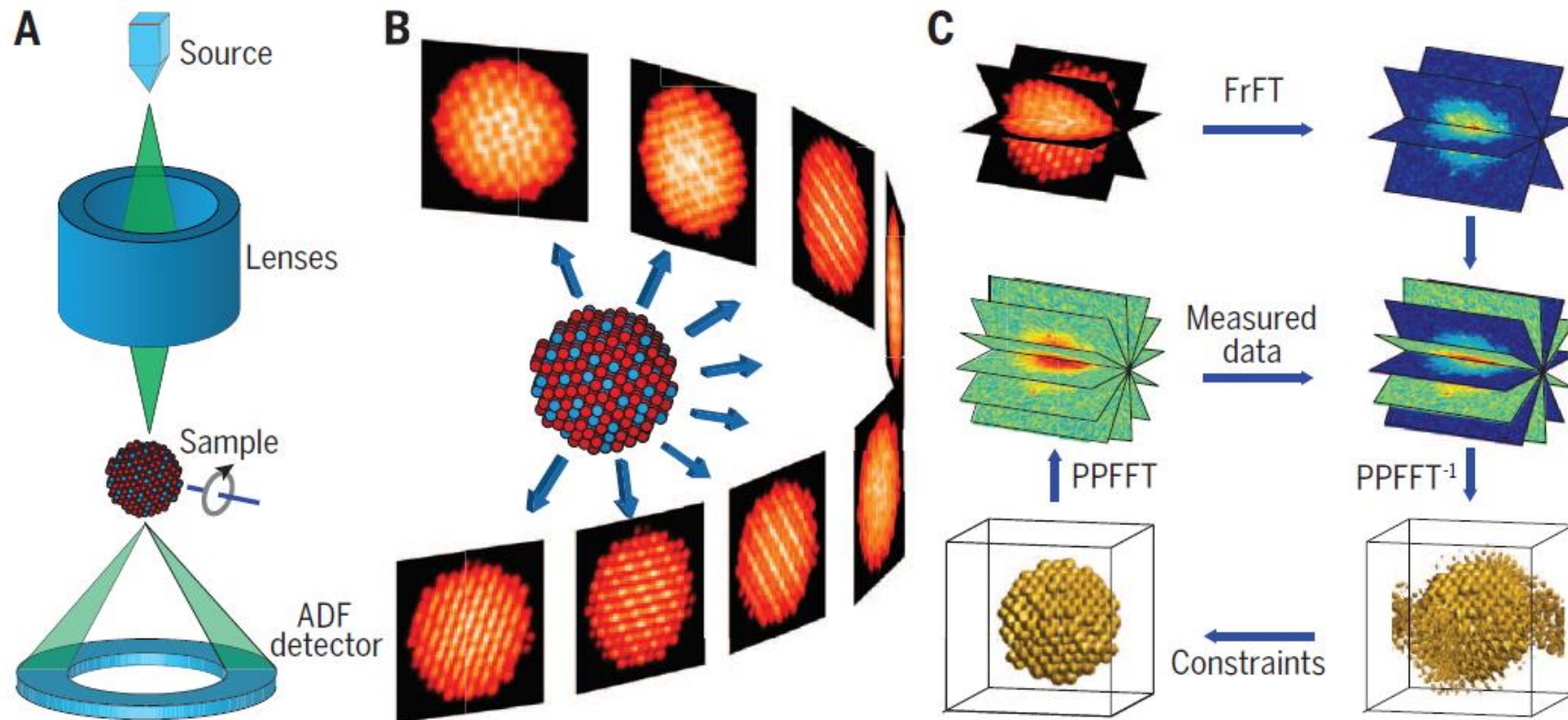
(b) Gate oxide stacks of  $\text{SiO}_2$  and HZO layers



Schematic of a transistor with a molybdenum disulfide channel and 1-nanometer carbon nanotube gate



TEM image of a cross section of the transistor



**Fig. 1. Schematic layout of AET.** (A) An electron beam is focused on a small spot and scanned over a sample to form a 2D image. The integrated signal at each scanning position is recorded by an ADF detector. (B) By rotating the sample around a tilt axis, a series of 2D images is measured at different tilt angles. (C) After preprocessing and alignment, the tilt series is inverted to Fourier slices by the fractional Fourier transform (FrFT). A 3D reconstruction is computed by using a Fourier-based iterative algorithm. From the 3D reconstruction, the coordinates of individual atoms are traced and refined to produce the 3D atomic model of the sample.