

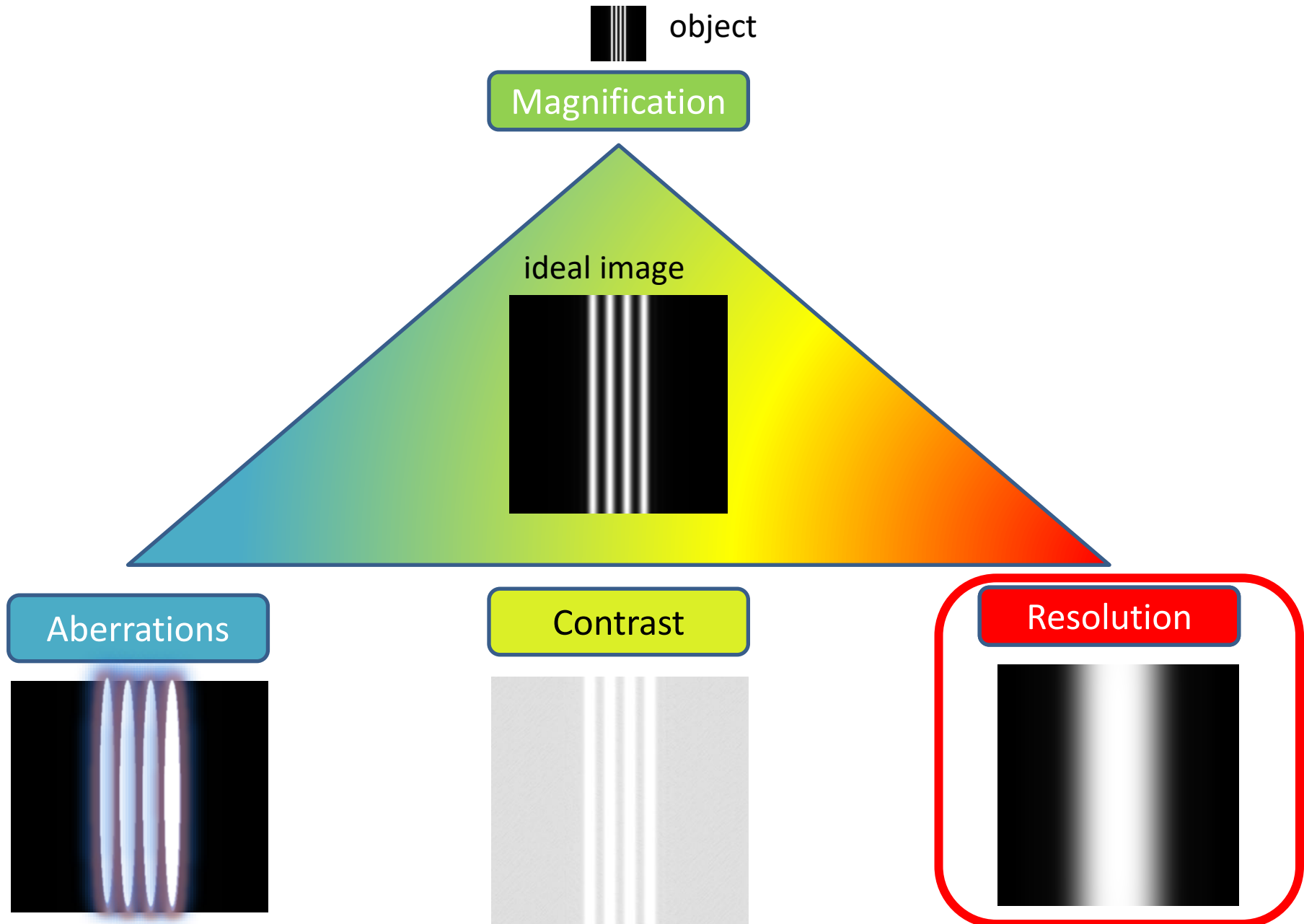
# MICRO-561

Fundamentals of Biomicroscopy

# Syllabus (tentative)

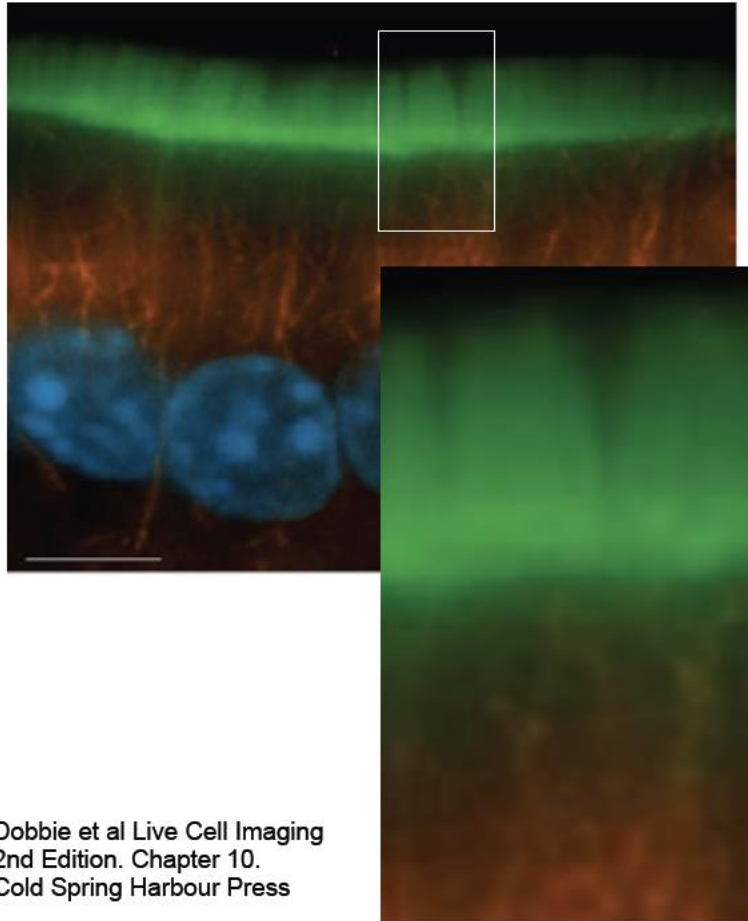
Lecture 1	Introduction & Ray Optics-1
Lecture 2	Ray Optics-2 & Matrix Optics-1
Lecture 3	Matrix Optics-2
Lecture 4	Matrix Optics-3 & Microscopy Design-1
Lecture 5	Microscopy Design-2
Lecture 6	Microscopy Design-3 & Resolution -1
Lecture 7	Resolution-2
Lecture 8	Resolution-3
Lecture 9	Contrast
Lecture 10	Fluorescence-1
Lecture 11	Fluorescence-2
Lecture 12	Fluorescence-3, Sources, Filters
Lecture 13	Detectors
Lecture 14	Bio-application Examples

# Important aspects for microscopy

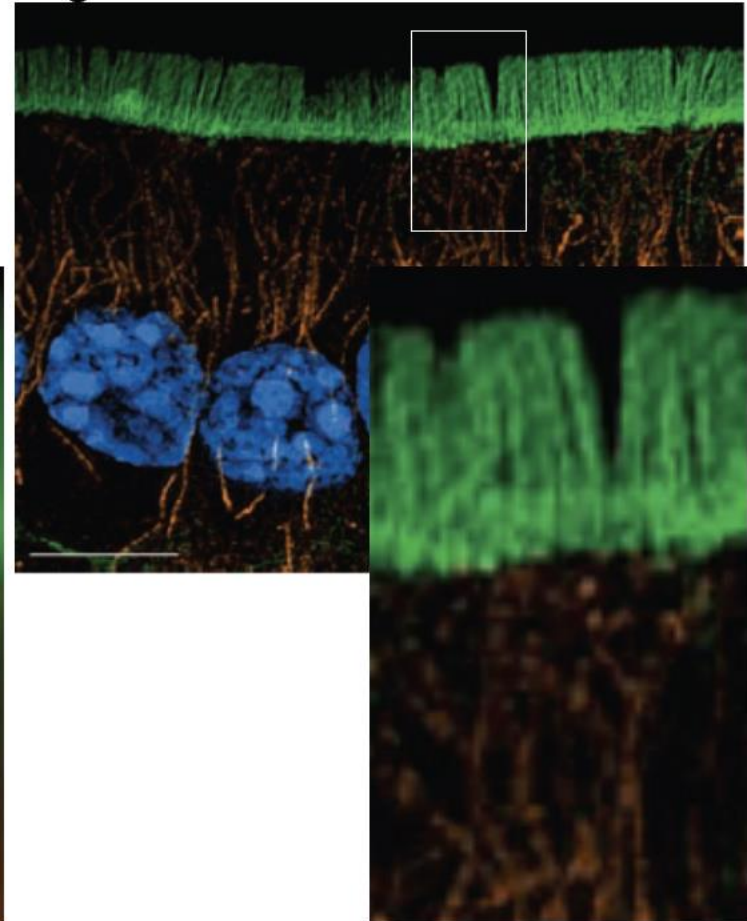


# Importance of Resolution in Microscopy

Normal resolution



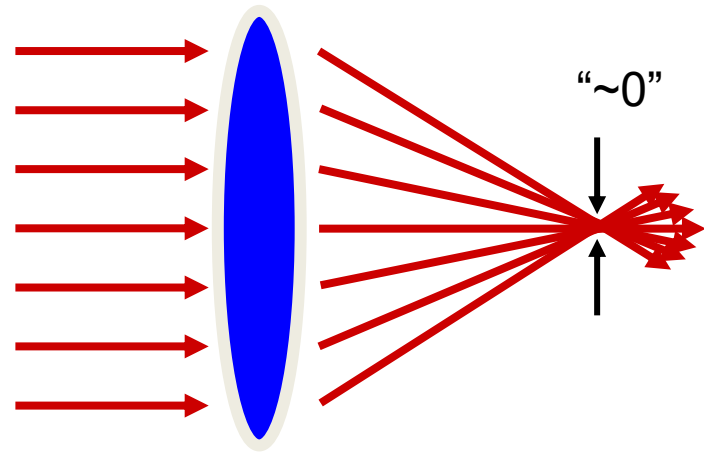
High resolution



Dobbie et al Live Cell Imaging  
2nd Edition. Chapter 10.  
Cold Spring Harbour Press

# Geometrical optics

Ray (geometrical) optics imply that we can focus a beam to an ideal “point” with zero size.



What does this imply for microscopy?



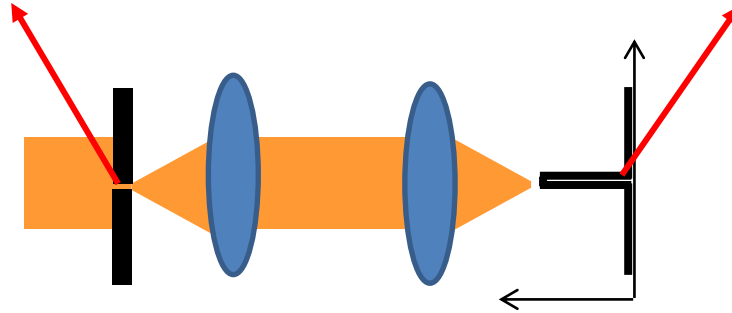
# Resolution in microscopy

## Object/Specimen:

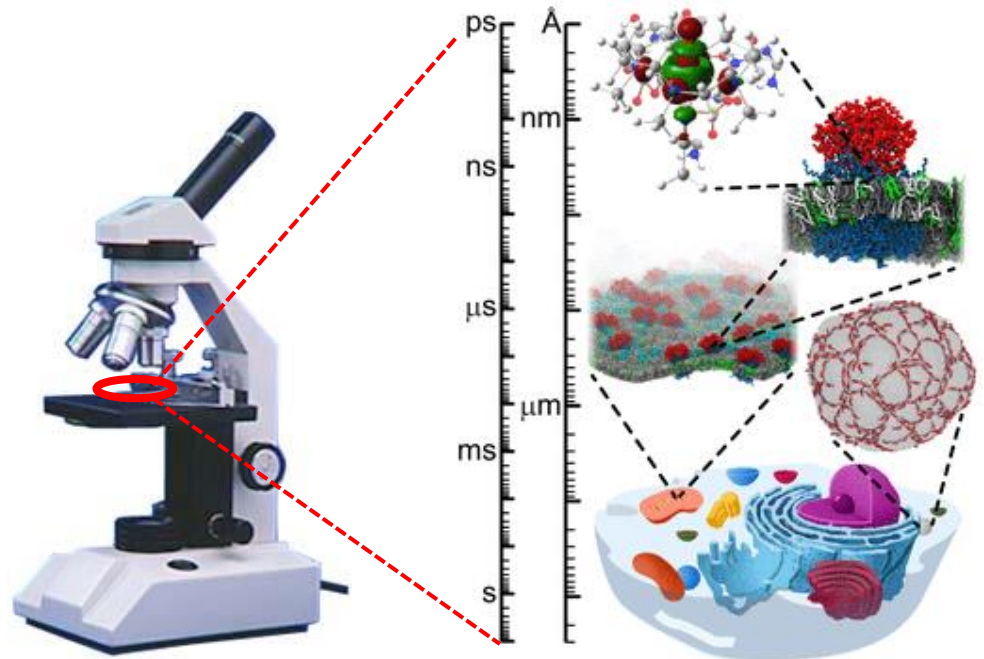
If object is a pin-hole  
(which represents a point source)

## Image:

Ideally the image will also a point source



Ability to image an exact copy of a point source (with ~zero dimension) implies that a microscope should be able to see **ANY** object, even those with extremely small dimensions such as: proteins, DNA, atoms ...

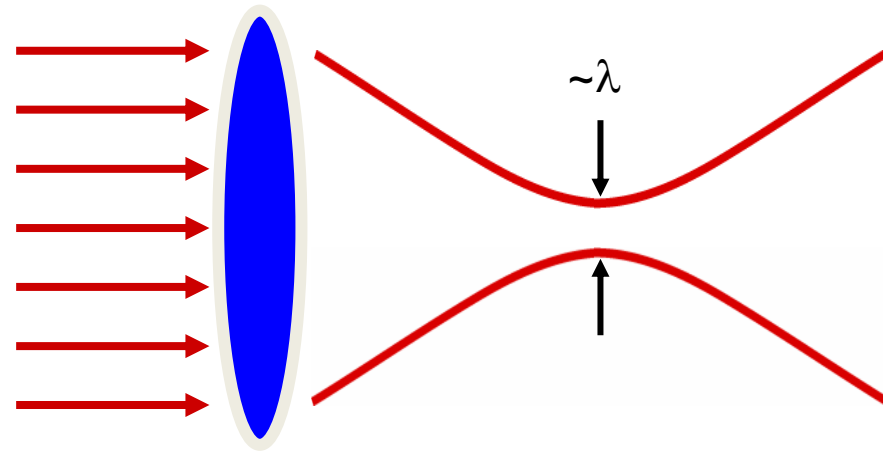
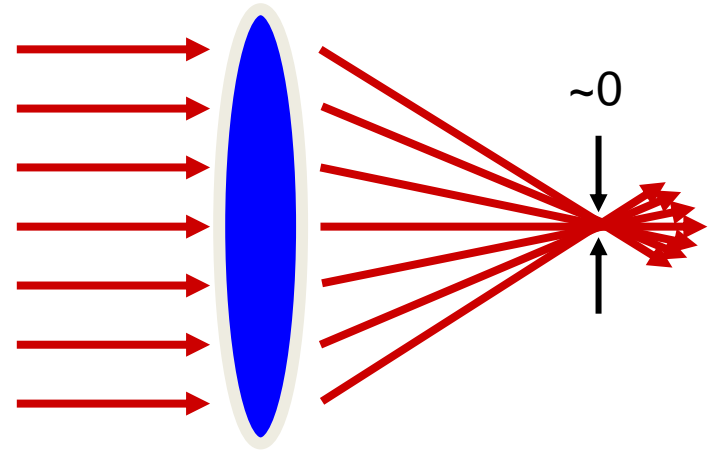


# Is geometrical optics the whole story?

Ray optics imply that we could focus a beam to an ideal “point” with zero size.

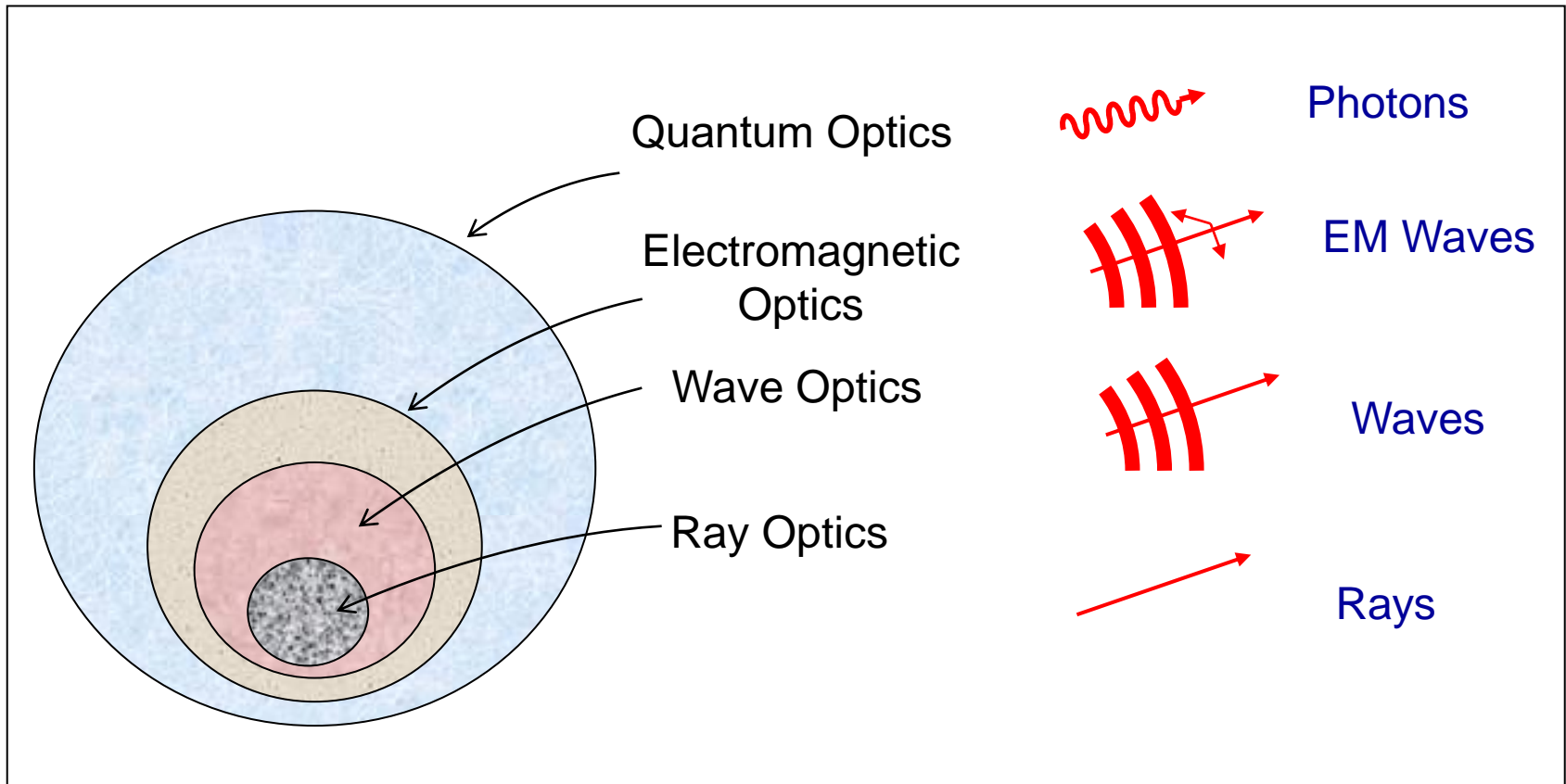
But, this does not hold in reality....

- The minimum possible “spot size” is about one wavelength,  $\lambda$ .
- The finite spot size is due to **diffraction**, which is not described by ray optics.
- This minimum spot size (i.e.  $\sim\lambda$ ) also gives the best spatial resolution one can achieve with a microscope.
- In practice, this implies an object having a dimension smaller than  $\lambda$ , will appear at best having a size of  $\lambda$ .



# Hierarchy of Theories in Optics

Ray optics cannot explain **diffraction** → we need to consider **wave optics**



# Outline

- **Wave optics**
  - **Wave equation**
  - **Wavelength, period, amplitude & phase of a wave**
  - **Monochromatic waves**
  - **Real waves and complex waves**
  - **Elementary waves: plane wave and spherical wave**

# Waves & Optics

**An optical wave is described mathematically by the **wavefunction**  $u(\mathbf{r}; t)$ , at position  $\mathbf{r} = (x, y, z)$  and time  $t$ .**

**Wavefunction satisfies the wave equation:**

$$\nabla^2 u - (1/v^2) \partial^2 u / \partial t^2 = 0$$

$$\nabla^2 = \partial^2 / \partial x^2 + \partial^2 / \partial y^2 + \partial^2 / \partial z^2$$

- “ $v$ ” is the speed of the wave.
- For E&M waves we use “ $c$ ”, for the speed of light in free space.
- Wave equation is general.
- It can be applied to E&M radiation at different spectrum (i.e. light waves, microwaves, radio waves) and also other waves such as water, acoustic, seismic..

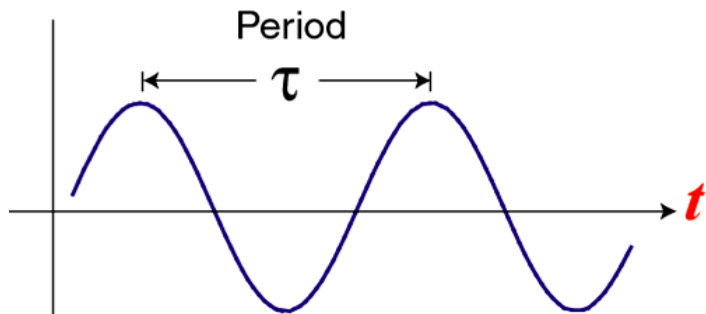
# Monochromatic waves

- A monochromatic wave is one of the possible solutions to wave equation.
- It has a harmonic dependence in time and space.

$$u(x,t) = A \cos[(k x - \omega t) - \theta]$$

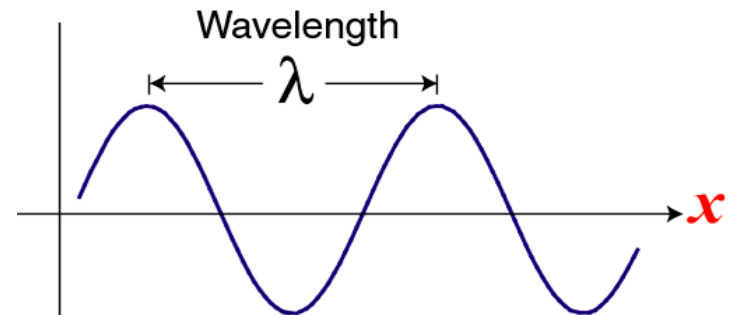
- Monochromatic wave function oscillates periodically in space & time.
- “One” wavelength defines the monochromatic wave (i.e. it is a “mono/single” color wave)

## Temporal quantities:



The angular frequency:  $\omega = 2\pi/\tau$   
The frequency:  $\nu = 1/\tau$

## Spatial quantities:

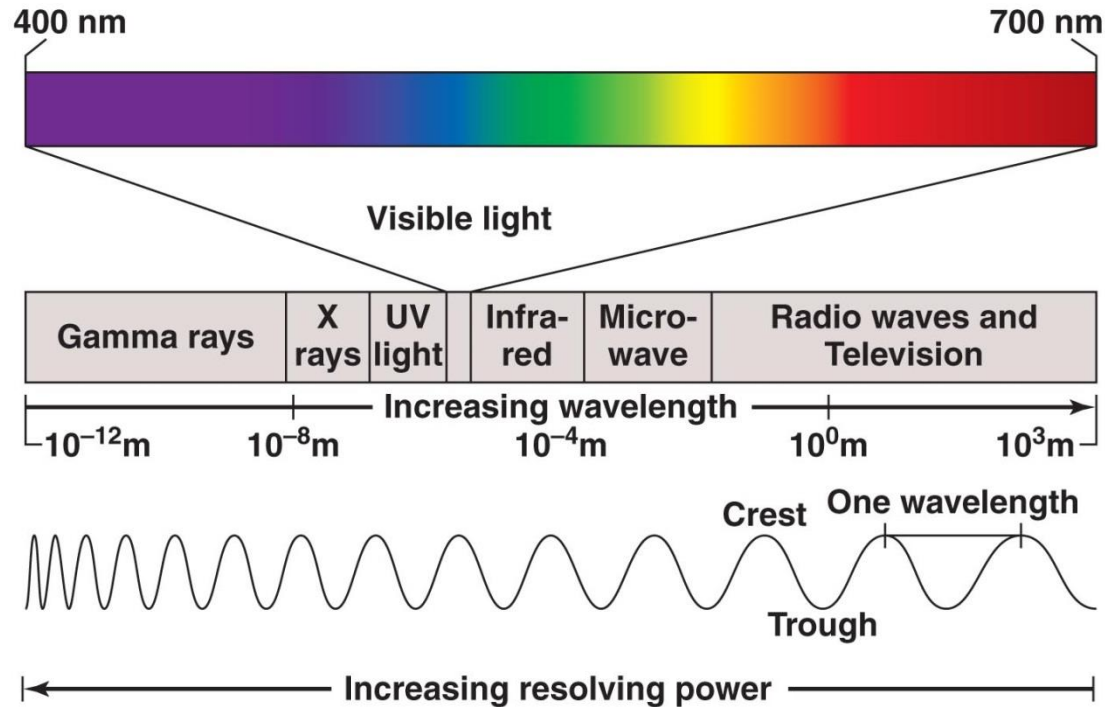


The k-vector:  $k = 2\pi/\lambda$   
The wave number:  $\kappa = 1/\lambda$

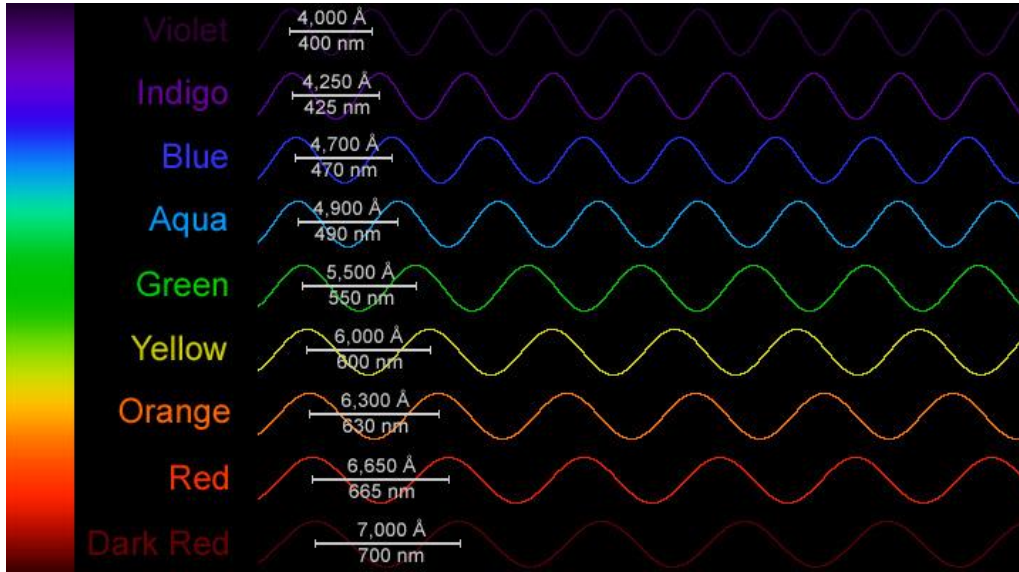
# Wavelength of Radiation

Wave radiation differs in **wavelength**

The electromagnetic spectrum

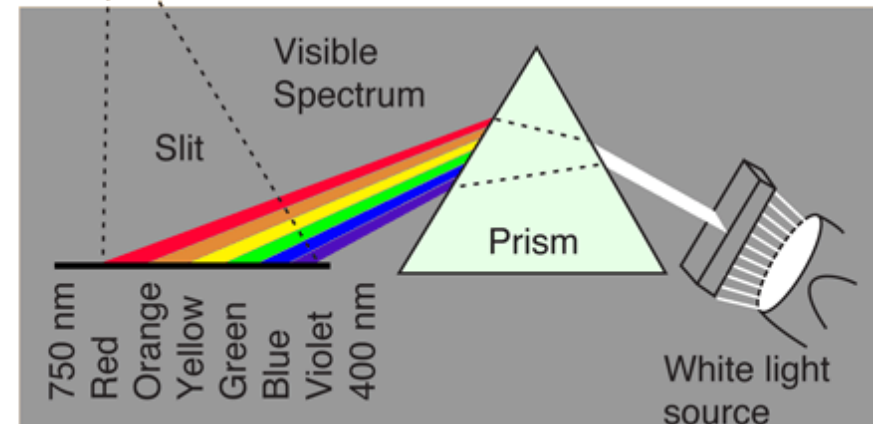


# Visible spectrum



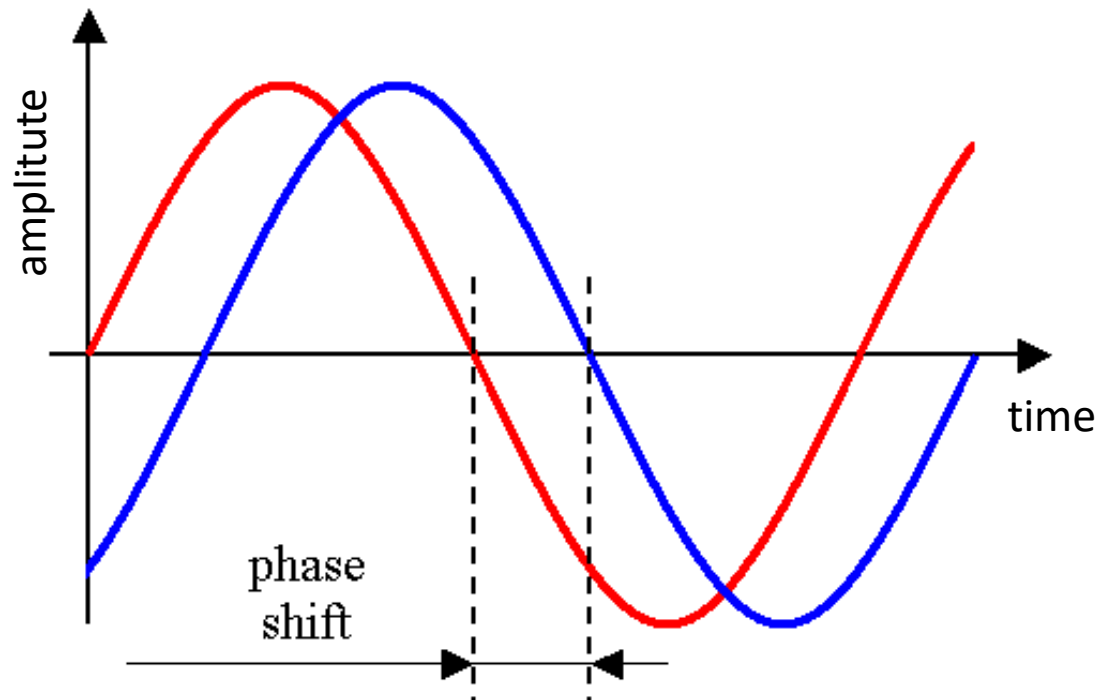
Radio	Far IR, Micro-wave	IR	UV	x-ray γ-ray
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White light is a combination of all wavelengths (colors, frequency ..)



# Characteristics of a wave - PHASE

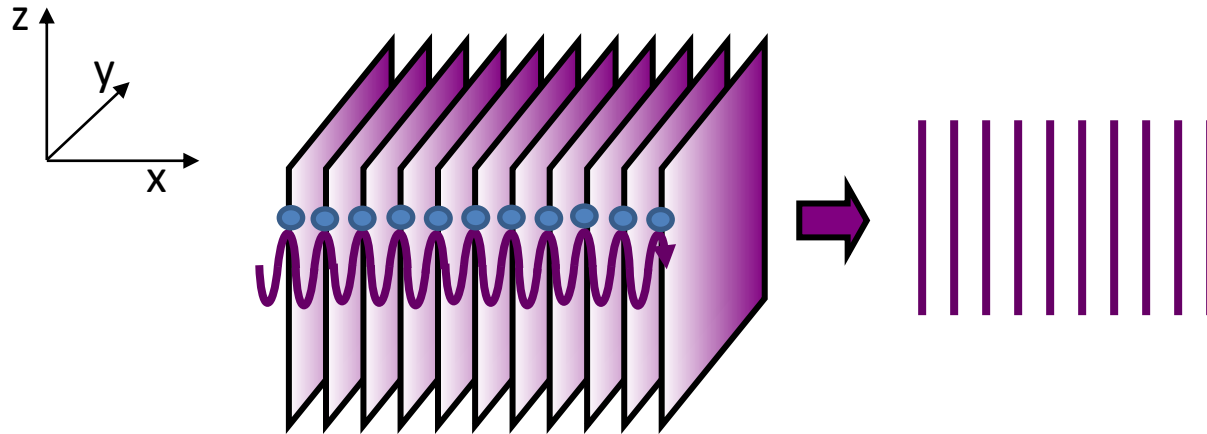
Below figure shows two waves with same amplitude & frequency (i.e. wavelength, period) but they differ in **PHASE**:



- Phase concept is introduced in wave optics.
- It does not exist in ray optics.

# Monochromatic “plane wave” function

This equation represents a monochromatic “plane” wave:  $u(x,t) = A \cos[(\omega t - kx) - \theta]$



Contours of maximum field, called **wave-fronts** or **phase-fronts**, are planes.

Usually we just draw lines

- They extend over all space (y-z).
- Wave-fronts are equally spaced: one wavelength apart.
- They're perpendicular to the propagation direction (x).

# Monochromatic wave function

This equation represents a monochromatic “**plane**” wave:  $u(x,t) = A \cos[(\omega t - k x) - \theta]$

If we generalize this equation to “**any**” monochromatic wave (planar, spherical, etc...), then the expression becomes:

$$u(r,t) = a(r) \cos[\omega t + \varphi(r)]$$

This represents a real wavefunction for a monochromatic wave with harmonic time dependence.

# Complex representation of a monochromatic wave function

In optics (and E&M), it is convenient to work with **complex wave** formalism

Real wave function is:  $u(r, t) = a(r) \cos[\omega t + \varphi(r)]$

**Complex wave** function representation of  $u(r, t)$  is:  $U(r, t) = a(r)e^{i\varphi(r)}e^{i\omega t}$

$$U(r, t) = U(r)e^{i\omega t}$$

Here,  $U(r)$  is called complex amplitude.

Mathematically,  $U(r, t)$  describes the wave completely, and the wavefunction  $u(r, t)$  is its real part:

$$u(r, t) = \text{Re}\{U(r, t)\} = \frac{1}{2}[U(r, t) + U^*(r, t)] = \frac{1}{2}[U(r)e^{i\omega t} + U^*(r)e^{-i\omega t}]$$

Here, sign \* signifies complex conjugation.

Complex wave function also satisfies the wave equation:  $\nabla^2 U - \frac{1}{c^2} \frac{\partial^2 U}{\partial t^2} = 0$

# Amplitude, intensity and phase of a monochromatic wave

$$u(r, t) = a(r) \cos [\omega t + \varphi(r)]$$

← this defines the real monochromatic wave,  $u(r, t)$

$$U(r, t) = a(r) e^{i\varphi(r)} e^{i\omega t}$$

← This is the complex wave function representation of  $u(r, t)$

$$U(r, t) = U(r) e^{i\omega t}$$

← Here,  $U(r)$  is called **complex amplitude**.

Real part of  $U(r)$  gives the **amplitude** of the wave:

$$|U(r)| = a(r)$$

**Intensity** of the wave is the square of its amplitude:

$$I(r) = |U(r)|^2 = a(r)^2$$

Imaginary part of  $U(r)$  gives the **PHASE** of the wave:

$$\arg\{U(r)\} = \varphi(r)$$

- “Phase” is introduced in wave optics. It does not exist in ray optics.
- “Phase” explains the “**interference & diffraction**” phenomena.

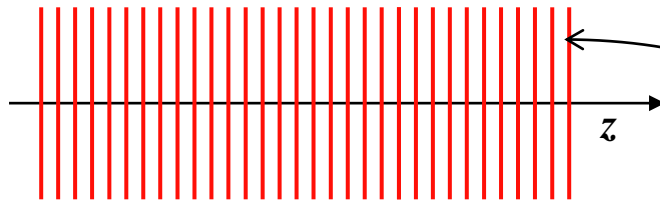
# 1<sup>st</sup> example for “elementary wave”: Plane wave

- Plane wave traveling in  $z$ -direction

$$u(\mathbf{r}; t) = a \cos[\omega t - kz + \varphi_0]$$

$$U(\mathbf{r}) = a \exp(i\varphi_0) \exp(-ikz)$$

$$U(\mathbf{r}) = A \exp(-ikz)$$



Wavefronts  $\varphi_0 - kz = 2\pi m$   
 $kz = 2\pi m + \text{constant}$

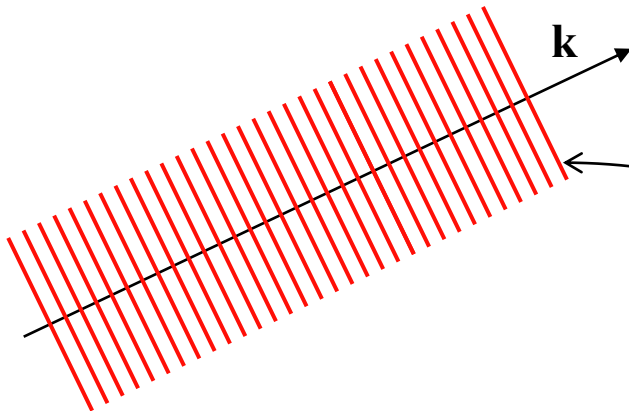
$$U(r) = Ae^{\pm ikz}$$

Propagation direction  
(left or right)

- Plane wave traveling in arbitrary direction

$$U(\mathbf{r}) = A \exp(-i\mathbf{k} \cdot \mathbf{r})$$

$$= A \exp[-i(k_x x + k_y y + k_z z)]$$

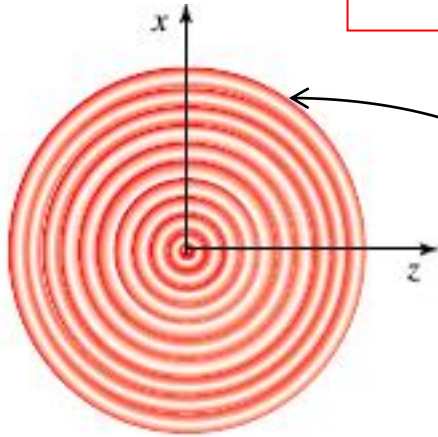


Wavefronts  $k_x x + k_y y + k_z z = 2\pi m + \text{constant}$

# 2<sup>nd</sup> example for “elementary wave”: spherical wave

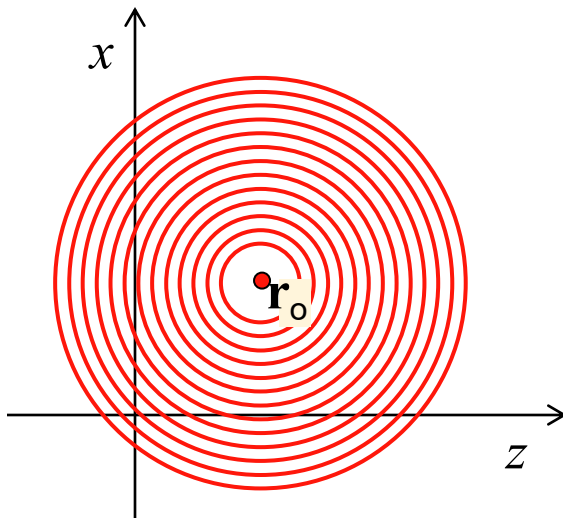
$$U(r) = \frac{A}{r} e^{\pm ikr}$$

Propagation direction  
(inward or outwards)



Wavefronts  $kr = 2\pi m + \text{constant}$

$$I(\mathbf{r}) = |A|^2 / r^2$$



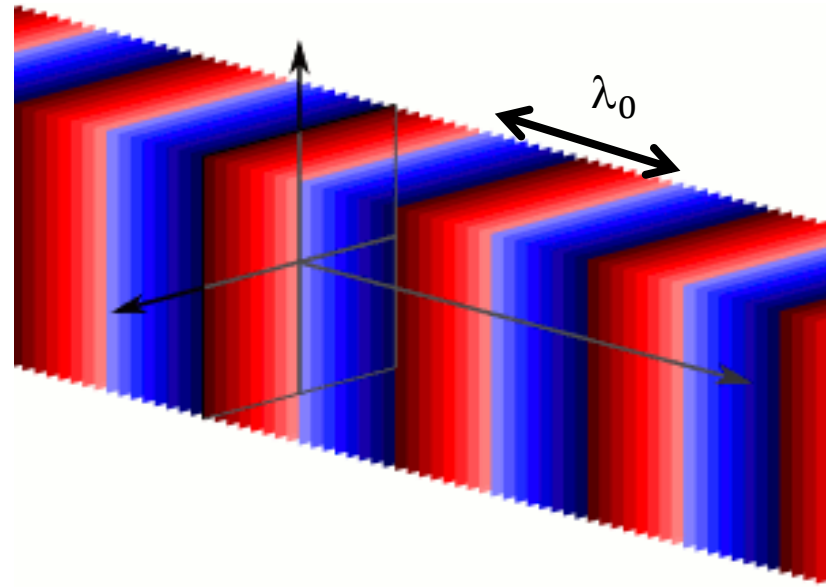
$$U(r) = \left( \frac{A}{|r - r_0|} \right) e^{-ik|r - r_0|}$$

Spherical wave centered at  $\mathbf{r}_0$

# Wave Propagation

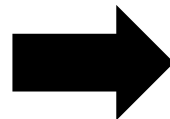
## Example:

Propagating plane wave in free space



- In free space, one wavelength ( $\lambda_0$ ) is the distance that light travels over one period in time ( $\tau$ ) at a speed of  $c$ :

Distance = Speed x Time



$$\lambda_0 = c \times \tau$$

- Velocity in vacuum (i.e. " $c$ ") =  $2.99792458 \cdot 10^8$  m/sec

# Wave Propagation in Free Space

- In free space, one wavelength ( $\lambda_0$ ) is the distance that light travels over one period in time ( $\tau$ ) at a speed of  $c$ :

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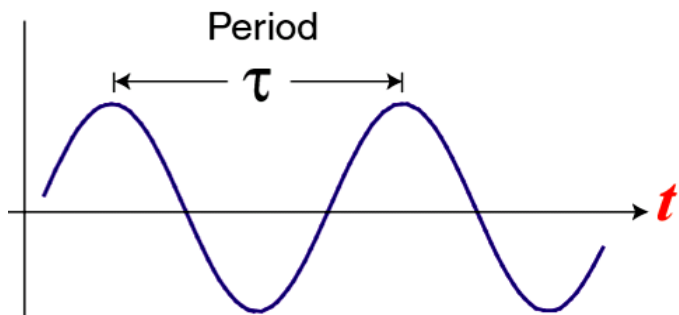
- Velocity in vacuum (i.e. “ $c$ ”) =  $2.99792458 \cdot 10^8$  m/sec

- Recall the definition of frequency & period in time:  $\nu = 1/\tau$

- Thus, velocity and wavelength of light are linked to each other as:

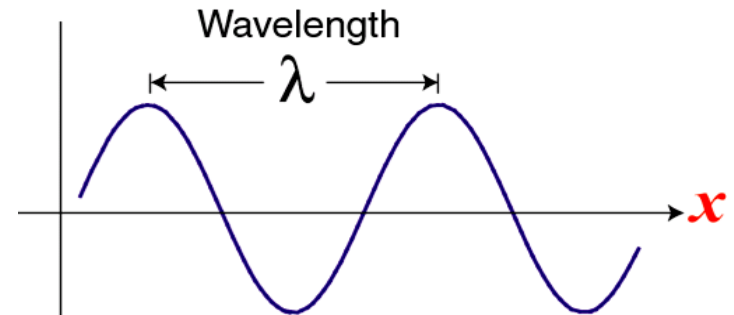
$$c = \nu \lambda_0$$

**Temporal** quantities:



The angular frequency:  $\omega = 2\pi/\tau$   
The frequency:  $\nu = 1/\tau$

**Spatial** quantities:



The k-vector:  $k = 2\pi/\lambda$   
The wave number:  $\kappa = 1/\lambda$

# Wave Propagation in Medium

- In free space:

$$c = \nu \lambda_0$$

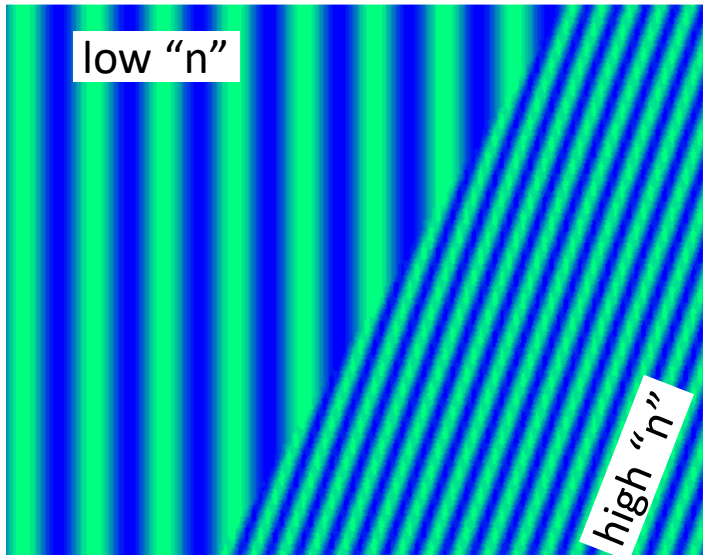
- When light enters into a medium (with index  $n$ ), its frequency (which is linked to energy) remains constant but its wavelength and speed ( $v$ ) change as follows:

$$v = \nu \lambda$$

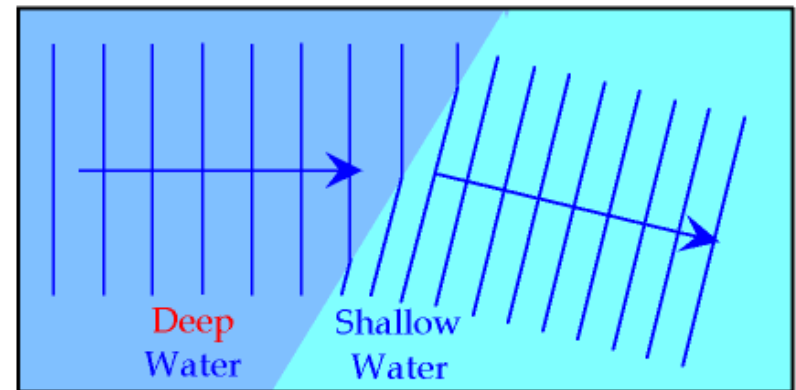
- Both wavelength and speed of light scale down linearly with the refractive index  $n$ :

$$v = c/n$$

$$\lambda = \lambda_0/n$$



Analogy with "water" waves



# Wave Propagation Through Objects

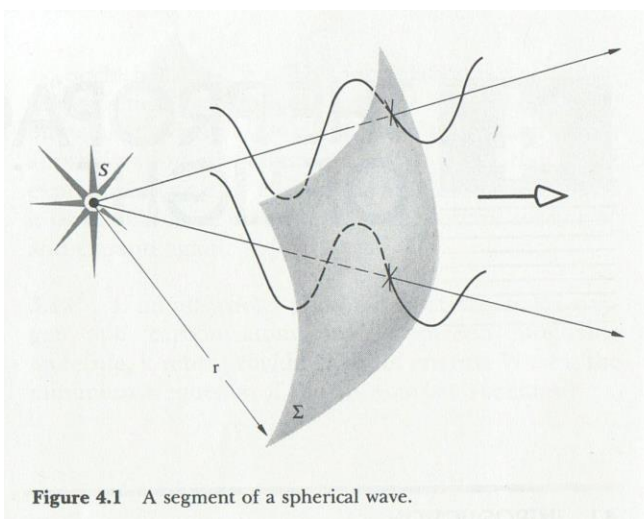


Figure 4.1 A segment of a spherical wave.

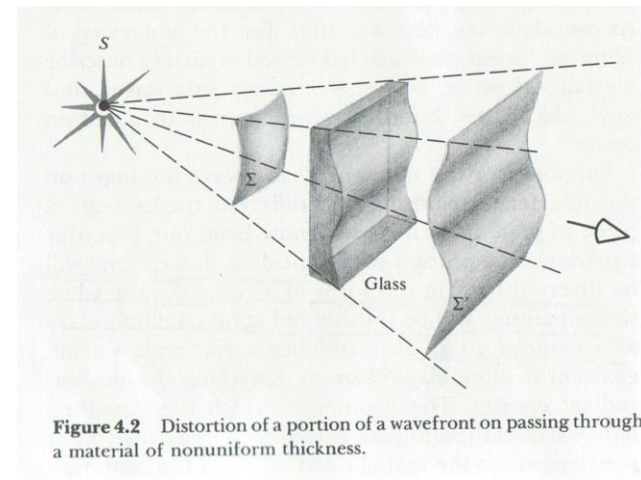
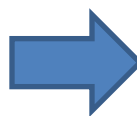


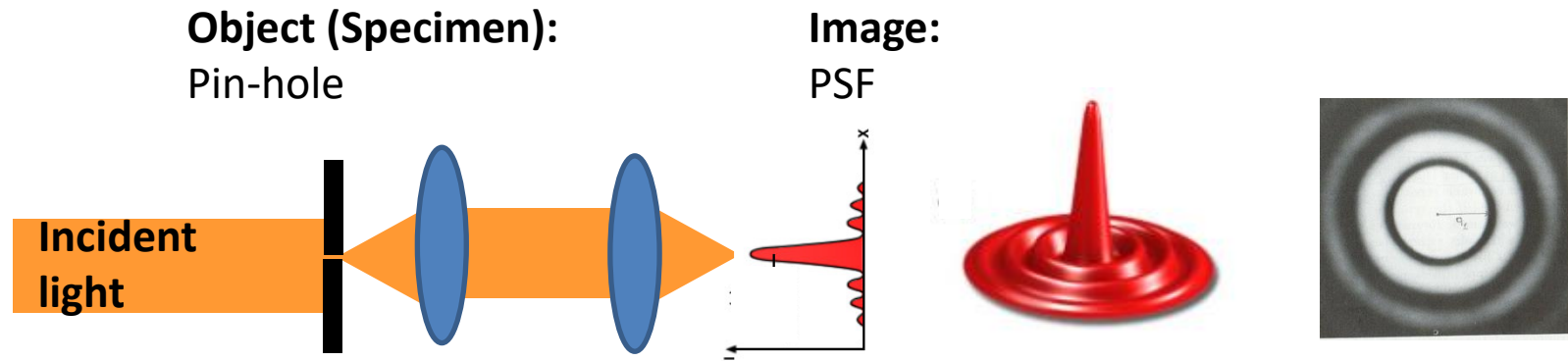
Figure 4.2 Distortion of a portion of a wavefront on passing through a material of nonuniform thickness.

consider a spherical wave interacts with an object (such as a lens, aperture, slit)



its wave fronts gets distorted with this interaction

# Wave phenomena & microscopy

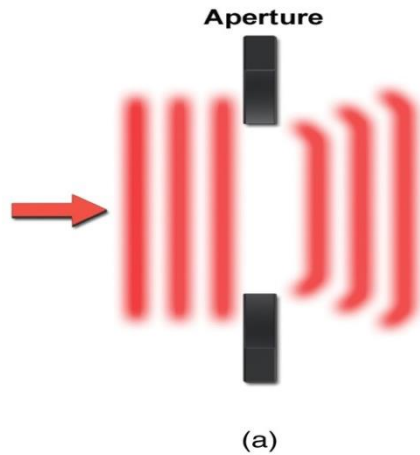


In a microscope:

- 1) Light from the illuminator is **diffracted** (i.e. spread) by the specimen (**object**)
- 2) Then, diffracted light is **collected** by the **objective lens**
- 3) And **focused** by the following optical components in the image plane where the propagating light waves **constructively & destructively interfere** to form the **image**

# Wave phenomena : diffraction

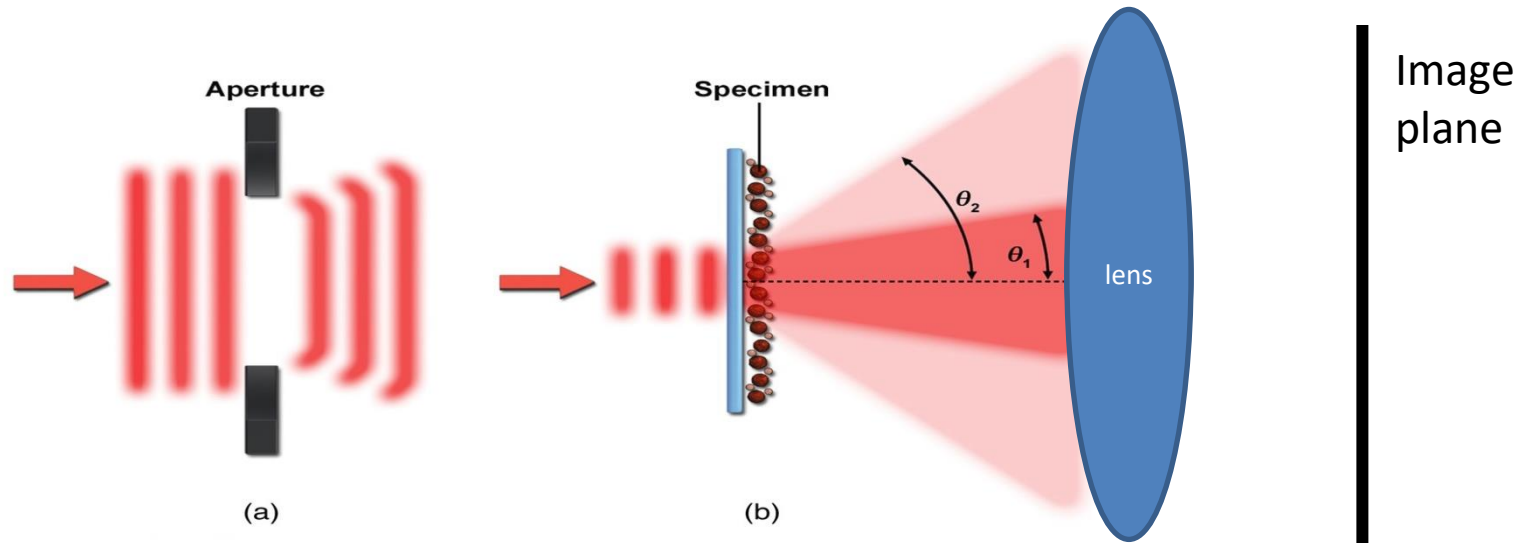
Depending on the specimen type, diffracted light can be perceived in different ways



**Example-1 (a):** When a beam of light is directed to an aperture, light appears to **bend around the edges**. The aperture could be the “stops” in the microscope (or the edges of the lenses)

# Wave phenomena : diffraction

Depending on the specimen type, diffracted light can be perceived in different ways



**Example-1 (a):** When a beam of light is directed to an aperture, light appears to **bend around the edges**. The aperture could be the “stops” in the microscope (or edges of the lenses)

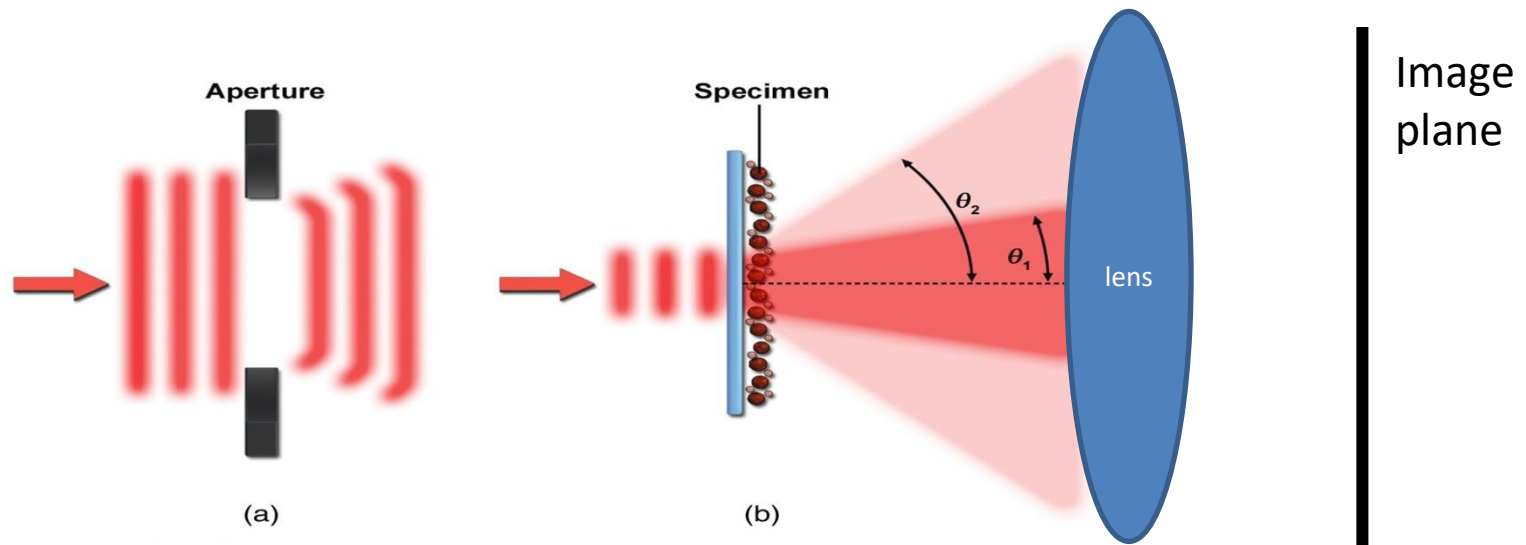
**Example-2 (b):** Diffraction also occurs when light illuminates a microscope slide covered with **small particles**. The amount of light **scattering** & the angle of spreading depend on the size & density of the diffracting particles on the slide.

In the right drawing, let's assume a mixture of 0.2  $\mu\text{m}$  & 2  $\mu\text{m}$  diameter particles as specimen.

The angle of spreading is inversely proportional to the particle size.

→ Larger angle ( $\theta_2$ ) corresponds to light diffraction by the smaller particles

# Wave phenomena : diffraction & interference



**Diffraction** - In microscopy, there are two primary sites for diffraction:

- 1- the specimen itself
- 2- the most limiting aperture of the microscope system

**Interference** can be seen as the combination of diffracted waves.

→ This is also the process responsible for creating images.

# Wave theory & wave phenomena (diffraction, interference ...)

- First deduced by Robert Hooke and mathematically formulated by Christiaan **Hyugens**
- Thomas **Young** demonstrated that wave theory best explained interference phenomenon.



Christiaan Hyugens  
1629-1695



Thomas Young  
1773-1829



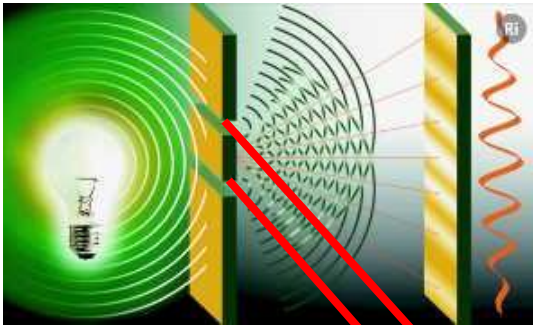
Augustin-Jean Fresnel  
1788-1827

- Augustin-Jean **Fresnel** synthesized the work of Hyugens and Young.
- Championed the wave theory in his 1818 memoir on diffraction.
- All this finally replaced by Newton's corpuscular (particle) theory .\*

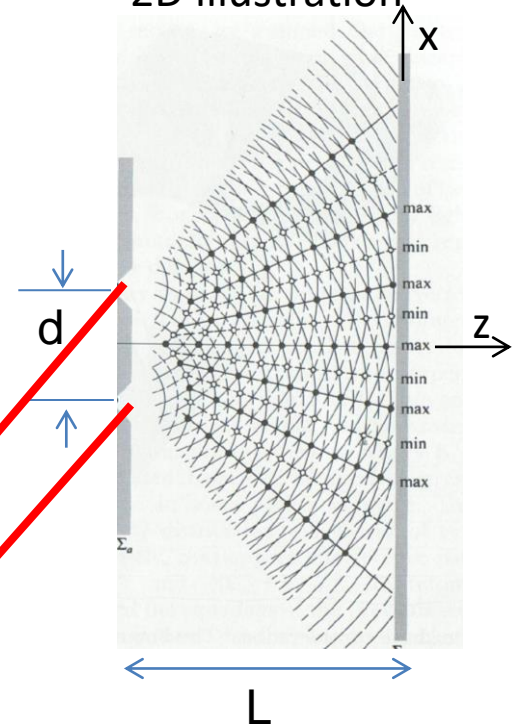
\* Until the 20<sup>th</sup> century (quantum era)

# Thomas Young's Double Slit Experiment

3D illustration



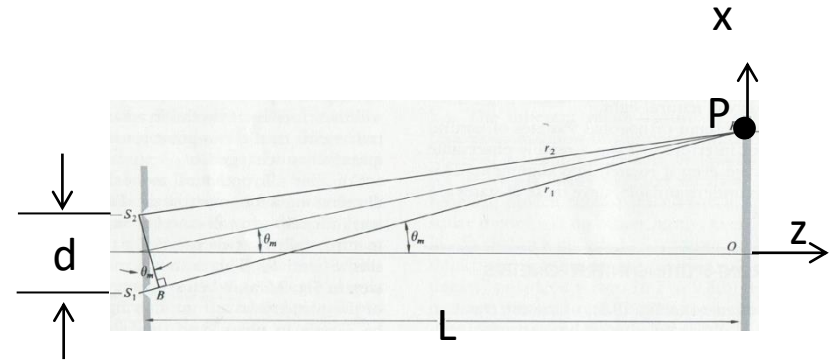
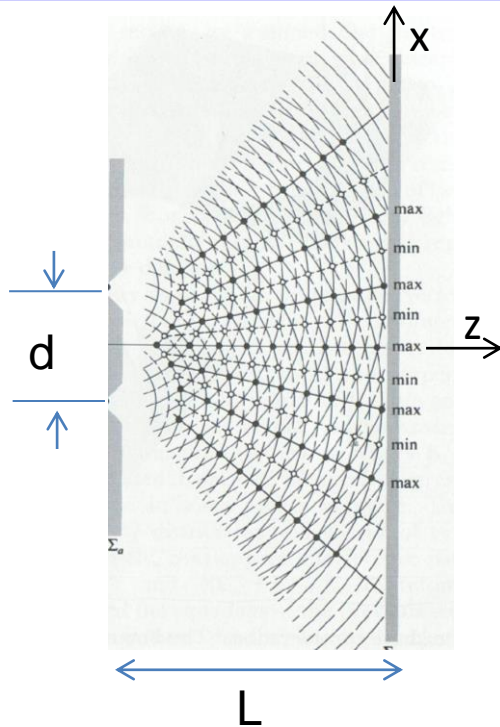
2D illustration



Consider having **two point sources** at the two slit positions (in cross-section view), as shown above.

- They are separated from each other by a distance  $d$  (right figure)
- The sources emit light waves with spherical wave-fronts
- These light waves (with spherical wave-fronts) interfere as they propagate in free space.

# Thomas Young's Double Slit Experiment



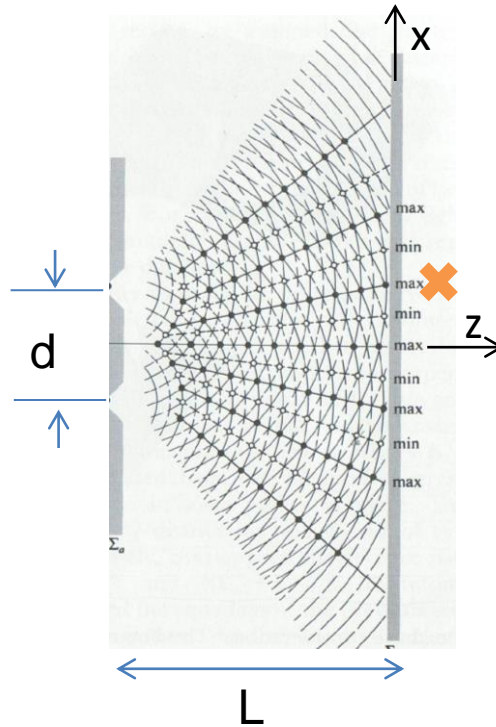
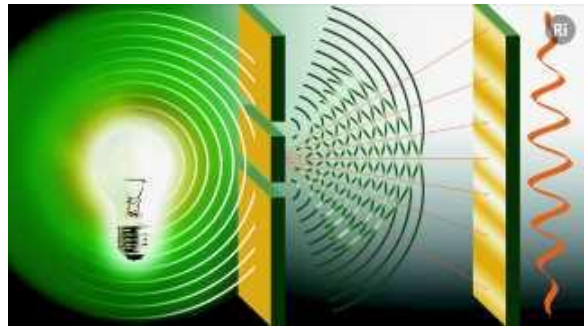
Consider having two point sources at the two slit positions, as shown above.

- Two sources emit waves with spherical wave-fronts.
- Let's put an observation screen "**L**" away from the double slit plane **at the z-axis**. Note that  $L \gg d, \lambda$
- **Along the x-axis** of the observation screen (lets take **point P**, on the right figure), there will be a **phase difference ( $\varphi$ )** between the two waves because each wave **travels over a different distance ( $r_1$  or  $r_2$ )** (see right figure).
- The interference of these waves with phase delay will led to an **interference pattern** on the observation screen
- The interference pattern will have **sinusoidal intensity fluctuations in the x-axis**.
- Intensity pattern is governed by the following interference equation:

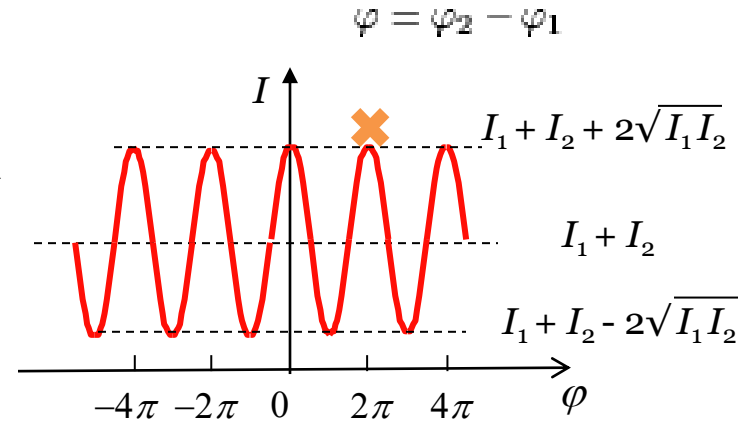
$$I = I_1 + I_2 + 2(I_1 I_2)^{1/2} \cos \varphi \quad \varphi = \varphi_2 - \varphi_1$$

- Amount of the **phase difference ( $\varphi$ )** dictates the final intensity at the observation plane.

# Thomas Young's Double Slit Experiment



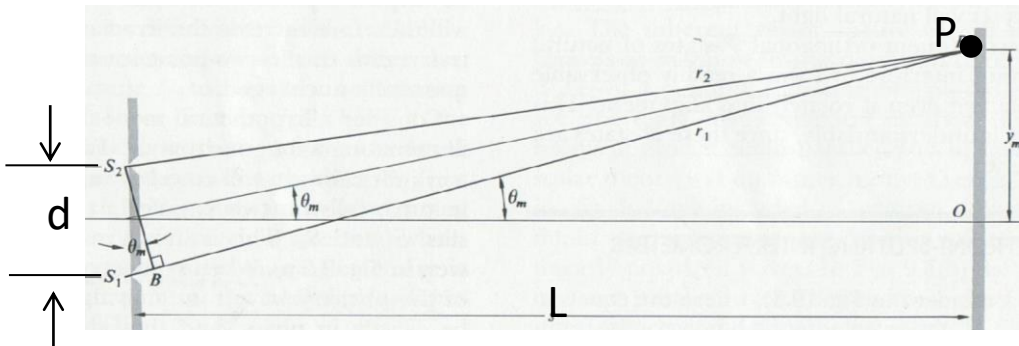
$$I = I_1 + I_2 + 2(I_1 I_2)^{1/2} \cos \varphi$$



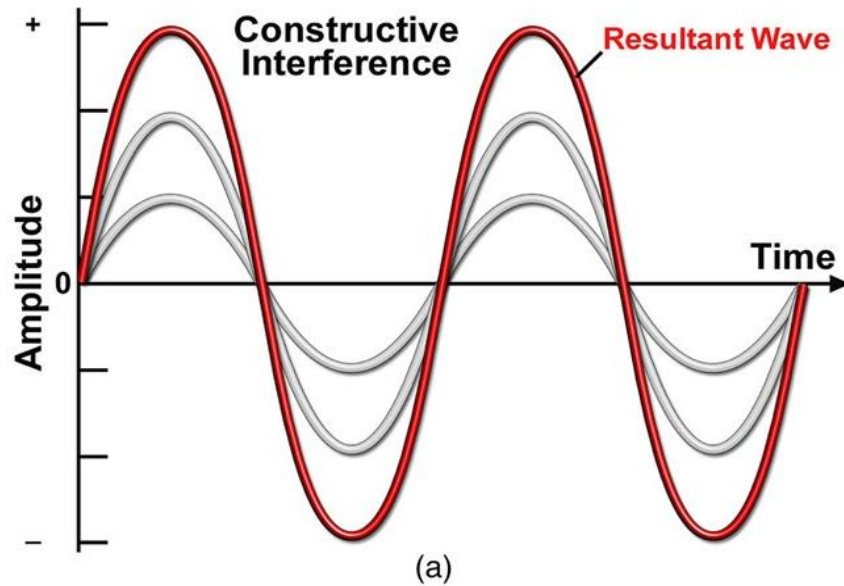
**1<sup>st</sup> max (m=+1):**

- Phase Difference ( $\varphi$ ) is  $+2\pi$
- Path Difference ( $S_1B$ ) is  $\lambda$

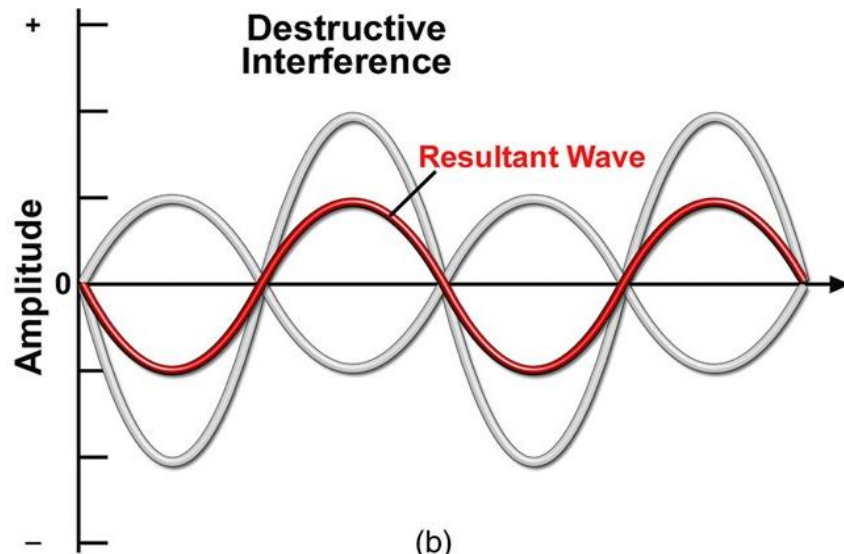
**→ Constructive Interference**



# Waves interference



(a)



(b)

Two waves are shown that oscillate in the plane of the page.

In these examples **(a)** and **(b)**:

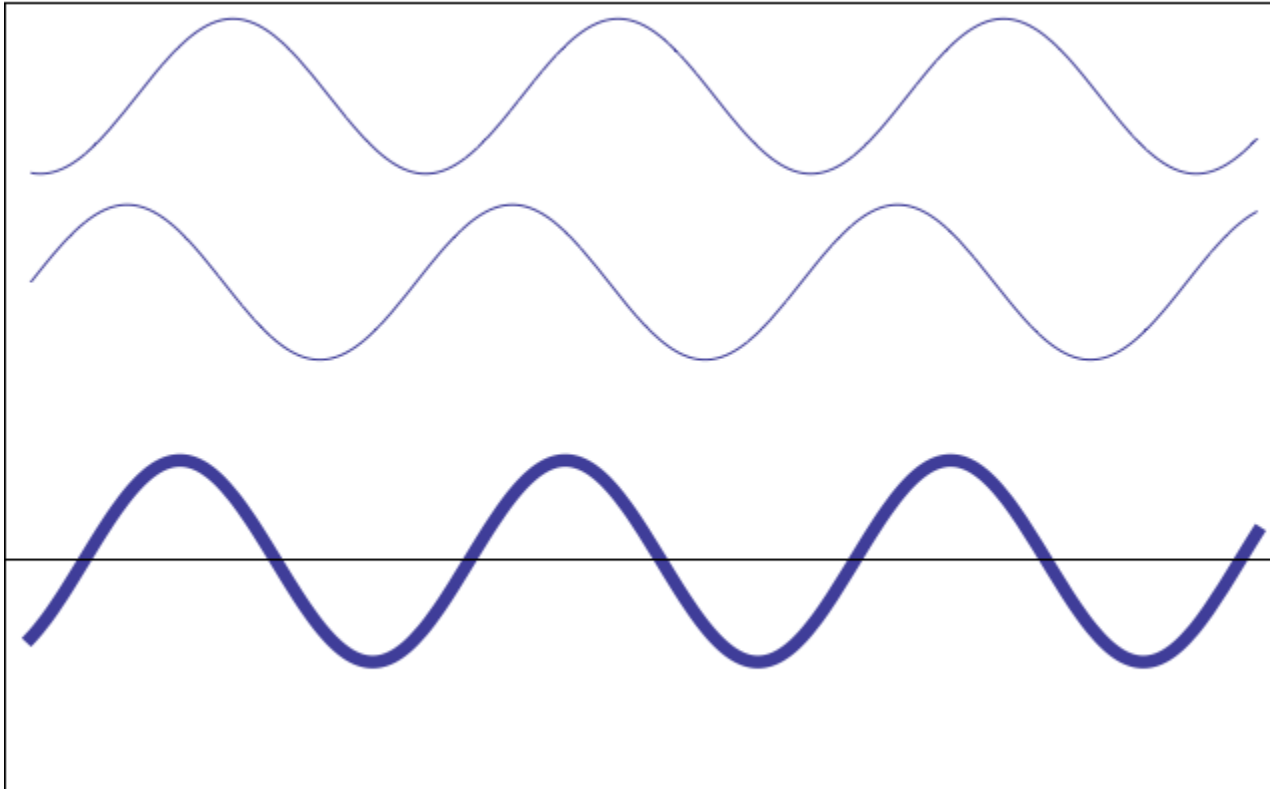
- The two waves (tinted gray) have the same wavelength, but vary in amplitude.
- The wave resulting from their interference is shown with red curve.

**(a) Constructive interference** occurs when these two waves have the same phase (a.k.a "in phase").

**(b) Destructive interference** occurs when these two waves are "out-of-phase".

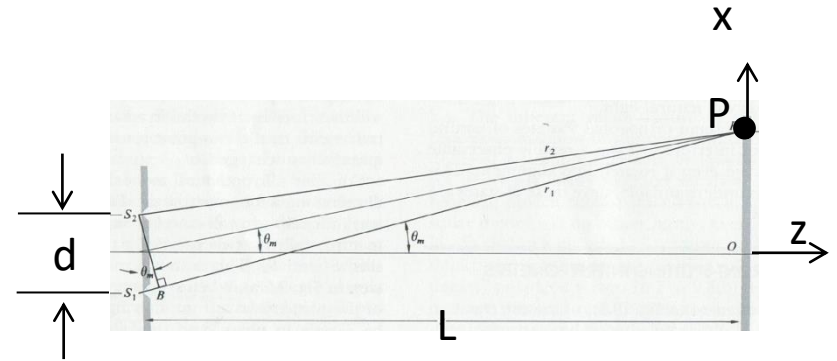
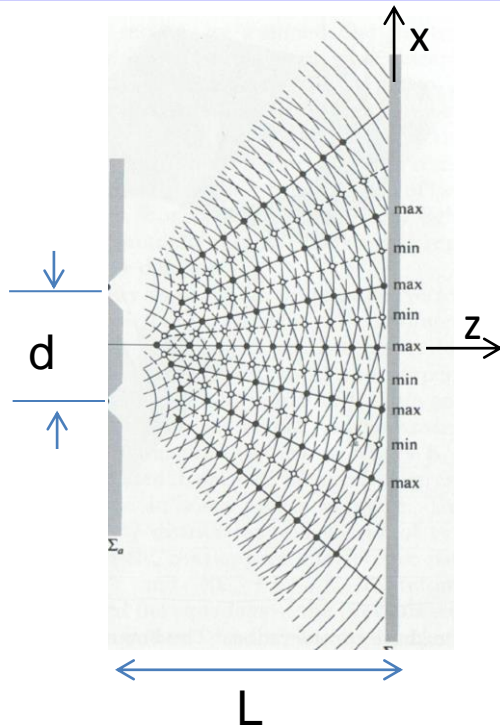
# Waves Interference

- Two waves (thin blue) are shown to oscillate in the plane of the page. They have same wavelength (frequency) and same amplitude.
- The wave resulting from their interference is shown with thick blue curve.



- **Constructive interference** occurs when these two waves have the same phase (a.k.a “in phase”). It leads to maximum intensity
- **Destructive interference** occurs when these two waves are “out-of-phase”. It leads to minimum intensity

# Thomas Young's Double Slit Experiment



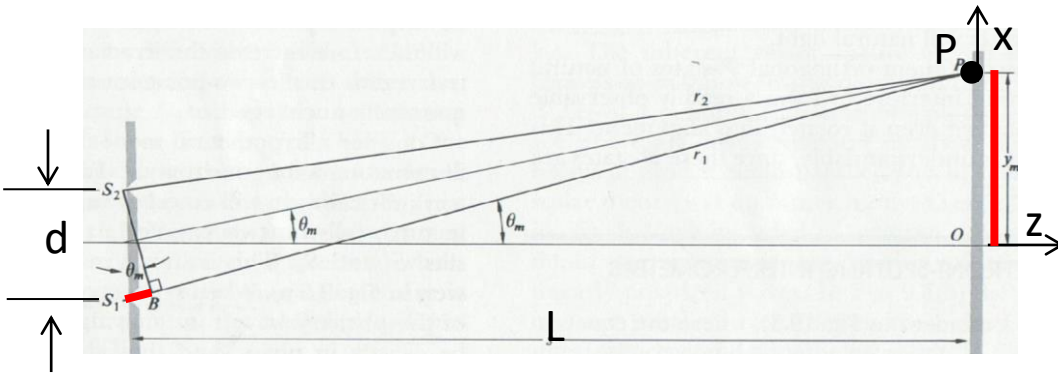
**For two point sources at the two slit positions & observation plane "L" away from the slits:**

- The interference of the waves with phase delays will lead to an **interference pattern** on the observation screen
- The interference pattern will have **sinusoidal intensity fluctuations in the x-axis**.
- Intensity pattern is governed by the following interference equation:

$$I = I_1 + I_2 + 2(I_1 I_2)^{1/2} \cos \varphi \quad \varphi = \varphi_2 - \varphi_1$$

- Amount of phase difference ( $\varphi$ ) dictates the final intensity at the observation plane.
- **Max** intensity points corresponds to **constructive** interference, where  $\varphi = 0, 2\pi, 4\pi, 6\pi \dots$  in the above equation
- **Min** intensity points corresponds to **destructive** interference, where  $\varphi = \pi, 3\pi, 5\pi, 7\pi \dots$  in the above equation

# Thomas Young's Double Slit Experiment



$$\text{Phase} \rightarrow \varphi_1 = kr_1 = \frac{2\pi}{\lambda} r_1$$

$$\varphi_2 = kr_2 = \frac{2\pi}{\lambda} r_2$$

Phase difference  $\rightarrow \varphi = \varphi_2 - \varphi_1$

Path difference  $\rightarrow |S_1B| = r_2 - r_1$

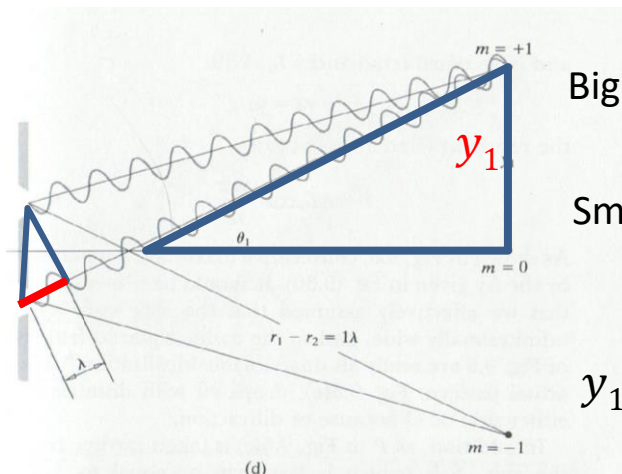
Location of the observation  $\rightarrow y_m$

For 1<sup>st</sup> max ( $y_1$ ):

- Constructive Interference
- Path difference is  $1 \times \lambda$

For m<sup>th</sup> max ( $y_m$ ):

- Path difference is  $m \times \lambda$
- Constructive Interference



Big triangle  $\rightarrow \sin\theta_1 \sim \frac{y_1}{L}$

Small triangle  $\rightarrow \sin\theta_1 = \frac{\lambda}{d}$

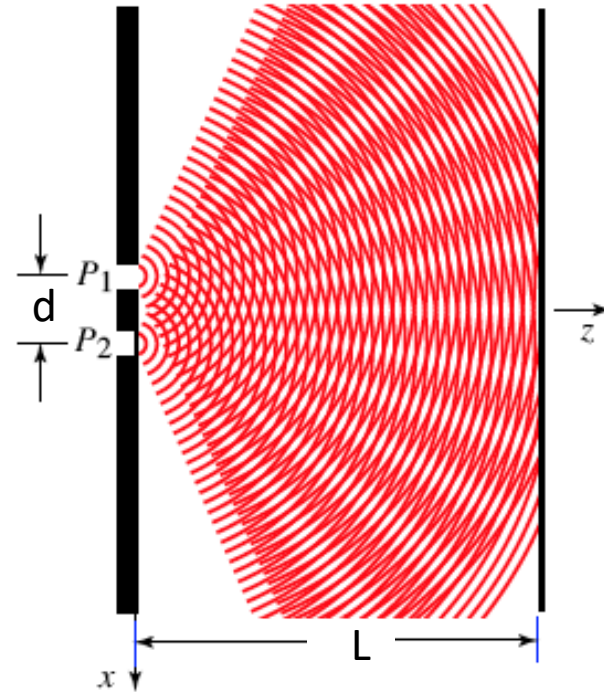
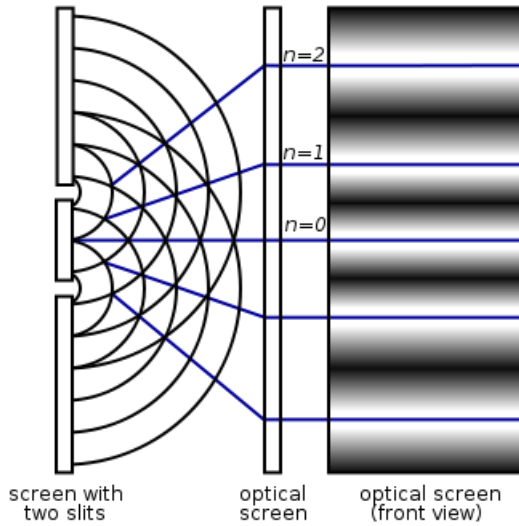
$$y_1 = \frac{\lambda L}{d} \iff \lambda = \frac{y_1 d}{L}$$

generalize

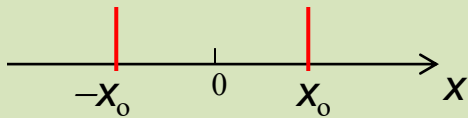
$$y_m = m \frac{\lambda L}{d}$$

$$\lambda = \frac{y_m d}{m \cdot L}$$

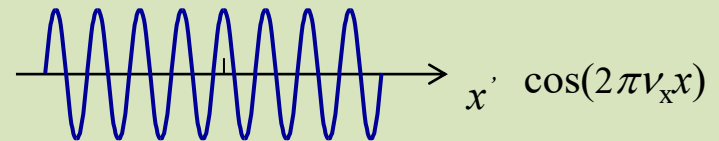
# Double Slit (with no width) Experiment



INPUT:



OUTPUT:



PS: Ignore the slit width