

The background image is a composite of two scenes. On the left, a traditional wooden windmill stands on a grassy bank next to a body of water, with high-voltage power lines and pylons in the background. On the right, a large industrial power plant is visible, featuring a prominent cooling tower and various structures. A semi-transparent red rectangular box is overlaid on the right side of the image, containing the course title in white text.

ME-251: Thermodynamics and energetics I

First Law for Open Systems

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2025 Fall Semester

Photo Credit: Trougnouf

- Estimate thermodynamic properties in special cases

- Incompressible substance
- Ideal gas

○ Reading: Moran 3.10-3.15

- Which of the following statements describes a supercritical fluid?
 - A. A substance at a temperature and pressure **both** below its critical point, exhibiting distinct liquid and gas phases.
 - B. A substance at a temperature above its critical point temperature **or** at a pressure above its critical pressure, where liquid and gas phases are indistinguishable.
 - C. substance at a temperature and pressure **both** above its critical point, where liquid and gas phases are indistinguishable.
 - D. A substance at its triple point, where solid, liquid, and gas phases coexist in equilibrium.

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luepfl

Additional rationale on $du = c_v dT$ for ideal gas expansion or compression

- For ideal gas, u is only a function of temperature
- For any process between state 1 and state 2 with varying specific volume $v_1 \neq v_2$, consider a process that brings state 1 to state 2', in which $v_1 = v_2'$ and $T_2 = T_2'$
- Note that for ideal gas, $u_2 = u_2'$, and $u_2 - u_1 = u_2' - u_1$
- The internal energy change does not depend on whether the volume is held constant

- Liquids and solids are often modeled as incompressible $v = \text{const}$

$$u = u(T) \quad c_v = \frac{du}{dT} = c(T) = c_p \quad \text{At constant } T, dh = vdp$$

- $\frac{pv}{RT} = \frac{p\tilde{v}}{\tilde{R}T} = Z \equiv 1$, the dilute limit of real gases

T : temperature [K]

\tilde{R} : universal gas constant [J/K/mol]

R : specific gas constant [J/K/kg]

\tilde{v} : molar volume [m³/mol]

$$u = u(T)$$

$$h = u + pv = u(T) + RT$$

$$c_v = \frac{du}{dT} = c(T)$$

$$c_p = c_v(T) + R$$

$$\text{Specific heat ratio } k = \frac{c_p}{c_v} = 1 + \frac{R}{c_v}$$

Additional rationale on $du = c_v dT$ for ideal gas expansion or compression

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- The internal energy change does not depend on whether the volume is held constant

Follows ideal gas law

$$c_v = \text{const}, \quad c_p = c_v + R$$

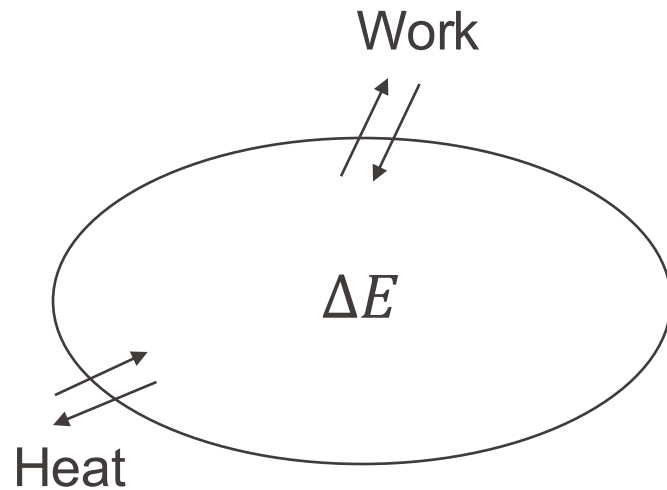
Common approximations:

Monoatomic $c_v = 1.5R$

Diatomic $c_v = 2.5 R$

Polyatomic: $c_v = 3R$

- Understand the difference between the analyses of closed and open systems
- Apply mass balance to open systems
- Apply energy balance to open systems
- Reading: Moran 4.1 - 4.5



Can be intuitively understood as energy conservation

Energy change = Energy going in – Energy going out

Energy flows can take the form of heat, work, ...

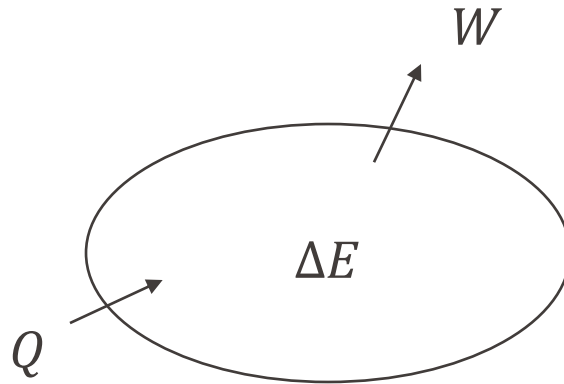
Change in the amount of energy contained within a system during some time interval

=

Net amount of energy transferred **into** the system by heat transfer during the time interval

–

Net amount of energy transferred **out of** the system by work during the time interval



$$\Delta E = E_2 - E_1 = Q - W$$

Everything is in [J]

Time Rate Form of First Law for Closed Systems

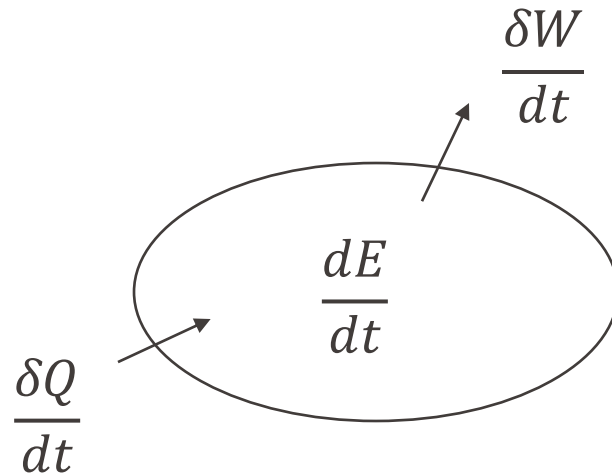
Time rate of change in the amount of energy contained within a system at time t

=

Net rate at which energy is transferred **into** the system by heat transfer at time t

-

Net rate at which energy is transferred **out of** the system by work at time t



$$\frac{dE}{dt} = \dot{Q} - \dot{W} \quad \text{Heat flow rate - Power}$$

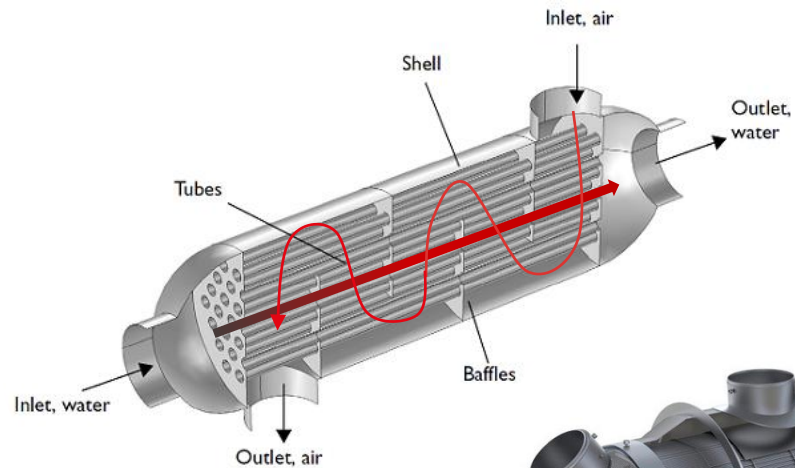
Everything is in [W]

- Type of work transfer
 - Expansion/compression (e.g., polytropic processes)
 - Electrical, spring, shaft,...

- Type of energy
 - Potential energy
 - Kinetic energy
 - Internal energy (thermal motion, latent, chemical, nuclear,...)

- Type of heat transfer (requires temperature gradient)
 - Conduction, convection, radiation

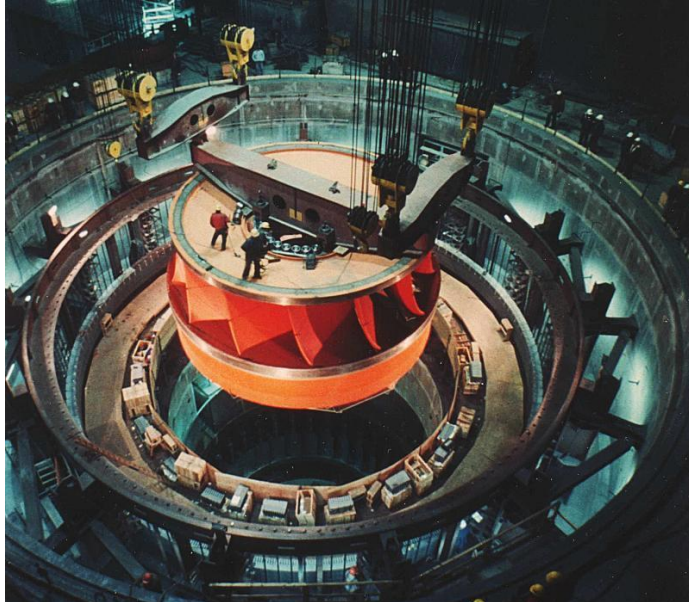
- Open systems: systems with mass transport across the boundary
 - Heat exchangers



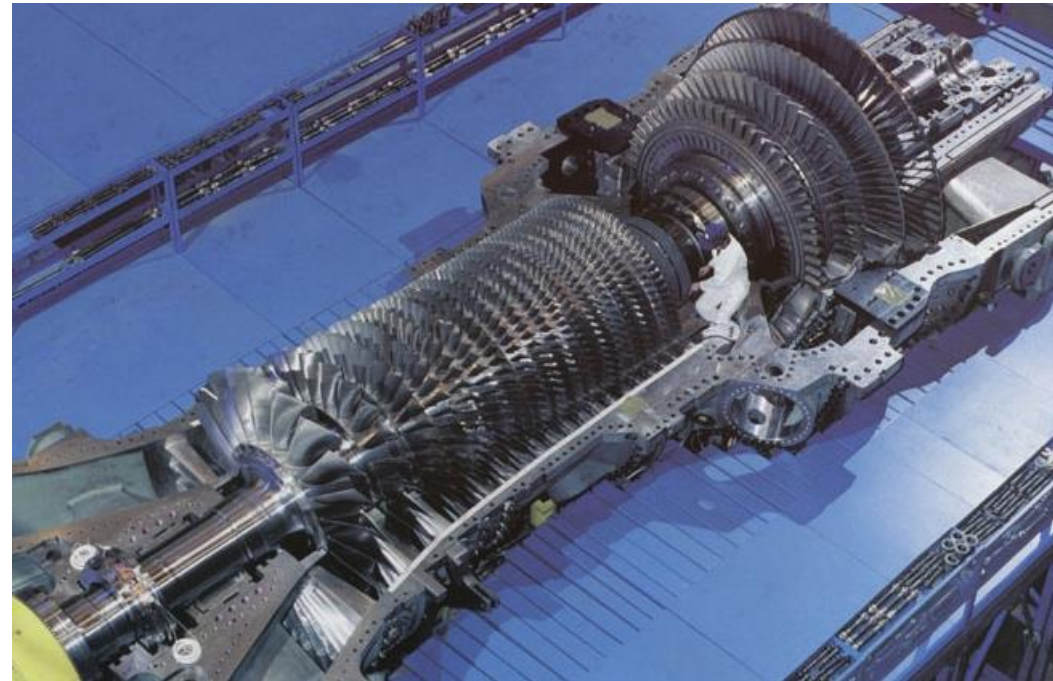
Alstom, Feedwater heater for nuclear power plant, 500 MW



- E.g., turbines

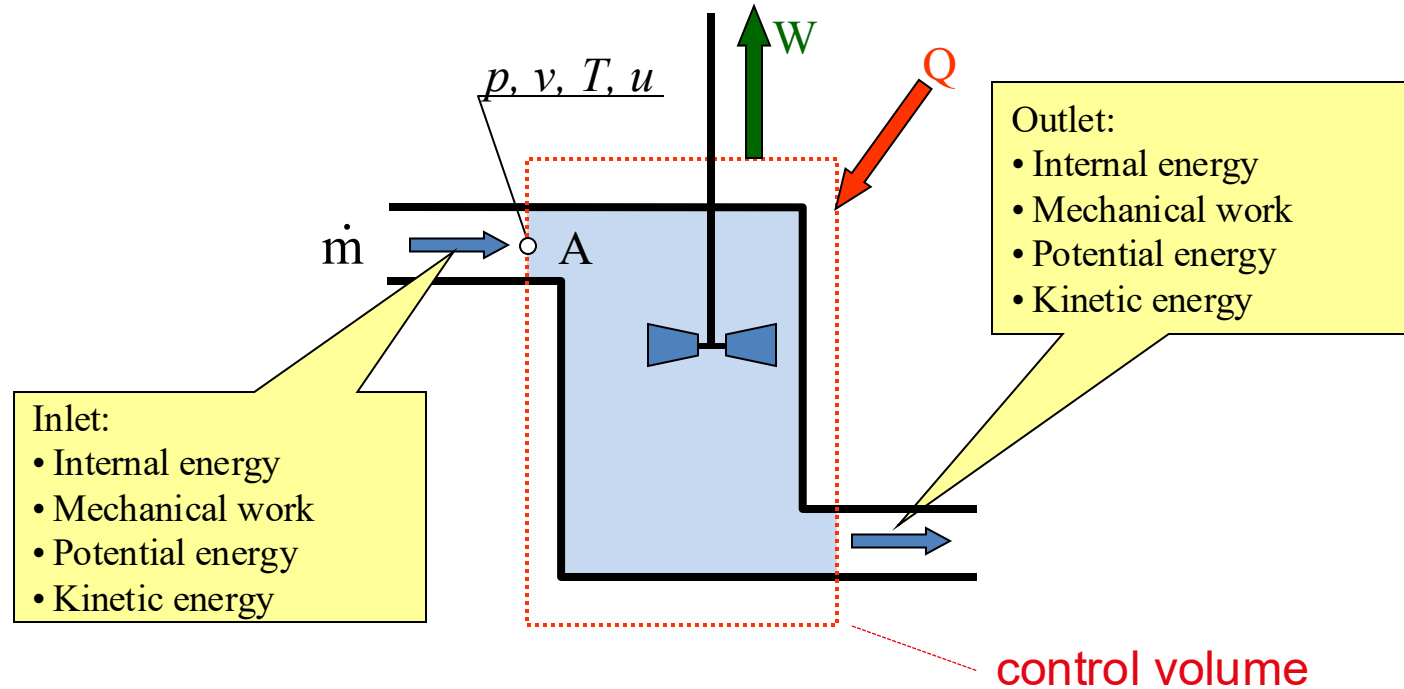


Francis turbine, Grand Coulee Dam,
biggest hydro plant in the US,
capacity 6800 MW, 33 turbines (different sizes)



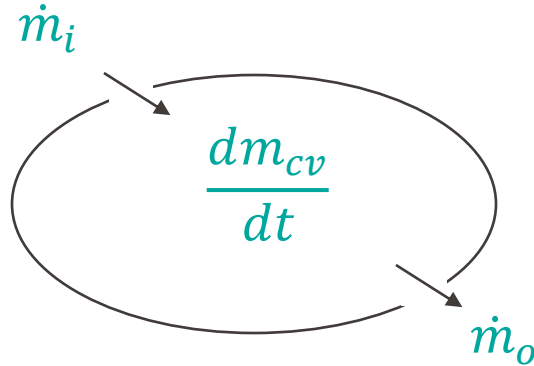
GE 9H gas turbine, 480 MW

Many real-world systems deal with flowing working fluid: turbines, engines, compressors, heat exchangers,...



Unlike closed systems, energy transfer can be due to incoming and outgoing mass flow

With opens systems, we consider both mass conservation and energy conservation



$$\frac{dm_{cv}}{dt} = \sum_i \dot{m}_i - \sum_o \dot{m}_o \quad [\text{kg/s}]$$

In general, there could be several inlets and outlets

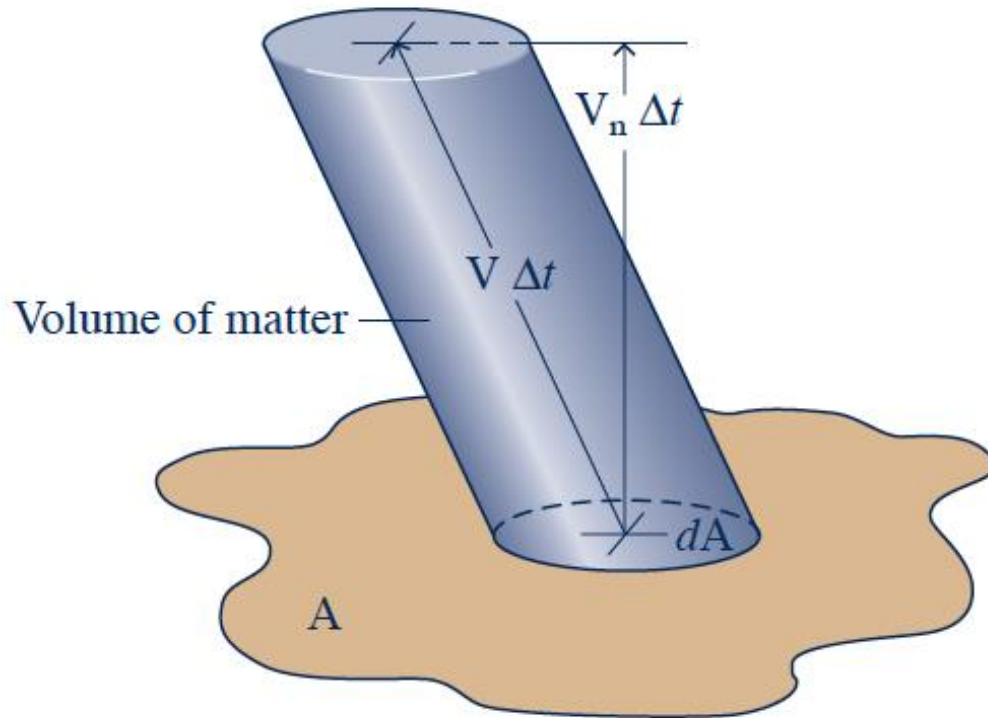
Time rate of change
in the amount of mass
contained within a
control volume at time t

=

Rate at which mass
is transferred **into** the
system at time t

–

Rate at which mass
is transferred **out of**
the system at time t



- Consider a small quantity of matter flowing with velocity V across an incremental area dA in a small time interval Δt

$$\left[\begin{array}{l} \text{amount of mass} \\ \text{crossing } dA \text{ during} \\ \text{time interval } \Delta t \end{array} \right] = \rho(V_n \Delta t) dA$$

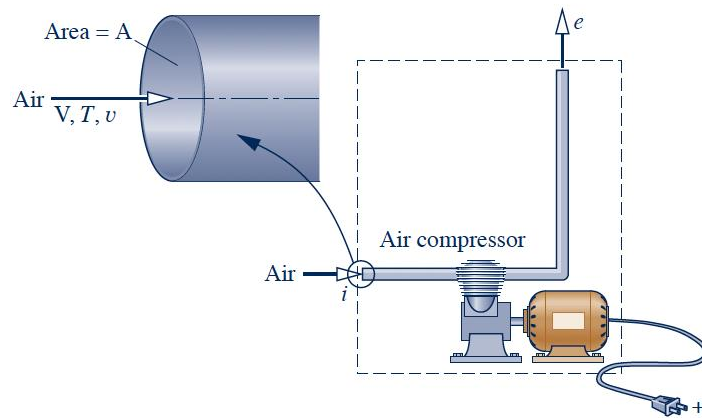
$$\left[\begin{array}{l} \text{instantaneous rate} \\ \text{of mass flow} \\ \text{across } dA \end{array} \right] = \rho V_n dA$$

$$\dot{m} = \int_A \rho V_n dA$$

In Moran, non-italic V represents velocity

- 1D form of the mass rate balance

- The flow is normal to the boundary at locations where mass enters or exits the control volume ($V = V_n$)
- Velocity and density are uniform over the specified inlet or exit area



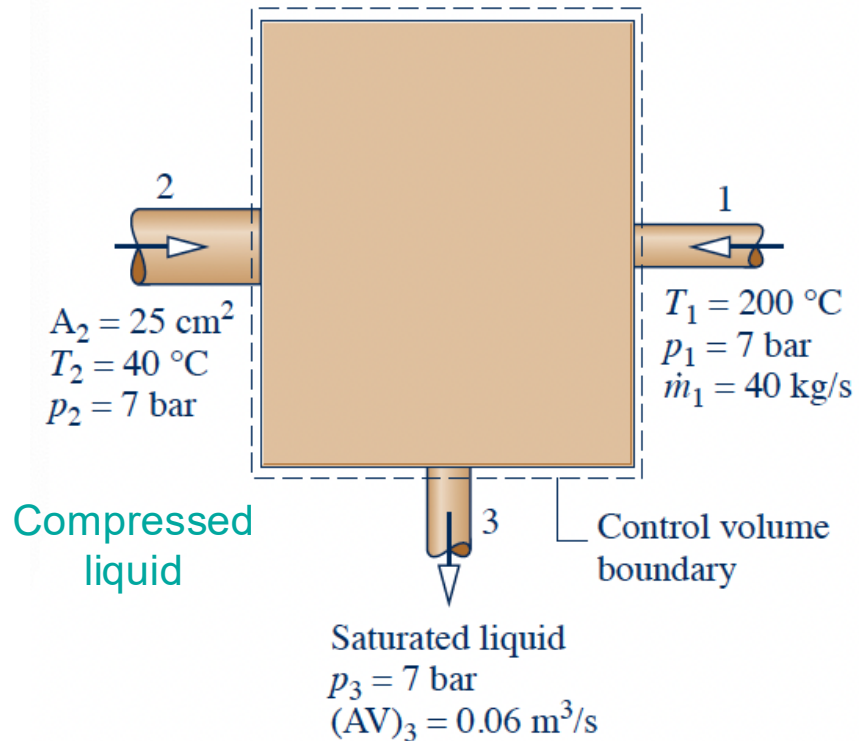
$$\dot{m} = \rho AV$$

$$\frac{dm_{cv}}{dt} = \sum_i \rho_i A_i V_i - \sum_o \rho_o A_o V_o$$

- Steady state: all properties in the control volume are invariant over time

$$\sum_i \dot{m}_i = \sum_o \dot{m}_o$$

- Steady mass in a control volume does not necessarily mean that the system is at a steady state (other properties can change)



A feedwater heater operating at steady state

Determine the mass flow rates at inlet 2 and at the exit, and the velocity V_2

$$v_3 = 1.108 \times 10^{-3} \text{ m}^3/\text{kg}$$

Approximate liquid water as incompressible substance

$$\frac{dm_{cv}}{dt} = \dot{m}_1 + \dot{m}_2 - \dot{m}_3$$

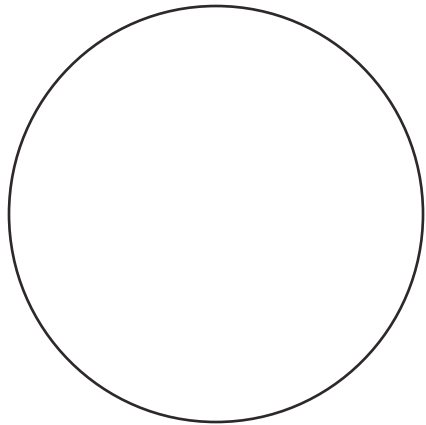
$$\dot{m}_2 = \dot{m}_1 - \dot{m}_3$$

$$\dot{m}_3 = \frac{(AV)_3}{v_3}$$

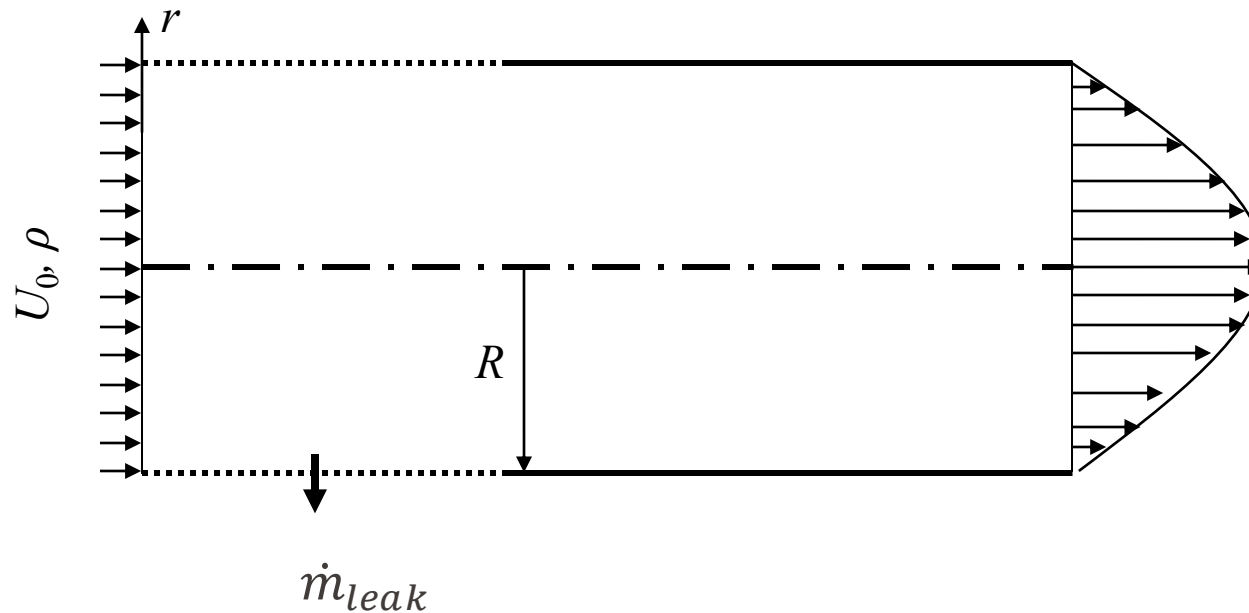
$$V_2 = \dot{m}_2 v_2 / A_2$$

$$v_2 = v_3$$

Example: steady-state pipe flow with steady leakage
(inlet and outlet velocities normal to the surface)

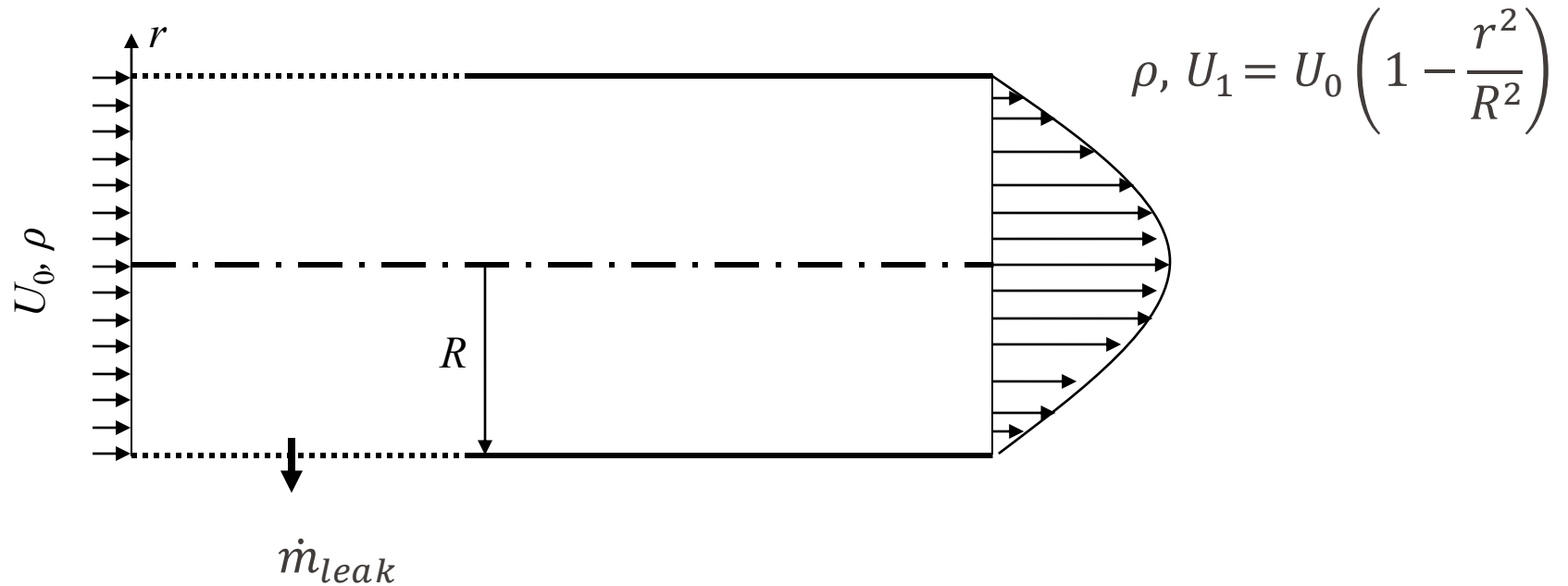
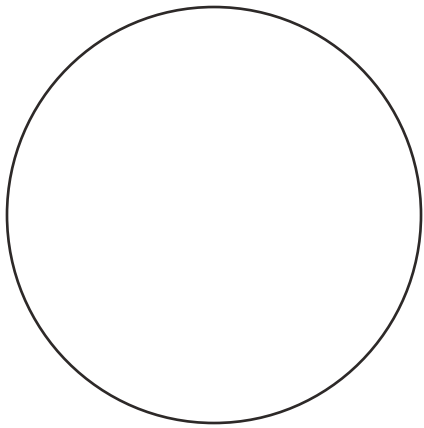


Cross-section



$$\rho, U_1 = U_0 \left(1 - \frac{r^2}{R^2} \right)$$

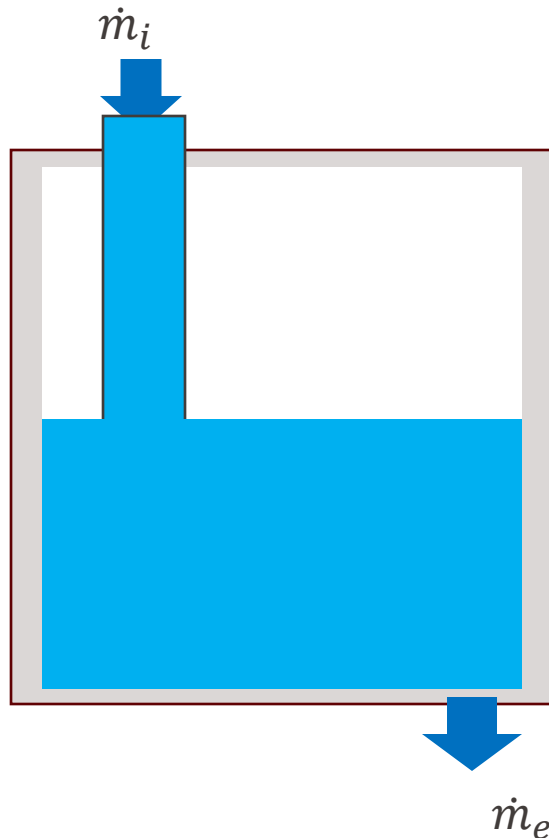
Calculate the mass flow rate leaking through the perforated pipe wall



$$\rho U_0 \pi R^2 = \dot{m}_{leak} + \int_0^R \rho U_1 2\pi r dr$$

$$\dot{m}_{leak} = \rho U_0 \pi R^2 - \int_0^R \rho U_1 2\pi r dr = \rho U_0 \pi R^2 - \frac{1}{2} \rho U_0 \pi R^2 = \frac{\pi}{2} \rho U_0 R^2$$

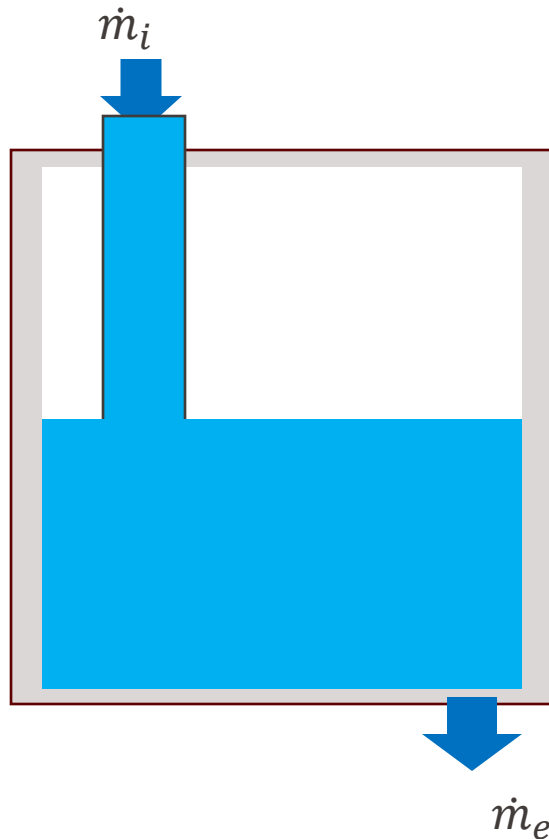
Example 4.2 in Moran (Transient Mass Balance)



- Water enters and exits an initially empty barrel.
 - Inlet flow rate is $\dot{m}_i = 7 \text{ kg/s}$
 - Outlet flow rate is proportional to the height of liquid inside $\dot{m}_e = 1.4L$, where L is the instantaneous liquid height (SI units)
 - Base area of the barrel is $A = 0.2 \text{ m}^2$ and density of water is $\rho = 1000 \text{ kg/m}^3$

Plot the variation of liquid height L as a function of time t

Example 4.2 in Moran (Transient Mass Balance)



$$\frac{dm_{cv}}{dt} = \dot{m}_i - \dot{m}_e$$

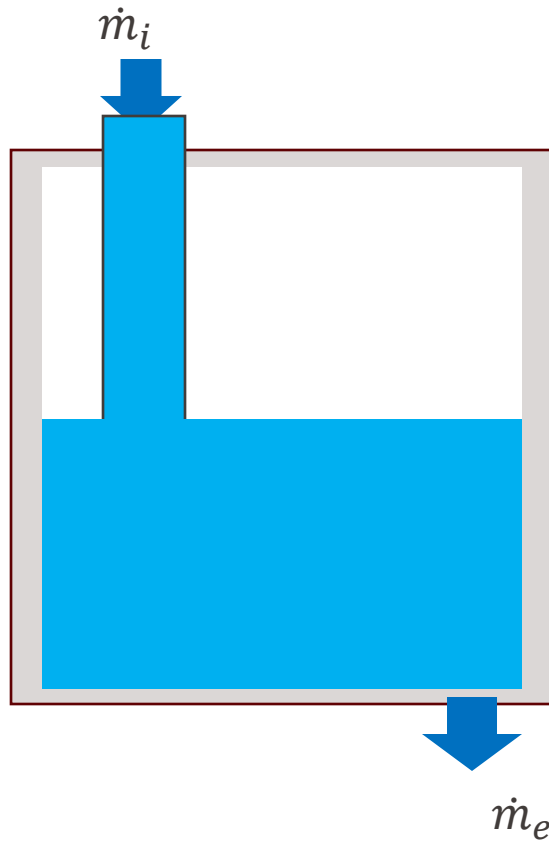
$$\frac{d(\rho AL)}{dt} = 7 - 1.4L \quad \text{SI units } \rho A = 200 \text{ [kg/m]}$$

$$\frac{dL}{dt} + 0.007(L - 5) = 0$$

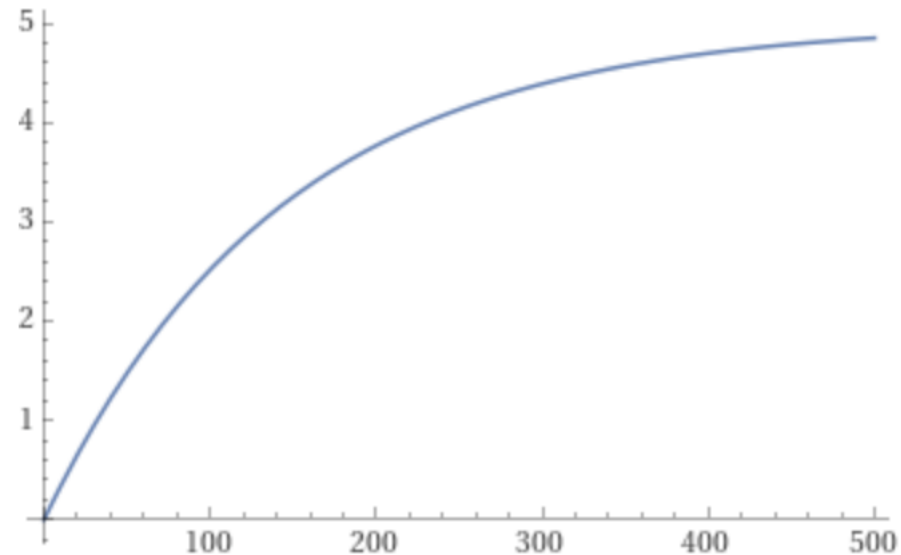
$$L = 5 + C \exp(-0.007t)$$

$$t = 0, L = 0 \Rightarrow C = -5 \quad L = 5[1 - \exp(-0.007t)]$$

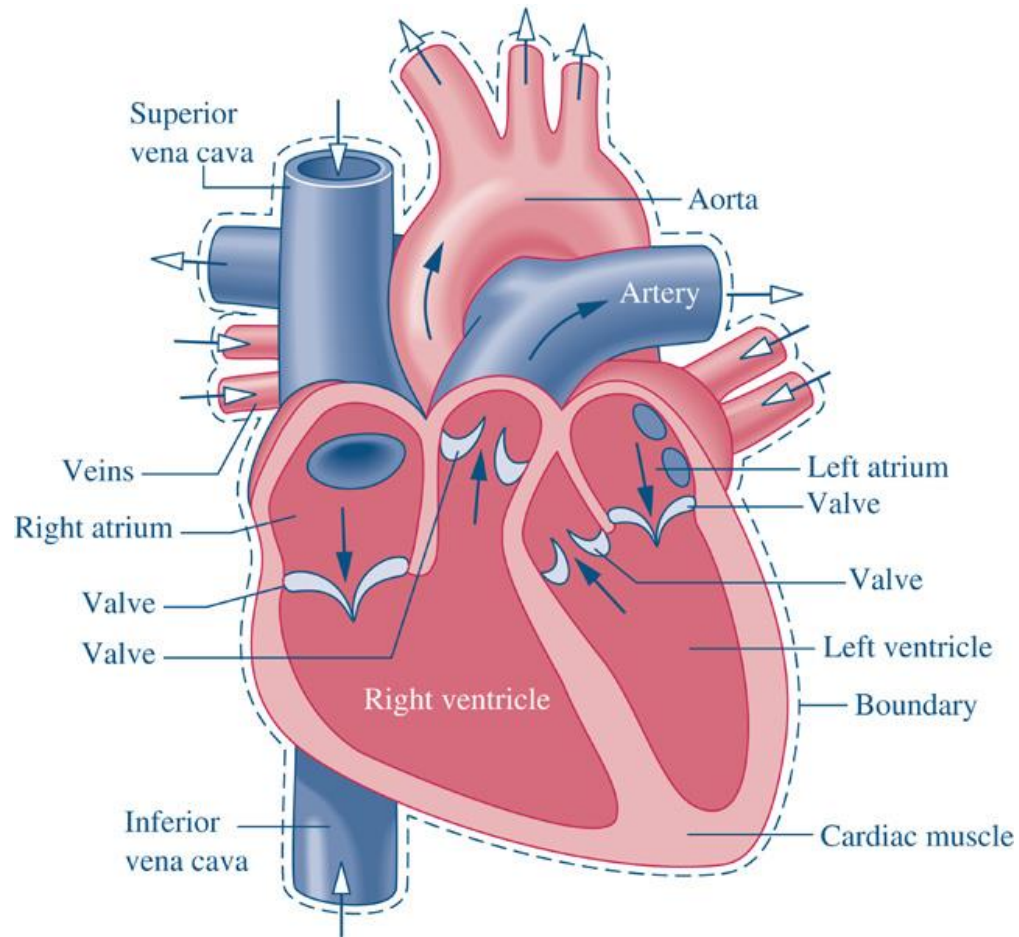
Example 4.2 in Moran (Transient Mass Balance)



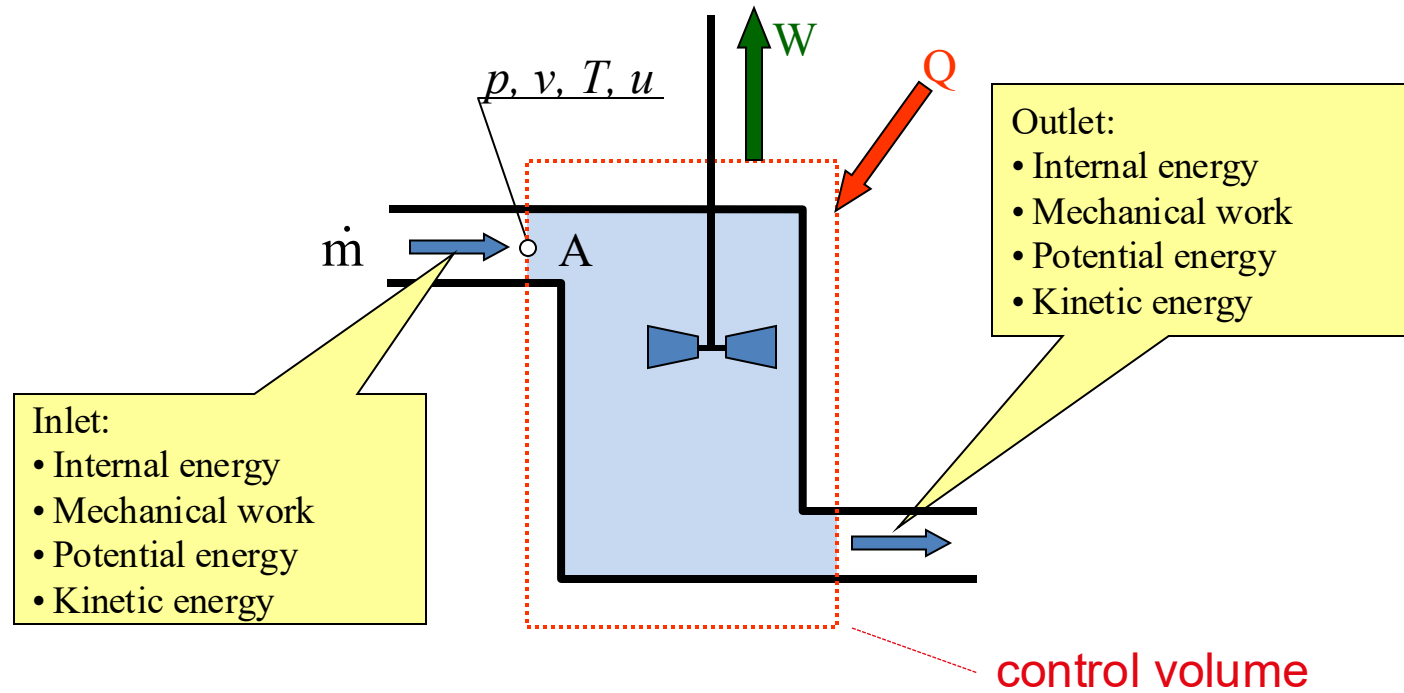
$$L = 5[1 - \exp(-0.007t)]$$



Filling height saturated around 5 m



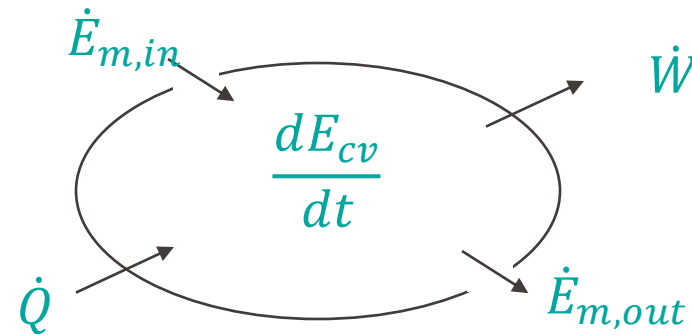
- Flow in a human heart is controlled by valves that intermittently allow blood to enter and exit as the heart muscles pump
- The boundary of this control volume is moving as the heart pulses



Unlike closed systems, energy transfer can be due to incoming and outgoing mass flow

- Energy rate balance

Time rate of change in the amount of energy contained within a control volume at time t



=

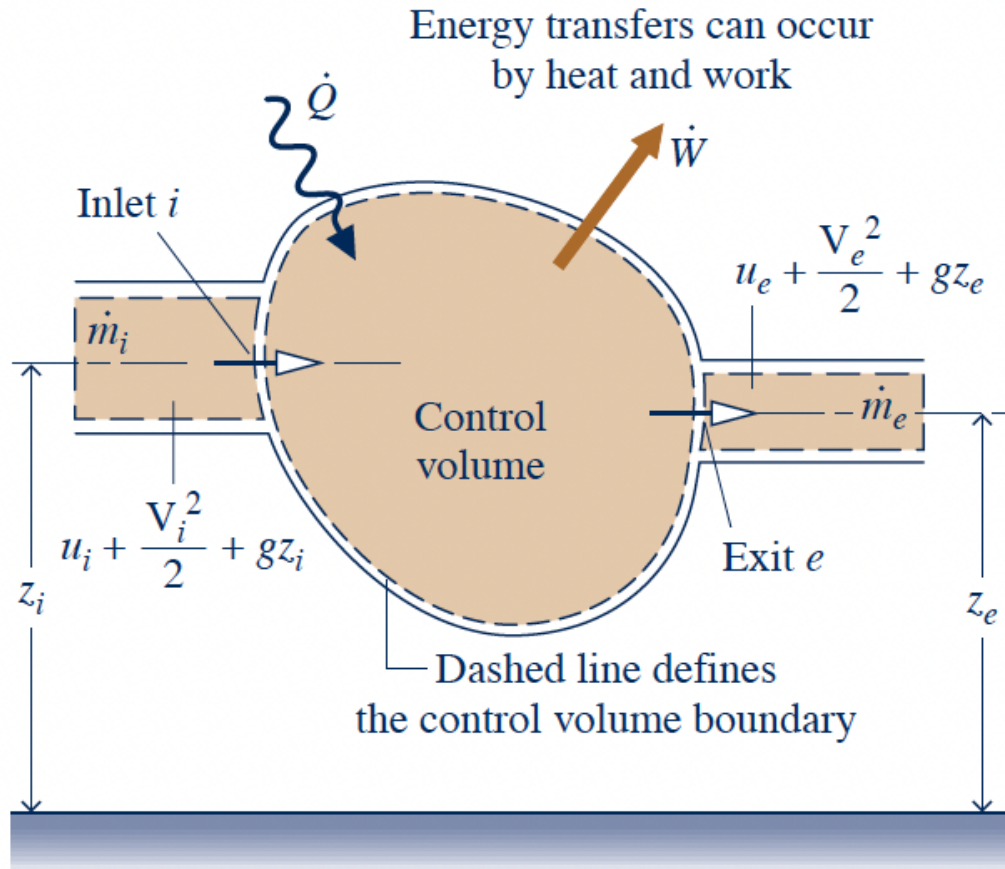
Net rate at which energy is transferred into the system by heat transfer at time t

-

Net rate at which energy is transferred out of the system by work at time t

+

Net rate at which energy is transferred into the system by accompanying **mass flow**



$$\frac{dE_{cv}}{dt} = \dot{Q} - \dot{W} + \dot{m}_i \left(u_i + \frac{V_i^2}{2} + gz_i \right) - \dot{m}_e \left(u_e + \frac{V_e^2}{2} + gz_e \right)$$

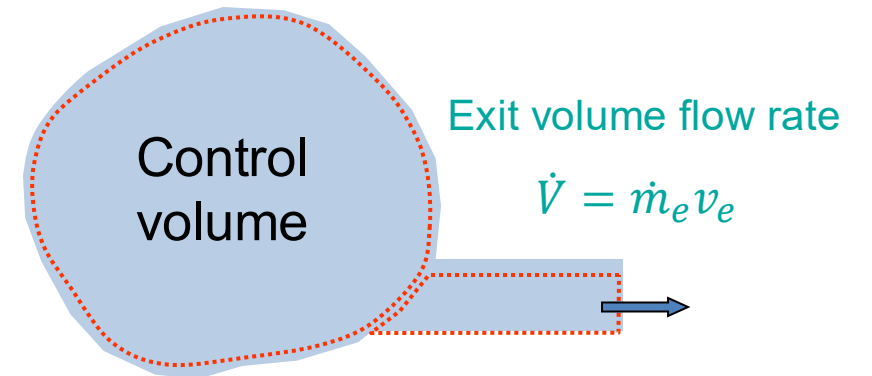
- Separate \dot{W} into two contributions
 - Work associated with fluid pressure as mass is introduced at inlets and removed at exits (**rate of flow work**)
 - All other works \dot{W}_{cv}

Rate of energy transfer by work **out of** control volume at exit

$$= p_e \dot{V} = p_e \dot{m}_e v_e$$

Rate of energy transfer by work **into** control volume at inlet

$$= \dot{m}_i (p_i v_i)$$



Total work $\dot{W} = \dot{W}_{cv} + \dot{m}_e (p_e v_e) - \dot{m}_i (p_i v_i)$

$$\frac{dE_{cv}}{dt} = \dot{Q} - \dot{W} + \dot{m}_i \left(u_i + \frac{V_i^2}{2} + gz_i \right) - \dot{m}_e \left(u_e + \frac{V_e^2}{2} + gz_e \right) \quad \text{One-inlet, one-outlet}$$

$$\frac{dE_{cv}}{dt} = \dot{Q}_{cv} - \dot{W}_{cv} + \dot{m}_i \left(u_i + p_i v_i + \frac{V_i^2}{2} + gz_i \right) - \dot{m}_e \left(u_e + p_e v_e + \frac{V_e^2}{2} + gz_e \right)$$

Recall enthalpy is defined as $h = u + pv$

$$\frac{dE_{cv}}{dt} = \dot{Q}_{cv} - \dot{W}_{cv} + \dot{m}_i \left(h_i + \frac{V_i^2}{2} + gz_i \right) - \dot{m}_e \left(h_e + \frac{V_e^2}{2} + gz_e \right)$$

More than one inlet and one outlet

$$\frac{dE_{cv}}{dt} = \dot{Q}_{cv} - \dot{W}_{cv} + \sum_i \dot{m}_i \left(h_i + \frac{V_i^2}{2} + gz_i \right) - \sum_e \dot{m}_e \left(h_e + \frac{V_e^2}{2} + gz_e \right)$$

$$E_{cv} = \int_{\mathcal{V}} \rho \left(u + \frac{\vec{V}^2}{2} + gz \right) d\mathcal{V}$$

$$\begin{aligned} \frac{d}{dt} \int_{\mathcal{V}} \rho \left(u + \frac{\vec{V}^2}{2} + gz \right) d\mathcal{V} = & \dot{Q}_{cv} - \dot{W}_{cv} + \sum_i \left[\int_A \left(h + \frac{\vec{V}^2}{2} + gz \right) \rho \vec{V} \cdot d\vec{A} \right]_i \\ & - \sum_e \left[\int_A \left(h + \frac{\vec{V}^2}{2} + gz \right) \rho \vec{V} \cdot d\vec{A} \right]_e \end{aligned}$$

Compare this to
$$\frac{dE_{cv}}{dt} = \dot{Q}_{cv} - \dot{W}_{cv} + \sum_i \dot{m}_i \left(h_i + \frac{\vec{V}_i^2}{2} + gz_i \right) - \sum_e \dot{m}_e \left(h_e + \frac{\vec{V}_e^2}{2} + gz_e \right)$$

$$\frac{dm_{cv}}{dt} = \sum_i \dot{m}_i - \sum_o \dot{m}_o$$

$$\frac{dE_{cv}}{dt} = \dot{Q}_{cv} - \dot{W}_{cv} + \sum_i \dot{m}_i \left(h_i + \frac{\vec{V}_i^2}{2} + gz_i \right) - \sum_e \dot{m}_e \left(h_e + \frac{\vec{V}_e^2}{2} + gz_e \right)$$

At steady state this reduces to

$$0 = \sum_i \dot{m}_i - \sum_o \dot{m}_o$$

$$0 = \dot{Q}_{cv} - \dot{W}_{cv} + \sum_i \dot{m}_i \left(h_i + \frac{\vec{V}_i^2}{2} + gz_i \right) - \sum_e \dot{m}_e \left(h_e + \frac{\vec{V}_e^2}{2} + gz_e \right)$$