



ME-251: Thermodynamics and energetics I

Thermodynamic Properties II

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- Enthalpy

$$H = U + pV$$

$$h = u + pv$$

$$\tilde{h} = \tilde{u} + p\tilde{v}$$

- Specific heat

$$c_v = \left(\frac{\delta Q}{\partial T} \right) \Big|_v = \left(\frac{\partial u}{\partial T} \right) \Big|_v$$

$$c_p = \left(\frac{\delta Q}{\partial T} \right) \Big|_p = \left(\frac{\partial h}{\partial T} \right) \Big|_p$$

- Evaluating thermodynamic properties with CoolProp

- Estimate thermodynamic properties in special cases

- Incompressible substance
- Ideal gas

○ Reading: Moran 3.10-3.15

- Liquids and solids are often modeled as incompressible $v = \text{const}$

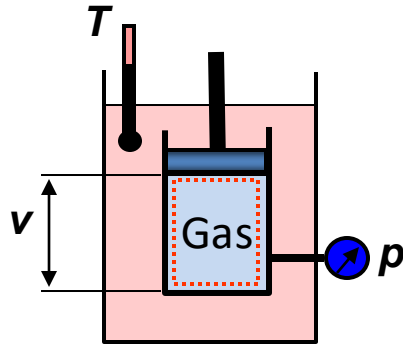
In general u is a function of v and $T \Rightarrow u = u(T)$ for incompressible substances

$$\Rightarrow c_v = \left(\frac{\partial u}{\partial T} \right) \Big|_v = \frac{du}{dT} = c_v(T) \text{ for incompressible substances}$$

$$\Rightarrow c_p = \left(\frac{\partial h}{\partial T} \right) \Big|_p = \left(\frac{\partial u + \partial pv}{\partial T} \right) \Big|_p = \frac{du}{dT} + p \frac{\partial v}{\partial T} = c_v \text{ ignoring thermal expansion}$$

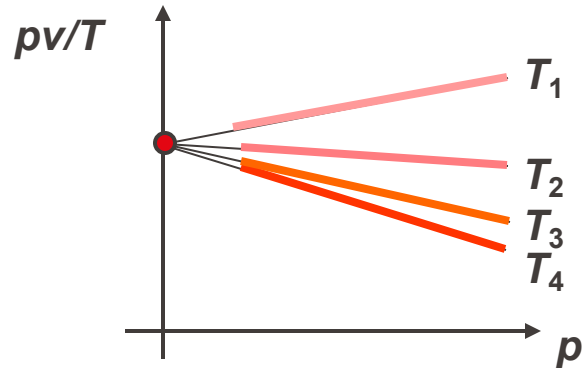
$$v = \text{const} \Rightarrow c_v = c_p = c(T)$$

$$h = u(T) + pv \Rightarrow \text{at constant } T, dh = d(pv) = v dp$$



- Consider gas confined in a piston-cylinder assembly and held at a constant temperature.
- Gas pressure is varied through piston movement
- At a constant temperature, the pressure and specific volume are measured.

$$\lim_{p \rightarrow 0} \frac{pv}{T} = R = \frac{\tilde{R}}{M}$$



M : molar mass [kg/mol]

T : temperature [K]

\tilde{R} : universal gas constant

R : specific gas constant

- $\tilde{R} = 8.314 \text{ J/K/mol}$
- \tilde{R} is the same for different gases, the specific gas constant R is different for different gases
- $\tilde{R} = k_B N_A$
 - The Boltzmann constant $k_B = 1.38 \times 10^{-23} \text{ [J/K]}$, converts temperature scale to thermal energy scale for individual particles
 - The Avogadro constant $N_A = 6.022 \times 10^{23} \text{ [1/mol]}$, historically defined to make one mole of hydrogen atoms weigh 1 gram
 - \tilde{R} converts temperature scale to thermal energy scale for one mole of particles

What is the SI unit for specific gas constant R

- A. [J]
- B. [J/K]
- C. [J/K/mol]
- D. None of above

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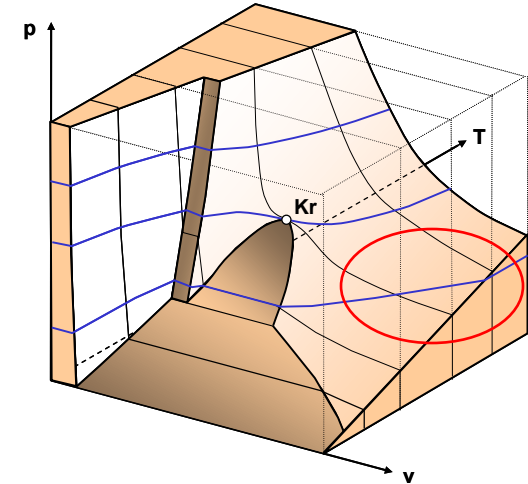
- $Z = \frac{pv}{RT} = \frac{p\tilde{v}}{\tilde{R}T}$, $\tilde{v} = vM = V/n$ molar volume
- $\lim_{p \rightarrow 0} Z = 1$
- In practice, Z is a function of temperature, pressure/specific volume, and gas species

$$Z = 1 + \frac{B(T)}{\tilde{v}} + \frac{C(T)}{\tilde{v}^2} + \frac{D(T)}{\tilde{v}^3} + \dots$$

$$Z = 1 + \hat{B}(T)p + \hat{C}(T)p^2 + \hat{D}(T)p^3 + \dots$$

Viral coefficients B , C , D (or \hat{B} , \hat{C} , \hat{D}) account for intermolecular interactions and can be derived from statistical mechanics

- Ideal gas model: $\frac{pv}{RT} = \frac{p\tilde{v}}{\tilde{R}T} = Z \equiv 1$, the dilute limit of real gases
 - No intermolecular interactions (other than elastic collision)
 - Volume of molecules ignored
- When we approximate real gases as ideal gas
 - In p-v diagram, $p \propto \frac{1}{v}$ for isotherms
 - In T-v diagram, $T \propto v$ for isobars



- Given a gas species, $u = u(T)$ (intermolecular interactions insignificant)
- $h = u + pv = u(T) + RT$, also just a function of temperature
- $c_v = \left(\frac{\partial u}{\partial T}\right) \Big|_v = \frac{du}{dT} = c_v(T)$
- $c_p = \left(\frac{\partial h}{\partial T}\right) \Big|_p = \frac{du}{dT} + R = c_v(T) + R$
- Specific heat ratio $k = \frac{c_p}{c_v} = 1 + \frac{R}{c_v}$
 - $c_v = \frac{R}{k-1}$; $c_p = \frac{kR}{k-1}$

- A process where $pv^n = \text{const}$
 - Isobaric process: $p = \text{const} \Rightarrow n = 0$
 - Isothermal process: $pv = RT = \text{const} \Rightarrow n = 1$
 - Isochoric processes: $v = \text{const} \Rightarrow n = +\infty$ ($p^{1/n}v = \text{const}$)
 - Isentropic process (not covered yet): $pv^k = \text{const} \Rightarrow n = k = c_p/c_v$

- Compression/expansion work (intensive form)

$$w = \int_1^2 p dv = \begin{cases} p_1 v_1 \ln\left(\frac{v_2}{v_1}\right), n = 1 \\ \frac{p_2 v_2 - p_1 v_1}{1 - n}, n \neq 1 \end{cases}$$

- Applying ideal gas law

$$w = \int_1^2 p dv = \begin{cases} RT \ln\left(\frac{v_2}{v_1}\right), n = 1, \text{ isothermal} \\ \frac{R(T_2 - T_1)}{1 - n}, n \neq 1 \end{cases}$$

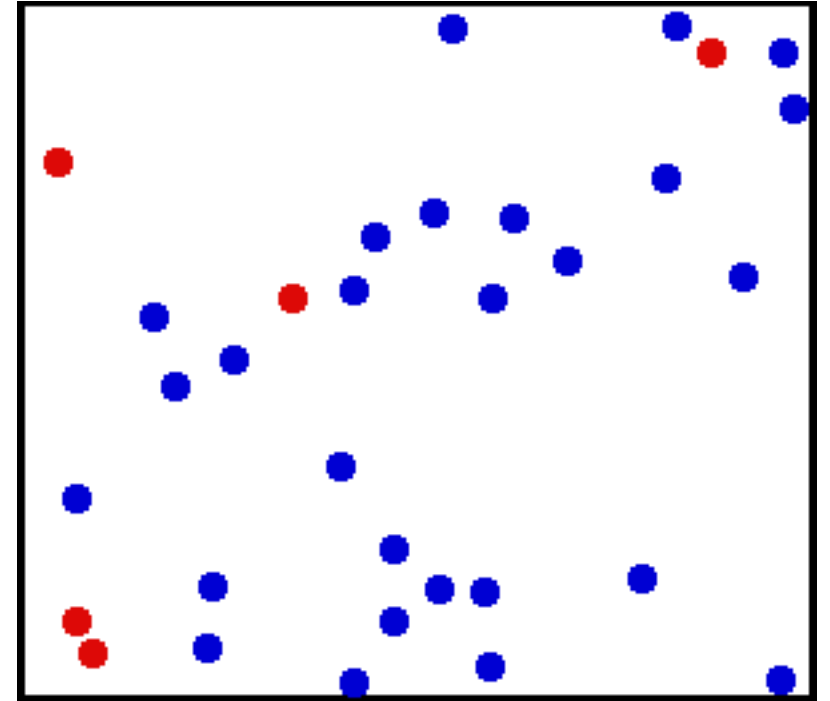
- Special case of ideal gas (also called calorically perfect)
 - Like an ideal gas and what's more, the specific heats are approximated as **independent of temperature**
 - $c_v = \text{const}$, $c_p = c_v + R$

Monoatomic $c_v = 1.5R$

Diatomic $c_v = 2.5 R$

Polyatomic: $c_v = 3R$

- Consider gas as a **large number of randomly moving particles** that collide with one another every now and then
- Collisions are elastic: **kinetic energy is conserved** before and after
- Between collisions, particles are not affected by any force field



Credit: A. Greg

- The kinetic theory defines temperature by considering the average kinetic energy of a particle in a system
- The average translational kinetic energy $e_{kt} = \frac{3}{2}kT$

- If an ideal gas system only contains monatomic molecules (Ar, He,..), one only needs to consider translational kinetic energy

$$U = \frac{3}{2} N_{sys} kT$$

- For monatomic ideal gases

$$c_v = \frac{1}{m_{sys}} \frac{dU}{dT} = \frac{3}{2} \frac{N_{sys}}{Mn} \frac{\tilde{R}}{N_A} = \frac{3}{2} R$$

- For diatomic and polyatomic molecules, internal energy can also come from rotation and vibration, still proportional to temperature.

