

Ergodic Theory

Solutions to Problem Sheet 9

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P1. Let (X, \mathcal{B}, μ, T) be a measure preserving system and let $f \in L^2(X)$. Let μ_f be the spectral measure of f with respect to U_T , and let $P_T f$ be the orthogonal projection onto the closed subspace of T -invariant functions in $L^2(X, \mu)$. Show that:

- (a) $\mu_f(\{0\}) = \|P_T f\|_2^2$.
- (b) If the system is ergodic, then: $\mu_f(\{0\}) = 0$ if and only if $\int f d\mu = 0$.
- (c) $\mu_f(\mathbb{T}) = \|f\|_2^2$.

(a) We want to compute

$$\mu_f(\{0\}) = \int_{\mathbb{T}} \mathbb{1}_{\{0\}}(x) d\mu_f(x).$$

Notice that

$$\mathbb{1}_{\{0\}}(x) = \lim_{N \rightarrow \infty} \frac{1}{N} \sum_{n=0}^{N-1} e(nx),$$

and by the Dominated convergence theorem, we get

$$\mu_f(\{0\}) = \lim_{N \rightarrow \infty} \frac{1}{N} \sum_{n=0}^{N-1} \int_{\mathbb{T}} e(nx) d\mu_f(x) = \lim_{N \rightarrow \infty} \frac{1}{N} \sum_{n=0}^{N-1} \langle T^n f, f \rangle.$$

Now, write $f = P_T f + f_{\text{erg}}$ where f_{erg} is the coboundary part of f . Using the mean ergodic theorem and orthogonality, we get that

$$\begin{aligned} \mu_f(\{0\}) &= \lim_{N \rightarrow \infty} \frac{1}{N} \sum_{n=0}^{N-1} \langle T^n P_T f, P_T f \rangle + \langle T^n f_{\text{erg}}, P_T f \rangle + \langle T^n P_T f, f_{\text{erg}} \rangle + \langle T^n f_{\text{erg}}, f_{\text{erg}} \rangle \\ &= \lim_{N \rightarrow \infty} \frac{1}{N} \sum_{n=0}^{N-1} \langle P_T f, P_T f \rangle + \langle T^n f_{\text{erg}}, f_{\text{erg}} \rangle \\ &= \langle P_T f, P_T f \rangle = \|P_T f\|_2^2. \end{aligned}$$

□

(b) Recall the mean ergodic theorem for the ergodic case. If T is ergodic, then $P_T f = \int f d\mu$. Thus

$$\mu_f(\{0\}) = \left\| \int f d\mu \right\|_2^2 = \left(\int f d\mu \right)^2,$$

from which the equivalence follows. □

(c) Using the spectral theorem we have that

$$\mu(\mathbb{T}) = \int_{\mathbb{T}} d\mu_f = \langle U_{\mathbb{T}}^0 f, f \rangle = \langle f, f \rangle = \|f\|_2^2.$$

□

In the following 3 exercises, \mathcal{H} is a separable Hilbert space, and for any $w \in \mathcal{H}$, we denote by \mathcal{H}_w the subspace generated by w , and we define $\chi_n(x) = e^{2\pi i n x}$. We define a unitary operator $U : \mathcal{H} \mapsto \mathcal{H}$. Recall that the spectral measure μ_w of w with respect to U is a finite Borel measure on $\mathbb{T} = \mathbb{R} \setminus \mathbb{Z}$, uniquely determined by the following property

$$\langle U^n w, w \rangle = \int_{\mathbb{T}} \chi_n(x) d\mu_w(x) \quad \forall n \in \mathbb{Z}. \quad (1)$$

Additionally, one can show that using this measure we obtain a unitary isomorphism $\mathcal{H}_w \cong L_{\mu_w}^2(\mathbb{T})$ where the action of U on \mathcal{H}_w corresponds to the action of M_{χ_1} on $L_{\mu_w}^2(\mathbb{T})$ and w corresponds to $\mathbb{1} \in L_{\mu_w}^2(\mathbb{T})$, and where M_{χ_1} is the unitary operator given by $M_{\chi_1} : f \in L_{\mu_w}^2(\mathbb{T}) \mapsto \chi_1 f \in L_{\mu_w}^2(\mathbb{T})$.

P2. Let $w, z \in \mathcal{H}$. Show that there exists a complex signed measure $\mu_{w,z}$ such that

$$\langle U^n w, z \rangle = \int_{\mathbb{T}} \chi_n d\mu_{w,z} \quad \forall n \in \mathbb{Z}.$$

Recalling the polarization identity: $\langle w, z \rangle = \frac{1}{4} \sum_{k=0}^3 i^k \|w + i^k z\|^2$, one can identify a suitable candidate for the sought-after measure $\mu_{w,z}$. Set $\mu_{w,z} = \frac{1}{4} \sum_{k=0}^3 i^k \mu_{w+i^k z}$.

By the spectral theorem, for $n \in \mathbb{Z}$ we obtain

$$\int_{\mathbb{T}} \chi_n d\mu_{w,z} = \frac{1}{4} \sum_{k=0}^3 i^k \int_{\mathbb{T}} \chi_n d\mu_{w+i^k z} = \frac{1}{4} \sum_{k=0}^3 i^k \langle U^n (w + i^k z), w + i^k z \rangle.$$

Next, we observe that for any k we have the following expression

$$\langle U^n (w + i^k z), w + i^k z \rangle = \langle U^n w, w \rangle + i^k \langle U^n z, w \rangle + i^{-k} \langle U^n w, z \rangle + \langle U^n z, z \rangle.$$

Note that $\sum_{k=0}^3 i^k = \sum_{k=0}^3 i^{2k} = 0$. Therefore, the first, second and fourth term drop out and we obtain

$$\int_{\mathbb{T}} \chi_n d\mu_{w,z} = \frac{1}{4} \sum_{k=0}^3 \langle U^n w, z \rangle = \langle U^n w, z \rangle,$$

as desired. □

P3. Let $\lambda \in \mathbb{T}$, $0 \neq w \in \mathcal{H}$. Show that w is an eigenvector of U for eigenvalue $\chi_1(\lambda)$ if and only if $\mu_w = \|w\|^2 \delta_\lambda$, where δ_λ is the Dirac measure at the point λ .

For the forward direction, assume that w is an eigenvector of U with eigenvalue $\chi_1(\lambda)$. By the unitary isomorphism, this implies that $\mathbb{1}$ is an eigenvector of M_{χ_1} with eigenvalue λ , which gives that

$$M_{\chi_1}^n \mathbb{1} = \lambda^n \mathbb{1}.$$

Note that $\chi_1^n = \chi_n$, so $M_{\chi_1}^n \mathbb{1} = \chi_n \mathbb{1}$, and thus μ_w -a.e., $\lambda^n = \chi_n$ holds. This implies that

$$\int_{\mathbb{T}} \chi_n d\mu_w = \lambda^n \int_{\mathbb{T}} d\mu_w = \lambda^n \|w\|^2.$$

Define the measure $\nu := \|w\|^2 \delta_\lambda$, from which we observe that $\lambda^n \|w\|^2 = \int_{\mathbb{T}} \chi_n d\nu$. Recalling that every character of \mathbb{T} is of the form χ_n , we have shown that the measures ν and μ_w agree on characters, which implies that they agree on trigonometric polynomials. By Stone-Weierstrass they agree on $C(\mathbb{T})$. By density, they agree on $L^2(\mathbb{T})$, which contains characters, and so $\mu_w = \|w\|^2 \delta_\lambda$, as desired.

For the other direction, we assume that $\mu_w = \|w\|^2 \delta_\lambda$ for $\lambda \in \mathbb{T}$. Since we have a point-mass measure, this implies that μ_w -a.e., $\chi_n = \chi_n(\lambda)$ for all n . In particular, $\chi_1 = \chi_1(\lambda)$ holds μ_w -a.e., which shows that $\mathbb{1}$ is an eigenvector of M_{χ_1} with eigenvalue $\chi_1(\lambda)$. By the unitary isomorphism $\mathcal{H}_w \cong L^2_{\mu_w}(\mathbb{T})$, this implies that w is an eigenvector of U with eigenvalue $\chi_1(\lambda)$, as claimed. \square

P4. In this exercise, we give a slightly different proof of Wiener's lemma.

Let μ be a finite Borel measure on \mathbb{T} and $p_n(\mu) := \int_{\mathbb{T}} \chi_n d\mu$.

(a) Show that μ has at most countably many atoms which we denote by x_1, x_2, \dots

This follows directly from μ being a finite measure as this implies that for any $\delta > 0$ the set of atoms x with $\mu(\{x\}) > \delta$ is finite, which in turn gives the claim (by applying this with $\delta = \frac{1}{n}$). \square

(b) We define the function $D_N \in C(\mathbb{T})$ by

$$D_N = \sum_{n=-N}^N \chi_n.$$

Show that D_N is real-valued and that it can also be written as

$$D_N(x) = \begin{cases} 2N + 1 & \text{if } x = 0 \in \mathbb{T} \\ \frac{\sin((N + \frac{1}{2})2\pi x)}{\sin(\pi x)} & \text{if } x \neq 0 \end{cases}$$

Moreover, show that it satisfies

$$\int_{\mathbb{T}} D_N(x) dx = 1.$$

The case $x = 0$ follows immediately from the definition of D_N .

For $x \neq 0$, we recognise that D_N is a geometric series and using the standard relation $\sin(\theta) = \frac{e^{i\theta} - e^{-i\theta}}{2i}$ we obtain our claim

$$\begin{aligned} D_N(x) &= \sum_{n=-N}^N (e^{2\pi i x})^n = e^{-2\pi i N x} (1 + \dots + (e^{2\pi i x})^{2N}) \\ &= e^{-2\pi i N x} \frac{(e^{2\pi i x})^{2N+1} - 1}{e^{2\pi i x} - 1} \\ &= \frac{e^{2\pi i (N+\frac{1}{2})x} - e^{-2\pi i (N+\frac{1}{2})x}}{e^{\pi i x} - e^{-\pi i x}} \\ &= \frac{\sin((N+\frac{1}{2})2\pi x)}{\sin(\pi x)}. \end{aligned}$$

Finally, for the integral we proceed as follows

$$\int_{\mathbb{T}} D_N(x) dx = \int_{\mathbb{T}} \sum_{n=-N}^N \chi_n(x) dx = \sum_{n=-N}^N \int_{\mathbb{T}} \chi_n(x) dx = \sum_{n=-N}^N \mathbb{1}_{\{0\}}(n) = 1.$$

□

(c) Show that any finite measure ν on \mathbb{T} satisfies the following

$$\lim_{N \rightarrow \infty} \frac{1}{2N+1} \sum_{n=-N}^N p_n(\nu) = \nu(\{0\}). \quad (2)$$

We have

$$\frac{1}{2N+1} \sum_{n=-N}^N p_n(\nu) = \frac{1}{2N+1} \sum_{n=-N}^N \int_{\mathbb{T}} \chi_n d\nu = \int_{\mathbb{T}} \frac{1}{2N+1} D_N d\nu.$$

By part (b), we know that $\frac{1}{2N+1} D_N$ converges to 0 for $x \neq 0$ as $N \rightarrow \infty$, it is bounded by 1 and is equal to 1 for $x = 0$. Therefore, by the Dominated Convergence Theorem, we obtain

$$\lim_{N \rightarrow \infty} \frac{1}{2N+1} \sum_{n=-N}^N p_n(\nu) = \int_{\mathbb{T}} \mathbb{1}_{\{0\}} d\nu = \nu(\{0\}).$$

□

(d) Denote by $\Delta = \{(t, t) \mid t \in \mathbb{T}\}$. Let ν be the push-forward of the product measure $\mu \times \mu$ on \mathbb{T}^2 under the map $(t_1, t_2) \in \mathbb{T}^2 \mapsto t_1 - t_2 \in \mathbb{T}$. Show that $\nu(\{0\}) = \mu \times \mu(\Delta)$ and $p_n(\nu) = |p_n(\mu)|^2$.

By definition, for any measurable set $A \subset \mathbb{T}$, we have $\nu(A) = \mu \times \mu(D^{-1}(A))$, where D is the difference map $D : (t_1, t_2) \in \mathbb{T}^2 \mapsto t_1 - t_2 \in \mathbb{T}$. Observe that $D^{-1}(\{0\}) = \Delta$, so the first claim follows immediately.

For the second claim, note that for any simple function $f = \mathbb{1}_A$ for $A \subset \mathbb{T}$ measurable,

the following identity is satisfied:

$$\int_{\mathbb{T}} f \, d\nu = \int_{\mathbb{T}^2} f \circ D \, d(\mu \times \mu).$$

Therefore, this identity also holds for any linear combinations of simple functions, and therefore for any L^1 -function, which follows by density. In particular, we can take $f = \chi_n$ and using Fubini's Theorem, we conclude

$$\begin{aligned} p_n(\nu) &= \int_{\mathbb{T}} \chi_n \, d\nu = \int_{\mathbb{T}^2} \chi_n \circ D \, d(\mu \times \mu) \\ &= \int_{\mathbb{T}} \int_{\mathbb{T}} \chi_n(t_1 - t_2) \, d\mu(t_1) \, d\mu(t_2) \\ &= \int_{\mathbb{T}} \int_{\mathbb{T}} \chi_n(t_1) \overline{\chi_n(t_2)} \, d\mu(t_1) \, d\mu(t_2) \\ &= p_n(\mu) \overline{p_n(\mu)} \\ &= |p_n(\mu)|^2. \end{aligned}$$

□

(e) Show that $\mu \times \mu(\Delta) = \sum_{i=1}^{\infty} \mu(\{x_i\})^2$ and deduce Wiener's lemma:

$$\lim_{N \rightarrow \infty} \frac{1}{2N+1} \sum_{n=-N}^N |p_n(\mu)|^2 = \sum_{i=1}^{\infty} \mu(\{x_i\})^2. \quad (3)$$

Using Fubini's Theorem once again, we have

$$\mu \times \mu(\Delta) = \int_{\mathbb{T}} \int_{\mathbb{T}} \mathbb{1}_{\Delta}(t_1, t_2) \, d\mu(t_1) \, d\mu(t_2) = \int_{\mathbb{T}} \mu(\{t_2\}) \, d\mu(t_2).$$

By part (a), the last integrand is non-zero on the countable set of atoms we identified, and so we obtain

$$\begin{aligned} \int_{\mathbb{T}} \mu(\{t_2\}) \, d\mu(t_2) &= \int_{\mathbb{T}} \sum_{i=1}^{\infty} \mathbb{1}_{\{x_i\}}(t_2) \mu(\{t_2\}) \, d\mu(t_2) \\ &= \sum_{i=1}^{\infty} \mu(\{x_i\})^2, \end{aligned}$$

as desired. Finally, using parts (c) and (d), we deduce Wiener's lemma

$$\begin{aligned} \lim_{N \rightarrow \infty} \frac{1}{2N+1} \sum_{n=-N}^N |p_n(\mu)|^2 &= \lim_{N \rightarrow \infty} \frac{1}{2N+1} \sum_{n=-N}^N p_n(\nu) \\ &= \nu(\{0\}) \\ &= \mu \times \mu(\Delta) \\ &= \sum_{i=1}^{\infty} \mu(\{x_i\})^2. \end{aligned}$$

□