

# 5.1 Microphones theory and design

# Introduction & Definitions

# Introduction

Applications, conditions of use

➔ different specifications:

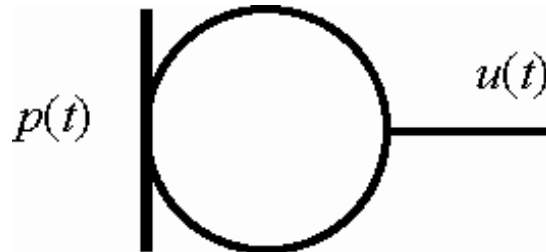
- Performances: sensitivity, directivity, non-linearities
- Mechanical, dimensions, solidity
- Electrical, impedance, polarization
- External influences
- Cost, useability

# Introduction

4 categories of microphones:

- Communication
- Professional broadcasting/recording
- General use
- Measurements

Symbol:

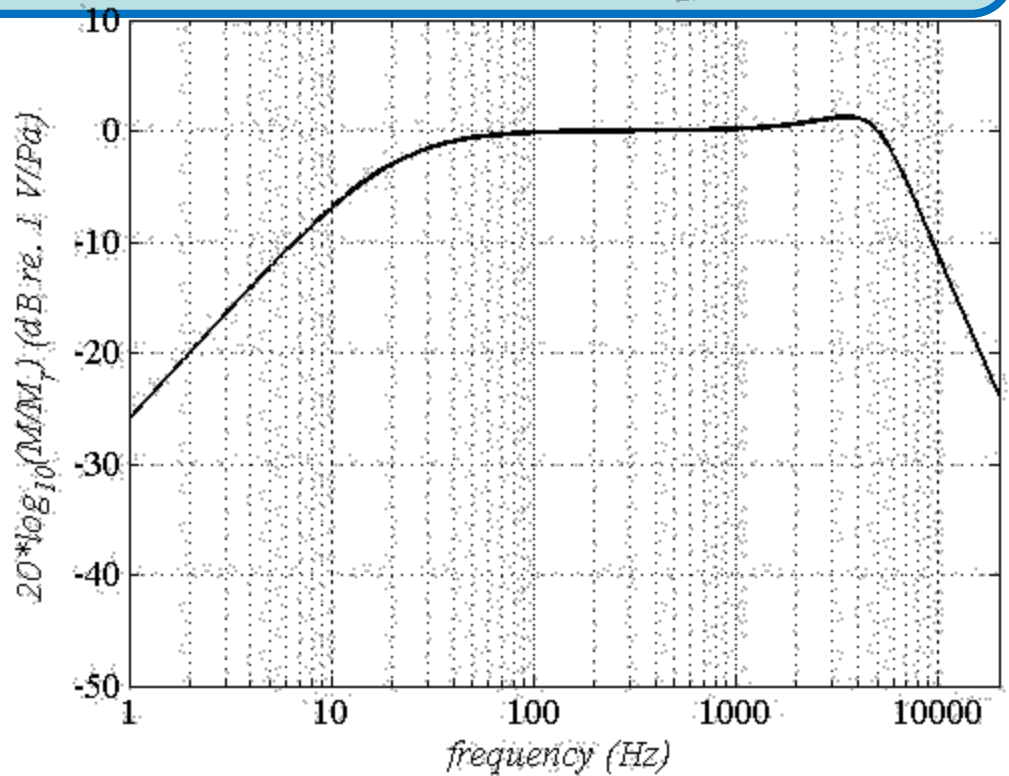


# Sensitivity

Sensibilité microphone

$$M = \frac{U}{\tilde{p}}$$

$$L_M = 20 \log_{10} \left( \frac{M}{M_r} \right)$$



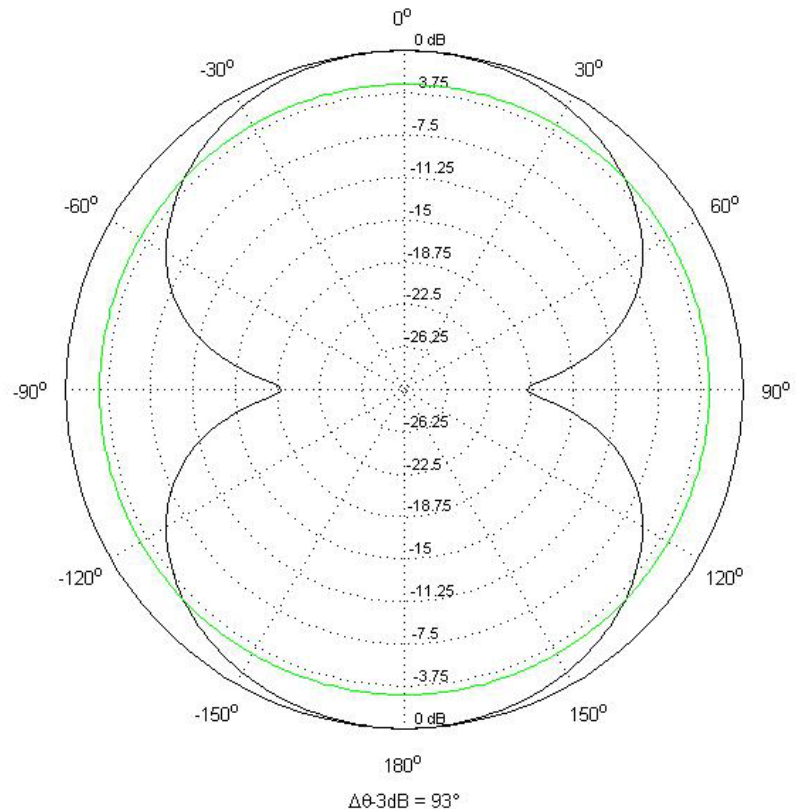
*Free field sensitivity:  $p$  at reference point, for progressive plane wave, before introducing the microphone (diffraction)*

# Directivity

Directivity: variation of  $M$  or  $L_M$  after incidence angle (in plane waves) w/ respect to a reference axis

$$M(\theta) = \frac{U(\theta)}{\tilde{p}}$$

$$L_M = 20 \log_{10} \left( \frac{M(\theta)}{M(\theta_0)} \right)$$



# Directivities

- Omnidirectional:

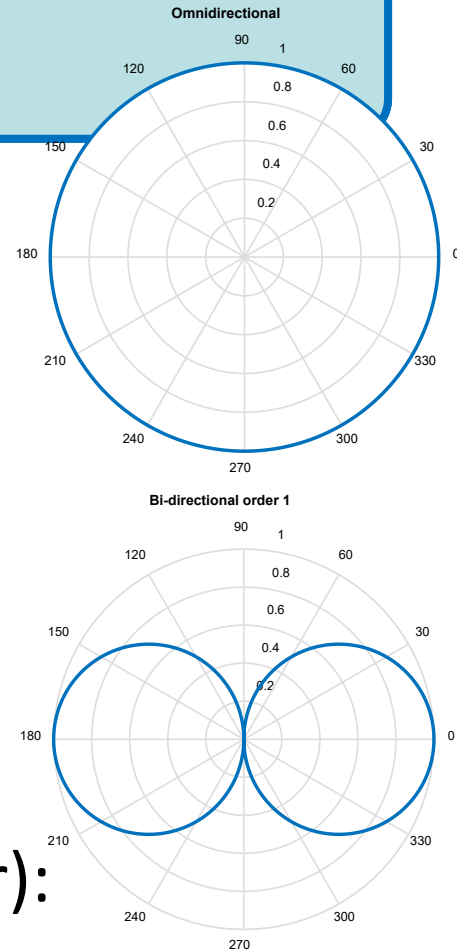
$$\frac{M(\theta)}{M_0} = 1$$

- Bidirectional of order n (positive integer):

$$\frac{M(\theta)}{M_0} = (\cos \theta)^n$$

- Unidirectional of order n (positive integer):

$$\frac{M(\theta)}{M_0} = \left( (1 - \mu) + \mu \cos \theta \right) (\cos \theta)^{n-1}$$

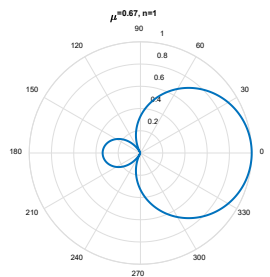
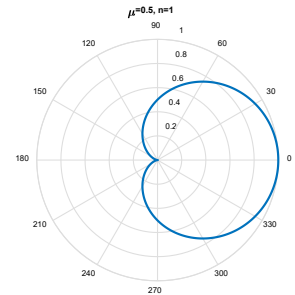
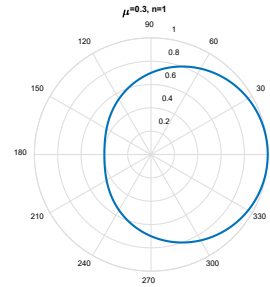


# Directivities

- Unidirectional of order n (positive integer):

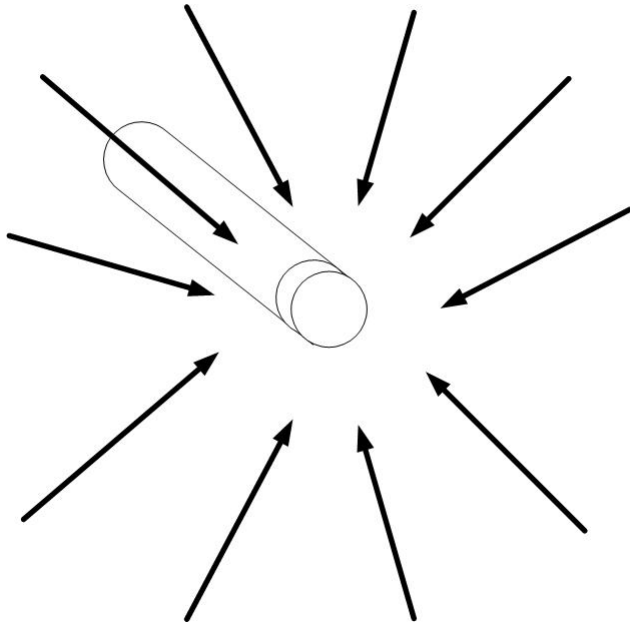
$$\frac{M(\theta)}{M_0} = \left\{ (1 - \mu) + \mu \cos \theta \right\} (\cos \theta)^{n-1}$$

- $\mu=0,3$ : infracardioid
- $\mu=0,5$ : cardioid
- $\mu=0,63$ : supercardioid
- $\mu=0,75$ : hypercardioid



# Sensitivities in diffuse field

Ideally diffuse sound field (isotropic)



$$M_d = \sqrt{\frac{1}{4\pi} \int_0^{4\pi} M^2(\theta, \phi) d\Omega}$$

where  $M(\theta, \phi)$  is the sensitivity in free field of the microphone.

# Directivity factor and index

- Hypothesis: microphone subject to:
  - prog. plane wave along  $\theta=0^\circ$ , rms value  $p$  (signal)
  - diffuse field of same rms value (noise)

Directivity factor = signal to noise ratio

$$\Delta = (M_o/M_d)^2$$

Directivity index:  $L_{\Delta} = 10 \cdot \log_{10}(\Delta)$

# Directivity factor and index

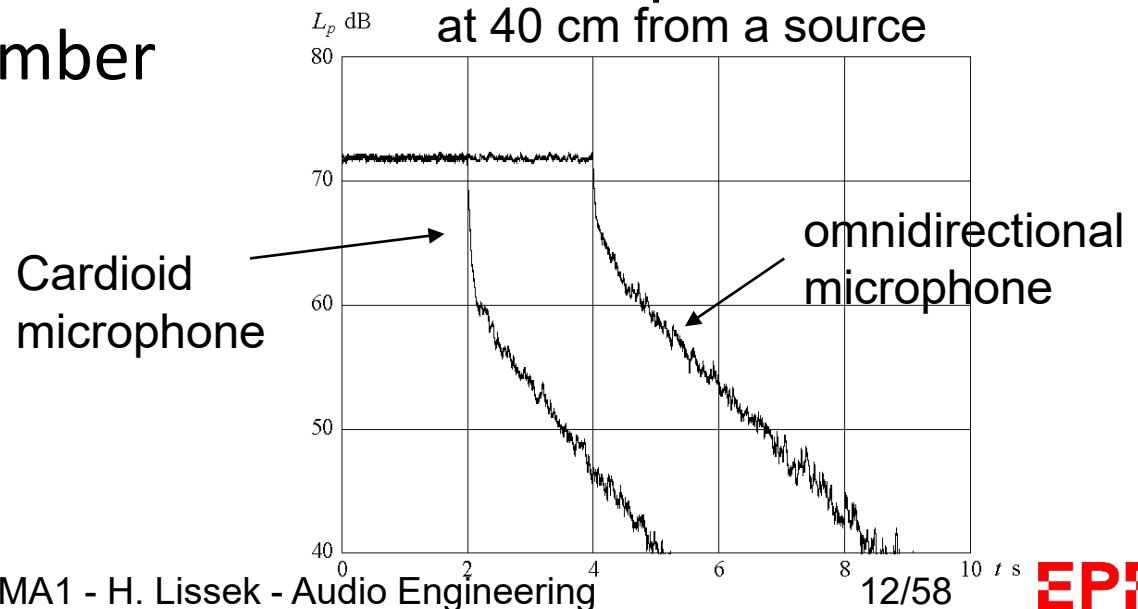
Tableau 8.3 Caractéristiques de directivités.

Microphone	$L_{\Delta}$ dB	$\theta_{-3}$ °	$L_M(90^\circ)$ dB	$L_M(180^\circ)$ dB	$\theta(M=0)$ °
omnidirectionnel	0	—	0	0	—
bidirectionnel $n = 1$	4,8	90	$-\infty$	0	$\pm 90$
cardioïde $n = 1$	4,8	131	-6	$-\infty$	180
cardioïde $n = 2$	8,8	76	$-\infty$	$-\infty$	$180, \pm 90$
hypercardioïde $n = 1$	6,0	105	-12	-6	$\pm 110$
supercardioïde $n = 1$	5,7	115	-8,6	-11,7	$\pm 126$

# Property: reverberant field reduction

These properties highlight the differences of pressure decreases measured with different types of microphones:

→  $L_{\Delta}$  correspond to a discontinuity in the decrease of diffuse field measured with the microphone in a reverberant chamber



# Property: diffraction

microphone=obstacle

→  $p$  different if microphone not « present »

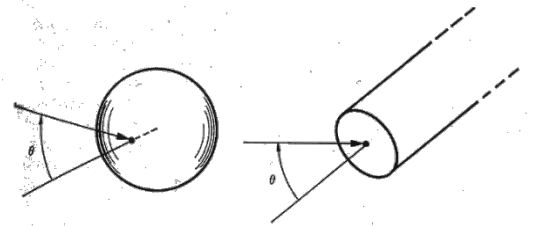
principle= see diffraction

diffraction = increase of directivity

when microphone dimensions  $\gg$  wavelength

→ Importance of microphone shapes and dimensions :

- spheres and cylinders: optimal
- lower dimensions → broadening of the band



# Sensitivity in pressure

*Sensitivity in pressure*  $M_p = \frac{\text{output voltage}}{\text{effective pressure } \tilde{p}_+}$

Progressive plane waves =  $p_+ > p \rightarrow M_p < M_d$  (free field)

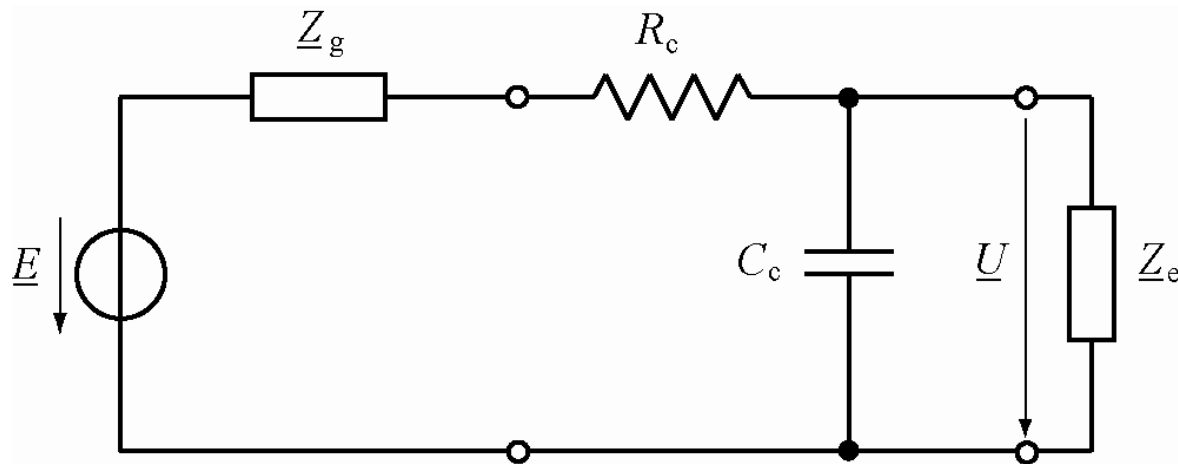
Sensitivity in pressure useful if diffraction can be neglected (microphone in a dispositive much wider, in an acoustical system)

Else: *sensitivity in free field*

# Microphone-equipment matching

After Thévenin, voltage  $U$  depends on load impedances:

- impedance matching ( $Z_e = Z_g^*$ )
- function in open circuit:  $Z_e > 5Z_g$  (resistance)



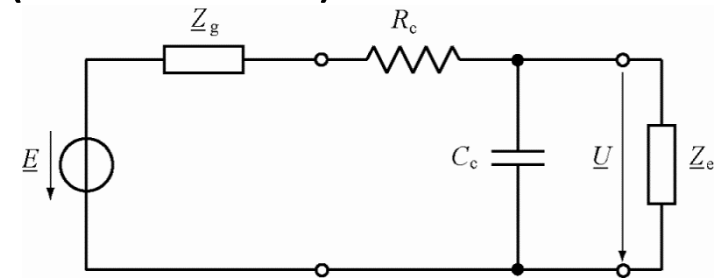
# Nominal values and conditions

Normalized specifications (CEI268-4):

- Nominal impedance  $Z_n = |Z_g|$  @ 1 kHz
- Nominal load impedance  $Z_{en} = |Z_e|$  @ 1 kHz
- Nominal output voltage  $U_n = U$  (nom. conditions)
- Nominal sensitivity in free field  $M_n$  (nom. cond.)

Nominal conditions:

- load impedance =  $Z_{en}$
- prog plane wave along  $\theta=0^\circ$ ,  $f=1$  kHz,  $L_p=80$  dB (without micro)
- nom. alimentation (active micro), or preamp.



# Proximity microphone

Proximity microphone designed to be close to the mouth

*Paraphonic sensitivity*: sensitivity of a proximity microphone, measured in conditions:

- Artificial mouth, at 50 mm of the microphone
- CEI:  $L_p=104$  dB at reference point of microphone
- CCITT (Consultative Committee for International Telegraph and Telephone):  $L_p=93$  dB

# Other specifications

Nominal matching values: 3 ranges:

- low impedance: 15 à 60  $\Omega$
- mid impedance : 100 à 10000  $\Omega$
- high impedance: 2 à 100 k $\Omega$

Output connectors:

- Symmetrical output (no connection to the mass)
- Asymmetrical output : 1 positive, 1 point to the mass

Limit characteristics:

- limit acoustic pressure(<harmonic distortion),
- max. allowed peak pressure,
- etc.

# Other specifications

- exterior influences
- « pop » protection (depends on the « action mode »)
- responses to impulsive sounds (non-linearities)
- background noise, self noise

# Microphones theory

# Description

Sound pressure  $p$  exerted on a diaphragm of surface  $S_d$

→ Force  $F = p \cdot S_d$

These forces induce a movement of the diaphragm  
(mechanical system)

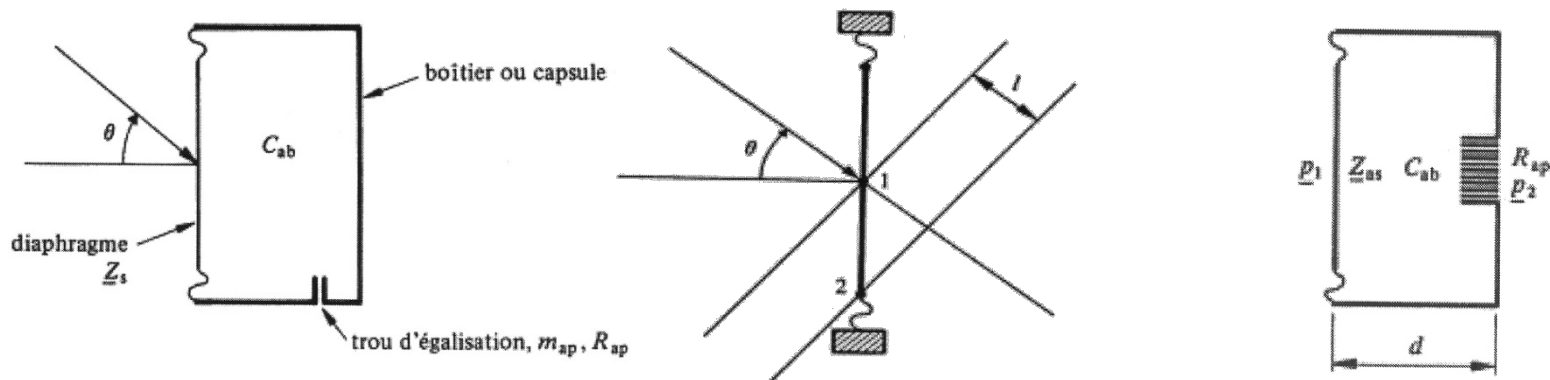
Transduction (generator behavior) converts mechanical  
quantity into electrical one

**What are the cascades allowing the respect of  
specifications (see flat response)?**

# Acoustic action modes

*Acoustic action modes* = ways the acoustic pressure impacts the diaphragm

- *pressure microphones*: pressure impacts only one face
- *pressure gradient microphones* : on the 2 faces
- *mixed mode microphones* : directly on 1 front face  
+ via one acoustical system connected to rear face



# Diaphragm movement

diaphragm = mechanical resonator (in all cases)

+ acoustic components

+ electric components

→  $Z_{mt}$

Acoustic field:

resulting applied forces  $F_s \Rightarrow v_d = \frac{F_s}{Z_{mt}}$

# Conversions in velocity vs. elongation

Electromechanical coupling equations

→ 2 types of behavior

- conversion in *velocity* :  $E = K_v v$

(electrodynamic, electromagnetic, magnetostrictive)

- conversion in *elongation* :  $E = K_\xi \xi$

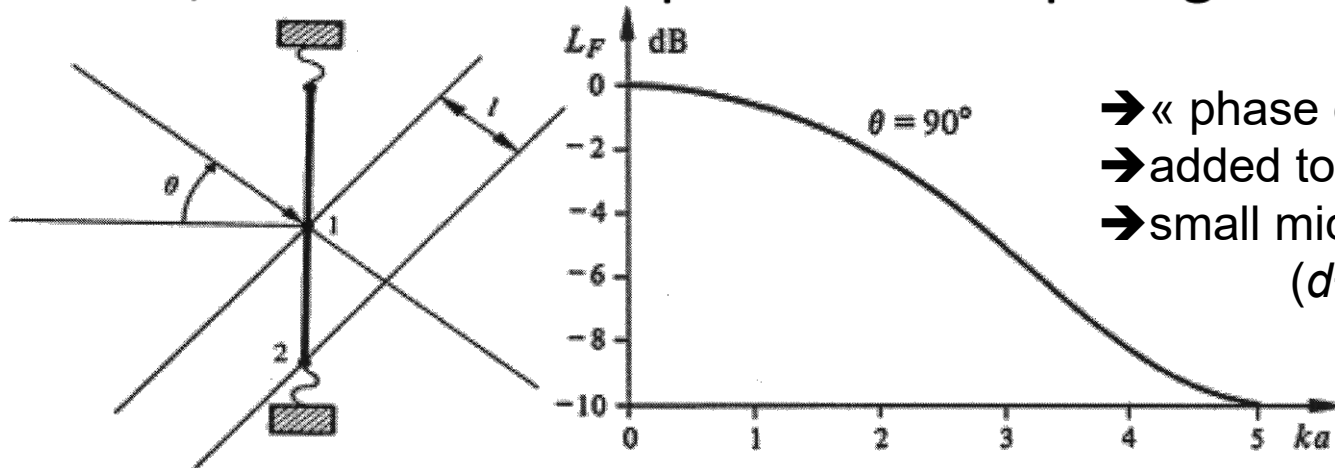
(electrostatic, piezoelectric)

# Force on one face

- Rigid diaphragm, plane disc of radius  $a$ , suspended, subject to plane wave, incidence  $\theta$
- dimensions  $\ll \lambda$

$\theta=0^\circ$ , diaphragm = wave plane  $\rightarrow F_s = S_d p$

If  $\theta \neq 0^\circ$ , : non uniform phase on diaphragm

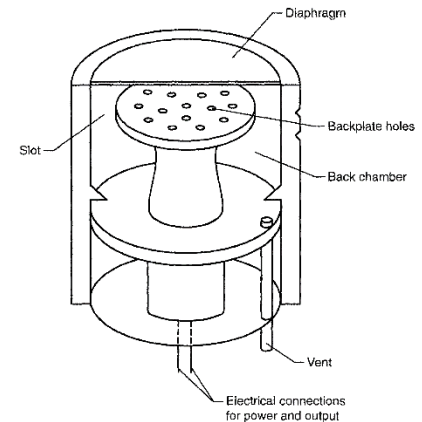


- $\rightarrow$  « phase effect »
- $\rightarrow$  added to diffraction
- $\rightarrow$  small microphone hypothesis ( $d < \lambda$ ) alleviates this

# Pressure microphone

Realization:

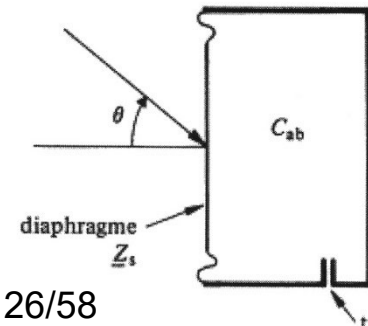
- Casing (compliance  $C_{ab}$ ) isolates rear face to sound field
- Equalization of static pressure :
  - small opening (mass  $m_{ab}$ , resistance  $R_{ab}$ )
  - so that  $\underline{q}_p$  (in the opening) is negligible (within microphone bandwidth)



Small microphone:

total force  $F_s = F_d = S_d p$  do not depend on  $\theta$ : omnidirectional  
 $F_s$  do not depend on  $f$  neither

**We want to have  $U$  independent on  $f$**



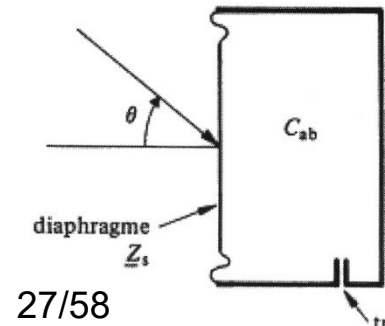
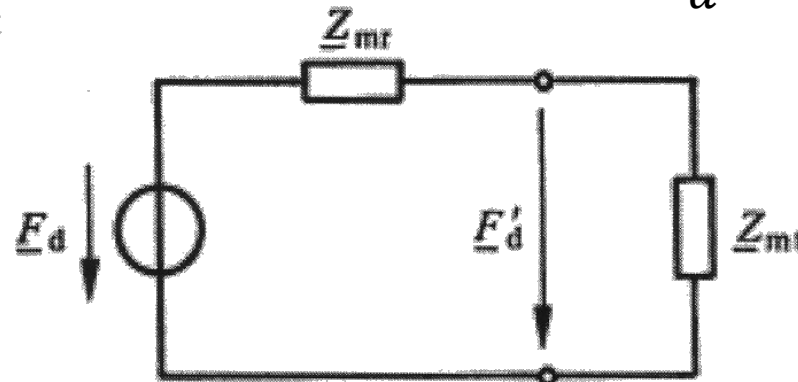
# Field's source impedance

cf Thévenin:

Acoustic field = source + source impedance

diaphragm = piston in a box  $\rightarrow$  self radiation  $Z_{mr}$

$\rightarrow$  Characteristic impedance  $Z_n = \frac{Z_{mt}}{S_d}$



# Pressure microphones

- Action in *pressure*, conversion in *velocity*:

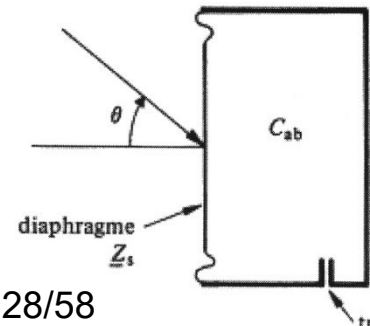
$$\underline{E} = K_v \underline{v} = K_v \frac{\underline{F}_s}{\underline{Z}_{mt}} = \frac{K_v S_d}{\underline{Z}_{mt}} \underline{p}$$

→ if we want  $\underline{M}$  independent on  $f$ :  $\underline{Z}_{mt} \triangleq R_{mt}$

- Action in *pressure*, conversion in *elongation*:

$$\underline{E} = K_\xi \underline{\xi} = \frac{K_\xi S_d}{j\omega \underline{Z}_{mt}} \underline{p}$$

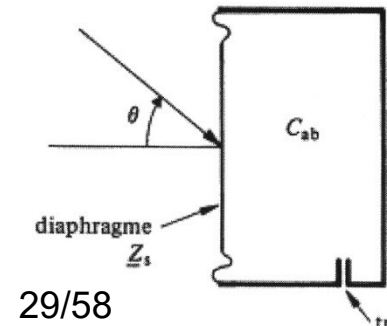
→ if we want  $\underline{M}$  independent on  $f$ :  $\underline{Z}_{mt} \triangleq \frac{1}{j\omega C_{mt}}$



# Pressure microphones

Conclusions: 2 fundamental types

- *conversion in velocity* and *resistive control*
- *conversion in elongation* and *compliant control*



# Pressure gradient microphones

$p$  exerted on 2 faces. Realization: diaphragm elastically suspended on annular surrounding

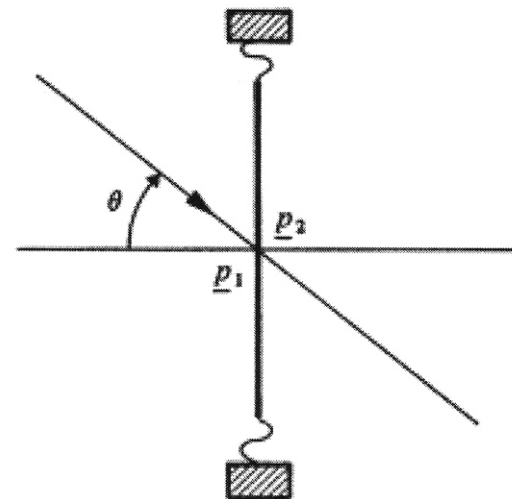
$$\text{Total force } F_s = F_1 - F_2 = S_d(p_1 - p_2)$$

$$(p_1 - p_2) = -\vec{\nabla} p \cdot \vec{n} \underbrace{l}_{e \cos \theta}$$

Locally

$$F_s = j\omega\rho S_d e \cos \theta v_1$$

→ velocity microphone



# Pressure gradient microphones

## Behavior under plane waves (far field)

$\underline{p} = \underline{Z}_s \underline{v}$ ,  $\underline{Z}_s$  specific impedance of **spherical wave**

- $kr \gg 1$ : spherical waves  $\equiv$  plane waves

$\underline{p} = \rho c \underline{v}$

$$\underline{F}_s = j \frac{\omega S_d e \cos \theta}{c} \underline{p}_1$$

depends on

- incidence  $\theta \rightarrow$  bidirectional
- $S_d$  and  $e \rightarrow$  big « volume » of diaphragm preferred
- frequency (slope +6dB per octave)



$$Z_s(r) = Z_c \frac{kr}{kr - j}$$
$$\begin{cases} R_s = Z_c \frac{(kr)^2}{1 + (kr)^2} \\ X_s = Z_c \frac{kr}{1 + (kr)^2} \end{cases}$$

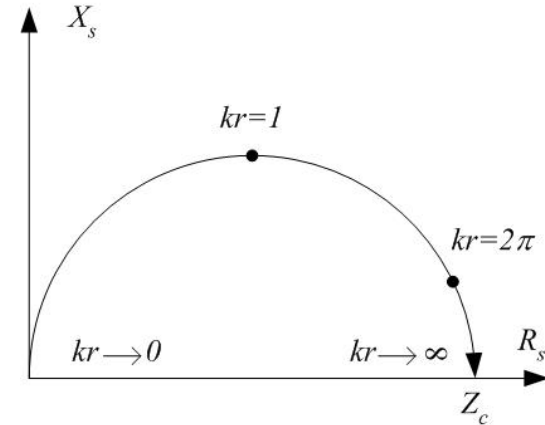
# Pressure gradient microphones

## Behavior under spherical waves (proximity)

$\underline{p} = \underline{Z}_s \underline{v}$ ,  $\underline{Z}_s$  specific impedance of **spherical wave**

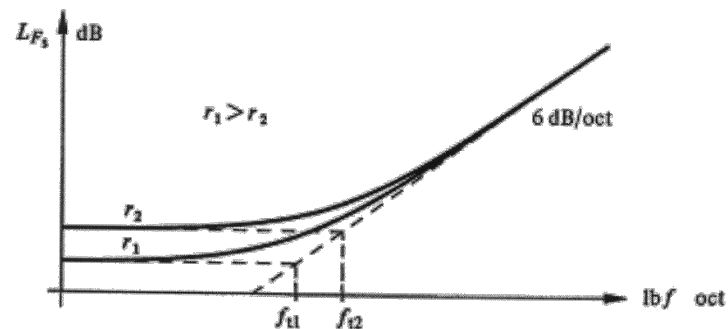
- $kr \gg 1$ : spherical waves  $\equiv$  plane waves

- $kr \ll 1$ :  $\underline{F}_s \cong \frac{S_d e \cos \theta}{r} \underline{p}_1$



Remark: getting closer to microphone, sensitivity increases + does not depend on  $f$

→ proximity effect



$$Z_s(r) = Z_c \frac{kr}{kr - j}$$

$$\begin{cases} R_s = Z_c \frac{(kr)^2}{1 + (kr)^2} \\ X_s = Z_c \frac{kr}{1 + (kr)^2} \end{cases}$$

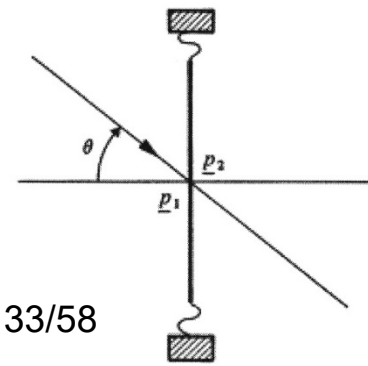
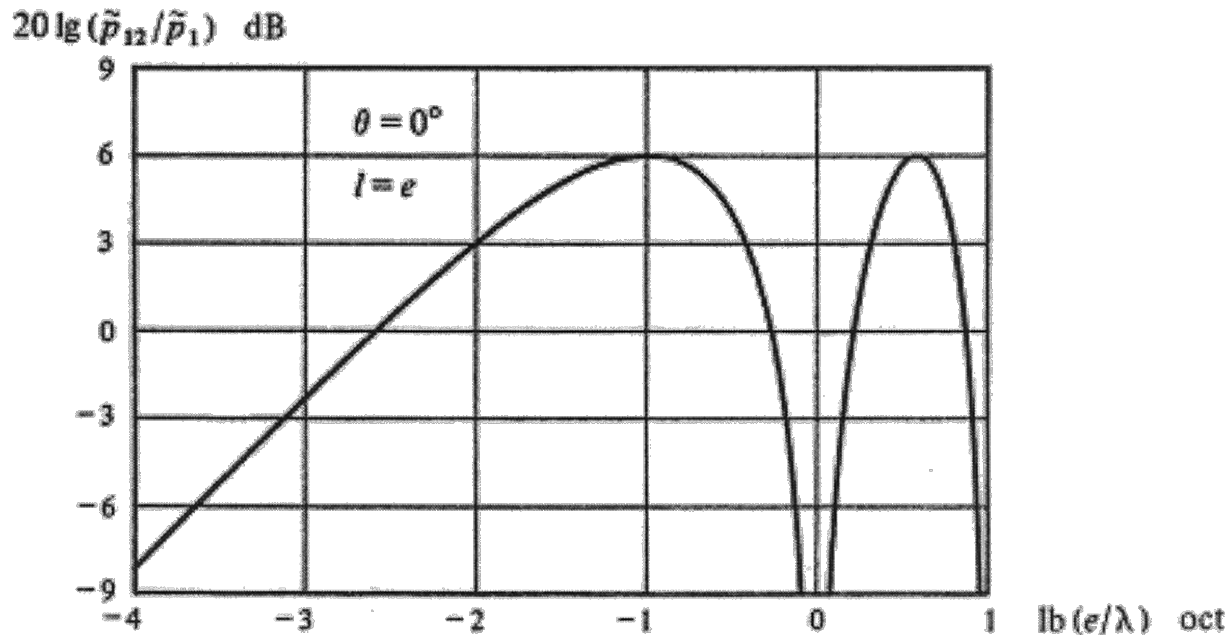
# Pressure gradient microphones

## Behavior at high frequencies

$\lambda$  becomes of same order of magnitude than  $l$

$$(p_1 - p_2) = -\nabla p \cdot \vec{n}l \quad \text{not valid}$$

→ new approximation:  $p_{12} = (p_1 - p_2) = p_1 (1 - e^{-jkl})$



# Pressure gradient microphones

## Increase of sensitivity with a screen

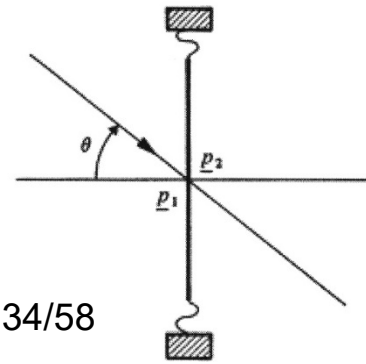
Velocity microphone :

$F_s <$  force exerted on one single face

→ sensitivity lower than pressure microphone

→ if  $F_s$  increases, sensitivity increases

screen:  $p_1 > p_2$  → increases  $F_s$



# Pressure gradient microphones

Condition of realization: far field

- Action in *velocity*, conversion in *velocity*:

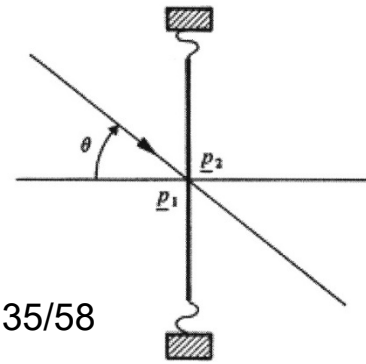
$$E = K_v v_1 = K_v \frac{F_s}{Z_{mt}} = \frac{j\omega K_v S_d e \cos \theta}{c Z_{mt}} p_1$$

if we want M independent on *f*:  $Z_{mt} \triangleq j\omega m_{mt}$

- Action in *velocity*, conversion in *elongation*:

$$E = K_\xi \xi = \frac{K_\xi S_d e \cos \theta}{c Z_{mt}} p_1$$

if we want M independent on *f*:  $Z_{mt} \triangleq R_{mt}$

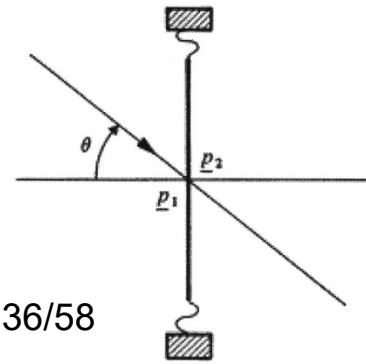


# Pressure gradient microphones

Condition of realization: far field

2 fundamental types :

- ***conversion in velocity*** and ***control by the mass***
- ***conversion in elongation*** and ***resistive control***



# Pressure gradient microphones

Condition of realization: proximity microphones

- conversion in velocity:

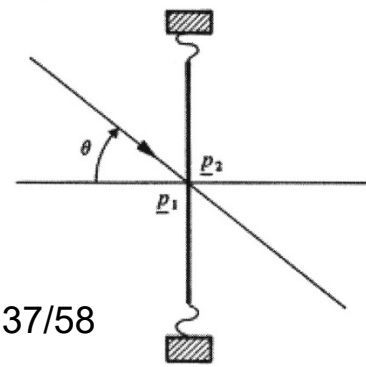
$$E = \frac{K_v S_d e \cos \theta}{r Z_{mt}} p_1$$

if we want M independent on  $f$ :  $Z_{mt} \triangleq R_{mt}$

- conversion in elongation:

$$E = \frac{K_\xi S_d e \cos \theta}{j\omega r Z_{mt}} p_1$$

if we want M independent on  $f$ :  $Z_{mt} \triangleq \frac{1}{j\omega C_{mt}}$



# Pressure gradient microphones

Condition of realization: proximity microphones

2 fundamental types :

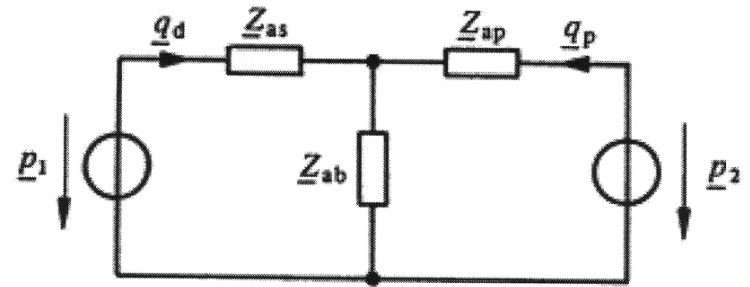
- ***conversion in velocity*** and ***resistive control***
- ***conversion in elongation*** and ***control by compliance***

# Mixed action microphones

Comprises one acoustic system: one or more openings (« rear access »)

equivalent acoustic scheme  
(diaphragm + acoustic system):

here,  $\underline{p}_2$  is defined on rear access;



Spherical waves: 
$$\underline{p}_2 = \underline{p}_1 \left[ 1 - \underbrace{d \cos \theta}_{\substack{\text{distance between} \\ \text{homolog points}}} \left( \underbrace{\frac{1}{r}}_{\substack{\text{proximity} \\ \text{effect}}} + jk \right) \right]$$

Plane waves: 
$$\underline{p}_2 = \underline{p}_1 [1 - jkd \cos \theta]$$

# Mixed action microphones

$$q_d = \frac{\left[ 1 + \left( Z_{ab} / Z_{ap} \right) (jkd \cos \theta) \right]}{Z_{as} + Z_{as} \left( Z_{ab} / Z_{ap} \right) + Z_{ab}} p_1$$

$$1 + \left( Z_{ab} / Z_{ap} \right) (jkd \cos \theta) = (1 - \mu) + \mu \cos \theta$$

$$\Rightarrow j\omega(d/c) Z_{ab} / Z_{ap} = \frac{\mu}{1 - \mu}$$

# Mixed action microphones

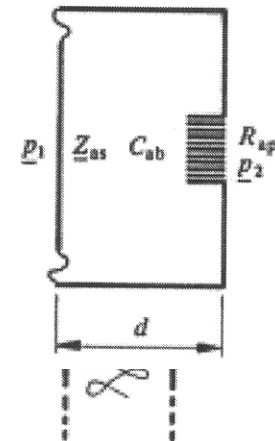
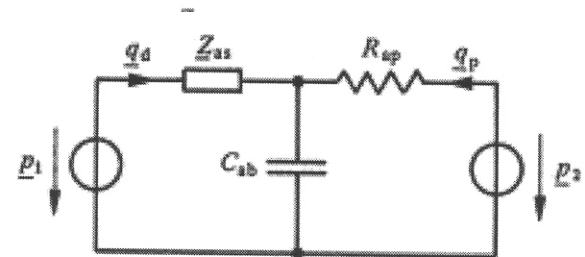
## Realization of a cardioid directivity ( $\mu=1/2$ )

$$\mu = \frac{1}{2} \Rightarrow \frac{\mu}{\mu-1} = 1 \quad \Rightarrow \quad Z_{ab} = \frac{Z_{ap}}{j\omega(d/c)}$$

2 realizations:

$$Z_{ap} = R_{ap} \Rightarrow Z_{ab} = \frac{1}{j\omega C_{ab}}$$

$$Z_{ap} = j\omega m_{ap} \Rightarrow Z_{ab} = R_{ab}$$



# Mixed action microphones

## Mastering the response curve

sensitivity independent on frequency

$$Z_{at} = Z_{as} + Z_{as} \underbrace{\frac{Z_{ab}}{Z_{ap}}}_{\mu/(1-\mu)j\omega(d/c)} + Z_{ab}$$

Small microphone, mixed action, and conversion in velocity

→  $|Z_{at}|$  constant

→ if conversion in velocity  $Z_{as} = j\omega m_{as}$  if acoustic system

$$R_{ap} C_{ab}$$

# Combined microphones

*Combined microphone* : constituted with 2 microphones in a single casing, coincident:

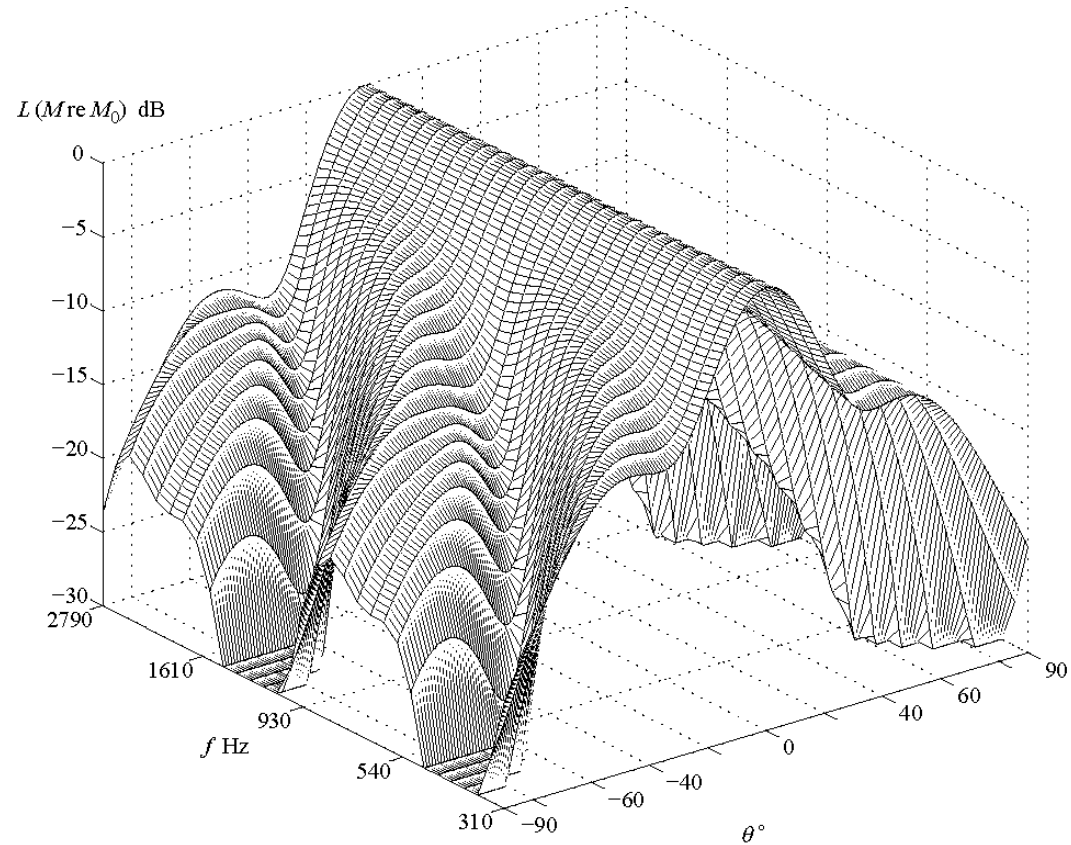
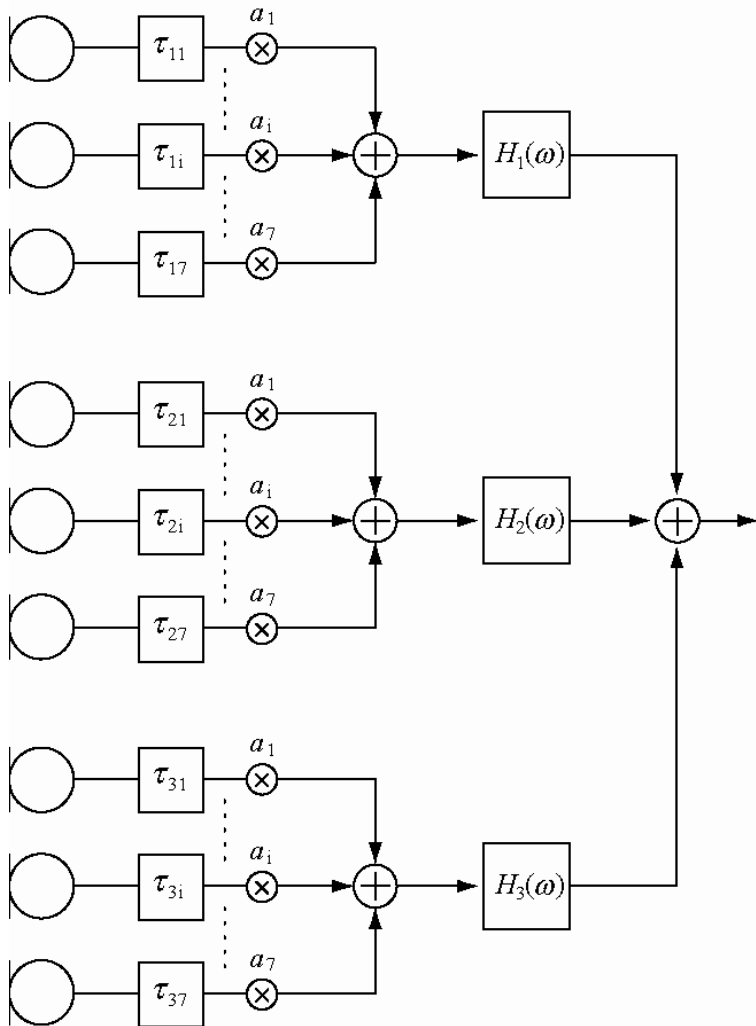
- 1 omnidirectional
- 1 bidirectional

weighted output voltages:  $\underline{E} = \alpha \underline{E}_1 + \mu \underline{E}_2$

Small microphones:

$$\underline{E} = (\alpha \underline{M}_1 + \mu \underline{M}_2 \cos \theta) \underline{p} = [(1 - \mu) + \mu \cos \theta] \underline{M}_2 \underline{p}$$

# Microphones arrays



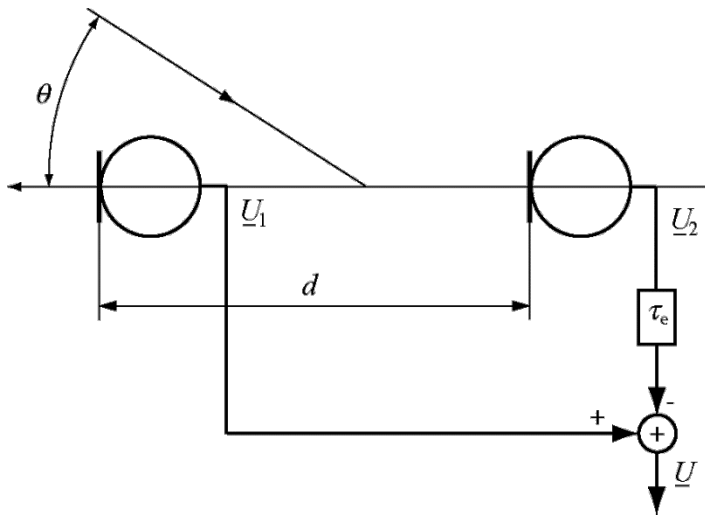
# Microphones arrays

2 small identical omni microphones, distant of  $d$  ( $kd < 1$ )

Plane wave, incidence  $\theta$ ;

electronic delay  $\tau_e$  between 2 output voltages

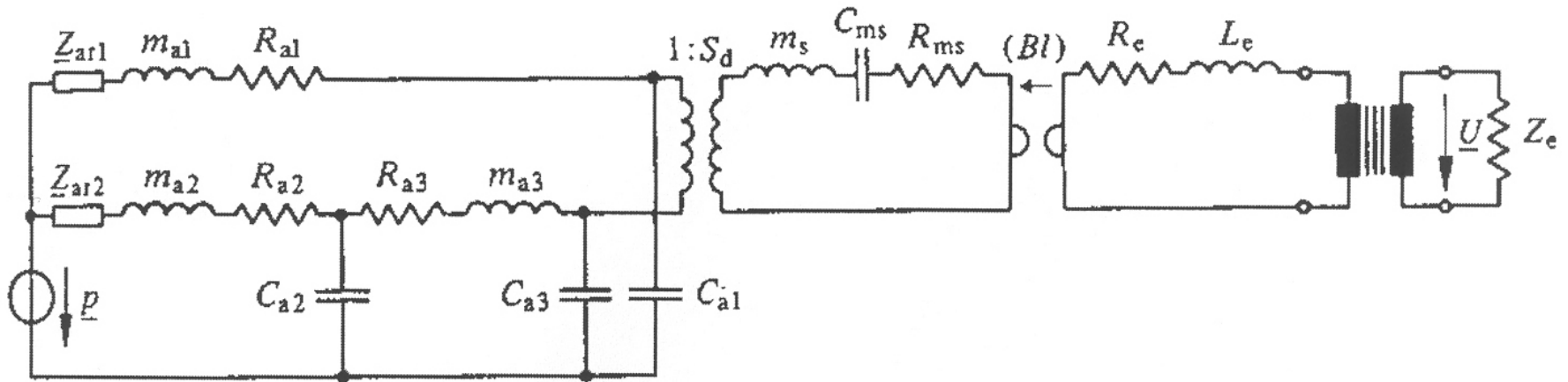
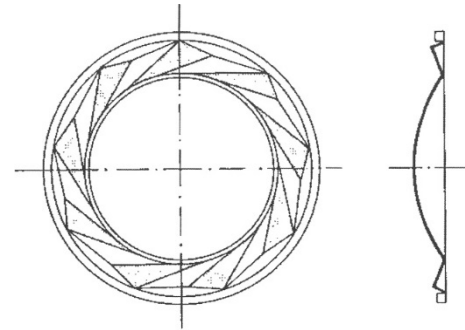
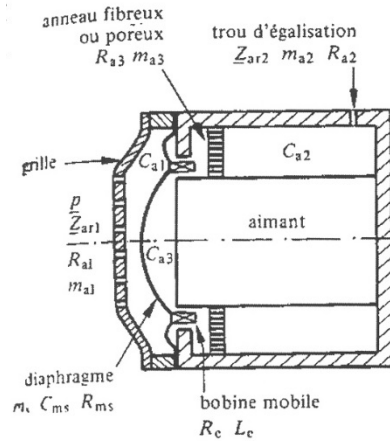
The 2 voltages are subtracted:



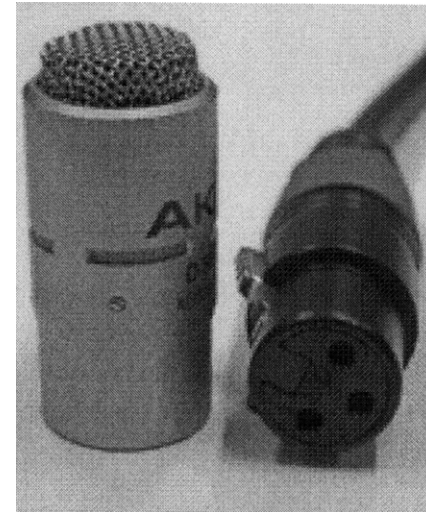
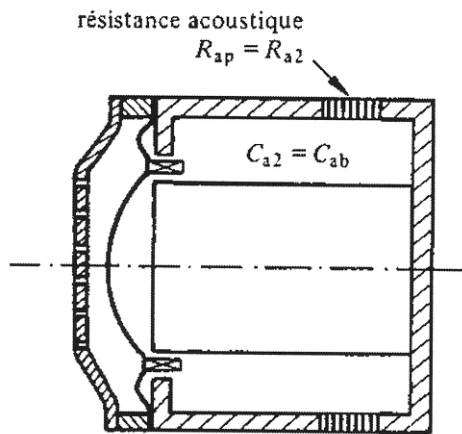
$$\begin{aligned}
 \underline{U} &= \underline{U}_1 - \underline{U}_2 \exp(-j\omega\tau_e) \\
 &= \underline{M}\underline{p}_1 \left\{ 1 - \exp \left[ -j\omega \left( \tau_e + \underbrace{\tau_a}_{d/c} \cos \theta \right) \right] \right\} \\
 &\cong \underline{M}\underline{p}_1 j\omega (\tau_e + \tau_a \cos \theta)
 \end{aligned}$$

# Microphones design

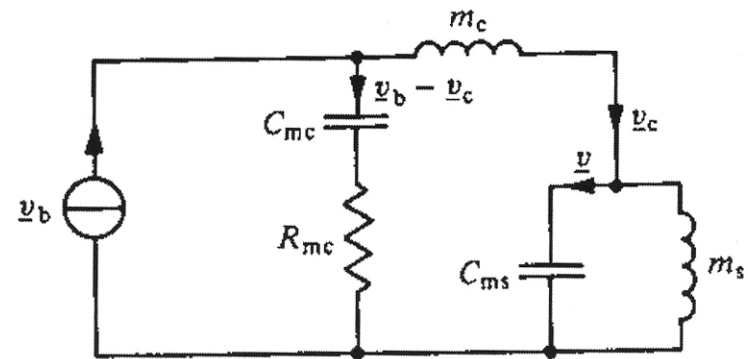
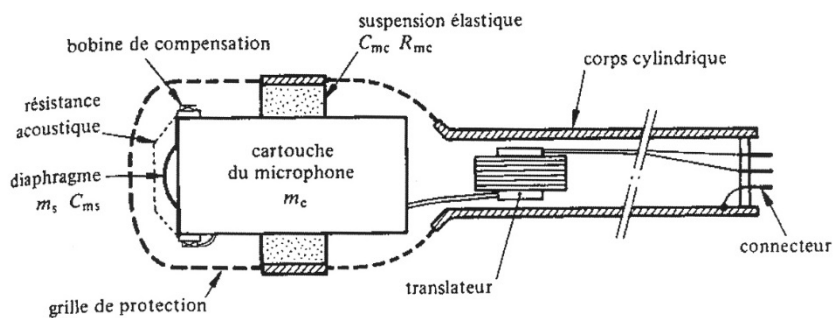
# Moving coil electrodynamic: omnidirectional



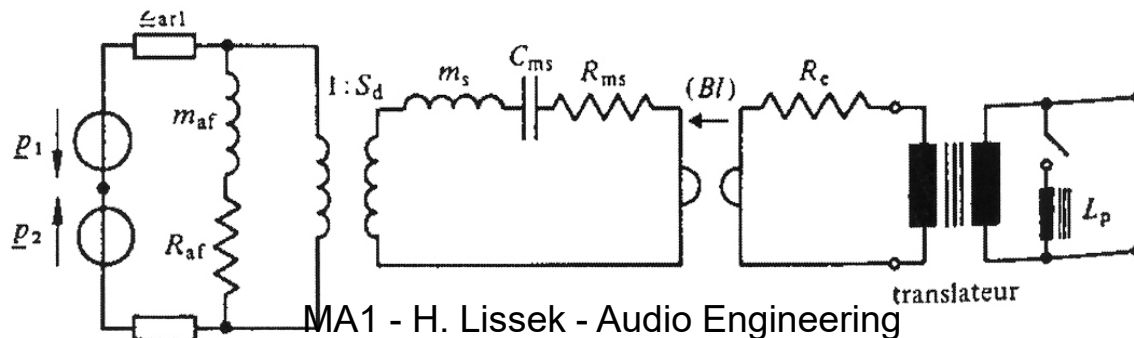
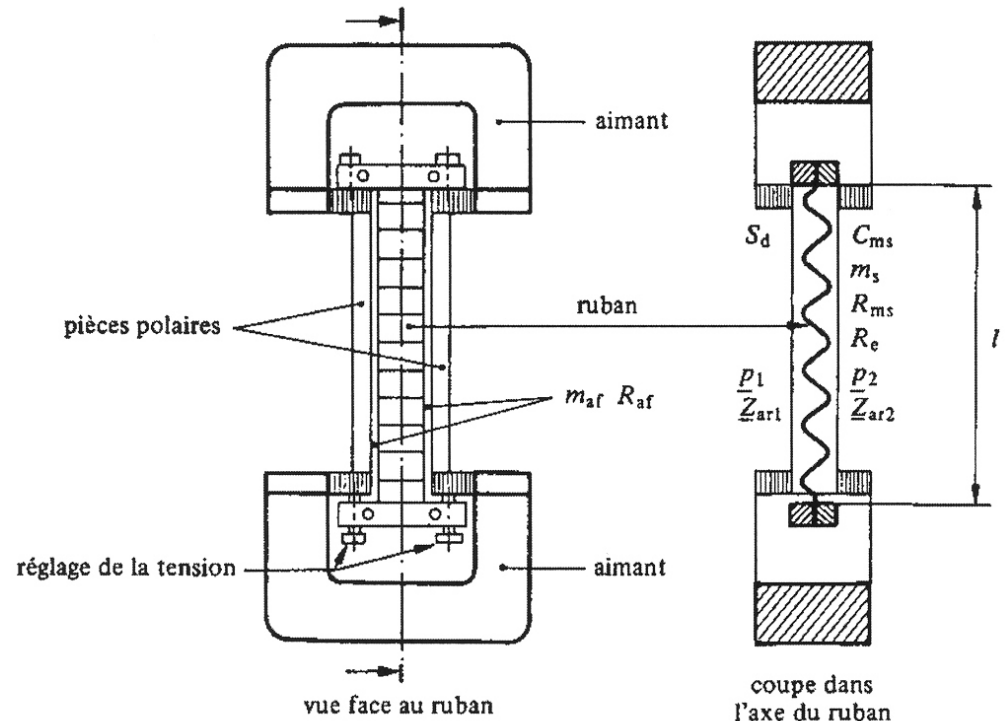
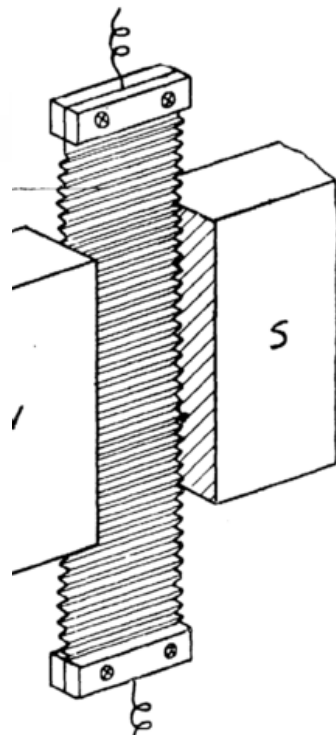
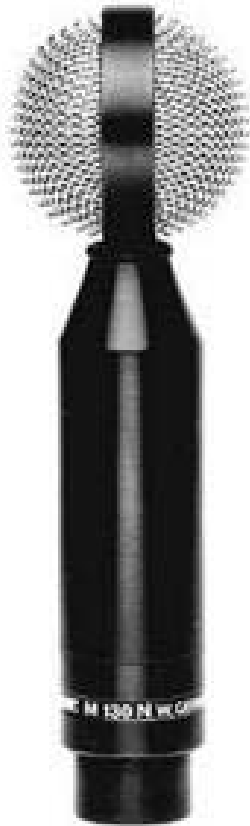
# Moving coil electrodynamic: hypercardioid



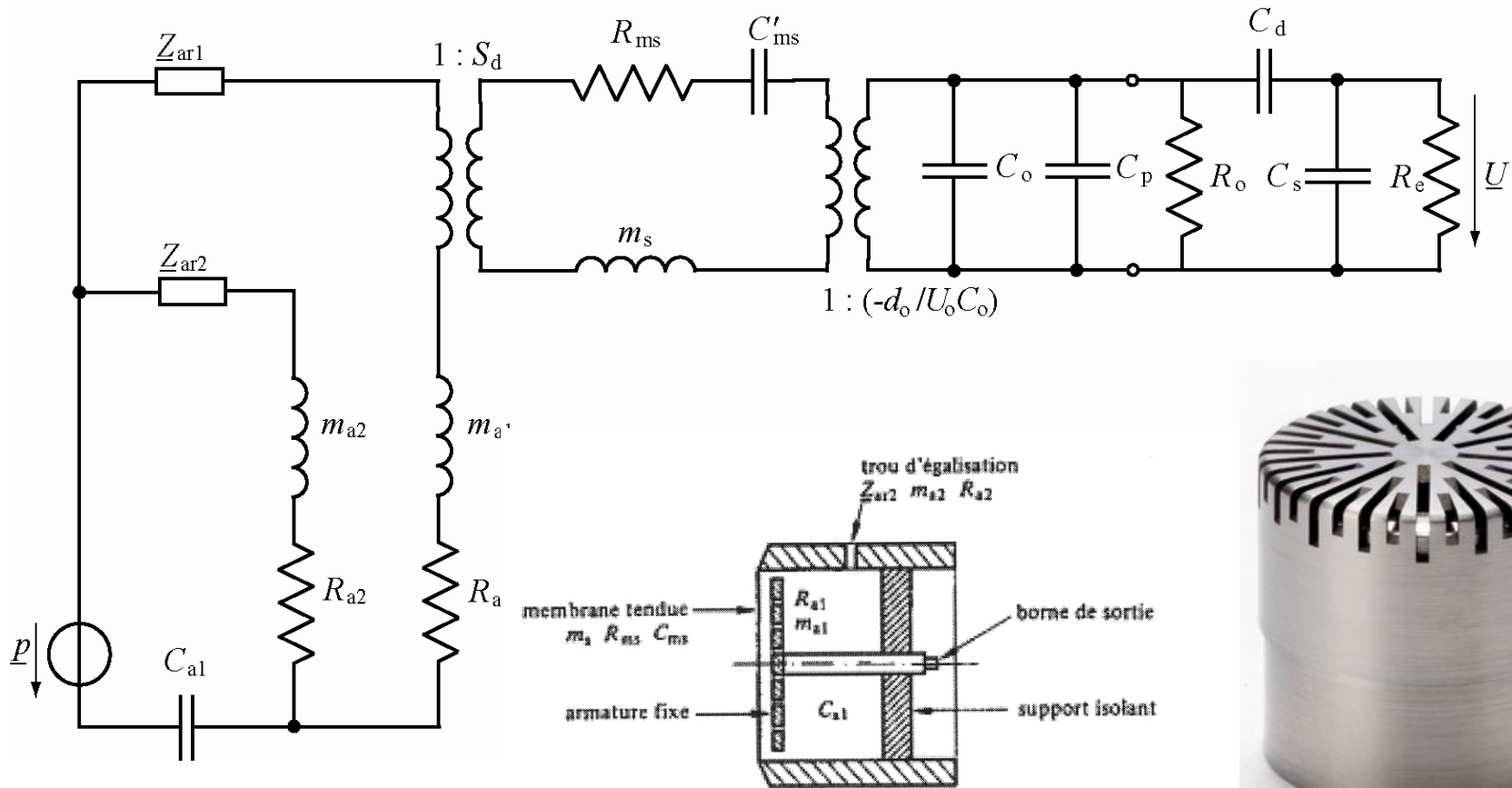
# Moving coil electrodynamic: mechanical protection



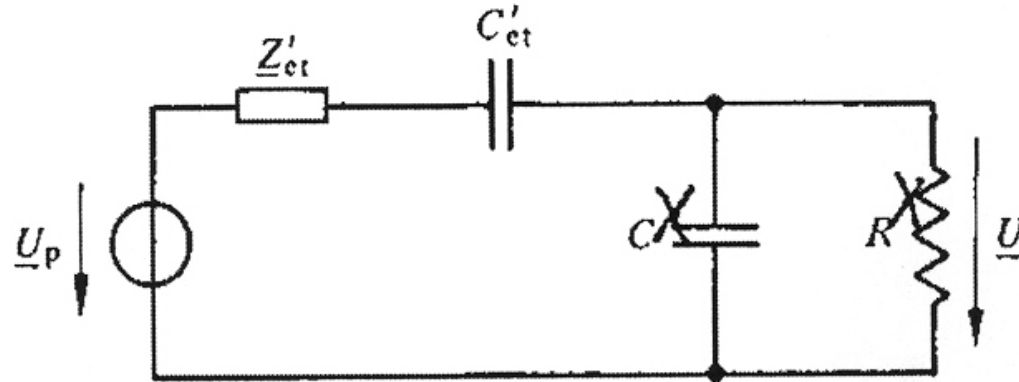
# Ribbon microphone (electrodynamical)



# Measurement electrostatic: omnidirectional

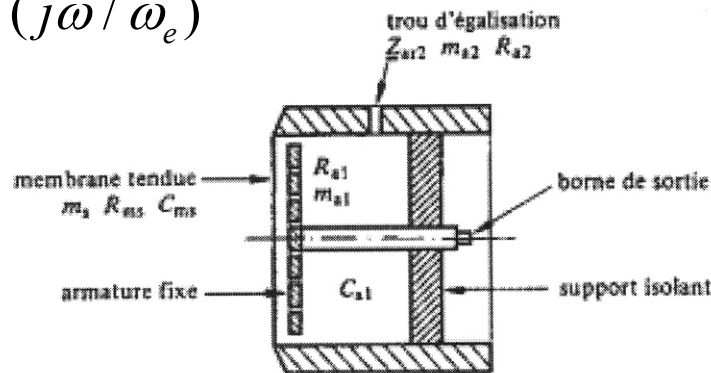


# Measurement electrostatic: omnidirectional

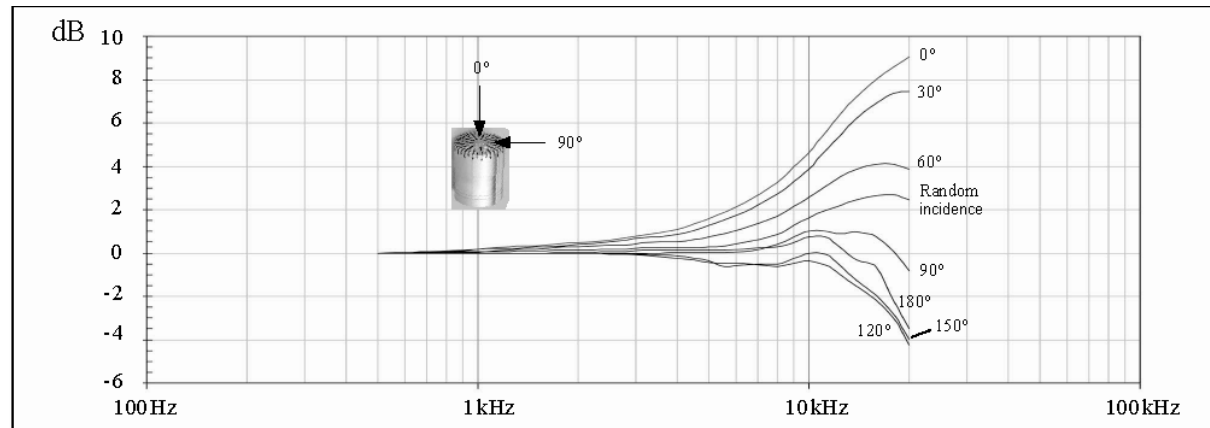
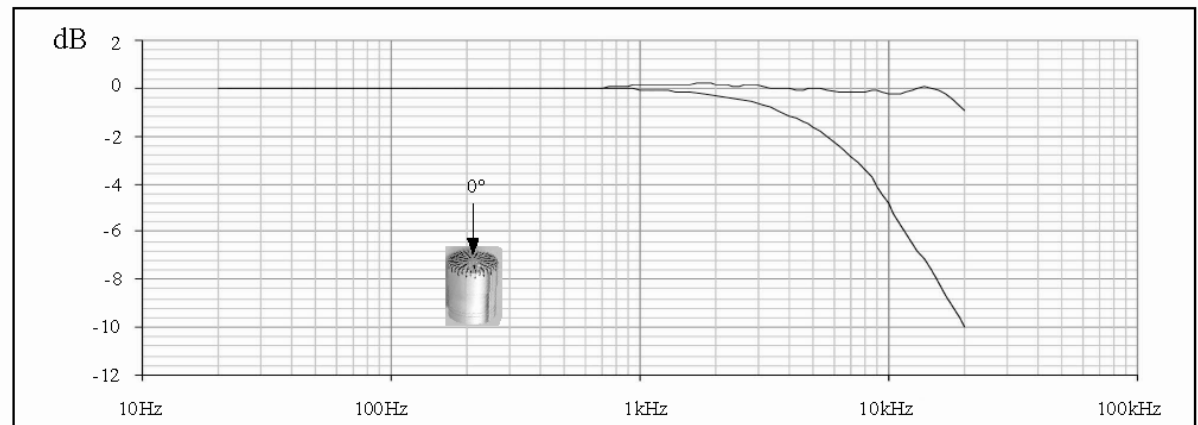


$$\frac{\underline{U}}{\underline{p}} = \frac{U_0 C'_{mt} S_d}{d_0} \cdot \frac{C_0}{C} \cdot \frac{(j\omega / \omega_e)}{1 + (j\omega / \omega_e)}$$

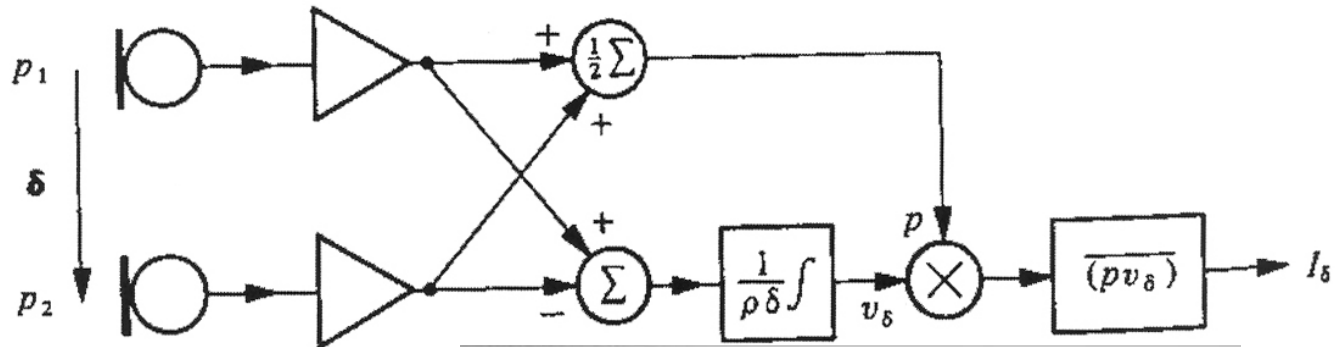
$$\omega_e = \frac{1}{RC}$$



# Measurement electrostatic: omnidirectional



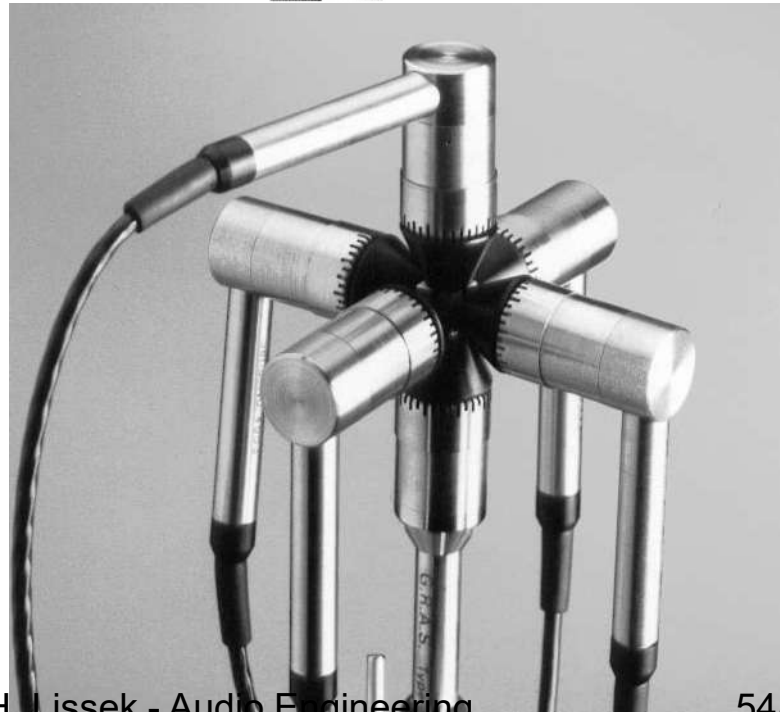
# Sound intensity measurement



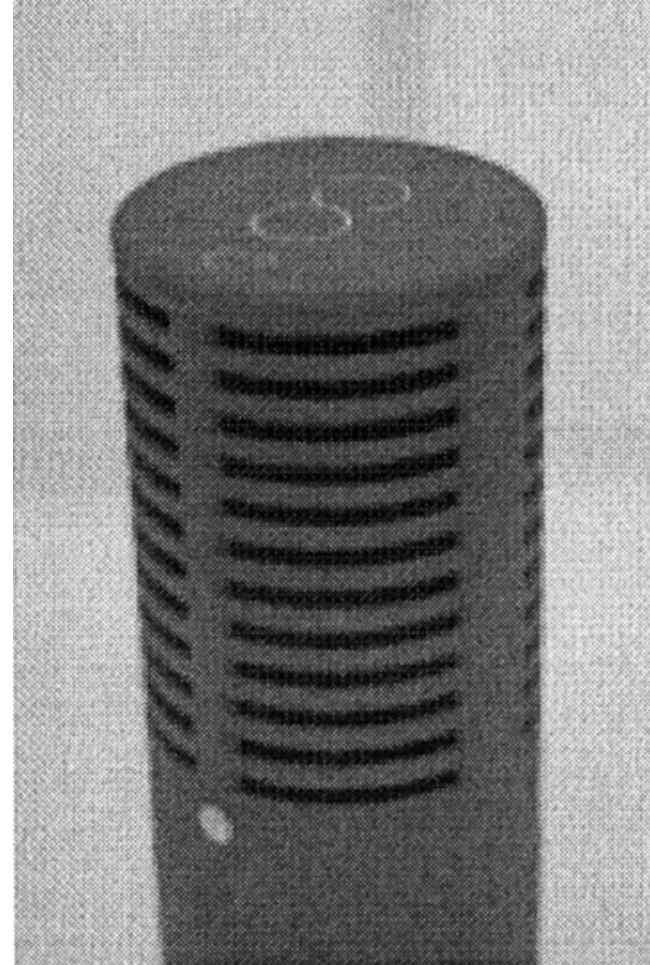
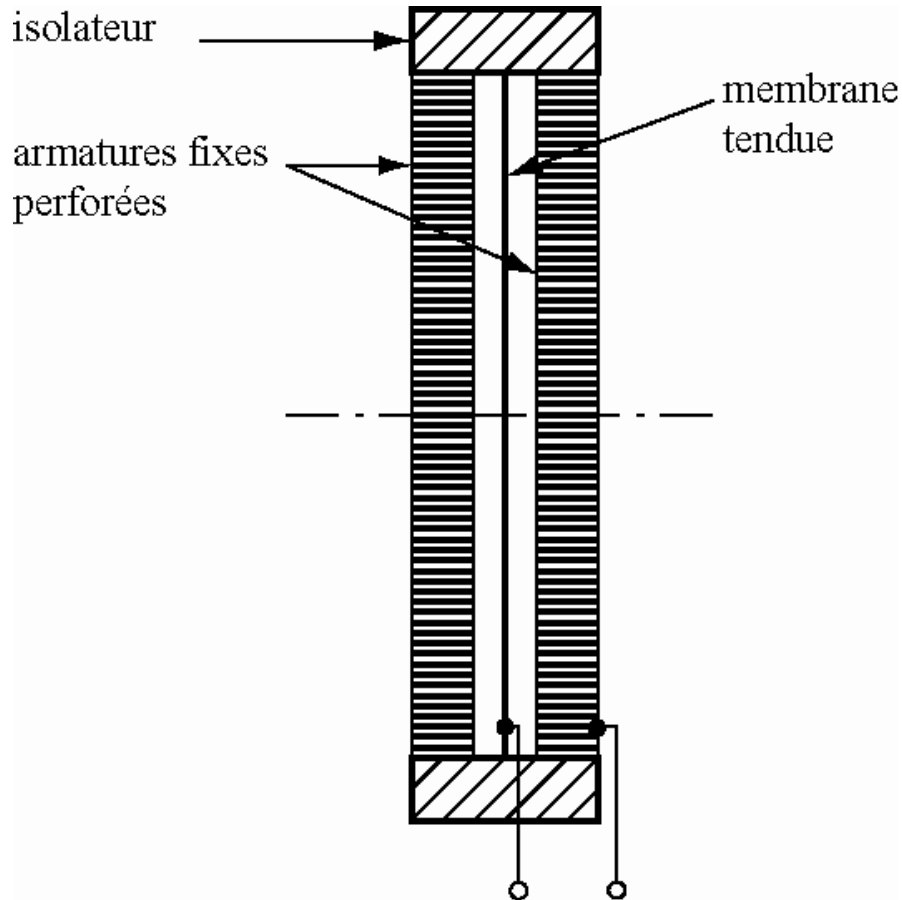
$$v_\delta = -\frac{1}{\rho} \int \partial_\delta p dt$$

$$= +\frac{1}{\rho \delta} \int (p_1 - p_2) dt$$

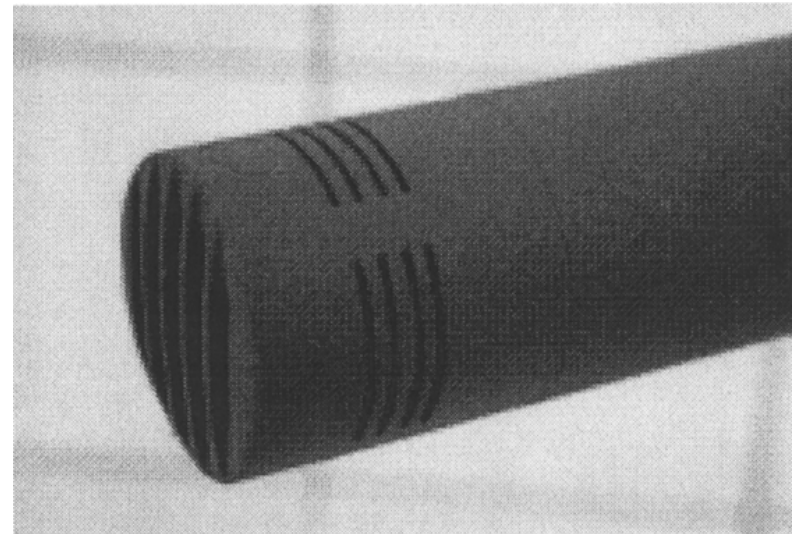
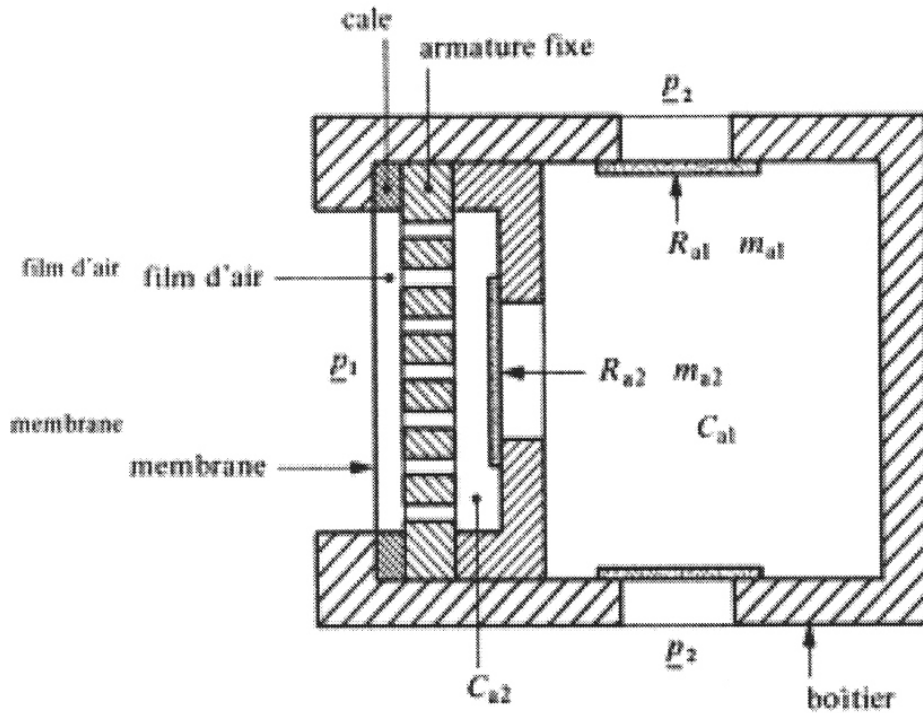
$$I_\delta = \frac{1}{2\rho\delta} \left[ \overline{(p_1 + p_2) \cdot \int (p_1 - p_2) dt} \right]$$



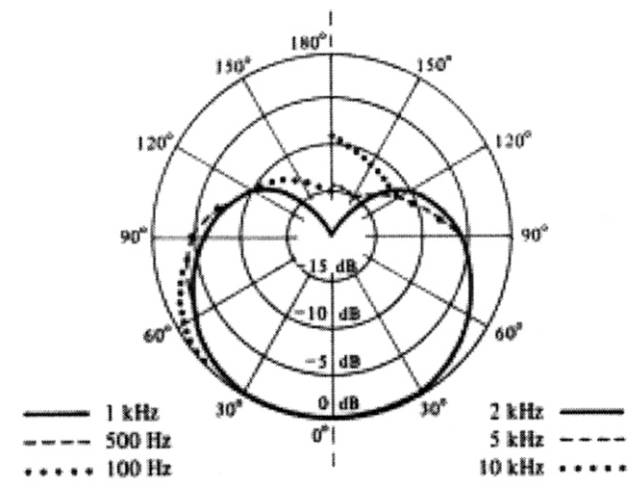
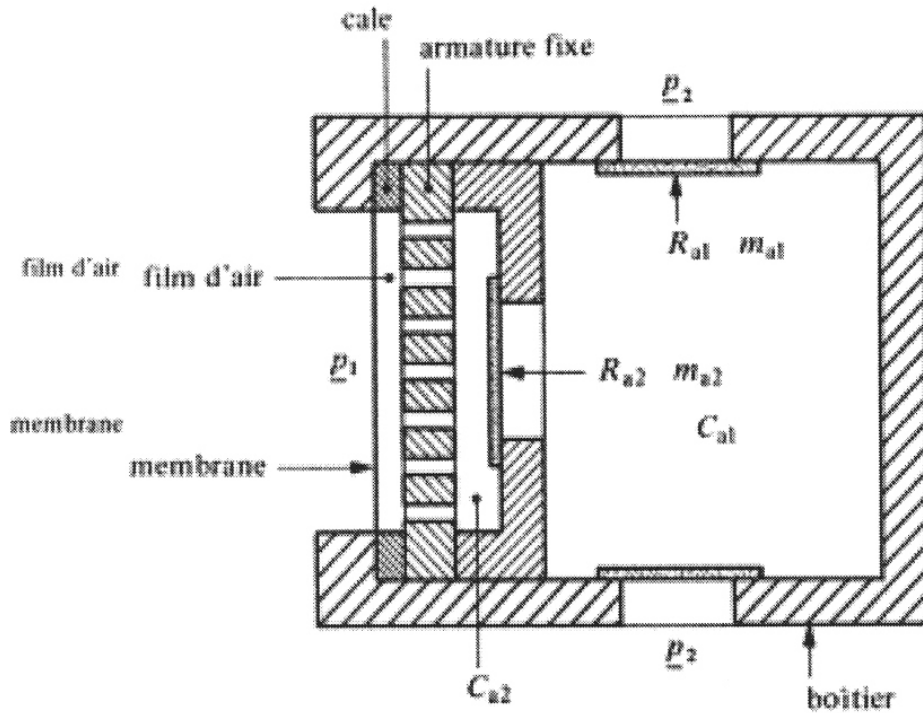
# Electrostatic: pressure gradient



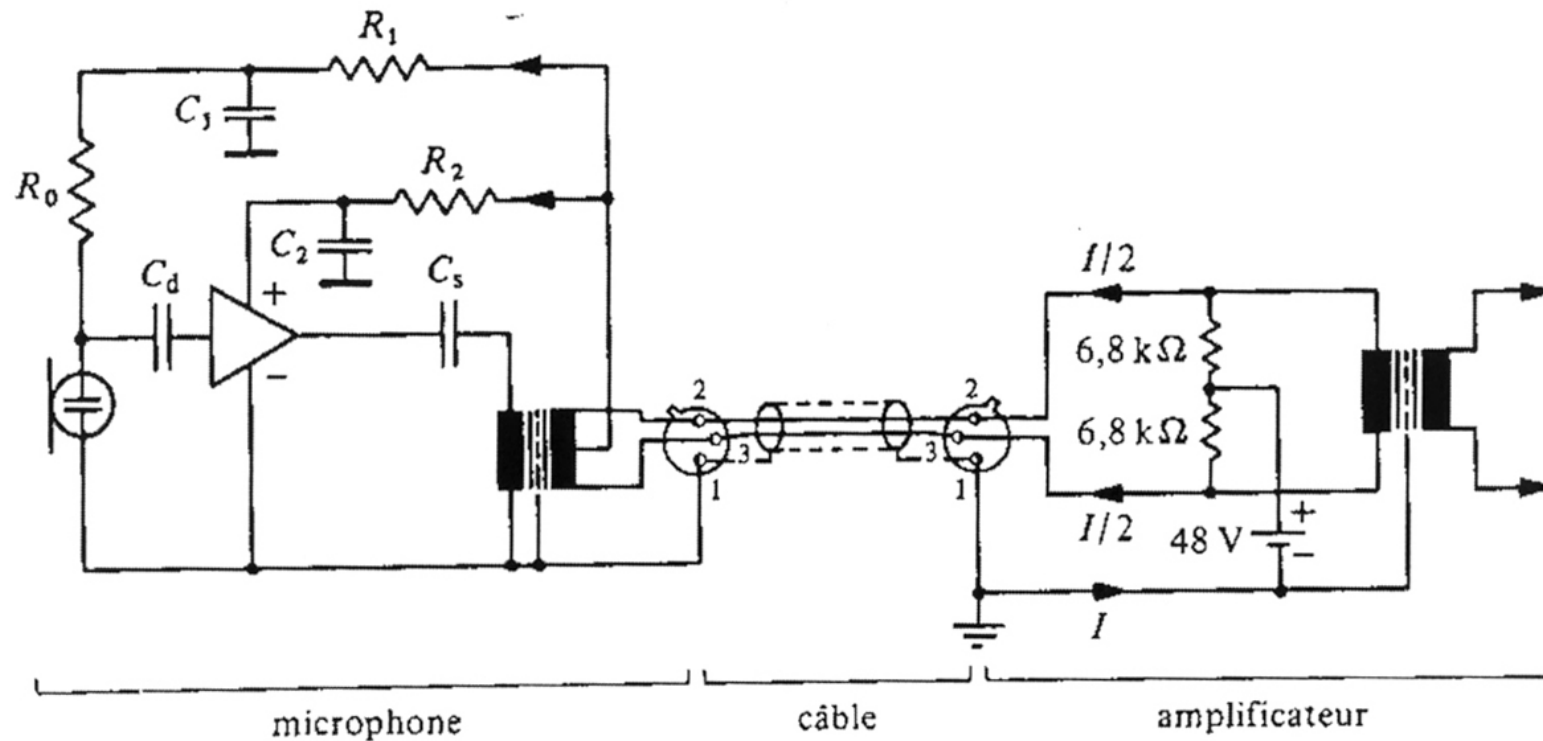
# Electrostatic: unidirectional



# Electrostatic: unidirectional



# Electrostatic: power supply



# Electret microphone

