

Parametrizing rotations

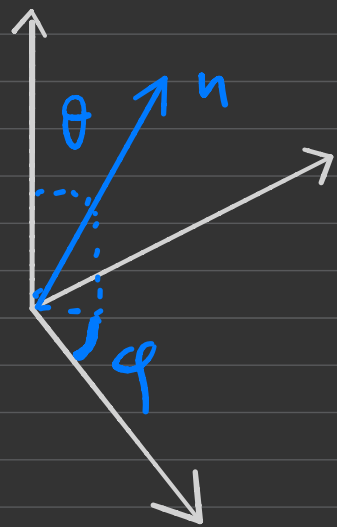
$$1) R(\underline{\theta}) = e^{-i \vec{\theta} \cdot \vec{J}}$$

$$\vec{\theta} = (\theta_1, \theta_2, \theta_3)$$

$$2) R(\varphi, \theta, \psi) = e^{-i\varphi J_3} e^{-i\theta J_2} e^{-i\psi J_3}$$

$\varphi, \theta, \psi \equiv$
Euler angles

2')



$$\vec{n} = (\sin\theta \cos\varphi, \sin\theta \sin\varphi, \cos\theta)$$

$$R(\vec{n}) \equiv R(\varphi, \theta, -\varphi) = e^{-i\varphi J_3} e^{-i\theta J_2} e^{i\varphi J_3}$$

$$\mathcal{R}(\vec{n})^{i3} = n^i$$

or

$$\boxed{\mathcal{R} \hat{z} = \vec{n}}$$

$$\mathcal{R}^{ij} z^j = \mathcal{R}^{i3}$$

$$\hat{z} = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$\hat{z}_i = \delta_{i3}$$

B) The helicity basis

Jacob J. Wick
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$$\vec{n} \equiv \frac{\vec{p}}{|\vec{p}|}, \quad \vec{p} = R(\vec{n}) \begin{pmatrix} 0 \\ 0 \\ p \end{pmatrix}$$

$$\begin{pmatrix} p^0 \\ 0 \\ 0 \\ p \end{pmatrix} = \begin{pmatrix} \sqrt{u^2 + p^2} \\ 0 \\ 0 \\ p \end{pmatrix} = e^{i\eta K_3} \begin{pmatrix} u \\ 0 \\ 0 \\ 0 \end{pmatrix}$$

$$p \equiv |\vec{p}| = \frac{tgh\eta}{\sqrt{1-tgh^2\eta}} u$$

$$\eta = tgh^{-1}\left(\frac{p}{p_0}\right)$$

$$\Lambda_{p\vec{p}}^h \equiv R(\vec{n}) e^{i\eta K_3}$$

$$(\vec{J} \cdot \vec{n}) U(R) = U(R) U(R)^{-1} (\vec{J} \cdot \vec{n}) U(R) = U(R) R_{ij} J^j n^i$$

$$\Rightarrow (\vec{J} \cdot \vec{n}) \Lambda_{\vec{P}}^h = \Lambda_{\vec{P}}^h J_3$$

$$\bullet |P, \sigma\rangle_h \equiv U(\Lambda_{\vec{P}}^h) |\bar{P}, \sigma\rangle$$

$$J \cdot n |P, \sigma\rangle = (J \cdot n) U(\Lambda_{\vec{P}}^h) |\bar{P}, \sigma\rangle = U(\Lambda_{\vec{P}}^h) J_3 |\bar{P}, \sigma\rangle$$

$$J \cdot n |P, \sigma\rangle = \sigma |P, \sigma\rangle$$

$$\begin{aligned}
 \textcircled{1} \quad R(\vec{n}) e^{i\eta K_3} &= e^{i\eta R(\vec{n}) K_3 R(\vec{n})^\dagger} \cdot R(\vec{n}) \\
 &= e^{i\eta \mathcal{R}_{i3} K_i} \quad R(\vec{n}) = e^{i\vec{\eta} \cdot \vec{K}} R(\vec{n})
 \end{aligned}$$

$$\begin{aligned}
 |P, \sigma\rangle_h &\equiv U\left(R(\vec{n}) e^{i\eta K_3}\right) |\vec{P}, \sigma\rangle = U\left(e^{i\vec{\eta} \cdot \vec{K}}\right) U(R(\vec{n})) |\vec{P}, \sigma\rangle \\
 &= U\left(e^{i\vec{\eta} \cdot \vec{K}}\right) |\vec{P}, \sigma'\rangle D_{\sigma' \sigma}^{(s)}(R(\vec{n})) \\
 &= |\vec{P}, \sigma'\rangle_s D_{\sigma' \sigma}^{(s)}(R(\vec{n}))
 \end{aligned}$$

Wigner Rotation in helicity basis

pure rotation $R \equiv R(\alpha, \beta, \gamma)$

$$\underline{R(\alpha, \beta, \gamma) \Lambda_I^h} = R(\alpha, \beta, \gamma) R(\varphi, \theta, -\varphi) e^{i\eta K_3}$$

$$= R(\varphi', \theta', \psi') e^{i\eta K_3}$$

$$= e^{-i\varphi' J_3} e^{-i\theta' J_2} e^{-i\psi' J_3} e^{i\eta K_3}$$

$$= \underbrace{e^{-i\varphi' J_3} e^{-i\theta' J_2} e^{i\varphi' J_3}}_{R(n')} e^{i\eta K_3} e^{-i(\varphi' + \psi') J_3}$$

$$= R(n') e^{i\eta K_3} e^{-i(\varphi' + \psi') J_3}$$

$$= \sum_{\vec{p}'}^h \overbrace{e^{-i(q' + \varphi')J_3}}$$

$W \equiv$ Wigner rotation

$$U(R(a, \vec{p}, \sigma)) |\vec{p}', \sigma\rangle_h = U(R) U(\Lambda_{\vec{p}}^h) |\vec{p}, \sigma\rangle$$

$$\vec{p}' \equiv R\vec{p}$$

$$= U(\Lambda_{R\vec{p}}) e^{-i(q' + \varphi')\sigma} |\vec{p}, \sigma\rangle$$

$$= e^{-i(q' + \varphi')\sigma} |R\vec{p}, \sigma\rangle$$

$$e^{i\alpha \vec{J} \cdot \vec{n}} = R(\vec{n}) e^{i\alpha J_3} R^{-1}(\vec{n})$$

$$|\vec{P}', \sigma\rangle_n = \underbrace{U(R(\vec{n}'))}_{\substack{\text{rotates } \vec{P}' \\ \text{into } \vec{P}}} e^{i\eta k_3} |\vec{P}, \sigma\rangle$$

$$e^{i\alpha \vec{J} \cdot \vec{n}} |\vec{P}', \sigma\rangle_n = \underbrace{U(R(\vec{n}')) e^{i\alpha J_3} U(R^{-1}(\vec{n}')) U(R(\vec{n}))}_{\substack{\text{rotates } \vec{P}' \\ \text{into } \vec{P}}} e^{i\eta k_3} |\vec{P}, \sigma\rangle$$

$$= U(R(\vec{n}')) e^{i\alpha J_3} e^{i\eta k_3} |\vec{P}, \sigma\rangle$$

$$= U(R(\vec{n}')) e^{i\eta k_3} |\vec{P}, \sigma\rangle e^{i\alpha \sigma}$$

$$= |\vec{P}, \sigma\rangle e^{i\alpha \sigma}$$

⊙ pure boost

$$\underline{H_h(P) \equiv U(\Lambda_{\vec{P}}^h)}$$

$$\Lambda_{\vec{P}}^h = R(n) e^{i\eta K_3}$$

$$W_h(\Lambda, P) \equiv H_h(\Lambda P)^{\dagger} U(\Lambda) H_h(P) \quad P' = \Lambda P$$

$$= e^{-i\eta' K_3} U^{\dagger}(R(n')) U(\Lambda) \underbrace{U(R(n)) e^{i\eta K_3}}$$

$$= U^{\dagger}(R(n')) \underbrace{e^{-i\vec{\eta}' \cdot \vec{K}} U(\Lambda) e^{i\vec{\eta} \cdot \vec{K}}}_{U(R(\vec{n}))} U(R(n))$$

$$\underline{U^{\dagger}(R(n'))} \quad \underline{W_s(\Lambda, P)} \quad \underline{U(R(\vec{n}))}$$

Last time

$$W_S(\Lambda, P) = \mathbb{1} + \frac{i(\vec{\ell}_\Lambda \wedge \vec{P}) \cdot \vec{J}}{p^0 + \omega} + \dots$$

$$\Lambda = \mathbb{1} + \omega \dots$$

$$\downarrow$$

$$\eta_\Lambda$$

• In helicity basis ($p^0 \gg \omega$)

$$\vec{n} = \frac{\vec{P}}{|\vec{P}|}$$

$$W_h = \mathbb{1} + \frac{i(\vec{\ell}_\Lambda \wedge \vec{n})_3 J_3}{1 + h_3} + \mathcal{O}\left(\frac{\omega}{p_0}\right)$$

$$U(\Lambda) |p, \sigma\rangle_h = U(\Lambda) H^h(p) |\bar{p}, \sigma\rangle$$

$$= H^h(\Lambda p) U(\omega_h(\Lambda, p)) |\bar{p}, \sigma\rangle$$

$$= H^h(\Lambda p) |\bar{p}, \sigma'\rangle D_{\sigma'\sigma}(\omega_h(\Lambda, p))$$

$$= |\Lambda p, \sigma'\rangle_h \left(\delta_{\sigma'\sigma} + i \frac{(\vec{n}_\Lambda \wedge \vec{n})_3}{1 + n_3} \sigma^3_{\sigma'\sigma} + \dots \right)$$

$$= |\Lambda p, \sigma\rangle_h \left(1 + i \frac{(\vec{n}_\Lambda \wedge \vec{n})_3}{1 + n_3} \cdot \sigma + \dots \right)$$

▲ In spin basis, action of pure rotation R , described by $W=R$, like in non-relativistic QM:

$$U(R)|\vec{p}, \sigma\rangle = |\vec{p}, \sigma'\rangle D(R)_{\sigma'\sigma}$$

$\downarrow \quad \quad \downarrow$
 $L \quad \quad S$

▲ In helicity basis spin is referred to $\vec{p}$ in lab frame: $W \neq R$

$$\text{II}_A) \quad \bar{P}^\mu = (K, 0, 0, K)$$

$$W^0 = \frac{1}{2} \varepsilon^{\alpha\mu\nu\rho} J_{\mu\nu} P_\rho = \frac{1}{2} \left(\varepsilon^{0123} J_{12} P_3 + \varepsilon^{0213} J_{21} P_3 \right)$$

$$= -J_3 \cdot K$$

$$W^3 = -J_3 K = \frac{1}{2} \varepsilon^{3120} J_{12} K \cdot 2$$

$$W^1 = K (K^2 - J^1)$$

$$[J_3, J_1] = iJ_2 \quad [J_3, J_2] = -iJ_1$$

$$W^2 = K (-K^1 - J^2)$$

$$[W^1, W^2] = \text{exercise} = -k^2 (iJ^3 - iJ^3) = 0$$

$$[W^3, W^1] = -ik W^2$$

$$[W^3, W^2] = +ik W^1$$

$$[J^3, W^1] = i W^2$$

$$[J^3, W^2] = -i W^1$$

$$\frac{W^3}{k} = J^3$$

$$[J^3, W^1] = i W^2$$

$$[J^3, W^2] = -i W^1$$

$$[W^1, W^2] = 0$$

ISO(2)

Convenient

$$W^{\pm} \equiv W' \pm i W^2$$

$$W^- = (W^+)^{\dagger}$$

$$\left[\begin{array}{l} [J^3, W^{\pm}] = \pm W^{\pm} \\ [W^{\pm}, W^{\pm}] = 0 \end{array} \right]$$

• $W^2 = W^{\mu} W_{\mu} \equiv \text{Casimir}$

$$= \cancel{(W^0)^2} - (W_1^2 + W_2^2 + \cancel{W_3^2}) = -W_1^2 - W_2^2 = -W^+ W^- \\ = -W^- W^+$$

$$\underline{W^2} \equiv \underline{-W^+ W^-} = -c \underline{1}$$

▲ $|\bar{P}, \sigma\rangle$ must host a unitary repr. of $ISO(2)$

• $ISO(2)$ is non compact and non semi-simple
 $\Rightarrow$ unitary irreps either $\dim=1$ or $\dim=\infty$

• $[J_3, W^\pm] = \pm W^\pm \Rightarrow W^\pm$ raising and lowering J_3

$$J_3 |2\rangle = 2|2\rangle \Rightarrow J_3 W^+ |2\rangle = (2+1) W^+ |2\rangle$$

$$J_3 W^- |2\rangle = (2-1) W^- |2\rangle$$

$$W^- W^+ = W^+ W^- = (W^+)^{\dagger} W^+ = C \cdot \mathbb{1} \quad C \geq 0$$

$$\langle \psi | A^\dagger A | \psi \rangle \geq 0$$

$$\langle \psi | A^\dagger A | \psi \rangle = \sum_n |\langle n | A | \psi \rangle|^2 \geq 0$$

• $c > 0$

$$(W^+)^+ W^+ = c \cdot \mathbb{1} \quad c > 0$$

$$\left(\frac{W^+}{\sqrt{c}}\right)^+ \left(\frac{W^+}{\sqrt{c}}\right) = \mathbb{1} = U^+ U$$

$$\frac{W^+}{\sqrt{c}} \equiv V = \text{unitary}$$

Toro Schuster

$$\Rightarrow \dots, \lambda-2, \lambda-1, \lambda, \lambda+1, \lambda+2, \dots$$

all eigenvalues !!

irrep is ∞ dimensional!

$$\Delta C = 0$$

$$(W^+)^+ W^+ = 0 \Rightarrow W^+ = W^- = 0$$

$$J_3 |2\rangle = 2 |2\rangle$$

$$W^\pm |2\rangle = 0$$

$\Rightarrow$ 1-dimensional irrep

▲ $C=0$

$$J^3 |\bar{p}, \lambda\rangle = \lambda |\bar{p}, \lambda\rangle$$

$$W^\pm |\bar{p}, \lambda\rangle = 0$$



$$W^\mu |\bar{p}, \lambda\rangle = -\lambda \bar{p}^\mu |\bar{p}, \lambda\rangle$$

$$\cdot \underline{W}^\mu = -\lambda \underline{\bar{p}}^\mu = -\lambda (\kappa, 0, 0, \kappa)$$

$$\cdot \lambda = \frac{\vec{J} \cdot \vec{\bar{p}}}{|\bar{p}|}$$

lossel

▲ Casimirs:

$$P^2 = W^2 = 0$$

$$\underline{W}^\mu = -\lambda \underline{\bar{p}}^\mu$$

$$\lambda \equiv \text{losseltz inu}$$

▲ Induced representation

$$\begin{pmatrix} k \\ 0 \\ 0 \\ k \end{pmatrix} \rightarrow \begin{pmatrix} k \\ k \end{pmatrix} \xrightarrow{\text{boost}} \begin{pmatrix} \cosh \eta & \sinh \eta \\ \sinh \eta & \cosh \eta \end{pmatrix} \begin{pmatrix} k \\ k \end{pmatrix}$$

$$= e^{\eta} \begin{pmatrix} k \\ k \end{pmatrix} \equiv \begin{pmatrix} P \\ P \end{pmatrix}$$

$$\eta_P \equiv \ln \frac{P}{k}$$

$$|\vec{P}, \lambda\rangle = \frac{U(\vec{R}(\vec{n})) e^{i\eta_P \cdot K_3}}{P^0} |\vec{P}, \lambda\rangle$$

$$\vec{P} = \vec{n} \cdot \vec{P}$$

$$P^0 \equiv P$$

$$\hookrightarrow U(\Lambda_P) |\bar{P}, \lambda\rangle$$

given

$$U^\dagger(\Lambda) P^\mu U(\Lambda) = \Lambda^\mu_\nu P^\nu$$

$$U^\dagger(\Lambda) W^\mu U(\Lambda) = \Lambda^\mu_\nu W^\nu$$

$$W^\mu = -\frac{1}{2} P^\mu$$

true for all P^μ

(of course $P^\mu P_\mu = 0$)

proof

$\Rightarrow$

$$W^\mu |P, \lambda\rangle = U(\Lambda_P) U(\Lambda_P)^\dagger W^\mu U(\Lambda_P) |\bar{P}, \lambda\rangle$$

$$p^\mu = \Lambda^\mu_\nu \bar{p}^\nu$$

$$= U(\Lambda_P) \Lambda^\mu_\nu \omega^\nu |\bar{P}, \lambda\rangle$$

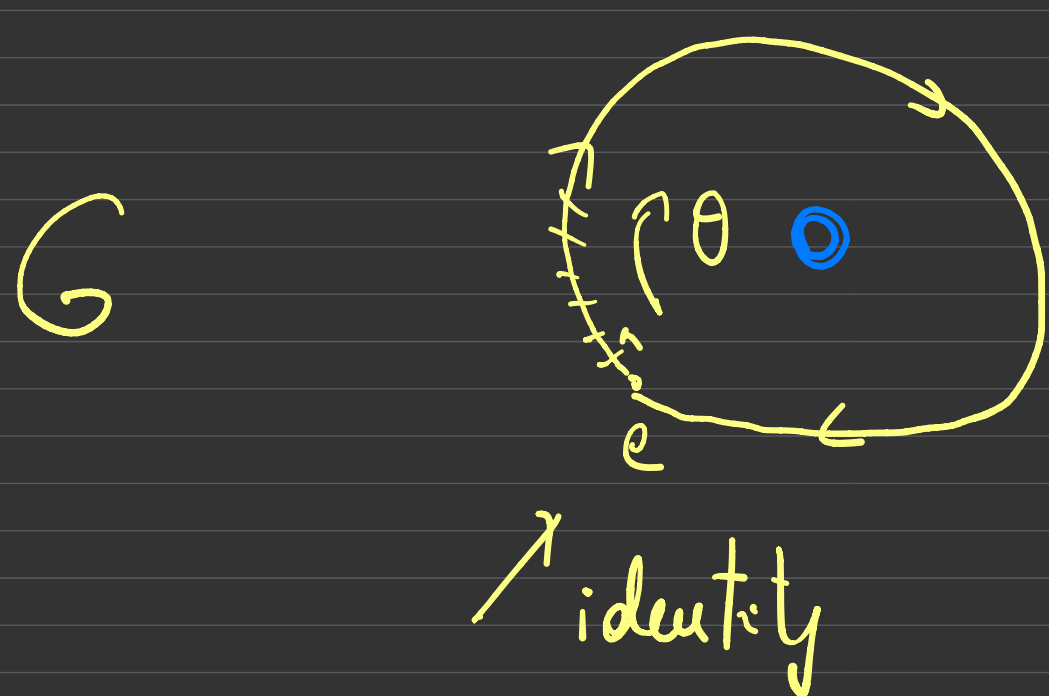
$$= U(\Lambda_P) \Lambda^\mu_\nu (-\lambda \bar{p}^\nu) |\bar{P}, \lambda\rangle$$

$$= -\lambda p^\mu \underbrace{U(\Lambda_P) |\bar{P}, \lambda\rangle}_{|P, \lambda\rangle}$$

$$= -\lambda p^\mu |P, \lambda\rangle$$

$\lambda \equiv$ Lorentz invariant

Quantization of λ



$$e^{i\theta X} \quad g_e$$

$$g(\theta) \quad g(0) = e \quad g(2\pi) = e$$

$$U(g(\theta)) = e^{i\theta X} = e^{i2\pi X} = e^{i\varphi} \quad U(e) = e^{i\varphi} \cdot 1$$

$$e^{i4\pi J_3} = 1 \quad e^{i4\pi J_2} \quad 4\pi J = 2\pi n$$

$\lambda = n/2$

CPT & anti-particles



$$\Theta \vec{j} \Theta = -\vec{j}$$

$$\Theta \vec{p} \Theta = \vec{p}$$

$$\Theta \vec{j} \cdot \vec{p} \Theta = -\vec{j} \cdot \vec{p}$$

$$\lambda \longrightarrow -\lambda$$

$$\lambda = \pm 1/2$$

$$\lambda = \pm 1$$

$$\lambda = \pm 2$$

$$\lambda = \pm 3/2$$

$\leftarrow ISO(3,1) + CPT$

Ex

{
photon
graviton
gravitino

massive $3/2$ $3/2, 1/2, -1/2, -3/2$