## Particle Physics 1 : Exercise 9

## 1) Properties of the chirality operator

Using the properties of the  $\gamma$ -matrices and the definition of the chirality operator  $\gamma^5 = i\gamma^0\gamma^1\gamma^2\gamma^3$ , show that

- (a)  $(\gamma^5)^2 = 1$
- (b)  $\gamma^{5\dagger} = \gamma^5$
- (c)  $\gamma^5 \gamma^\mu = -\gamma^\mu \gamma^5$

## 2) Chiral projection operators

Show that the chiral projection operators  $P_R = \frac{1}{2}(1 + \gamma^5)$  and  $P_L = \frac{1}{2}(1 - \gamma^5)$  are indeed projection operators, i.e. they satisfy the following properties:

$$P_R + P_L = 1$$
,  $P_R P_R = P_R$ ,  $P_L P_L = P_L$ ,  $P_R P_L = 0$ 

## 3) Differential cross section for $e^-\mu^- \to e^-\mu^-$ scattering

Using helicity amplitudes, calculate the differential cross section for  $e^-\mu^- \to e^-\mu^-$  scattering in the following steps

(a) From the Feynman rules of QED, show that the lowest-order QED matrix element for  $e^-\mu^- \to e^-\mu^-$  is

$$\mathcal{M}_{fi} = -\frac{e^2}{(p_1 - p_3)^2} g_{\mu\nu} \left[ \overline{u}(p_3) \gamma^{\mu} u(p_1) \right] \left[ \overline{u}(p_4) \gamma^{\nu} u(p_2) \right],$$

where  $p_1$  and  $p_3$  are the four-momenta of the initial- and final-state  $e^-$ , and  $p_2$  and  $p_4$  are the four-momenta of the initial- and final-state  $\mu^-$ .

(b) Working in the centre-of-mass frame, and writing the four-momenta of the initial and final-state  $e^-$  as  $p_1^{\mu} = (E_1, 0, 0, p)$  and  $p_3^{\mu} = (E_1, p \sin \theta, 0, p \cos \theta)$  respectively, show that the electron currents for the four possible helicity combinations are

$$\overline{u}_{\downarrow}(p_3)\gamma^{\mu}u_{\downarrow}(p_1) = 2(E_1c, p_s, -ip_s, p_c) 
\overline{u}_{\uparrow}(p_3)\gamma^{\mu}u_{\downarrow}(p_1) = 2(m_s, 0, 0, 0) 
\overline{u}_{\uparrow}(p_3)\gamma^{\mu}u_{\uparrow}(p_1) = 2(E_1c, p_s, ip_s, p_c) 
\overline{u}_{\downarrow}(p_3)\gamma^{\mu}u_{\uparrow}(p_1) = -2(m_s, 0, 0, 0)$$

where m is the electron mass,  $s = \sin \theta/2$  and  $c = \cos \theta/2$ .

(c) Explain why the effect of the parity operator  $\hat{P} = \gamma^0$  is

$$\hat{P}u_{\uparrow}(\mathbf{p},\theta,\phi) = u_{\downarrow}(\mathbf{p},\pi-\theta,\pi+\phi).$$

Hence, or otherwise, show that the muon currents for the four helicity combinations are

$$\overline{u}_{\downarrow}(p_4)\gamma^{\mu}u_{\downarrow}(p_2) = 2(E_2c, -p_s, -ip_s, -p_c) 
\overline{u}_{\uparrow}(p_4)\gamma^{\mu}u_{\downarrow}(p_2) = -2(M_s, 0, 0, 0) 
\overline{u}_{\uparrow}(p_4)\gamma^{\mu}u_{\uparrow}(p_2) = 2(E_2c, -p_s, ip_s, -p_c) 
\overline{u}_{\downarrow}(p_4)\gamma^{\mu}u_{\uparrow}(p_2) = 2(M_s, 0, 0, 0),$$

where M is the muon mass and  $E_2 = \sqrt{\mathbf{p}^2 + M^2}$  is the muon energy in the centre-of-mass frame.

(d) For the relativistic limit where  $E \gg M$ , show that the matrix element squared for the case where the incoming  $e^-$  and incoming  $\mu^-$  are both left-handed is given by

$$|\mathcal{M}_{LL}|^2 = \frac{4e^4s^2}{(p_1 - p_3)^4},$$

where  $s = (p_1 + p_2)^2$ . Find the corresponding expressions for  $|\mathcal{M}_{RL}|^2$ ,  $|\mathcal{M}_{RR}|^2$  and  $|\mathcal{M}_{LR}|^2$ .

(e) In this relativistic limit, show that the differential cross section for unpolarised  $e^-\mu^- \to e^-\mu^-$  scattering in the centre-of-mass frame is

$$\frac{d\sigma}{d\Omega} = \frac{2\alpha^2}{s} \cdot \frac{1 + \frac{1}{4}(1 + \cos\theta)^2}{(1 - \cos\theta)^2}.$$