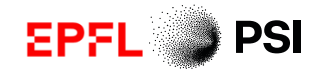


Materials Science at Large Scale Facilities

Sources

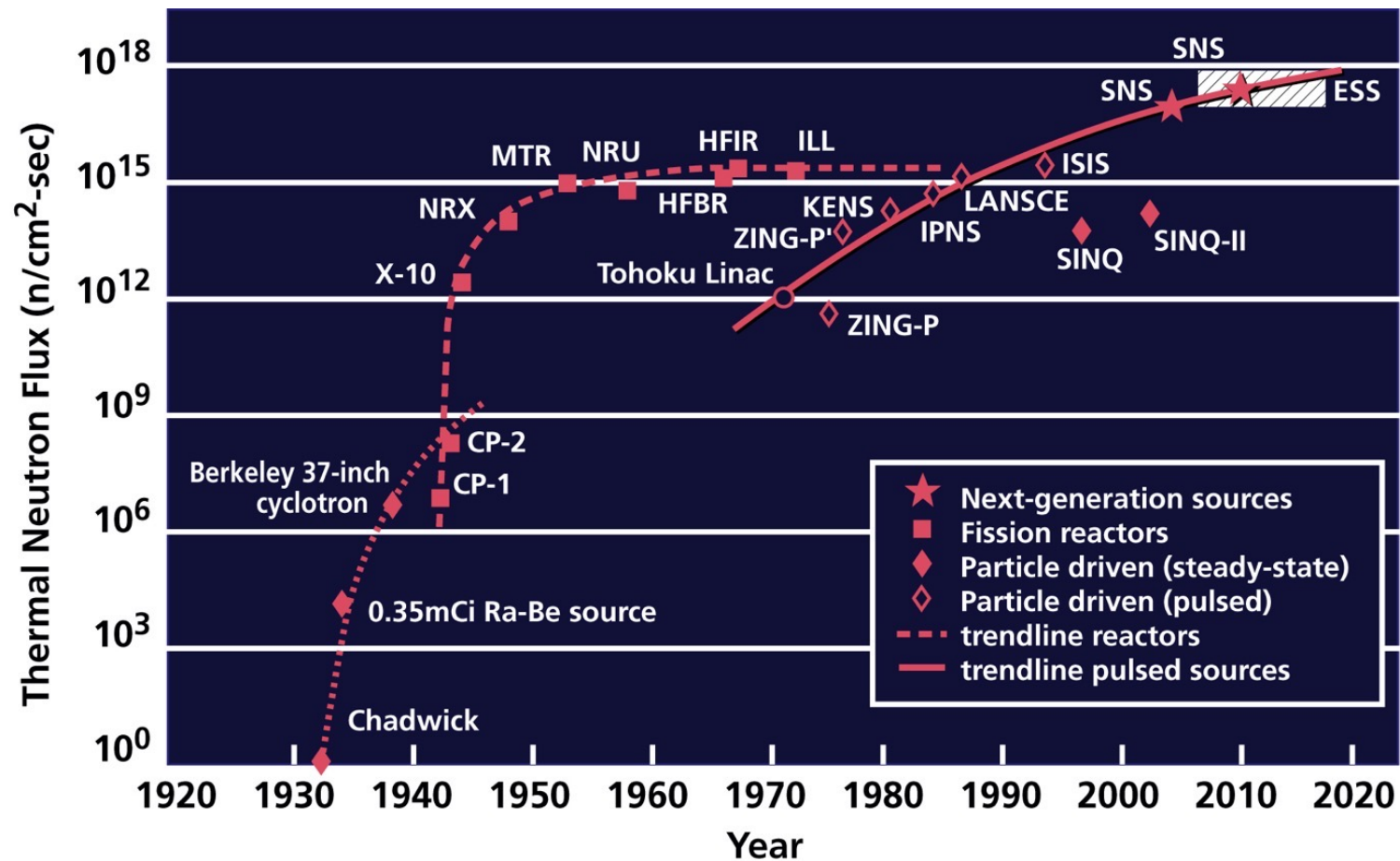
Steven Van Petegem
EPFL course MSE-435

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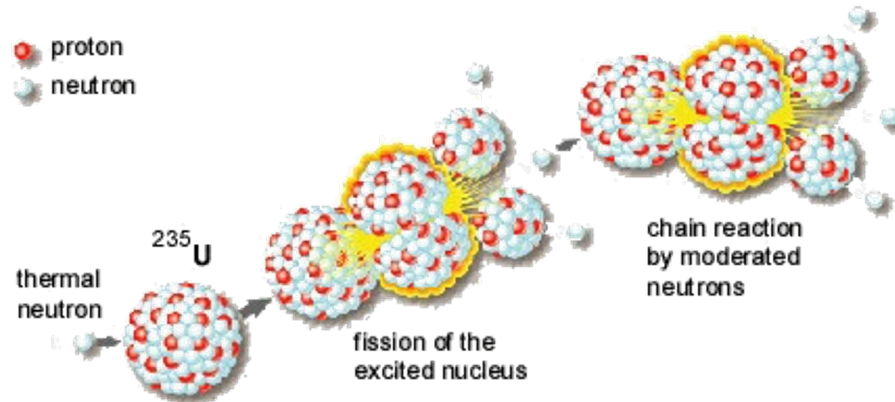
- **Neutron sources**
 - **Neutron reactors**
 - **Spallation neutron sources**
- Muon source
- Synchrotrons
- X-ray Free Electron Laser

Development of Neutron Research Facilities



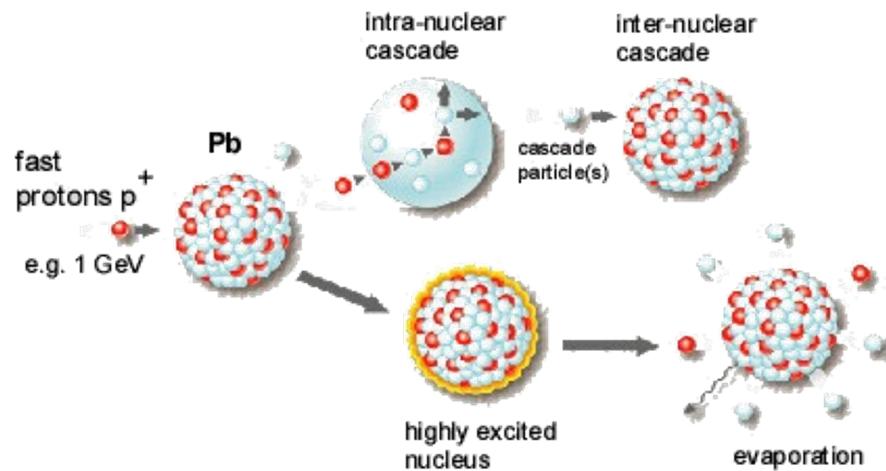
(Updated from *Neutron Scattering*, K. Skold and D. L. Price: eds., Academic Press, 1986)

Neutron production



Fission

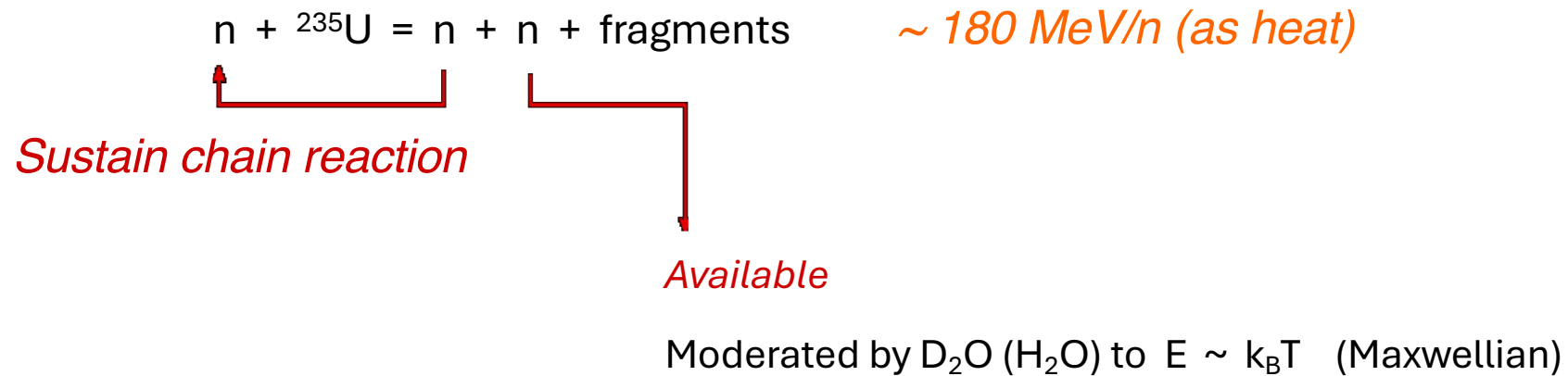
- Chain reaction
- Continuous flow
- ~ 1 neutron/fission



Spallation

- No chain reaction
- Accelerator driven
- Pulsed operation
- ~ 30 neutrons/proton

Neutron production - fission

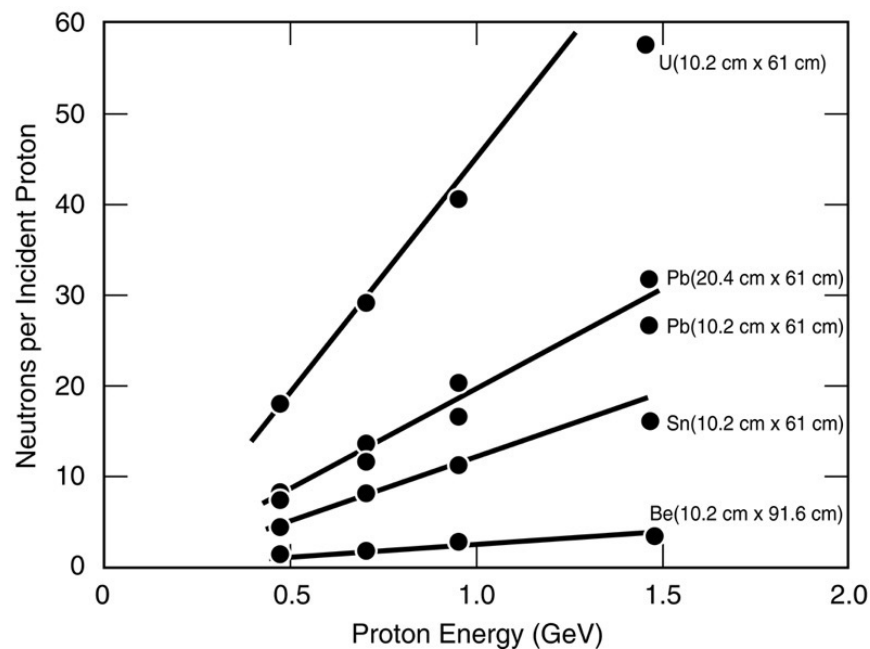


The chain reaction is started by inserting a neutron source, such as Californium-252 or Plutonium-238

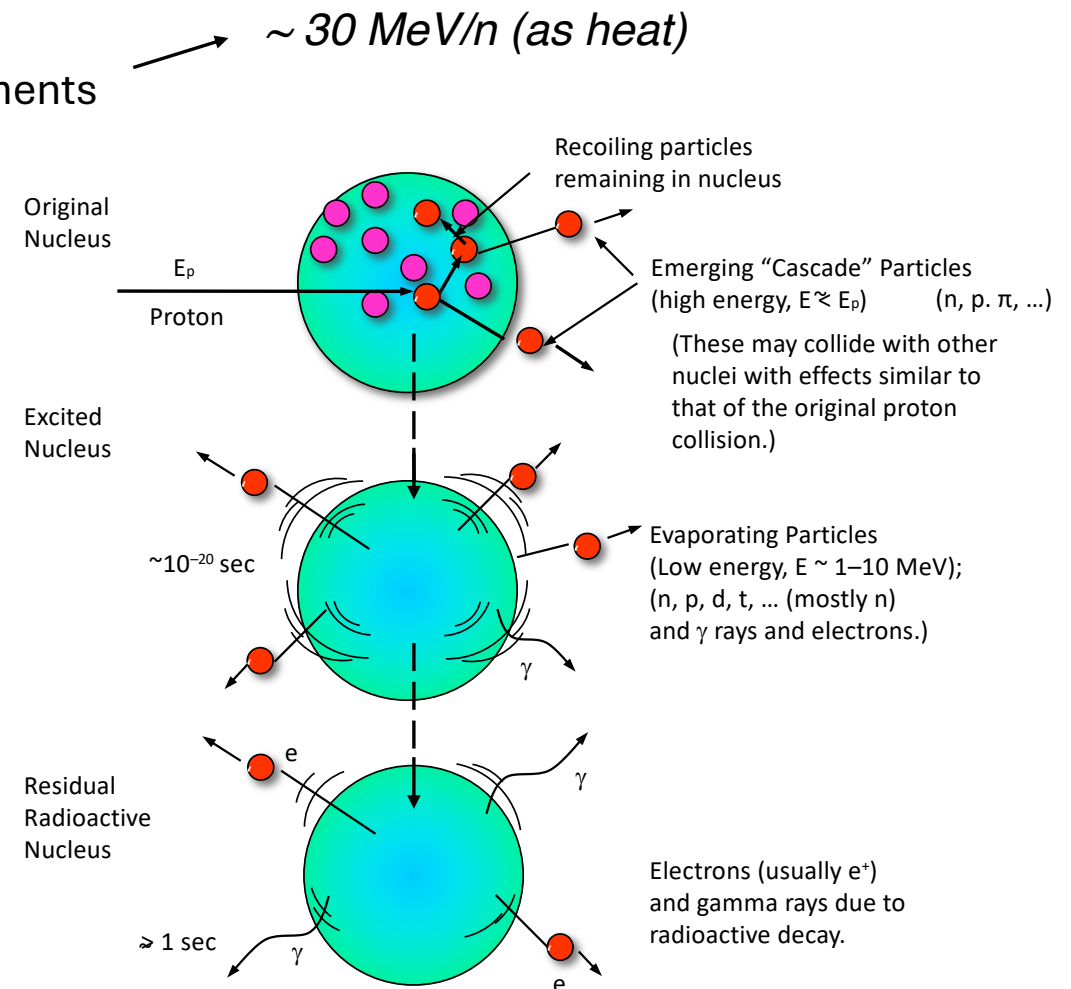
Neutron production - spallation

$p + \text{heavy nucleus} = 20 \sim 30 n + \text{fragments}$

1 GeV e.g. W, Pb, U



John M. Carpenter, IPNS, SNS, 26 September 2008



Neutron reactors

- Reactor e.g., HFR at ILL, Grenoble, France.

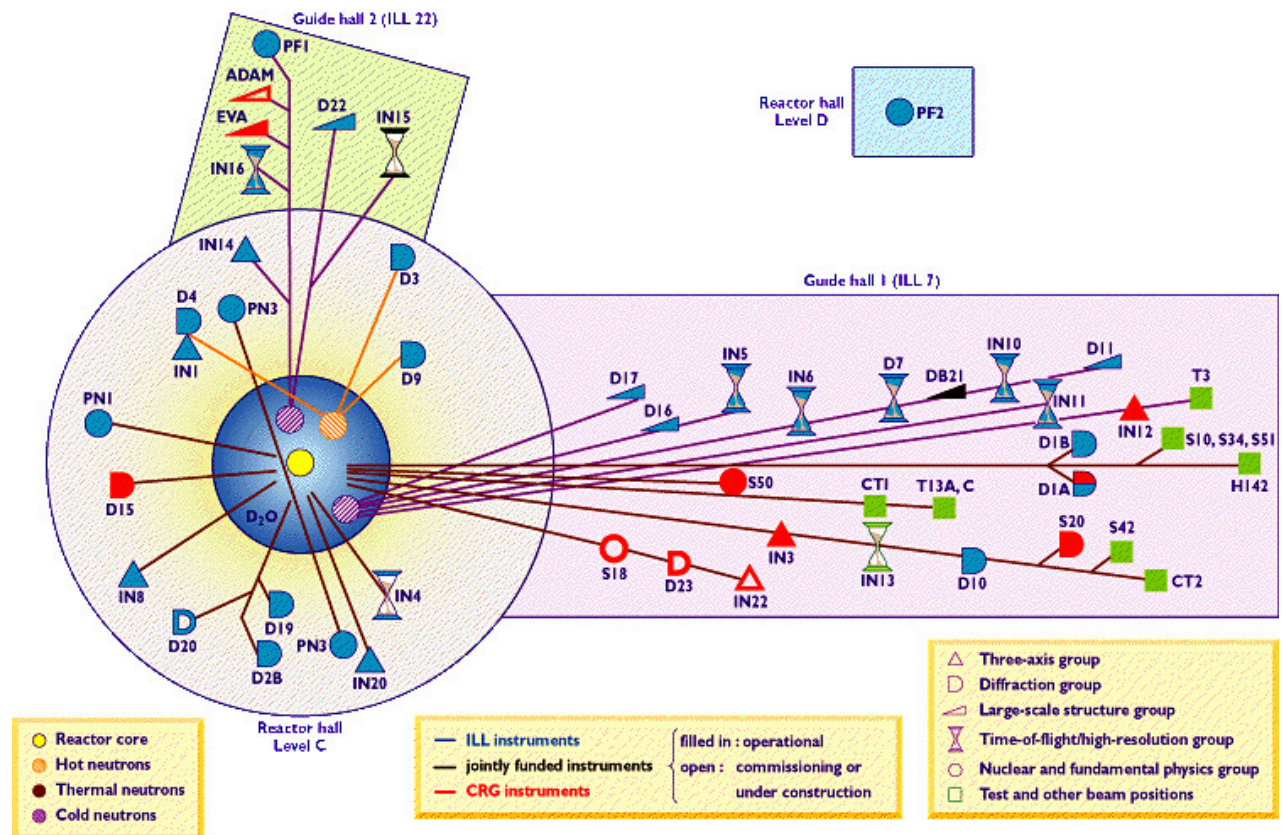
$$\sim 1.5 \times 10^{15} \text{ n/cm}^2/\text{s}$$

Advantages

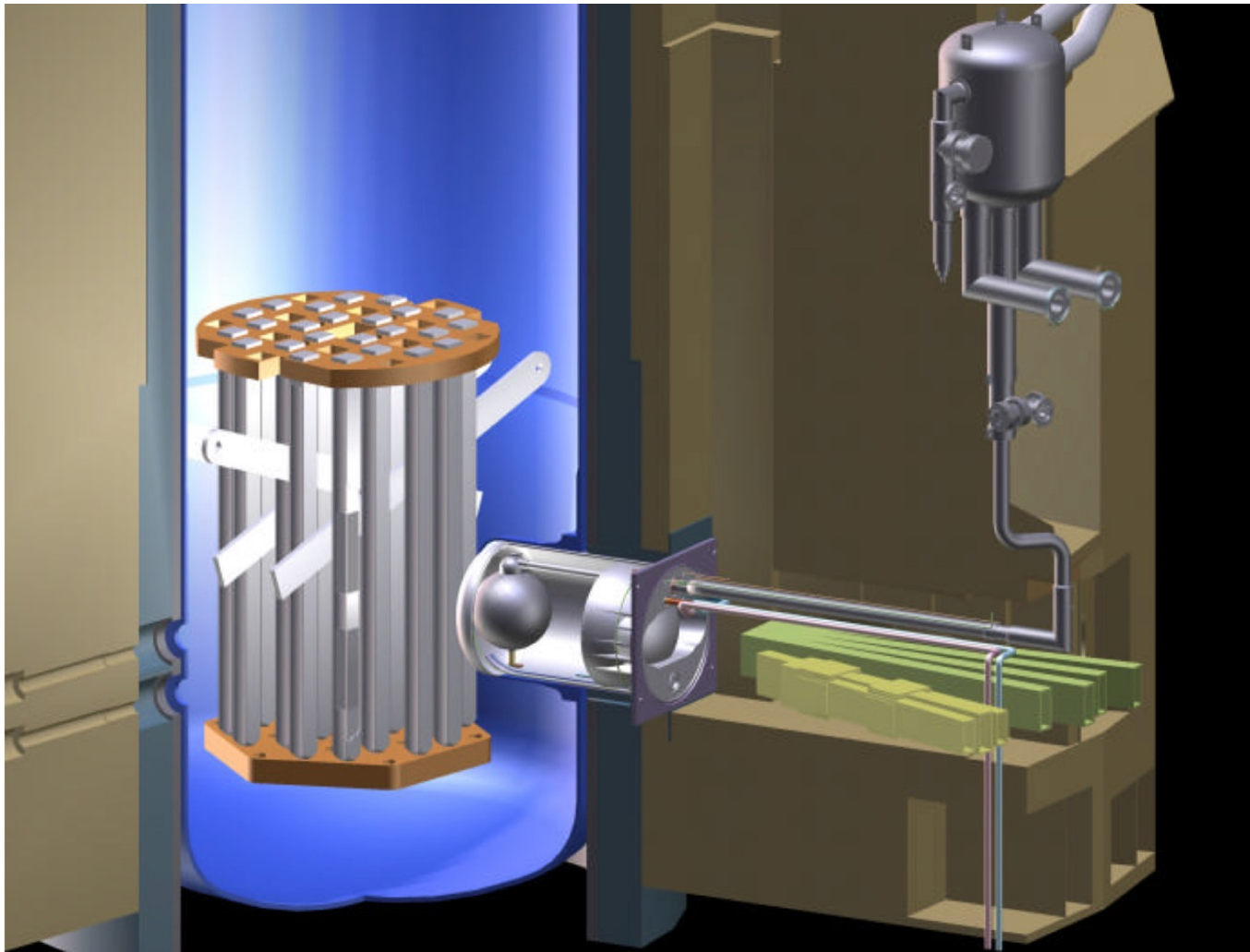
- High time averaged flux.
- Mature technology
- Very good for cold neutrons

Drawbacks

- Licensing (cost/politics).
- No time structure.



Neutron reactors



The National Institute of Standards and Technology (NIST) Reactor is a 20 MW research reactor with a peak thermal flux of 4×10^{14} N/sec. It is equipped with a unique liquid-hydrogen moderator that provides neutrons for seven neutron guides

Spallation sources

Pulsed spallation sources e.g., IPNS, ISIS, LANSCE, SNS

Advantages

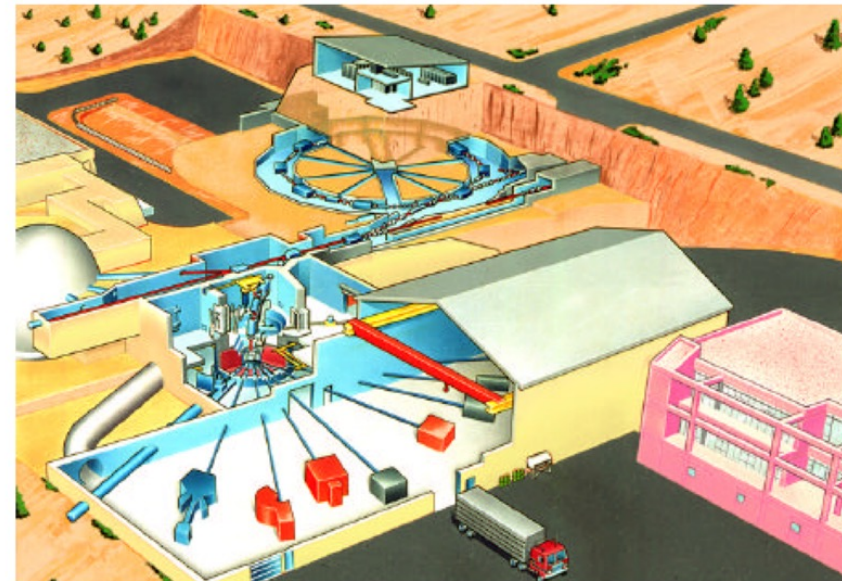
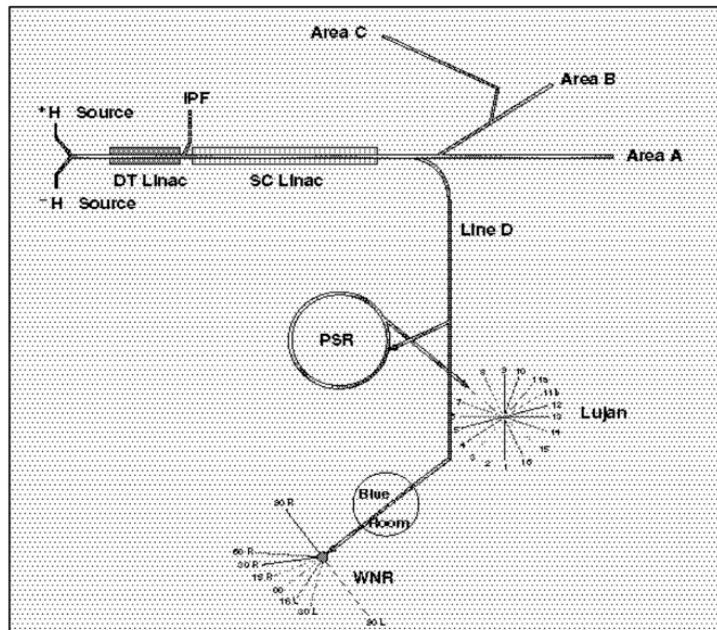
- High peak flux.
- Politically acceptable.
- Advantageous time structure for many applications.
- Accelerator based – politics simpler than reactors.
- Technology rapidly evolving.

Disadvantages

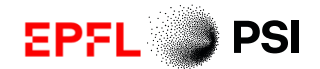
- Low time averaged flux.
- Not all applications exploit time structure.

Spallation sources - LANSCE

- Linac produces 20 H^- (a proton + 2 electrons) pulses per second
 - 800MeV, $\sim 800 \mu\text{sec}$ long pulses, average current $\sim 100 \mu\text{A}$
- Each pulse consists of repetitions of 270 nsec on, 90 nsec off
- Pulses are injected into a Proton Storage Ring with a period of 360 nsec
 - Thin carbon foil strips electrons to convert H^- to H^+ (i.e. a proton)
 - $\sim 3 \times 10^{13}$ protons/pulse ejected onto neutron production target



Spallation sources



- **CW spallation source** e.g., SINQ at Paul Scherrer Institut (PSI).
0.85 mA, 590 MeV, 0.9 MW
 1×10^{14} n/cm²/s average flux

Advantages

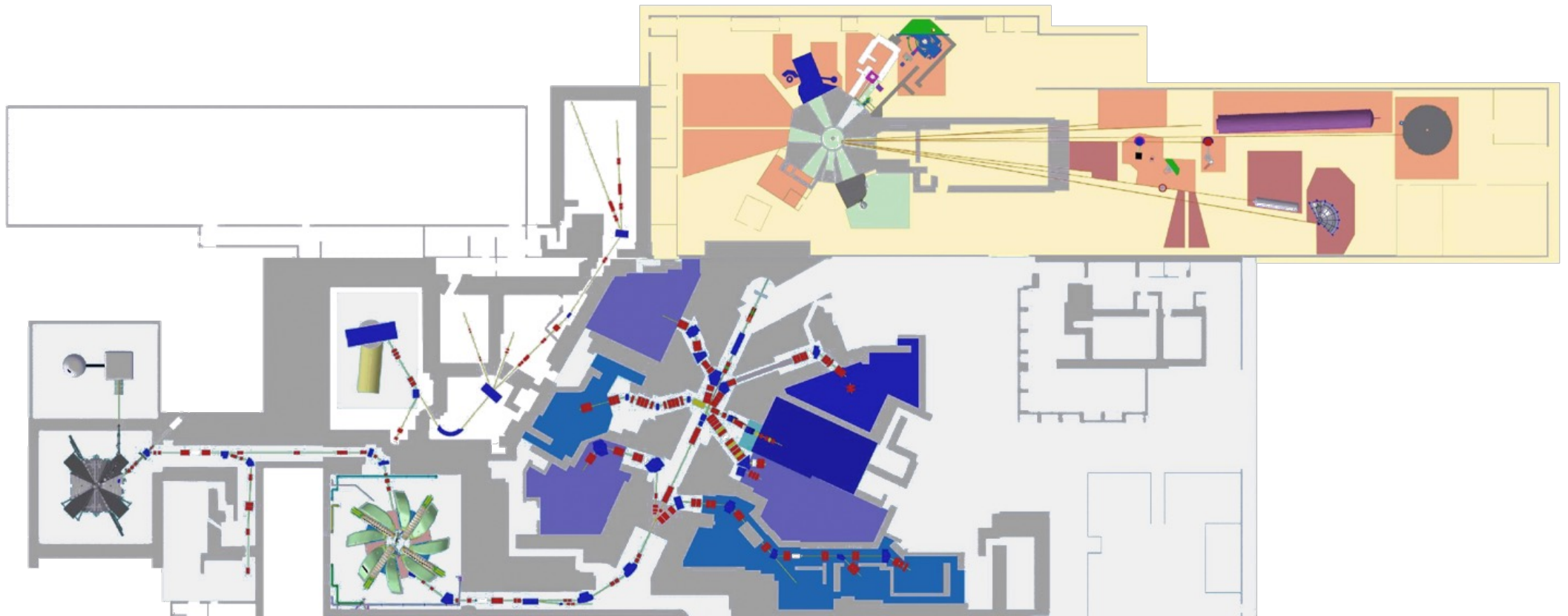
- High time averaged flux.
- Uses reactor type instrumentation (mature technology).
- Politically acceptable.
- piggy-backed on existing accelerator.

Disadvantages

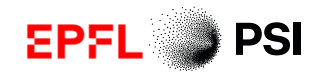
- No time structure.

SINQ – Swiss Spallation Neutron Source

EPFL  PSI



Proton accelerator



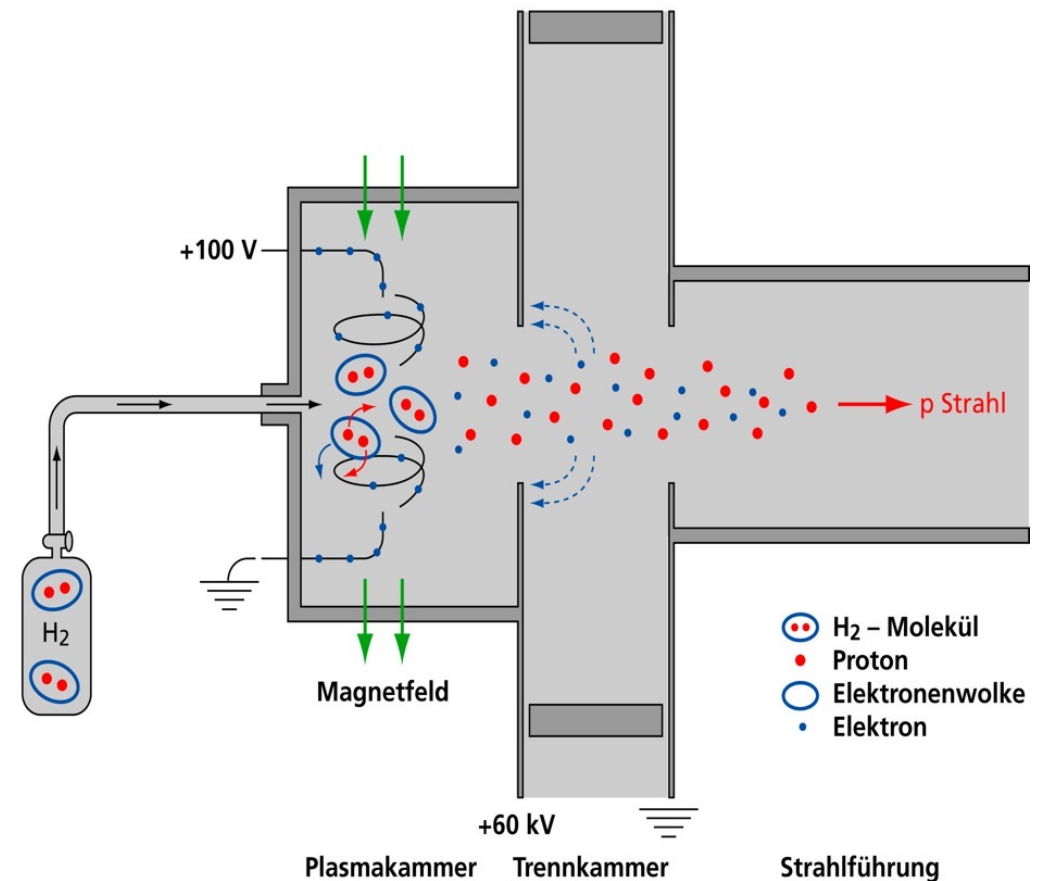
Movie on cyclotron

Generation of a proton beam

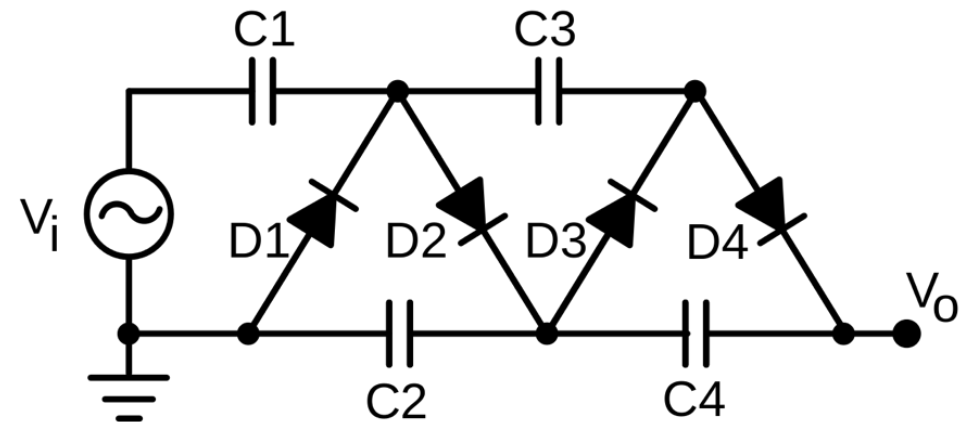
In the plasma chamber hydrogen atoms are separated into a mixture of protons and electrons by collisions with free electrons.

The protons are extracted from this plasma by a 60kV electric field and directed to the accelerator tube, which forms them into a beam with a particle velocity of 5% of the speed of light.

1 electron volt is the energy that a single positively charged particle possesses after it has passed through a potential difference of 1 volt. It is therefore about 0.16 trillionth of a joule.



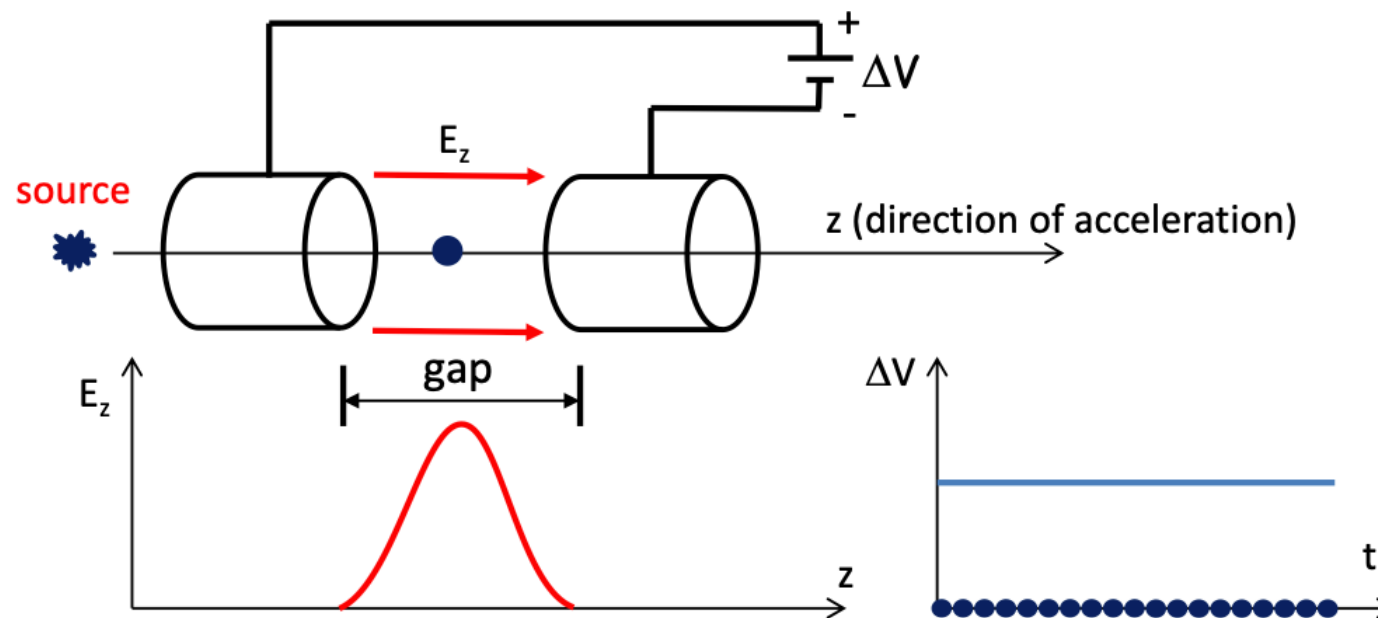
Cockcroft-Walton accelerator



- Voltage multiplier
- Basic principle:
<https://www.youtube.com/watch?v=ljUmDU4Uzp8>
- Maximum $\sim 10\text{MV}$
- @PSI: 0.87MV

Accelerators – basic principles

Consider the acceleration between two electrodes in DC

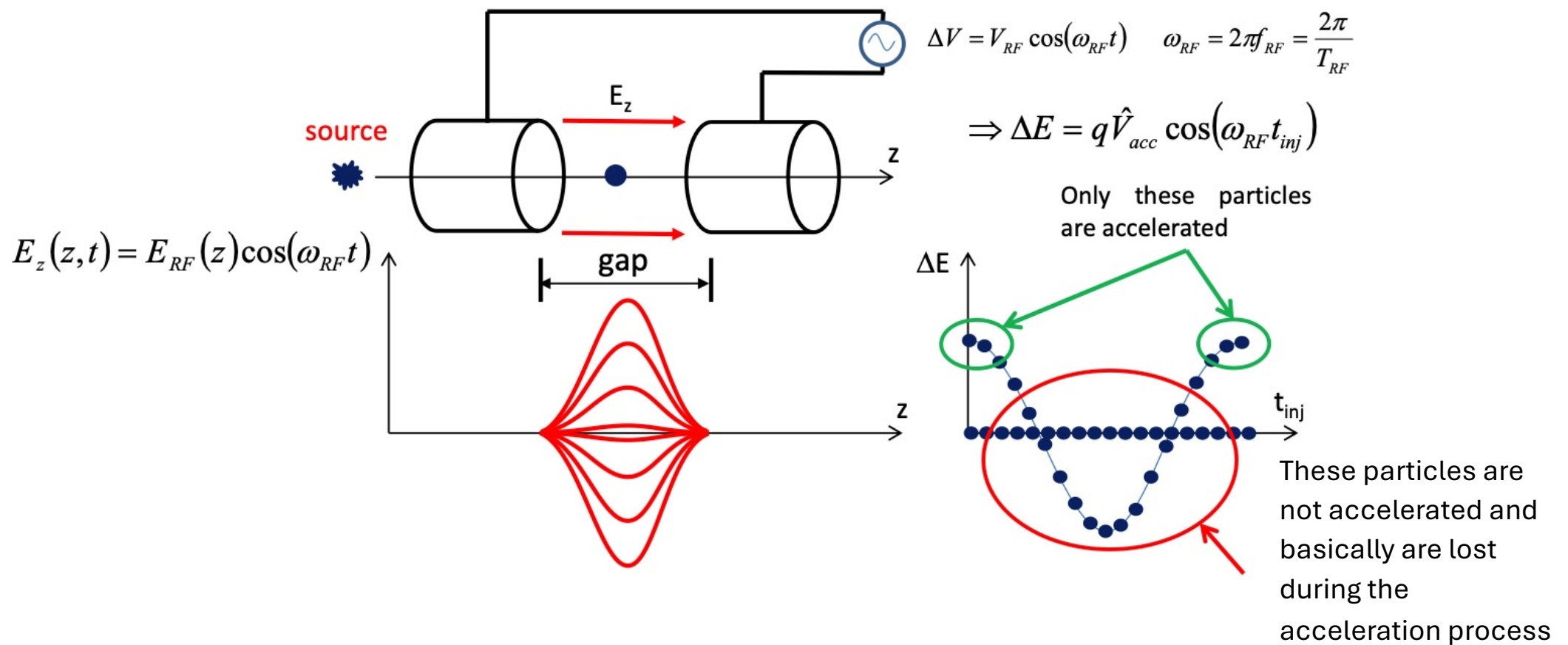


Energy gain per electrode: $\Delta E = q\Delta V$

<https://cas.web.cern.ch/sites/default/files/lectures/constantia-2018/alesinilinearaccelerators.pdf>

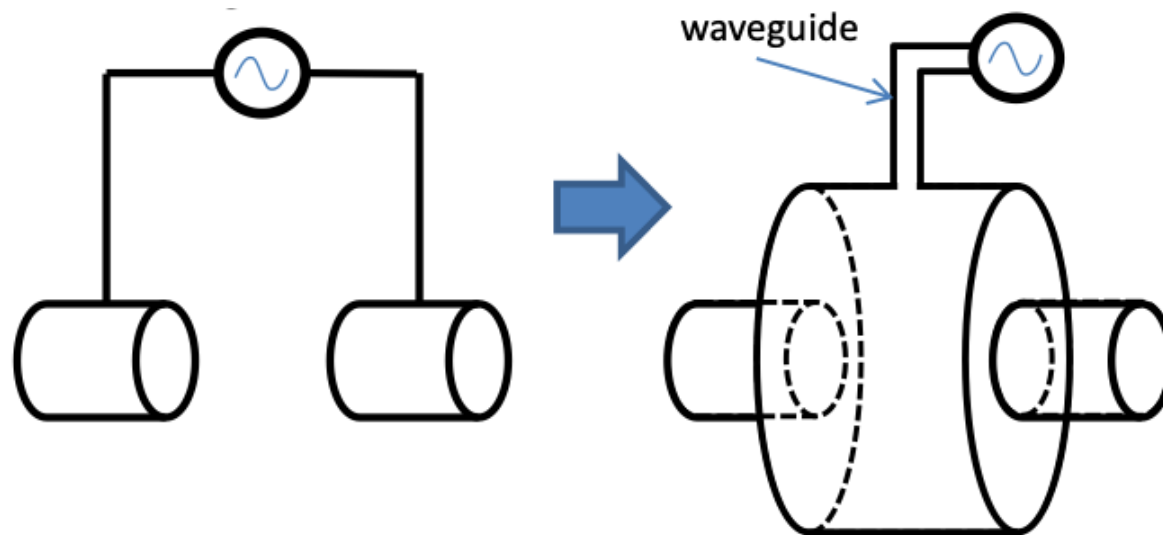
Accelerators – basic principles

Consider the acceleration between two electrodes fed by an RF generator



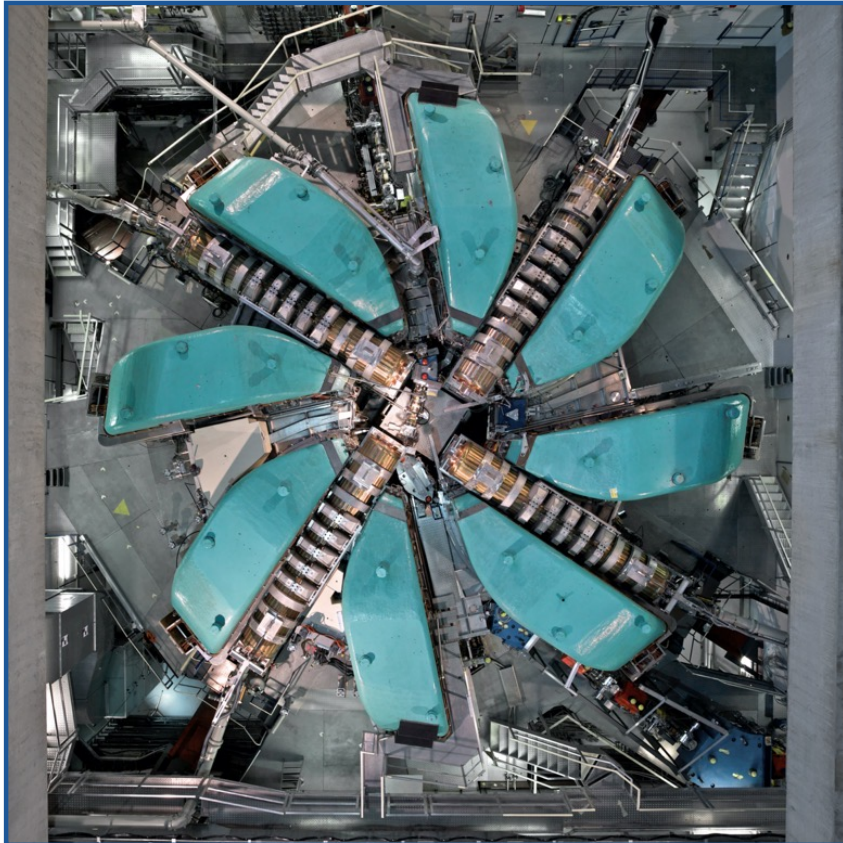
Accelerators – basic principles

In praxis: use a cavity which resonant frequency matches the RF generator frequency. Each cavity can be independently powered from the RF generator



- The cavities are metallic closed volumes where the electromagnetic fields have a particular spatial configuration (resonant modes).
- The modes are excited by RF generators that are coupled to the cavities.
- The resonant modes are called Standing Wave (SW) modes (spatial fixed configuration, oscillating in time).

Cyclotron



The main accelerator consisting of eight separate sector magnets and five high-frequency cavities, specially designed to accelerate a high-intensity proton beam.

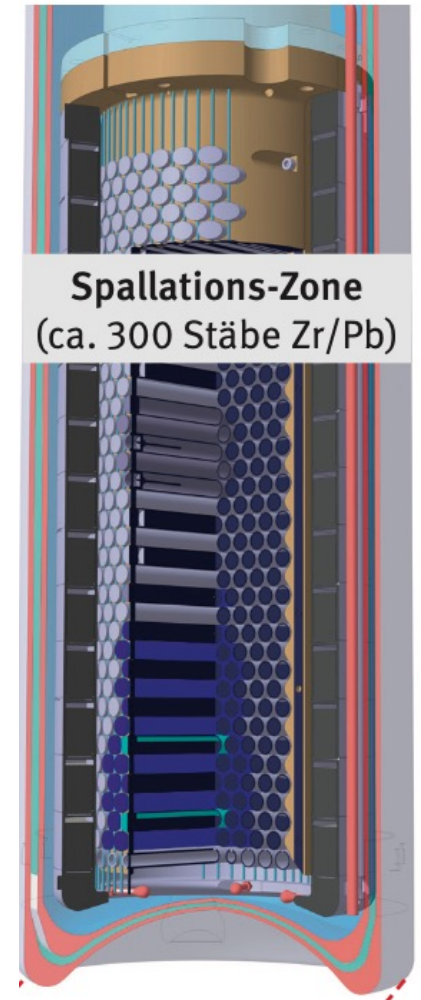
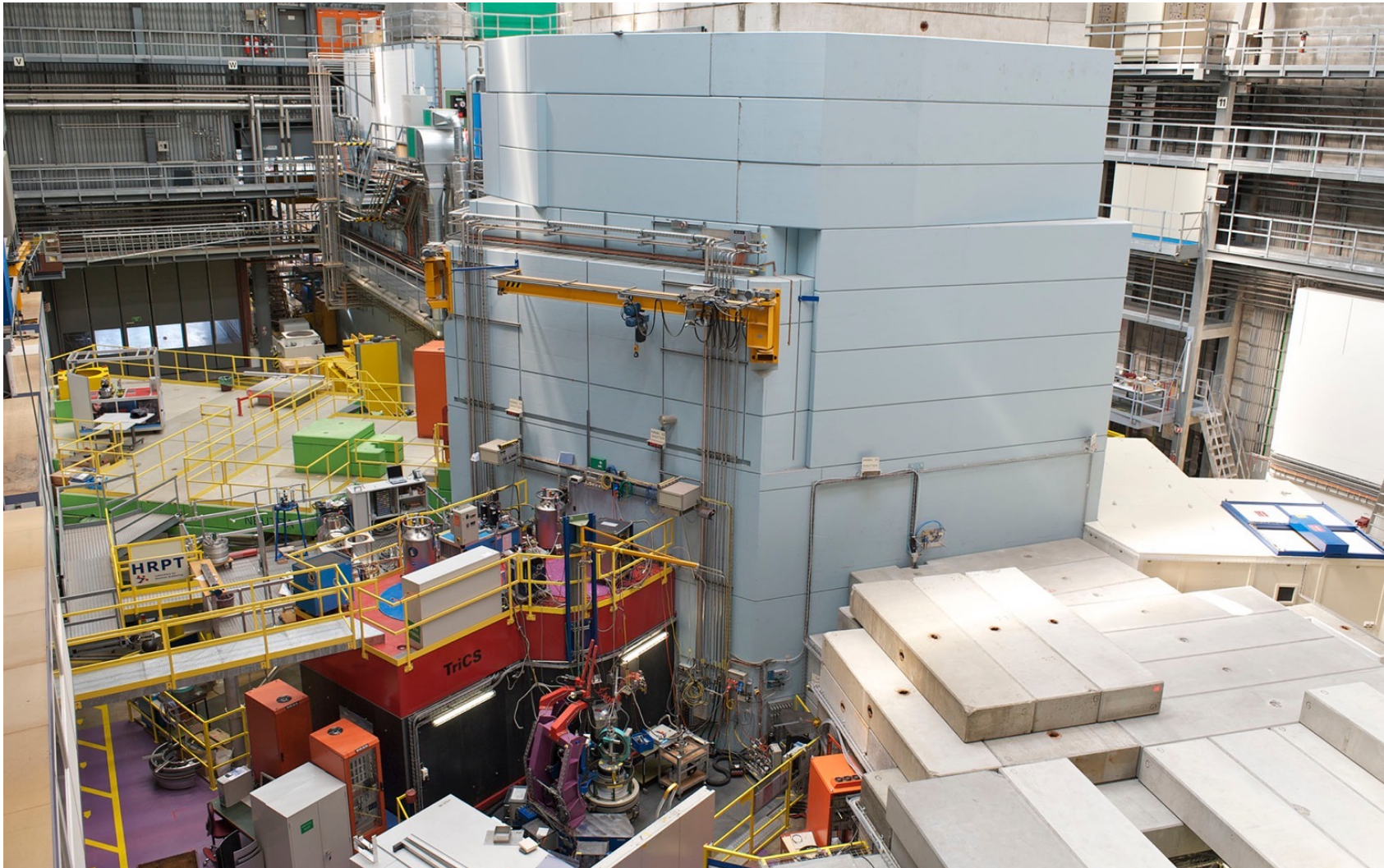
Injection energy:	72 MeV
Extraction energy:	590 MeV
Beam current:	1.6 mA
Time between two packets:	19.75 ns

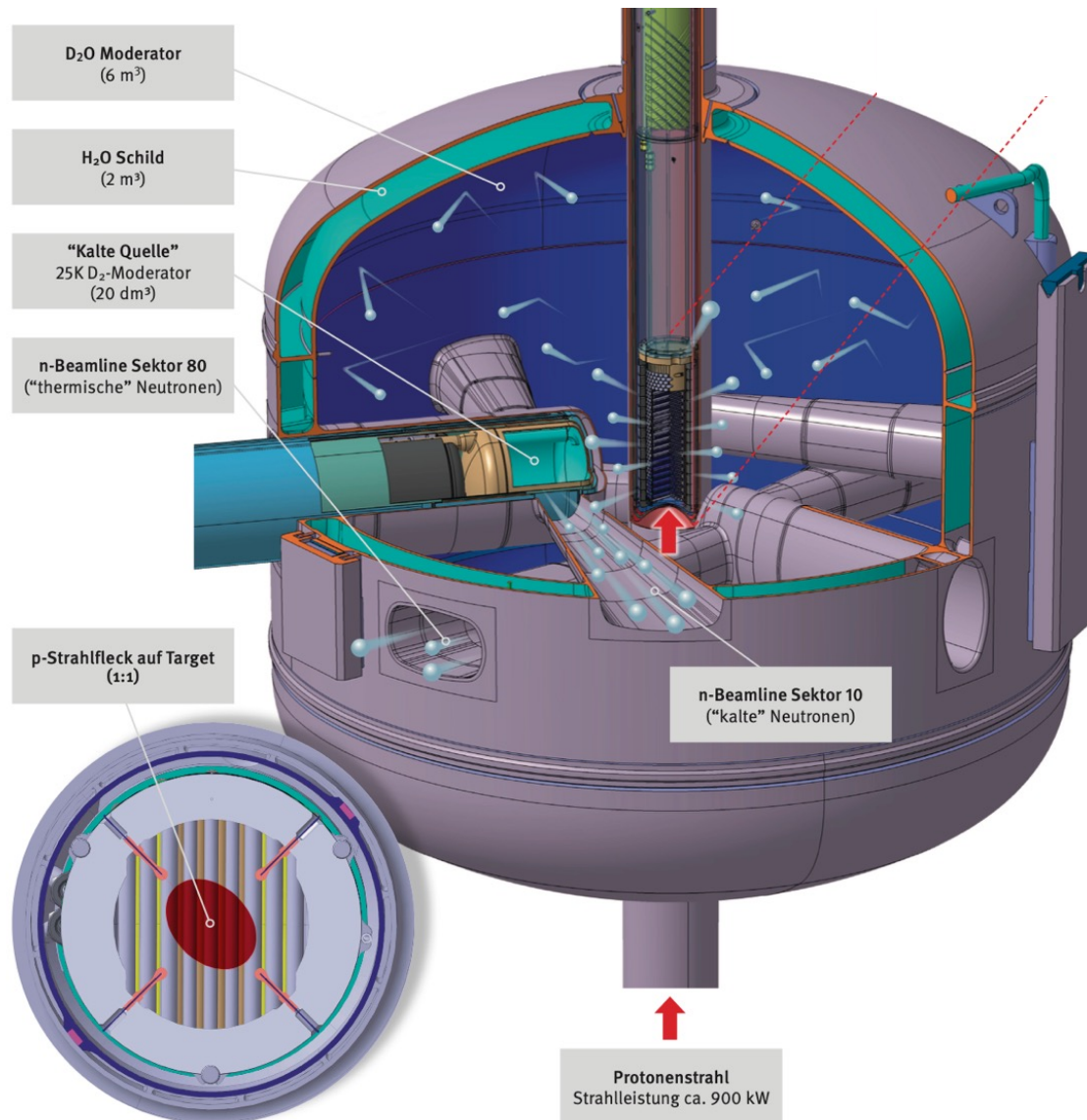
In the cyclotron, the proton beam is brought to a final energy of 590 MeV (82% of the speed of light). This means that a proton would fly around the earth approx. 6 times in one second (circumference of the earth approx. 40,000 km).

The proton beam current is approx. 2.4 mA, resulting in a beam power of over 1.3 MW.

Swiss Spallation Neutron Source

EPFL  PSI

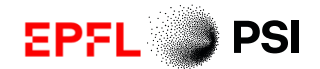




The neutrons that are released leave the target at a speed of about 20'000 km/s. Since this is too fast for materials science purposes, they are slowed down. This is done by atomic collisions in a moderator material that surrounds the target. After that, the neutrons still have a speed of about 2200 m/s, they are called thermal neutrons. If the neutrons are additionally decelerated in a very cold medium (liquid heavy hydrogen at -250°C), they become even slower. We then speak of cold neutrons, which still have a velocity of about 700 m/s.

Very fast neutrons that cannot be slowed down in the moderator have to be absorbed by the 5 m thick shielding made of iron and special concrete. In the process, their kinetic energy is converted into heat. The inner part of the shielding must therefore be cooled with water.

Neutron sources - flux



Reactors

DR3	Risø	$2 \times 10^{14} \text{ n/cm}^2/\text{s}$
ILL	Grenoble	$1.5 \times 10^{15} \text{ n/cm}^2/\text{s}$
FMR-II	Garching	$8 \times 10^{14} \text{ n/cm}^2/\text{s}$

Spallation sources

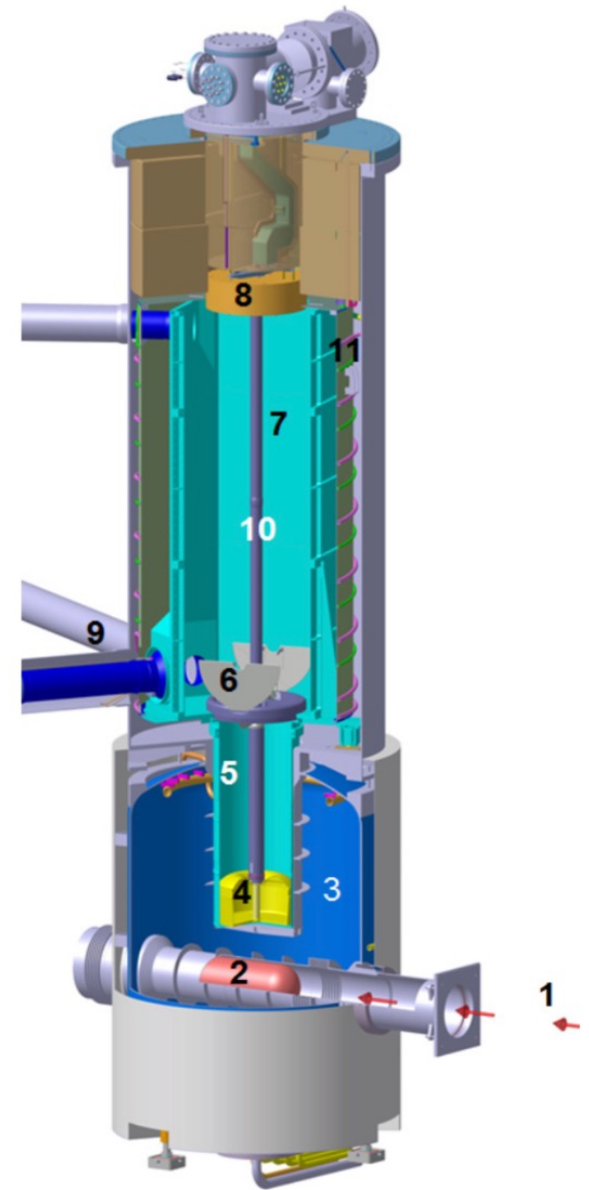
ISIS @ 160 kW	average	$1.2 \times 10^{13} \text{ n/cm}^2/\text{s}$
	peak	$6.0 \times 10^{15} \text{ n/cm}^2/\text{s}$
SNS @ 2 MW	average	$4.0 \times 10^{13} \text{ n/cm}^2/\text{s}$
	peak	$3.0 \times 10^{16} \text{ n/cm}^2/\text{s}$

Cooling down

Neutrons	Energy range	Wavelength [$\text{\AA}$]	Velocity [m/s]
ultra cold	$\leq 300 \text{ neV}$	≥ 500	≤ 8
very cold	300 neV - 0.12 meV	52.2 – 26.1	7.5 – 152
cold	0.12 meV - 12 meV	26.1 – 2.6	152 – 1515
thermal	12 meV - 100 meV	2.6 - 0.9	1515 - 4374
epithermal	100 meV - 1eV	0.9 - 0.28	4374 - $13.8 \cdot 10^3$
intermediate	1eV - 0.8MeV		
fast	$> 0.8\text{MeV}$		

Ultra cold neutron source

- The full intensity proton beam (1) – 1.5×10^{16} protons/s – impinges for up to 8s onto the spallation target (2) built out of 760 lead-filled Zircalloy tubes.
- Subsequently, neutrons are thermalized in the surrounding ultraclean heavy water (D₂O) (3) operating at a temperature of ~ 300 K. This water is also used to cool the spallation target.
- Solid deuterium at 5 K inside the D₂ vessel (4) allows further cooling of the neutrons, resulting in a cold neutron flux of a few times 10^{13} n/cm²/s.
- These neutrons are finally down-scattered to become ultracold via phonon interaction with the solid D₂.
- The UCN emanating at the top lose this boost energy in the 1 m rise in the vertical guide (5) and enter the ~ 2 m³ large UCN storage vessel (7) coated with diamond-like carbon.



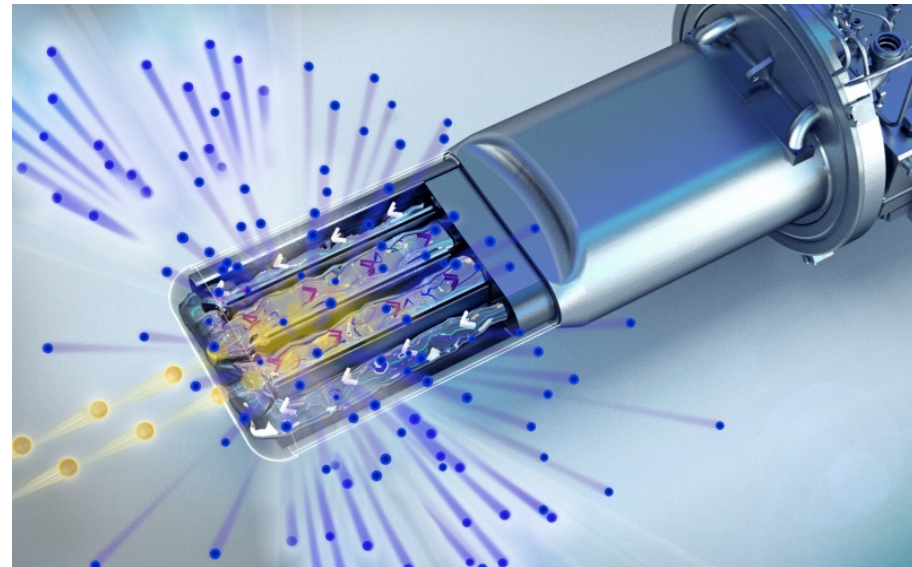
Spallation Neutron Source



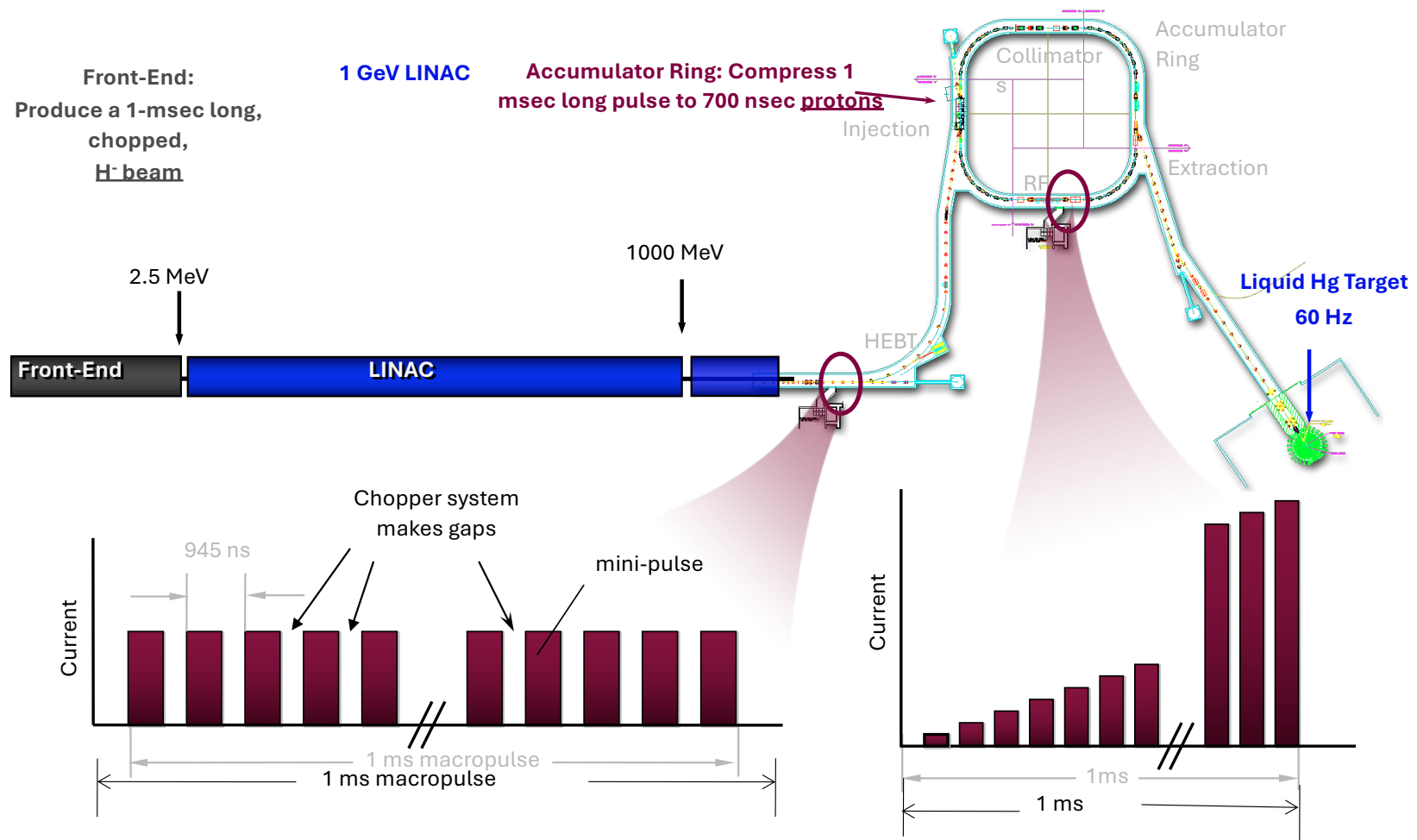
Oak Ridge, USA

Currently most intense pulsed neutron source

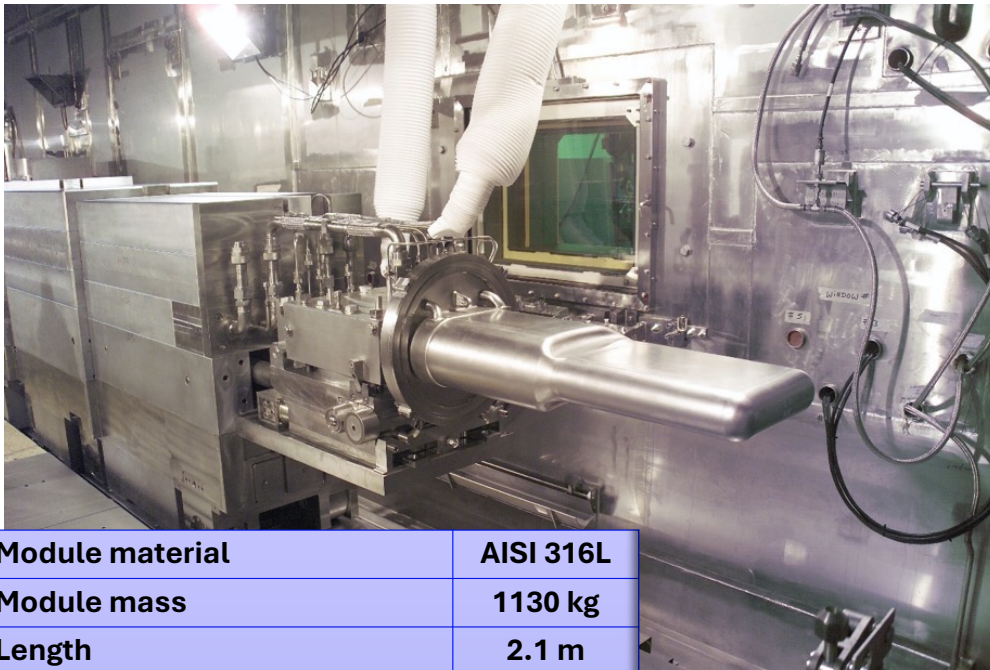
Cost: ~1.4 billion dollar



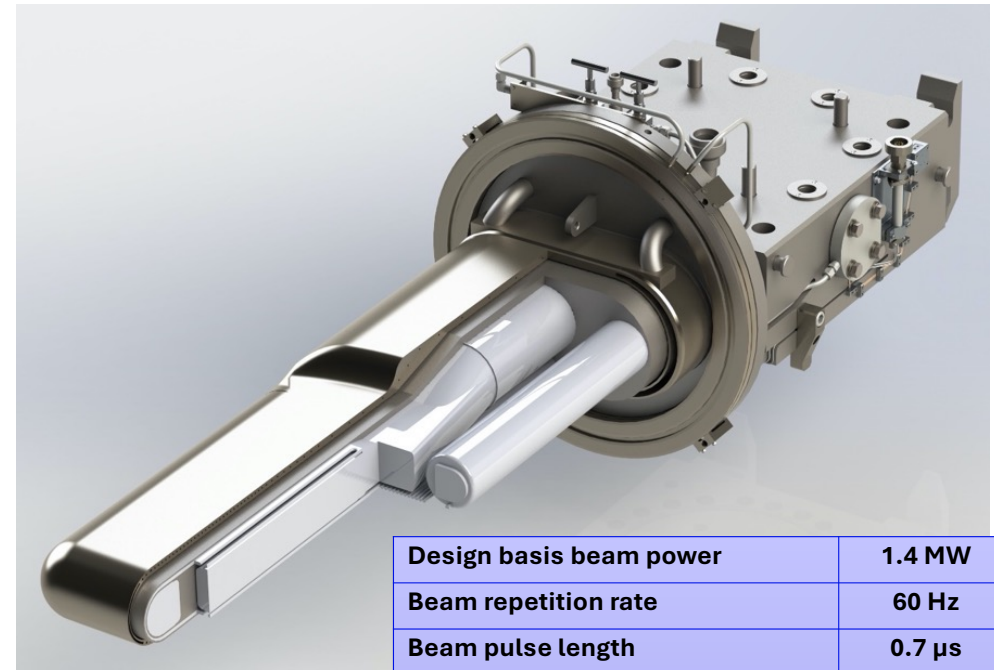
Spallation Neutron Source



Spallation Neutron Source

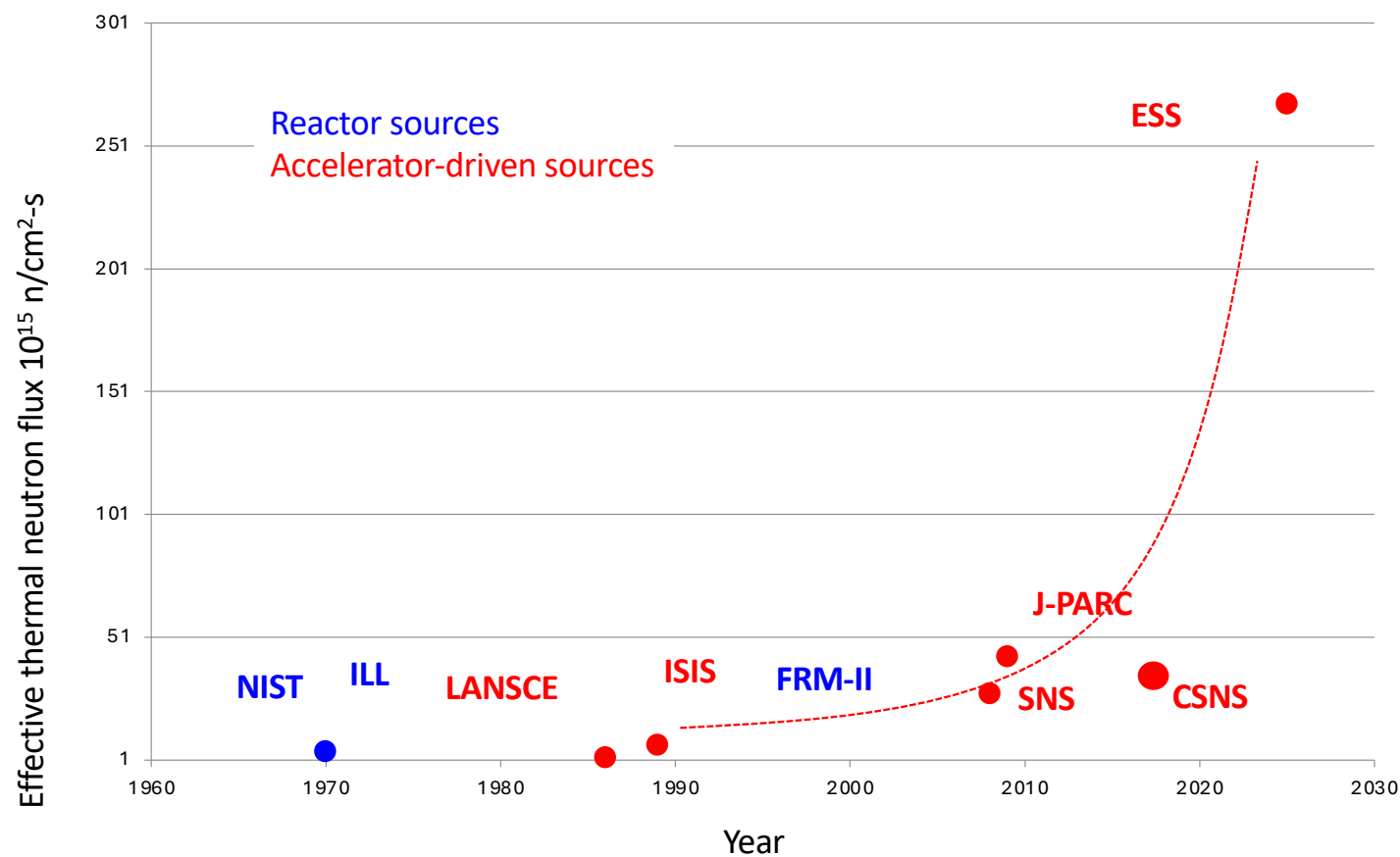


Module material	AISI 316L
Module mass	1130 kg
Length	2.1 m
Mercury mass inside module	794 kg
Mercury mass flow rate	19.4 t/min
Radiation damage limit	12 dpa

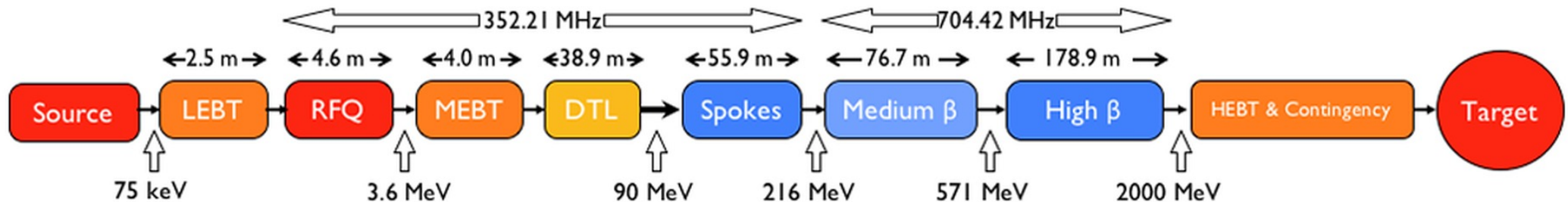


Design basis beam power	1.4 MW
Beam repetition rate	60 Hz
Beam pulse length	0.7 μ s
Energy per pulse	23 kJ
Proton energy	1 GeV

European Spallation Source



European Spallation Source



- The proton beam is transported through a Low Energy Beam Transport (LEBT) section to the Radio Frequency Quadrupole (RFQ), where it is bunched and accelerated up to 3.6 MeV.
- In the Medium Energy Beam Transport (MEBT) section the transverse and longitudinal beam characteristics are diagnosed and optimised for further acceleration in the Drift Tube Linac (DTL).
- After approximately 50 meters the protons have gained enough speed so they can be accelerated by superconducting cavities. These cavities are cooled by liquid helium to -271 °C.
- The protons reach 96% of the speed of light

European Spallation Source

The target is a 2.6 m diameter stainless steel disk containing bricks of a neutron-rich heavy metal: tungsten. It weighs almost five tons. The wheel rotates at 23.3 RPMs, in time with the arrival of the proton beam painted across the exterior of the wheel shroud. The unit is cooled by a flowing helium gas system interfaced with a secondary water system.

The neutrons, which, following spallation, travel at 10% of the speed of light, are then slowed down. This is achieved using a para-hydrogen and water-based moderator and a beryllium-lined reflector. The moderator-reflector system is housed in a replaceable plug, and also includes cryogenic hydrogen and water-cooling systems.

Once moderated, the neutrons are delivered to the instruments through 42 beam ports radiating from the Target Station.

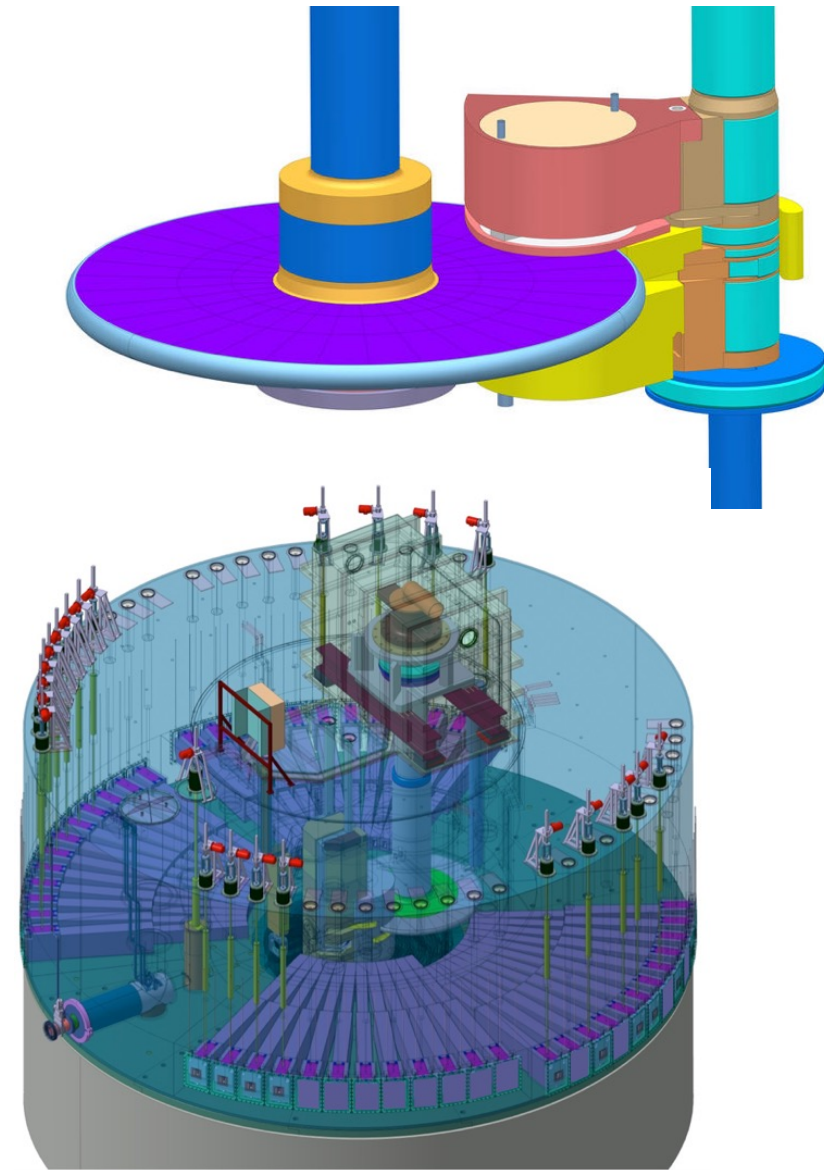
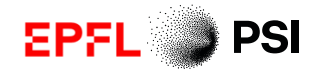
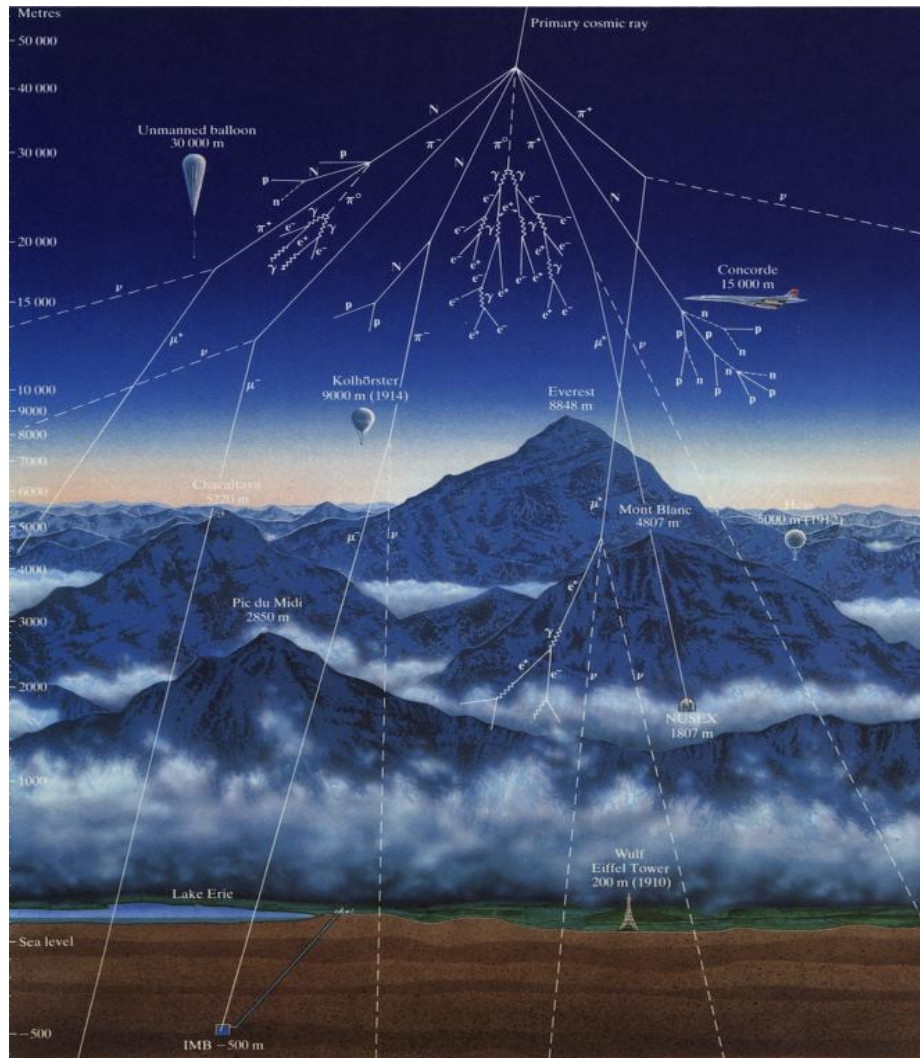


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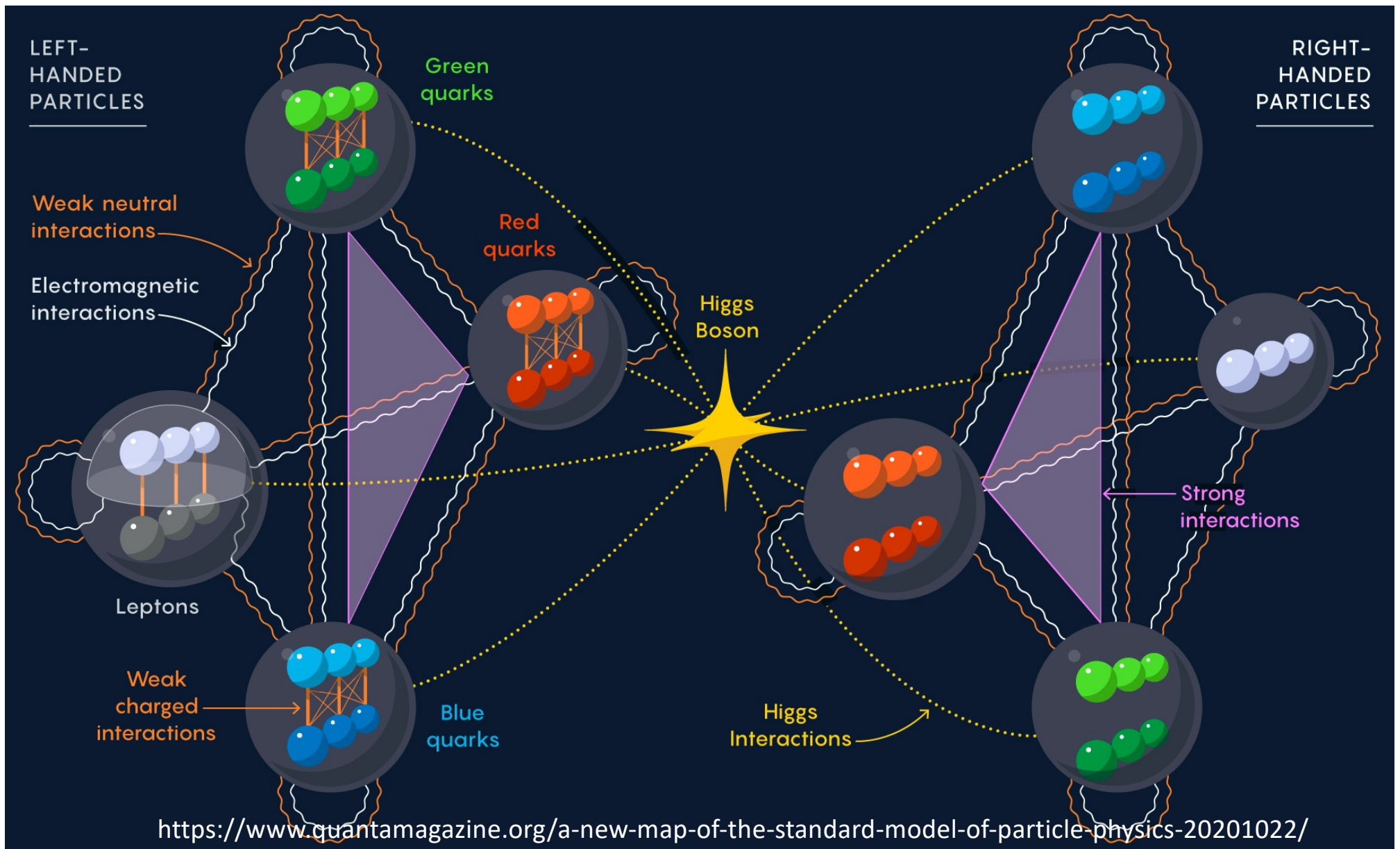


- Neutron sources
 - Neutron reactors
 - Spallation neutron sources
- **Muon source**
- Synchrotrons
- X-ray Free Electron Laser

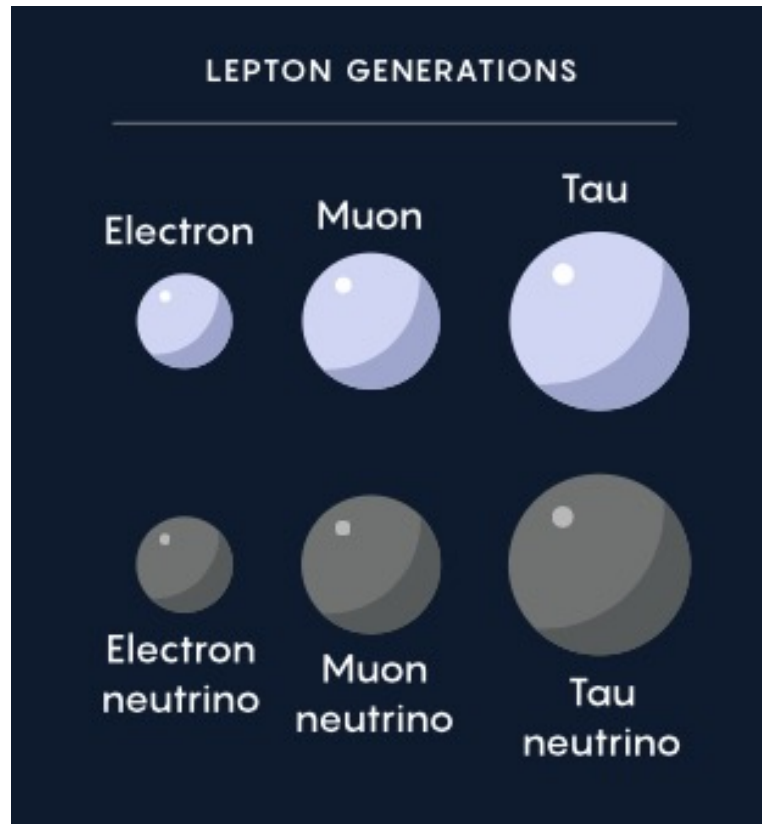
Myons



- Muons arriving on the Earth's surface are created indirectly as decay products of collisions of cosmic rays with particles of the Earth's atmosphere.
- When a cosmic ray proton impacts atomic nuclei in the upper atmosphere, pions are created. These decay within a relatively short distance into muons and muon neutrinos.
- Muons were discovered by Carl D. Anderson and Seth Neddermeyer at Caltech in 1936



Myons



Myon properties

Elementary particles:

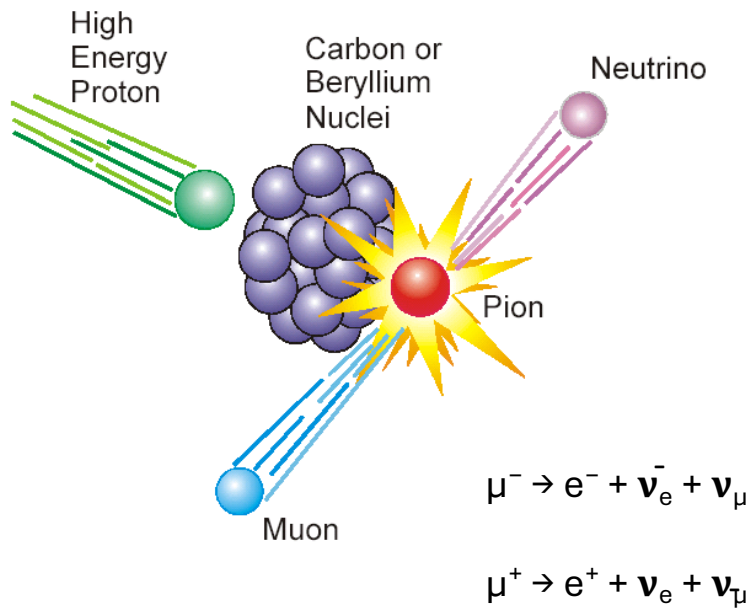
Mass:	$\sim 200 \times$ electron mass
	$\sim 1/9 \times$ proton mass
Charge:	+ e, oder -e
Spin:	$\frac{1}{2}$
Magnetic moment :	$3.18 \times m_p$
Lifetime:	2.2 μ sec

<https://www.quantamagazine.org/a-new-map-of-the-standard-model-of-particle-physics-20201022/>

Myon in the lab

$$\pi^+ \rightarrow \mu^+ + \nu_\mu$$

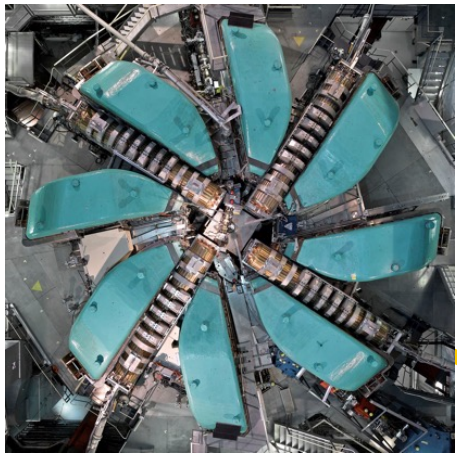
$$\pi^- \rightarrow \mu^- + \bar{\nu}_\mu$$



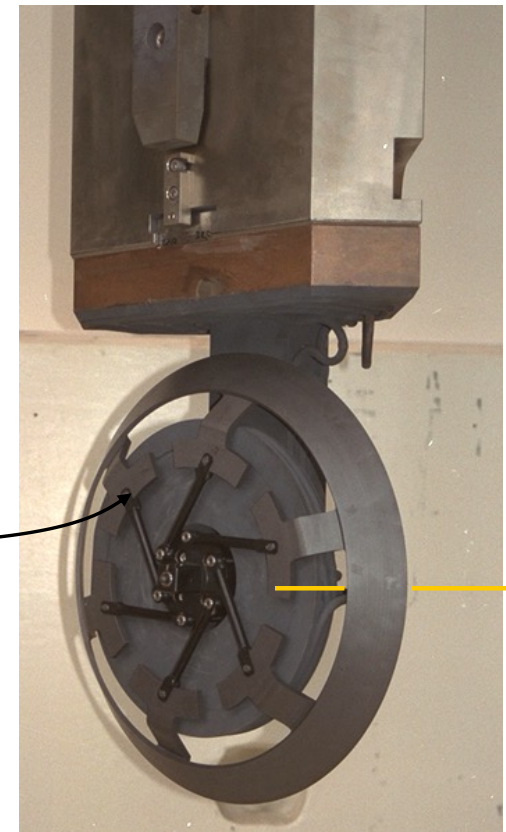
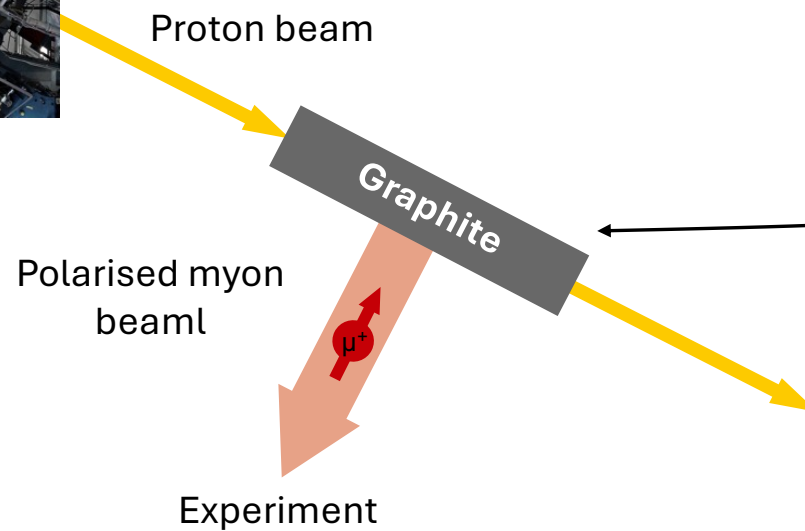
At PSI:

- $\approx 10^{16}$ Protonen / sec @ 600 MeV
- ~ 1 MW power on $5 \times 5 \text{ mm}^2$, 40 kW/mm^2
- Energy consumption of 1000 households and melts Fe in 0.1 ms
- $\approx 10^7 - 10^8 \mu^+/\text{sec}$, 100 % pol., $\approx 4 \text{ MeV}$
- Strongest myon source in the world

Myon source @ PSI

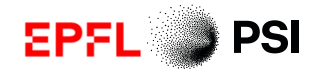


Proton
accelerator



Target

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- Neutron sources
 - Neutron reactors
 - Spallation neutron sources
- Muon source
- **Synchrotrons**
- X-ray Free Electron Laser

Synchrotrons



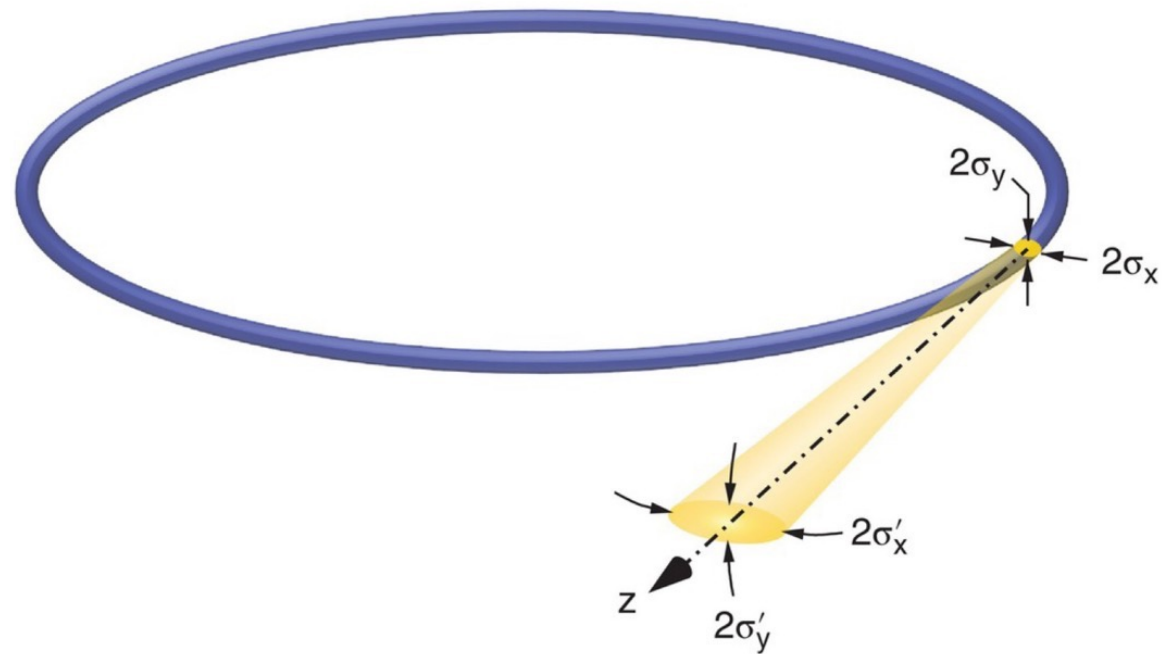


Emittance

The product of the linear source size σ and the beam divergence σ' in the same plane is known as ϵ , the emittance in that plan

$$\epsilon_x = \sigma_x \sigma'_x$$

$$\epsilon_y = \sigma_y \sigma'_y$$



Brilliance

The appropriate figure-of-merit for light sources is the brilliance.

It takes into account:

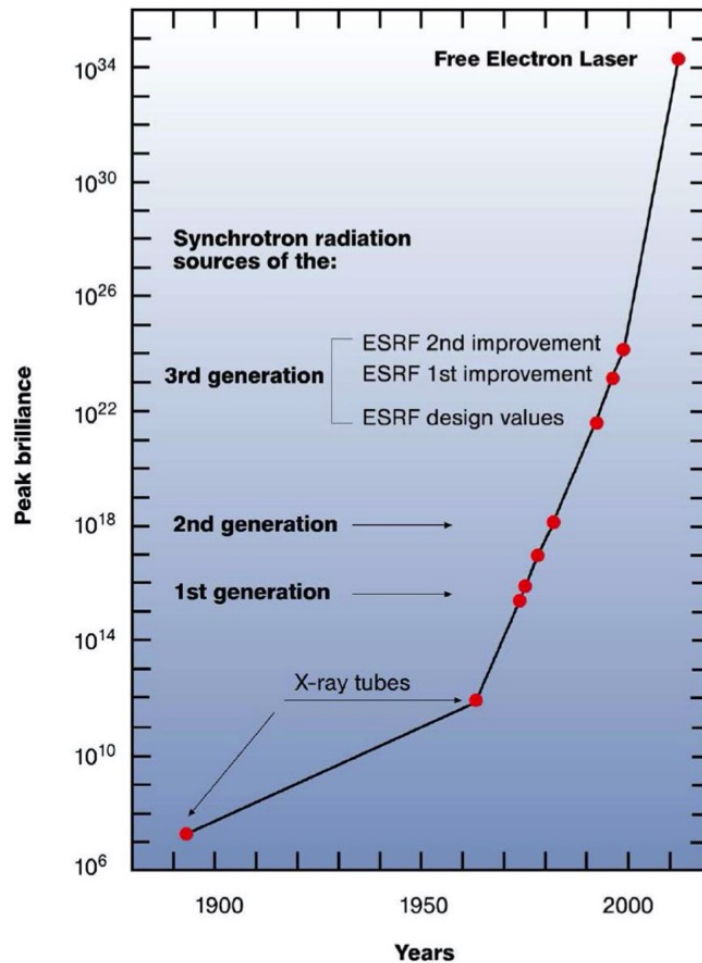
- the number of photons produced per time unit (use second)
- the angular divergence of the photons (in mrad^2)
- the cross-sectional area of the photon beam (use mm^2)
- the number of photons falling within a certain bandwidth (BW) of the central wavelength (0.1% are customary)

The unit for brilliance is:

$\text{photons} / \text{s} / \text{mm}^2 / \text{mrad}^2 / 0.1\% \text{ (BW)}$

or, flux per total emittance

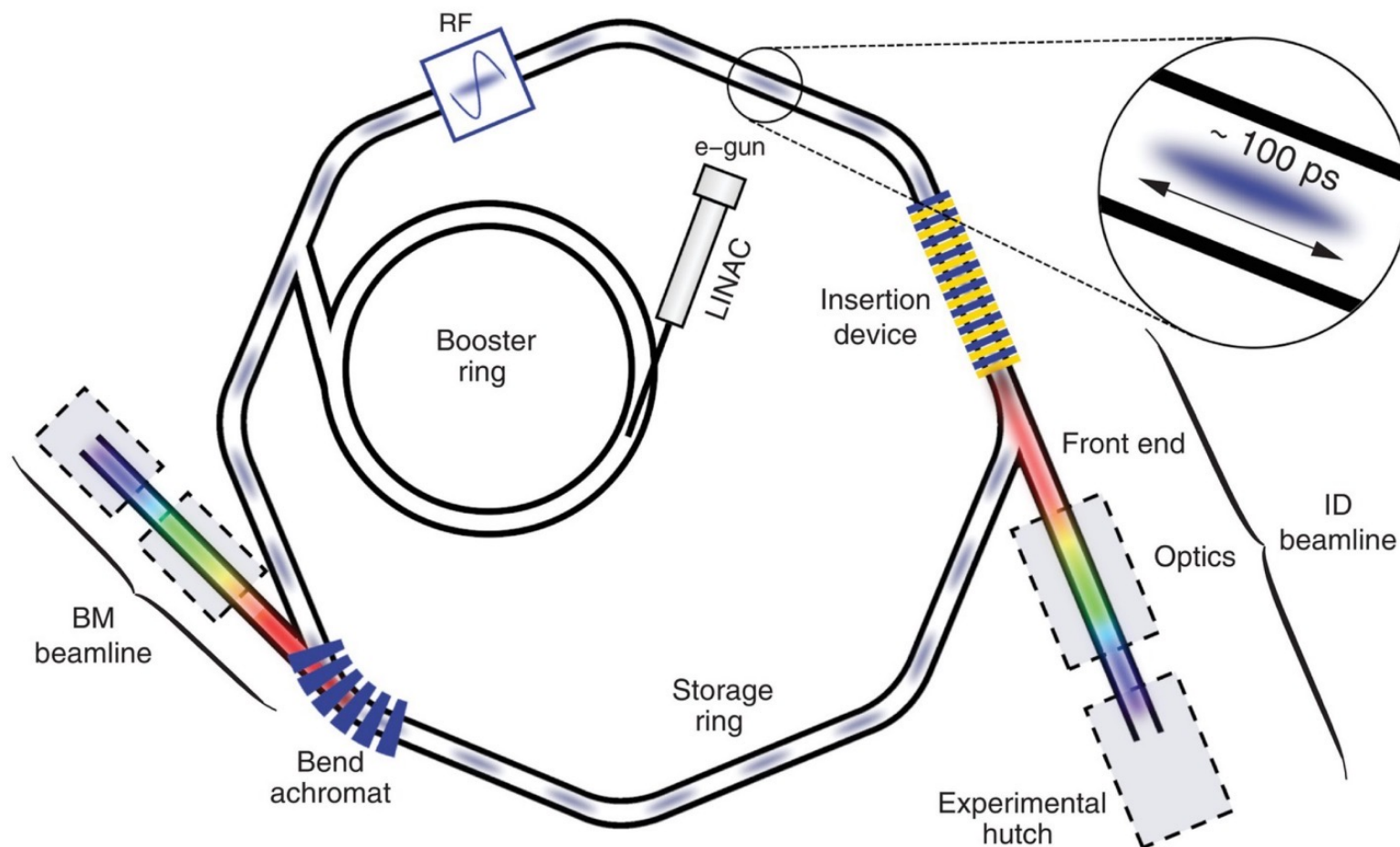
Brilliance



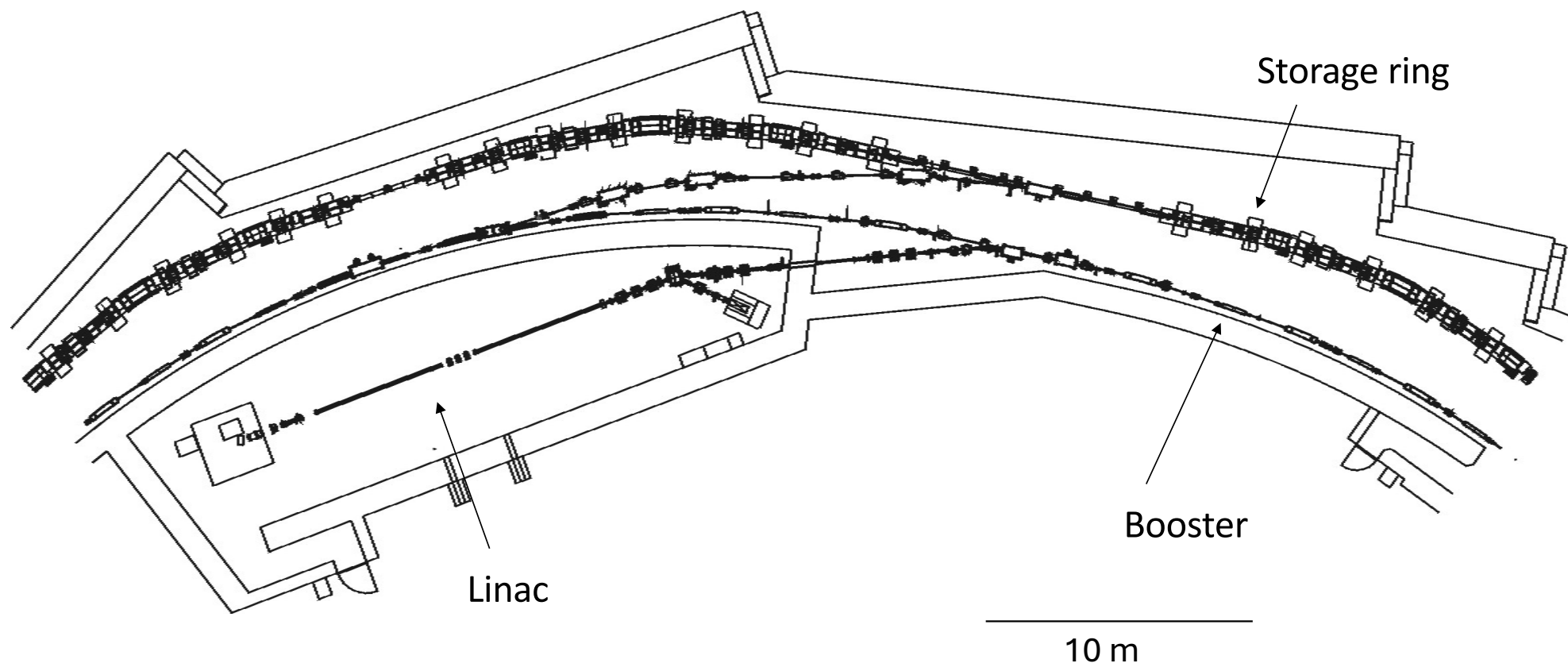
Classification of synchrotron light sources:

- **1st generation:** parasitic use of synchrotrons built for particle physics
- **2nd generation:** dedicated synchrotrons / storage rings built for photon science
- **3rd generation:** dedicated storage rings optimized for operation with insertion devices (wigglers and undulators)
- **4th generation:** diffraction-limited storage rings

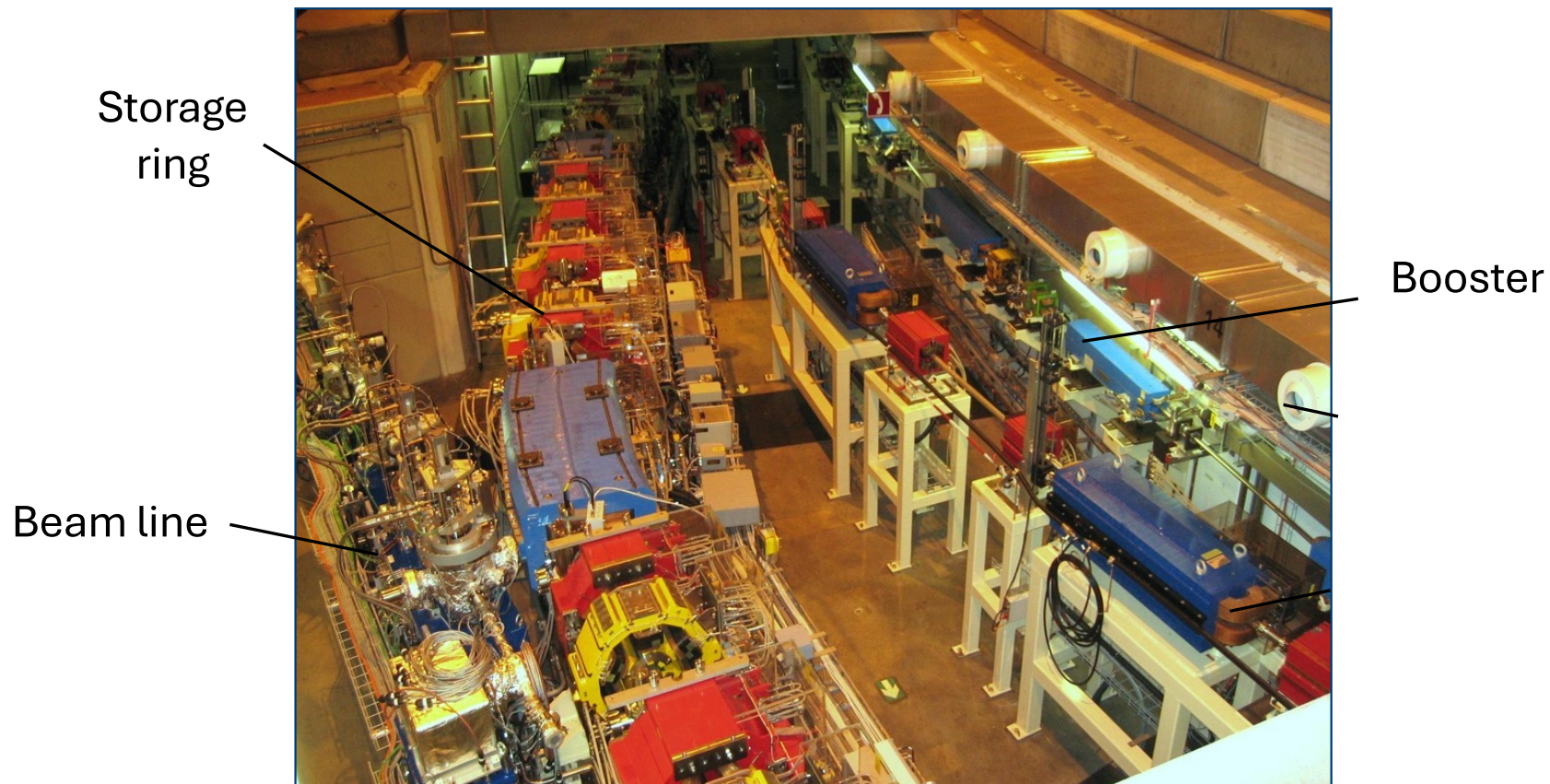
Synchrotron layout



Injector

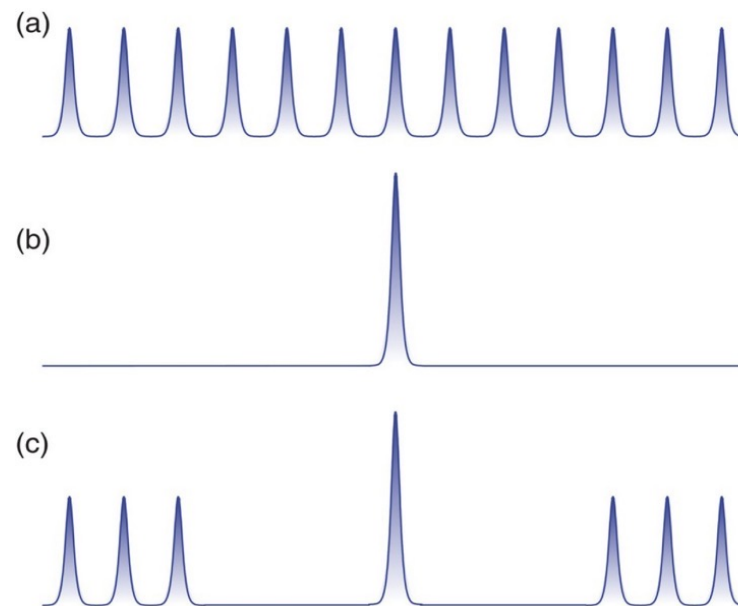


Swiss Light Source

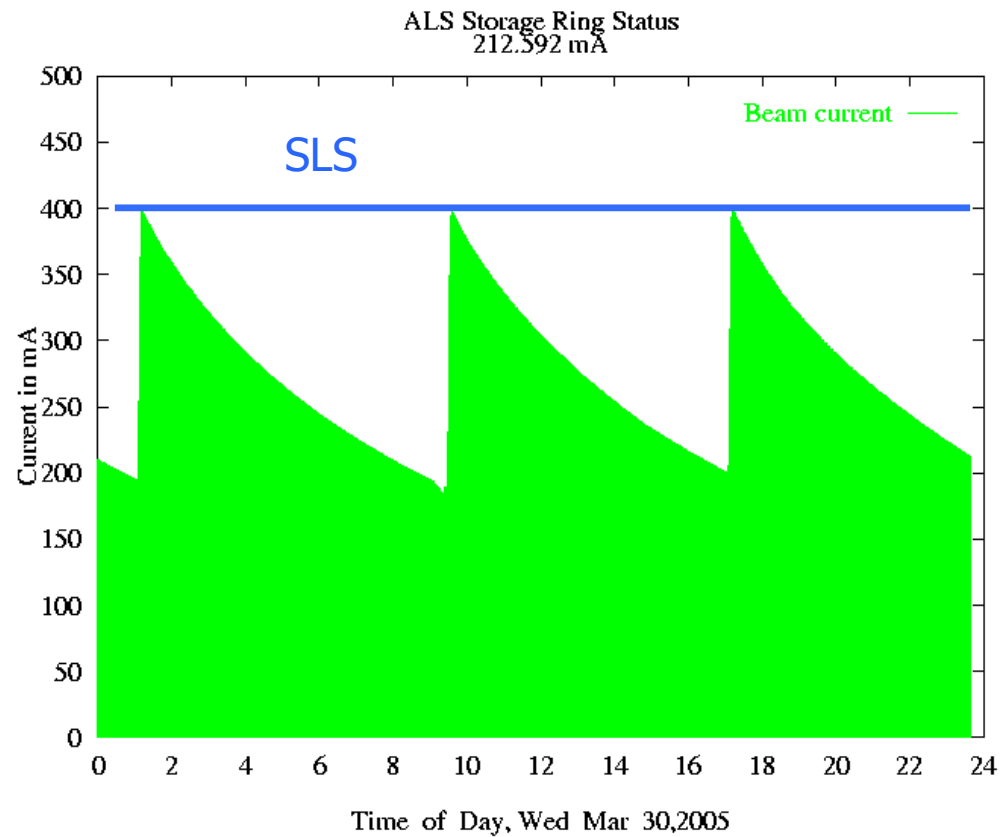


Filling modes

Normally, the ring is filled with bunches of electrons equally spaced from one another by a few nanoseconds. Some time-resolved experiments may require that only a single bunch of electrons is in the storage ring. In the hybrid ‘camshaft’ mode, a bunch (often containing more charge) is isolated by approximately ± 100 ns by dropping bunches on either side.



Filling modes



ALS (Berkeley , California):

Lifetime ~ 10 h ,

Refilled every 8 h

Beam current: 400 mA => 200 mA

SLS:

Lifetime ~8 h ,

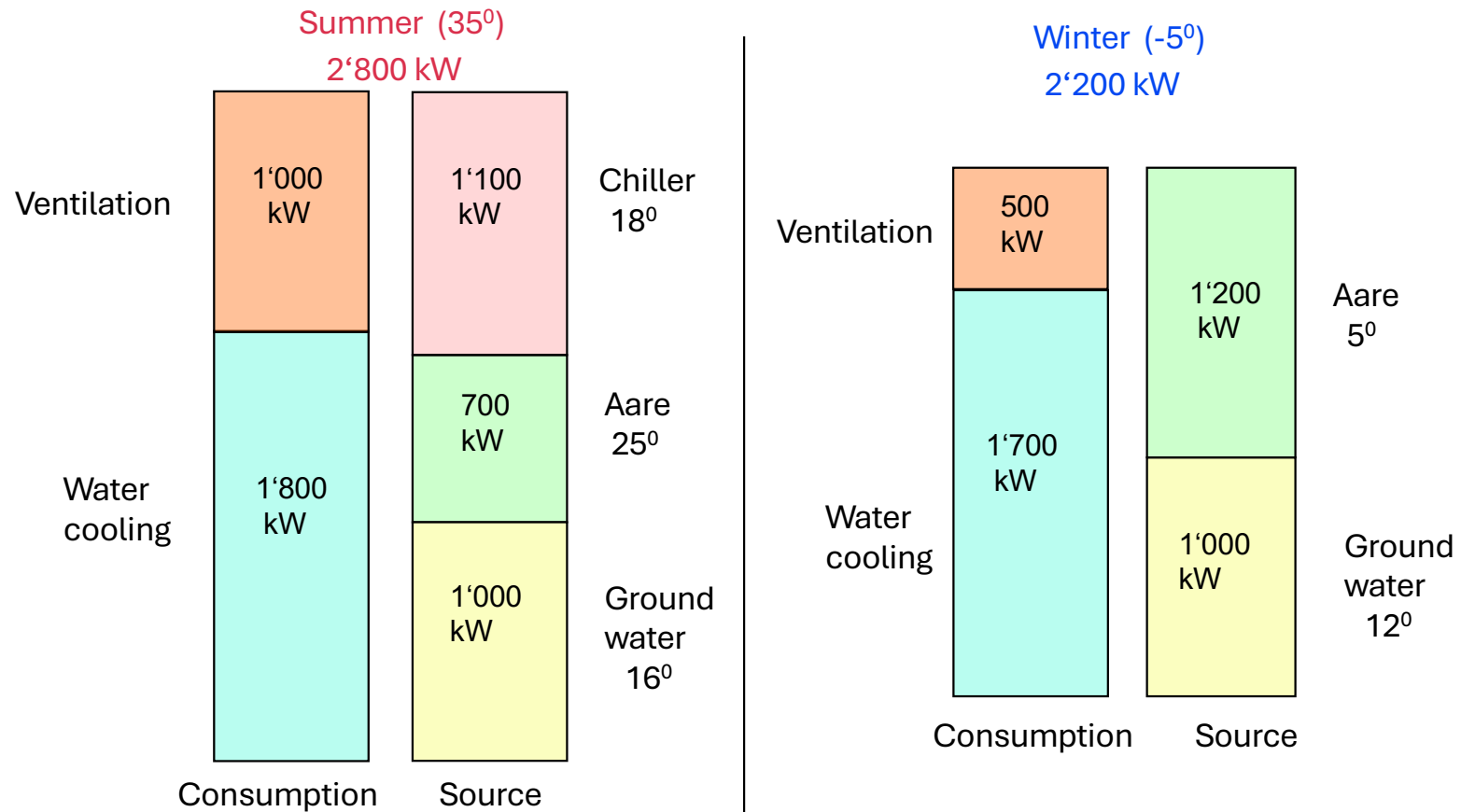
Top-up every 1½ min.

Beam current: 401 mA => 400 mA

Filling modes @ ESRF

Beam mode & characteristics	Intensity	Lifetime
Uniform filling mode , 992 bunches are equally distributed around the whole circumference of the storage ring. Refill time: ~3 minutes	200 mA	60 h
2*1/3 filling mode , 2 times one third of the storage ring is filled. The 2 one thirds are separated by an empty gap of 1/6th of the ring. Top up time: ~3 minutes	200 mA	55 h
7/8 + 1 filling mode , A train of 868 bunches (7/8 of the Storage Ring circumference) filled with 200 mA (0.23 mA / bunch). Both edges of the train are filled with 1 mA single bunch. The remaining 1/8 gap is filled in its center with a cleaned 2 mA-single bunch Filling time: ~ 5 minutes	200 mA	55 h
Hybrid mode 24*8 + 1 filling mode , One clean 4mA single bunch diametrically opposed to a ~ 196 mA multi-bunch beam composed of 24 groups of bunches spread over 3/4 of the storage ring circumference. Filling time: ~ 10 minutes.	200 mA	25 h
16 bunch filling mode , 16 highly populated and equally spaced bunches. Purity < 10 ⁻⁷ Filling time: ~5 mins.	90 mA	10 h
4*10 filling mode , 4 equidistant bunches _ 10 mA/bunch Filling time: ~ 5 minutes	40 mA	6 h

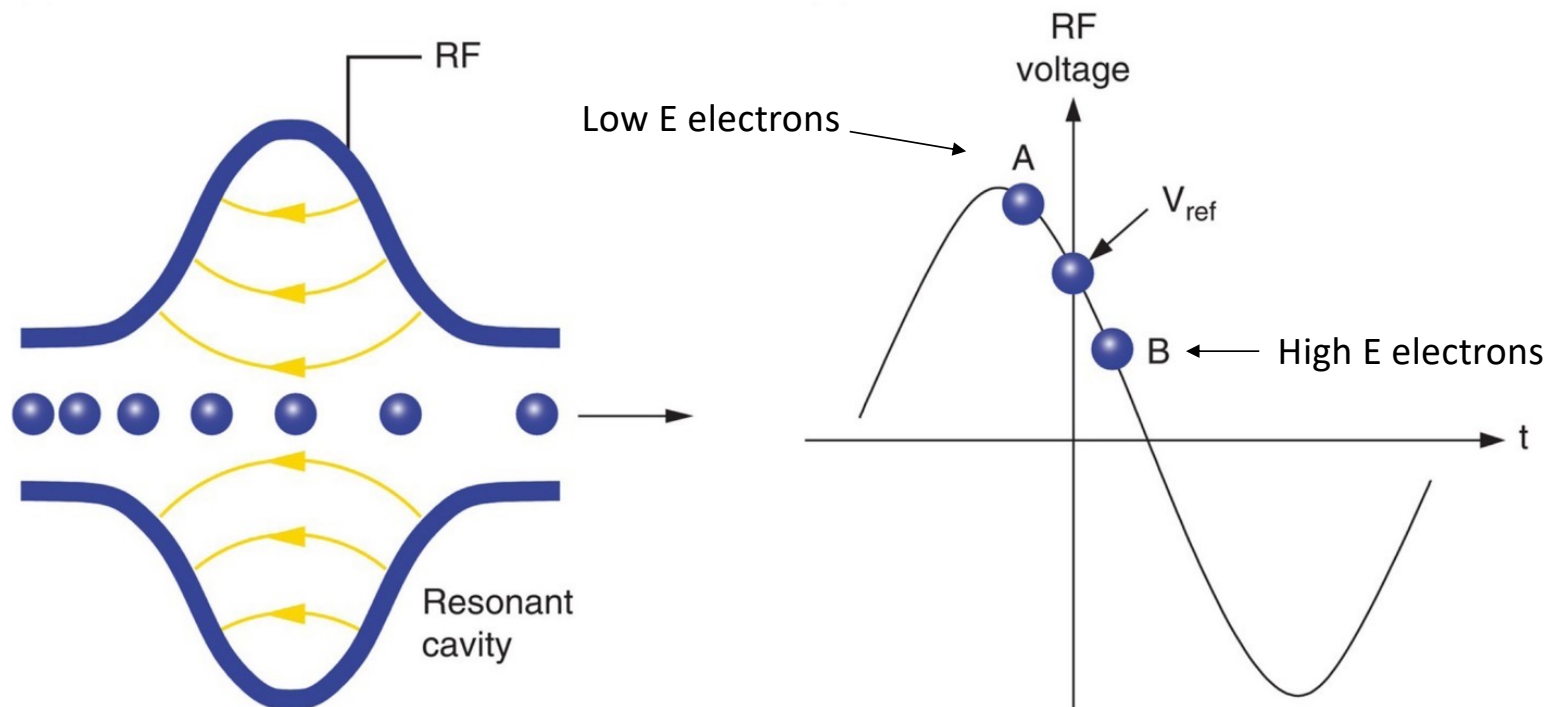
Cooling - SLS



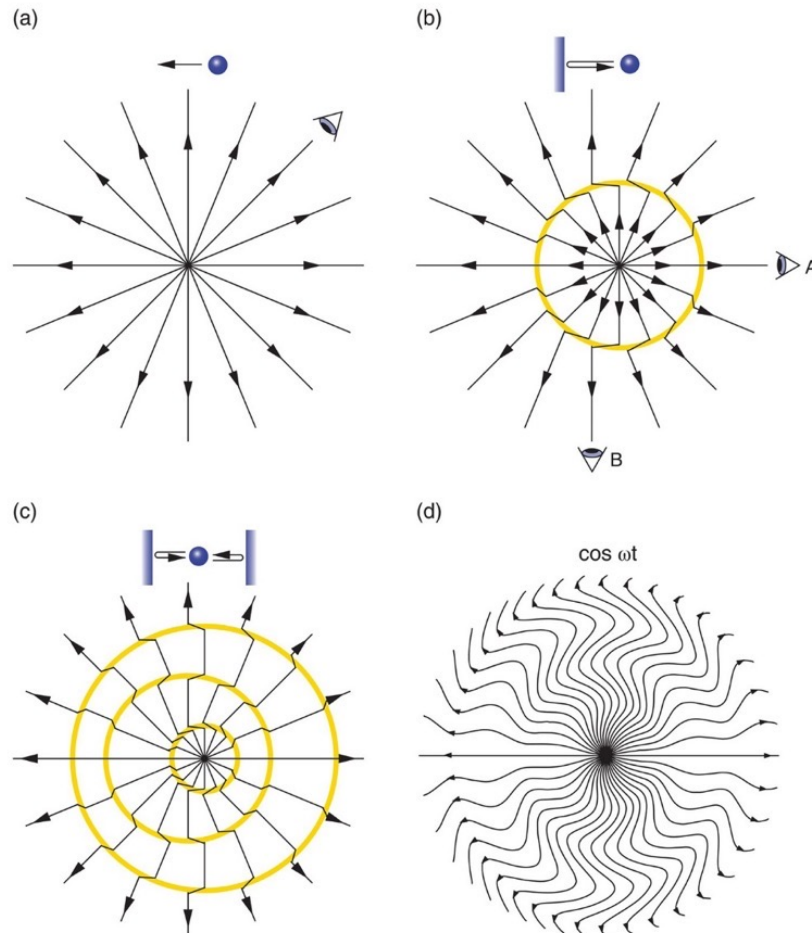
Acceleration

Electrons entering the resonant cavity at the right moment in its voltage cycle are accelerated by the electric field within the cavity generated by the klystron source.

‘Slow’ electrons entering the RF cavity at *A* will be given an extra boost, while ‘fast’ electrons (at *B*) receive less energy.



Generation of electromagnetic radiation (EMR)



a) At rest or constant velocity, no component of the field in any transverse direction

b) Single abrupt acceleration

- Distortion of electric field lines
- Amplitude of radiation $\sim \cos \theta$

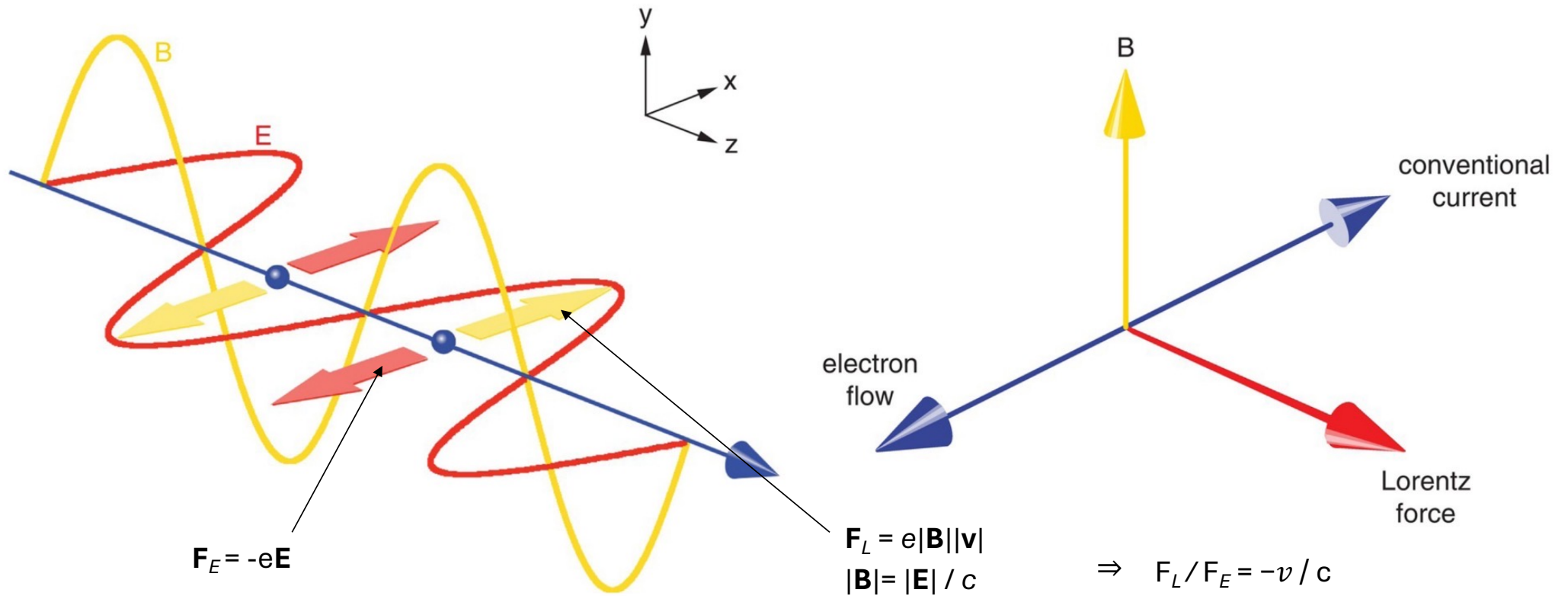
c) Multiple abrupt accelerations

- Set of pulses at regular intervals
- Multiple frequencies

d) Simple harmonic

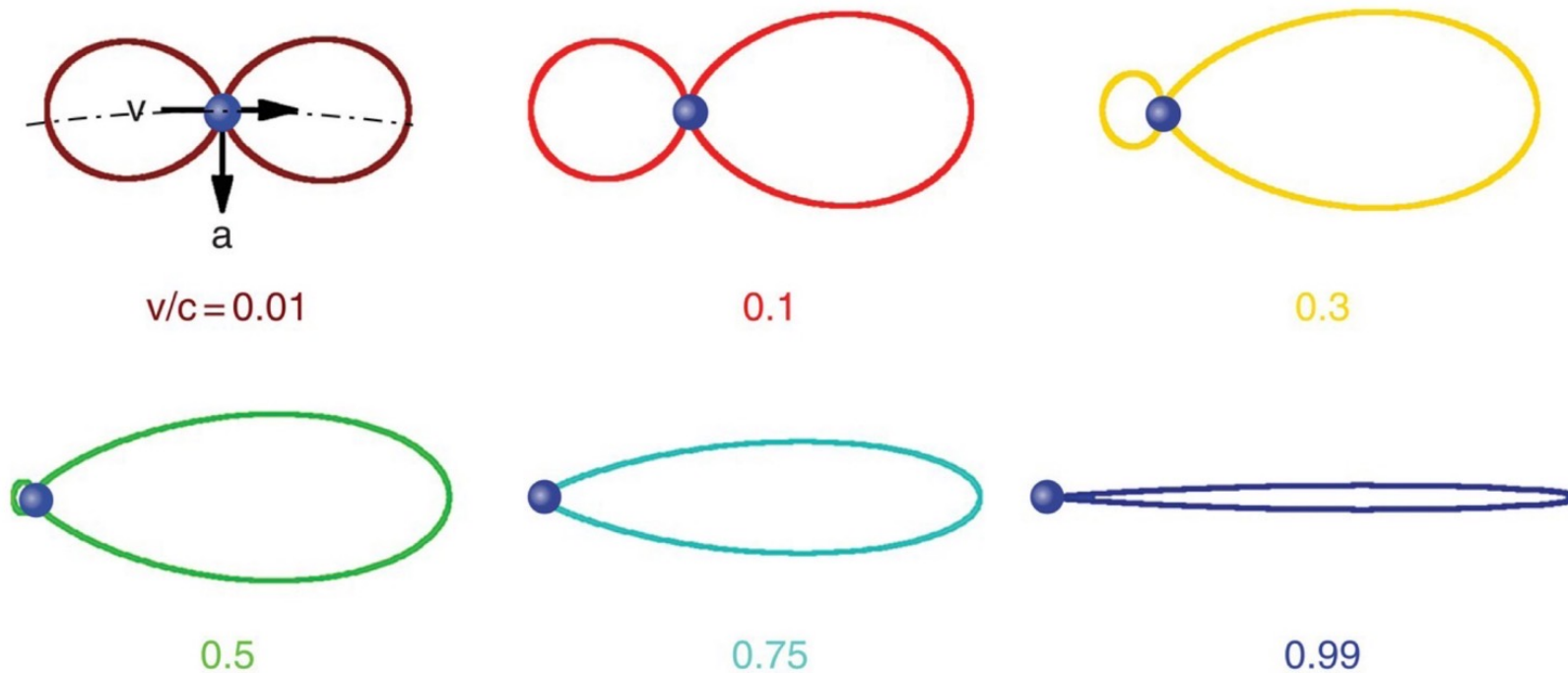
- Radiation with single frequency

Forces acting on charged particle by EMR



F_L only becomes nonnegligible compared to F_E for electrons travelling at relativistic velocities.

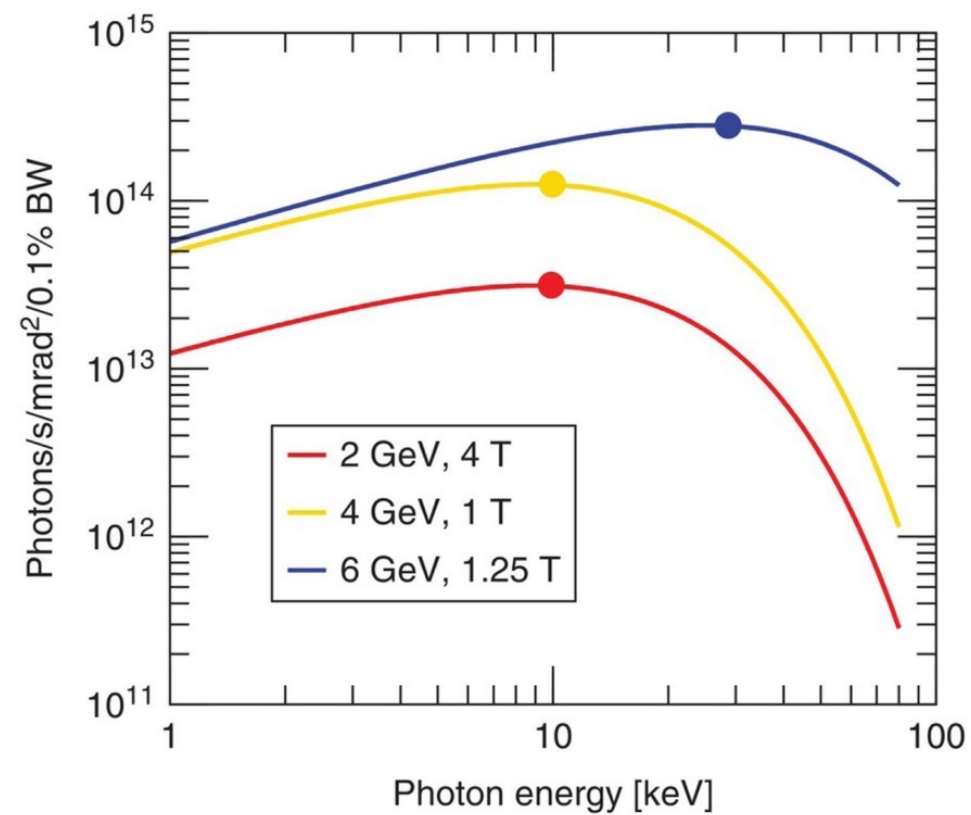
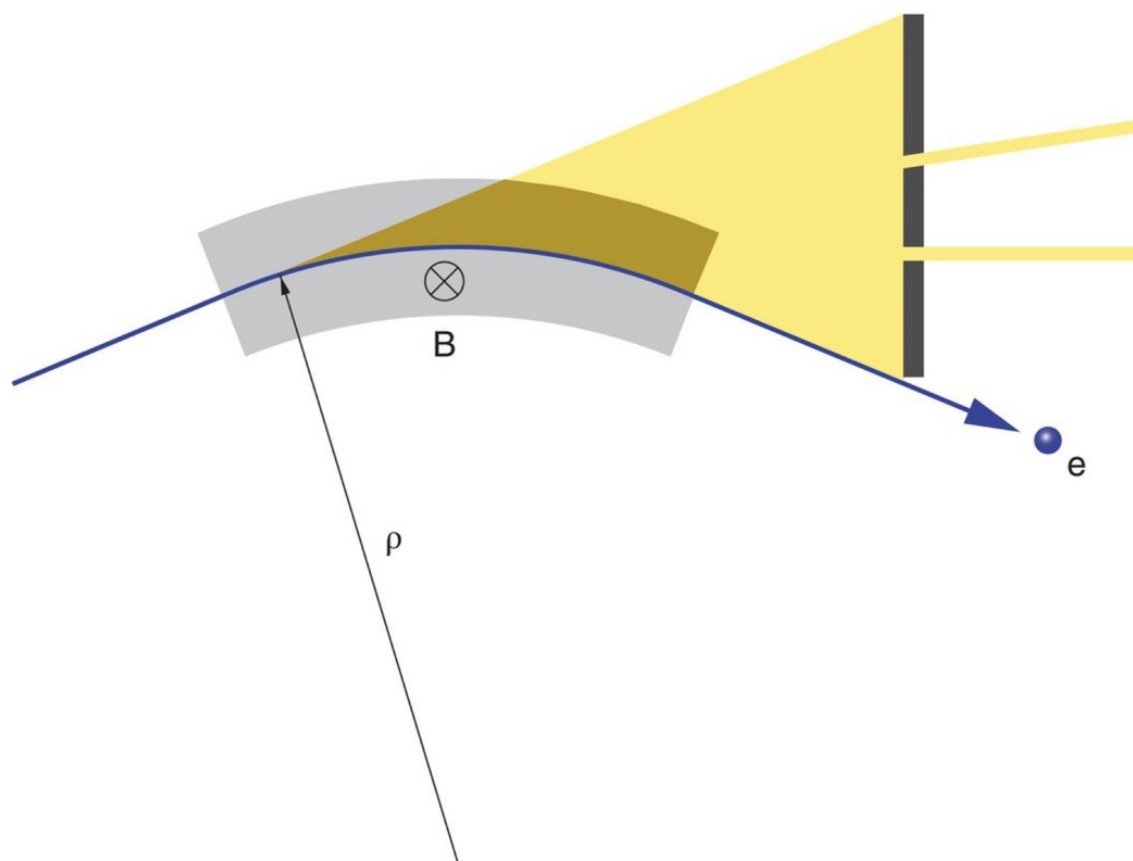
Radiation of an accelerating charged particle



$$P = \frac{2}{3} \frac{e^4 c^3}{(m_e c^2)^4} \mathcal{E}^2 B^2 = 6.763 \times 10^{-18} \mathcal{E}^2 B^2$$

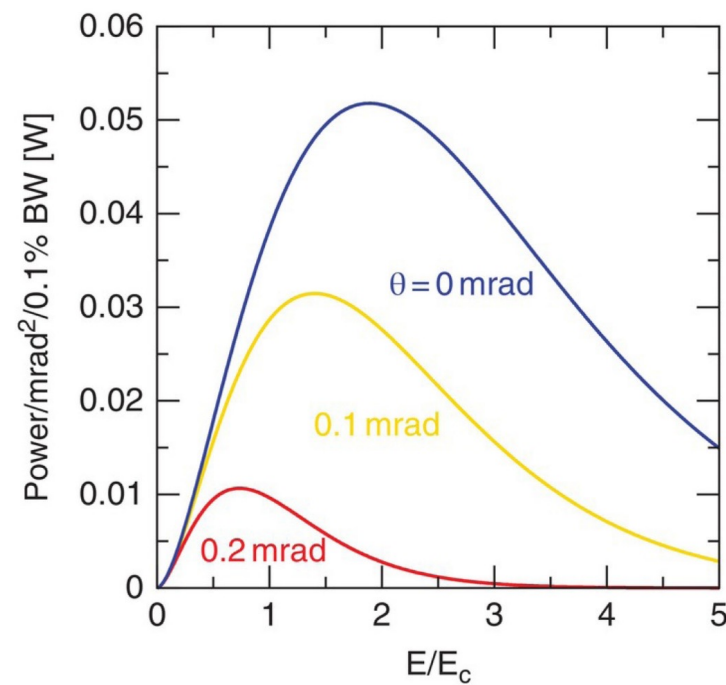
E.g. 3 GeV synchrotron storing a current of $I = 250$ mA and using bending magnets with field strengths of 1.4 T equates to a radiative power output due to circulation of the electrons alone of 186 W.

Bending magnet

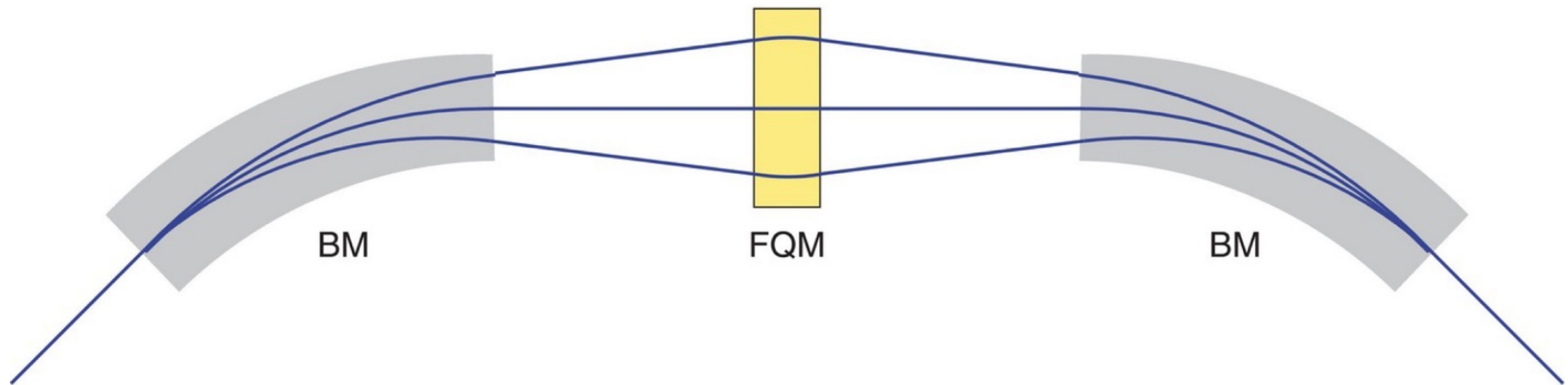


Bending magnet

The power spectra per mrad² and 0.1% bandwidth of bending-magnet radiation as a function of energy for different angular positions out of the plane of the electron orbit. The power spectrum falls off and the positions of the spectrum maxima move to lower energies with increasing angle from the axis.

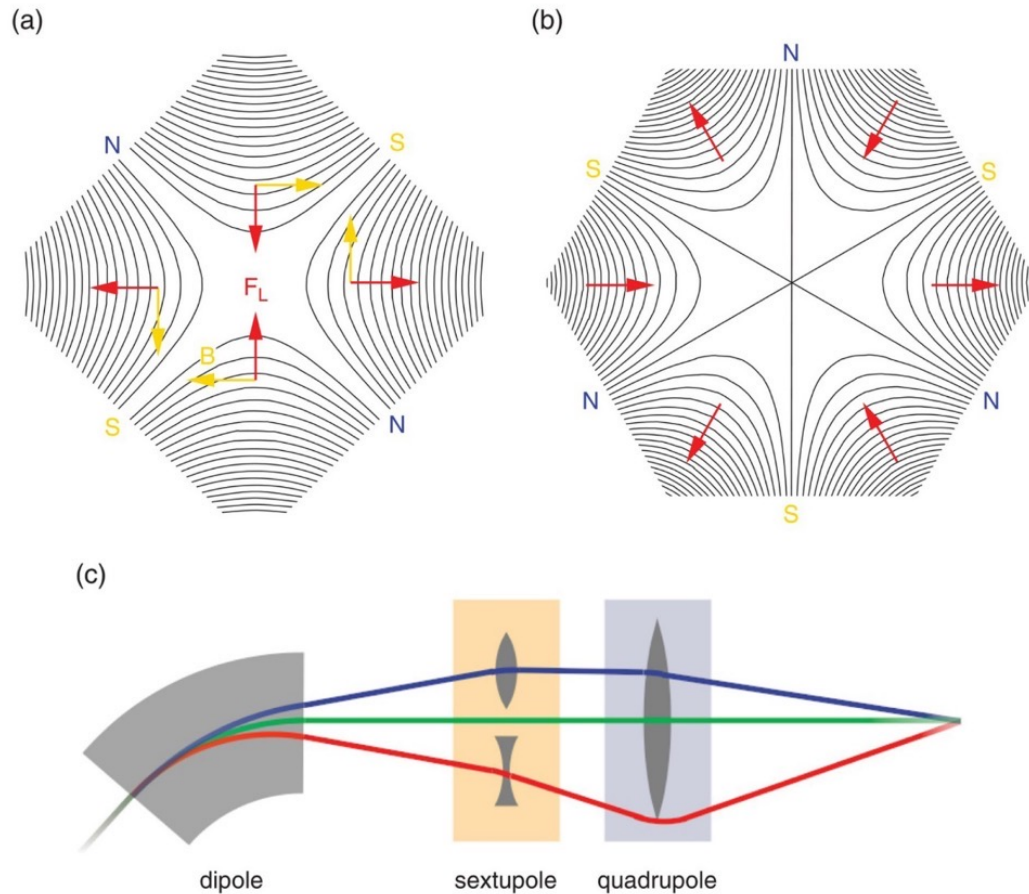


Electron beam path correction



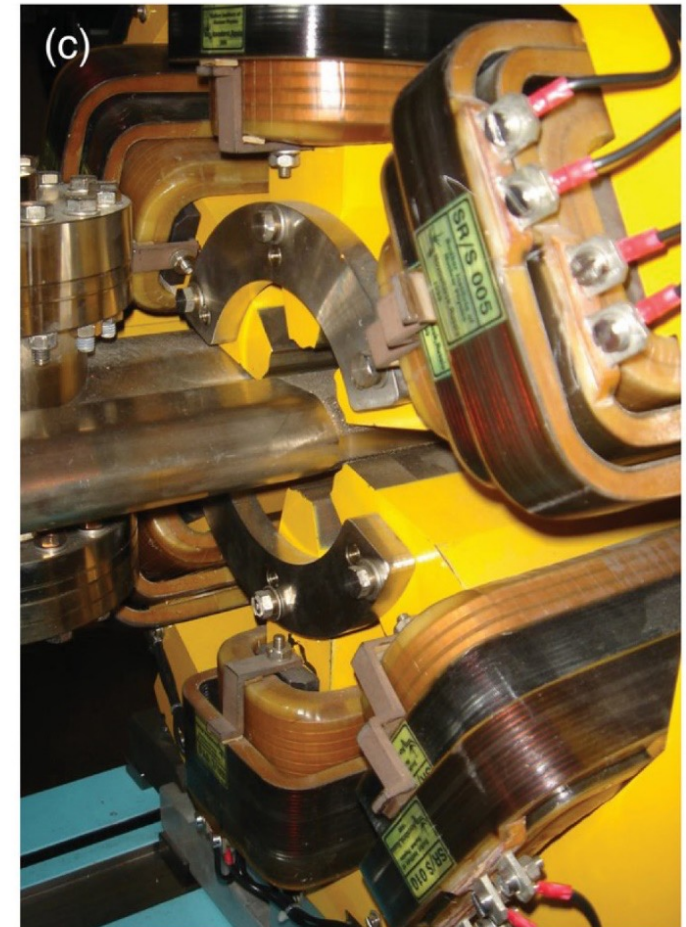
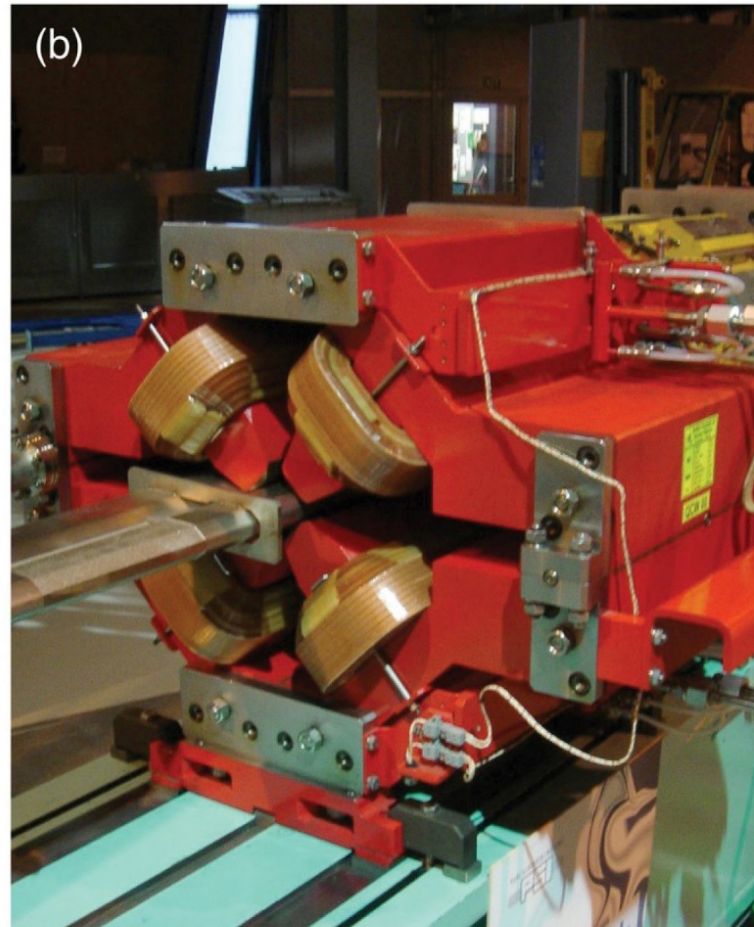
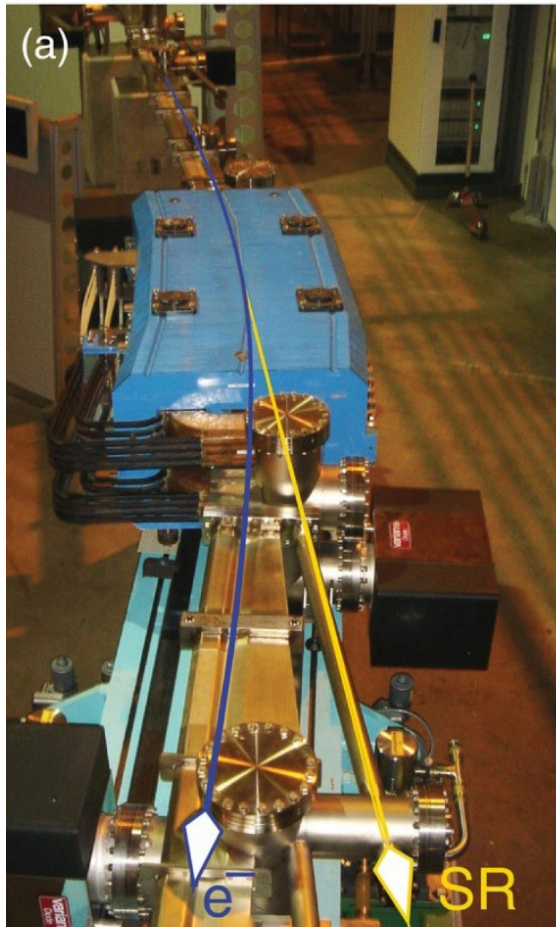
Angular dispersion of electrons of different energies as they pass through a bending magnet and the consequent increase in the electron-beam emittance can be corrected in a double-bend achromat by placing a focussing quadrupole magnet (FQM) symmetrically in between two identical dipole bends (BM). Lower-energy electrons are bent through larger angles by the dipoles than are those of higher energy.

Electron beam path correction

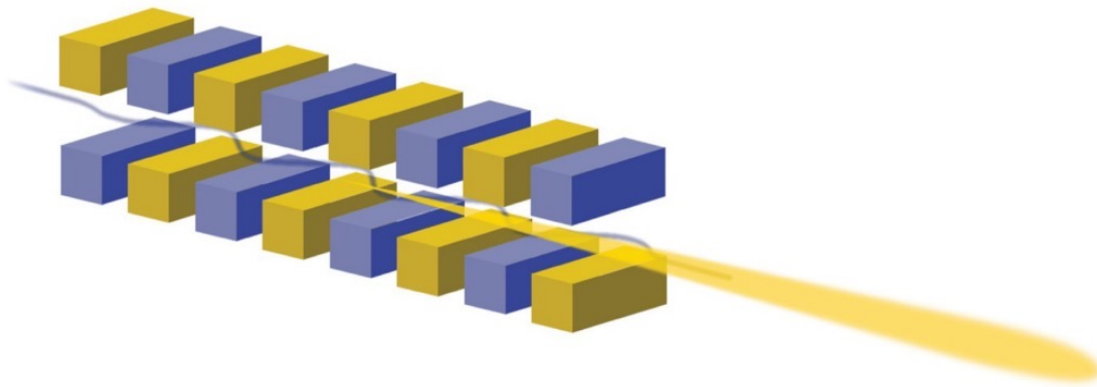


- In a quadrupole the Lorentz forces are such that in one plane (here, vertical) they focus the electron beam, while in the orthogonal plane (here, horizontal), the electrons are defocused. Therefore one could place two quadrupoles subsequently.
- Sextupoles have focal lengths that are inversely proportional to the distance of the electrons from the central axis
- The dispersion of the electrons' energies are sorted for their momentum by the sextupole in the plane perpendicular to its axis, which corrects for the quadrupole chromatic focussing error.

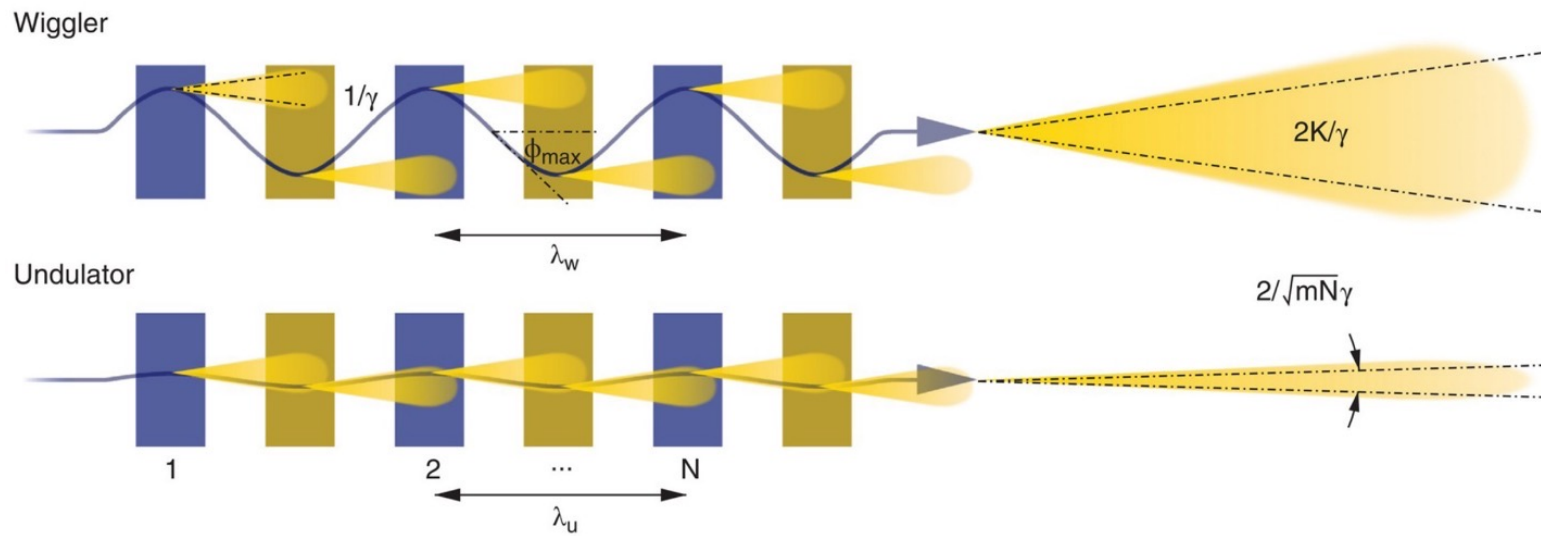
Magnetic lattice



Insertion devices



The lower magnet array of a U15 (15 mm period) undulator.



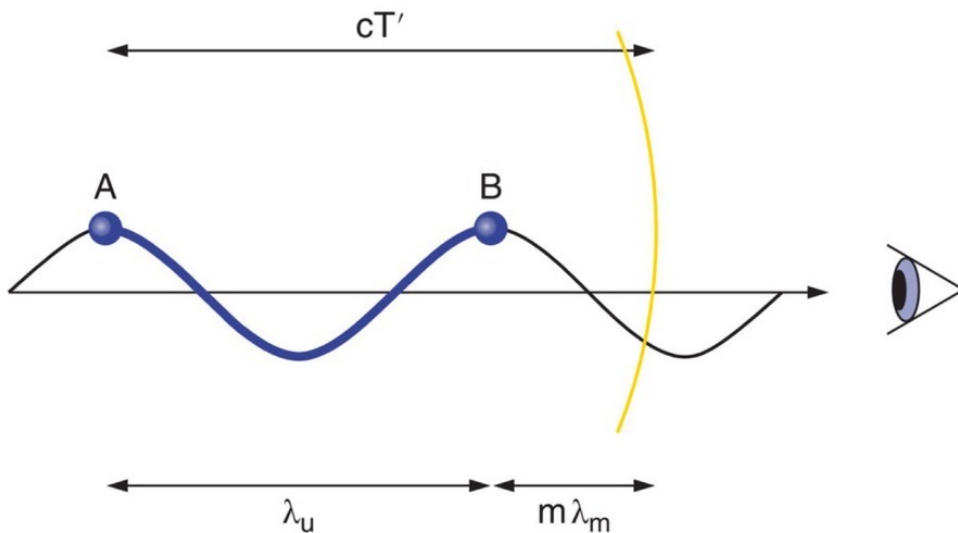
Undulator



Permanent magnets:
62 periods à 56mm

Undulators

Constructive interference between electromagnetic wavefronts emanating from equivalent points on the undulations occurs when their separation is an integer multiple of the radiation wavelength. The curved path S taken by the electron travelling at the relativistic velocity v between A and B is highlighted in blue



$$m\lambda_m = cT' - \lambda_u$$

$$\Delta E = \frac{2hc\gamma^2}{\lambda_u(1 + K^2/2)}$$

$$\phi_{\max} = K/\gamma$$

γ is the relativistic Lorentz factor

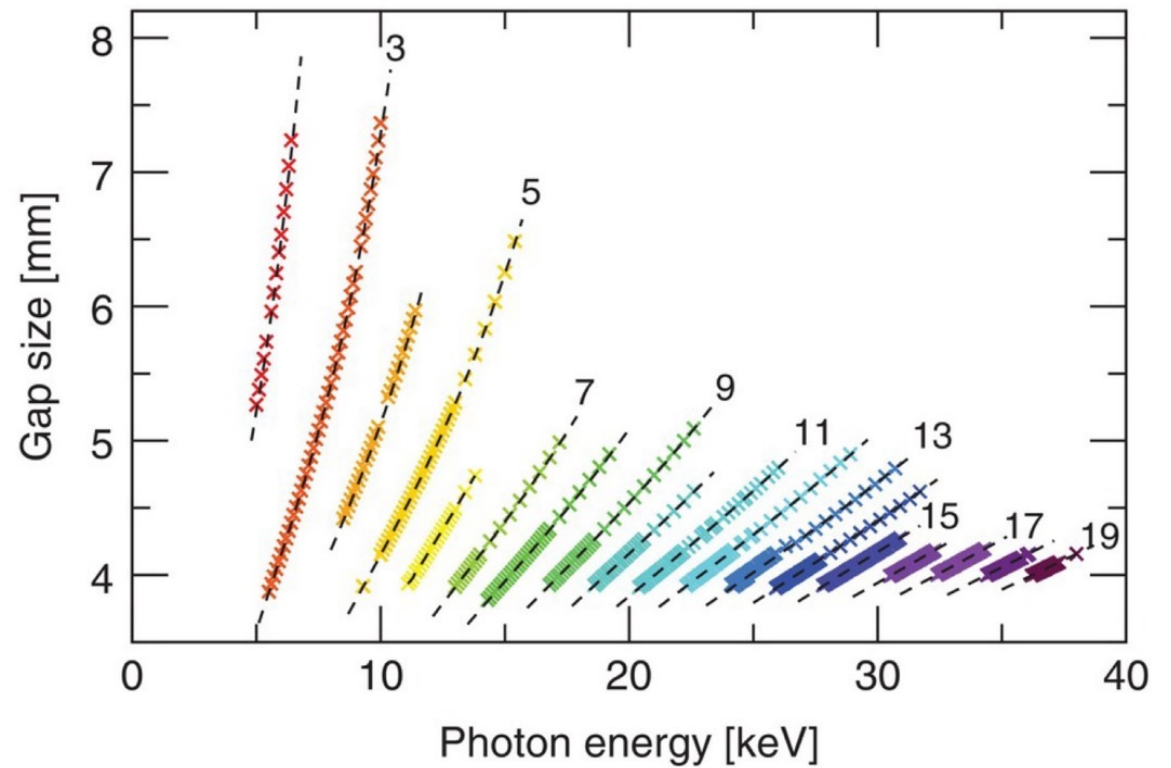
λ_u undulator spatial period

K magnetic deflection parameter

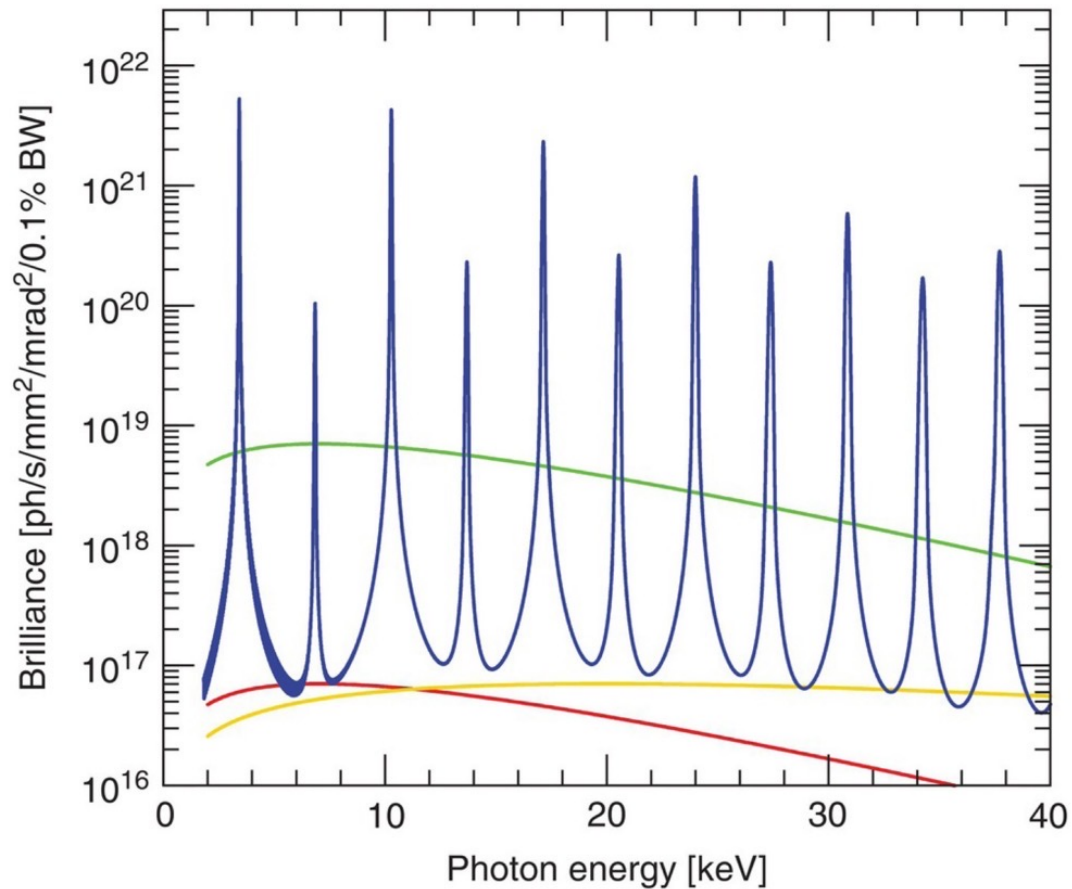
$\phi_{\max}$ = maximum angular deviation of electron oscillations

Undulator

The spectrum can be tuned by varying K . This is achieved by changing the gap between the two sets of magnetic poles and thereby the magnetic-field strength B_0 .

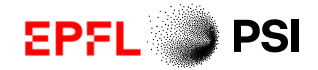


Comparison insertion devices



Comparison of brilliances at a 3 GeV DLSR running at 400 mA between a U14 undulator with $K = 1.6$ (blue), a bending magnet with $B = 1.41$ T (red), a superbend with $B = 4$ T (yellow), and a wiggler with the same field strength as the bending magnet and 100 periods (green).

Diffraction limited storage rings



Technological improvements in magnetic lattice technology has led to so-called diffraction limited storage rings (DLSR).

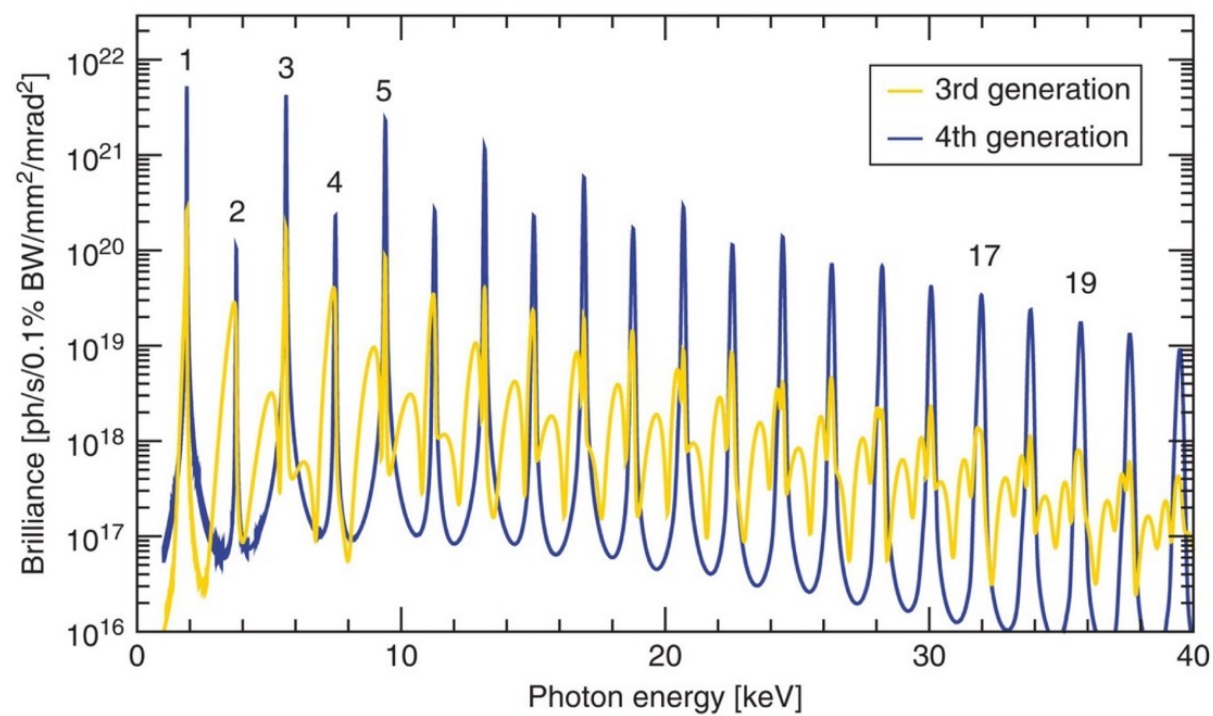
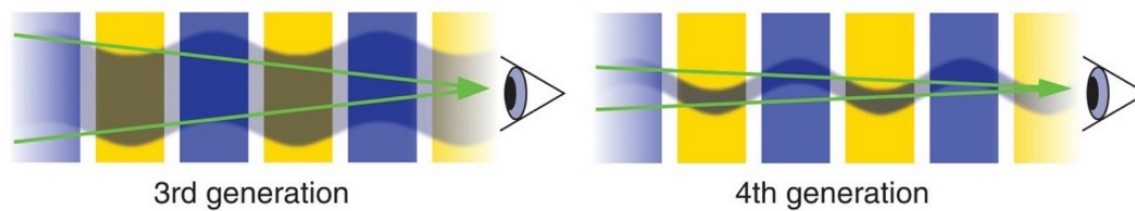
DLSR have 40x smaller horizontal emittance

- In third-generation facilities, the horizontal electron-beam width is, at approximately $100\text{ }\mu\text{m}$, two orders of magnitude larger than the oscillation amplitude A ($\sim 1\text{ }\mu\text{m}$)
- In DLSR the horizontal electron-beam width is of the order of $10\text{ }\mu\text{m}$.

Enhanced brilliance at spectral peaks

‘Cleaner’ energy spectra

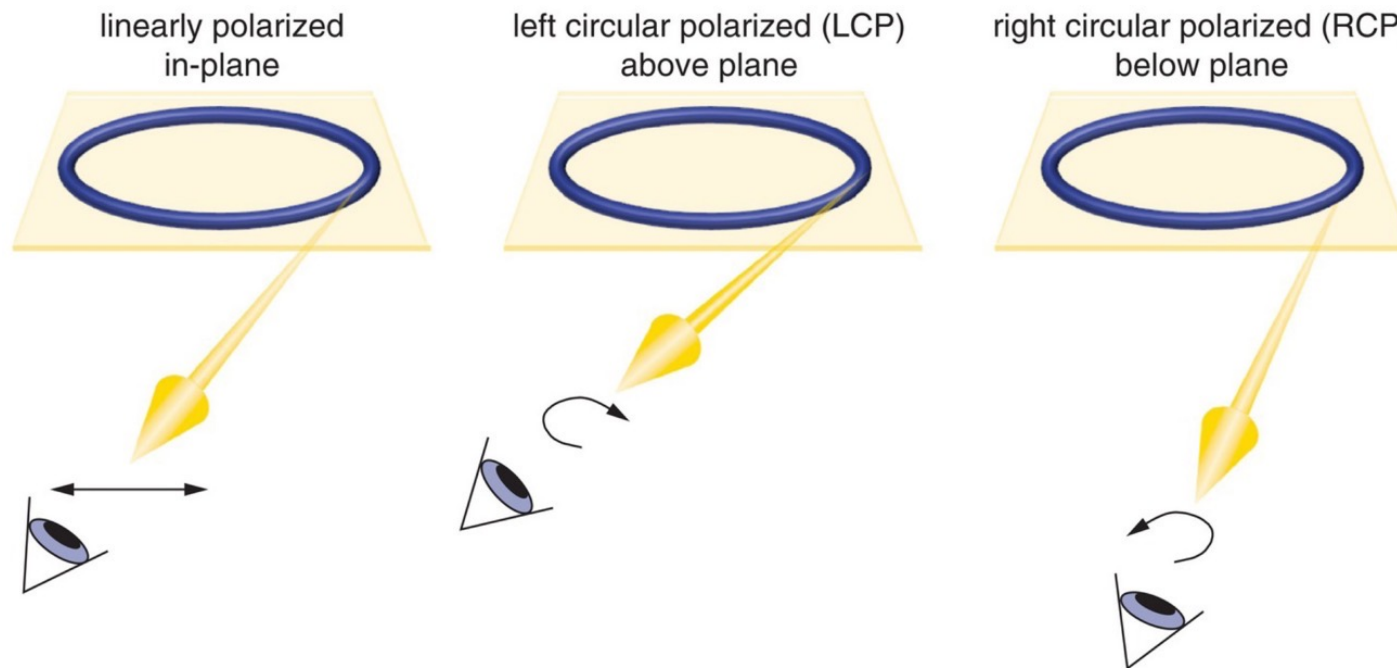
Diffraction limited storage rings



Polarization

Viewed in the electrons' orbital plane, the polarization of the radiation is linear in that plane, as the electrons appear to oscillate in the horizontal plane.

If the observer is above the plane of the storage ring, the electrons will appear to execute an elliptical orbit in a clockwise direction. This angular momentum from the observer's frame of reference is transferred to the emitted photons, which are left-circularly polarized

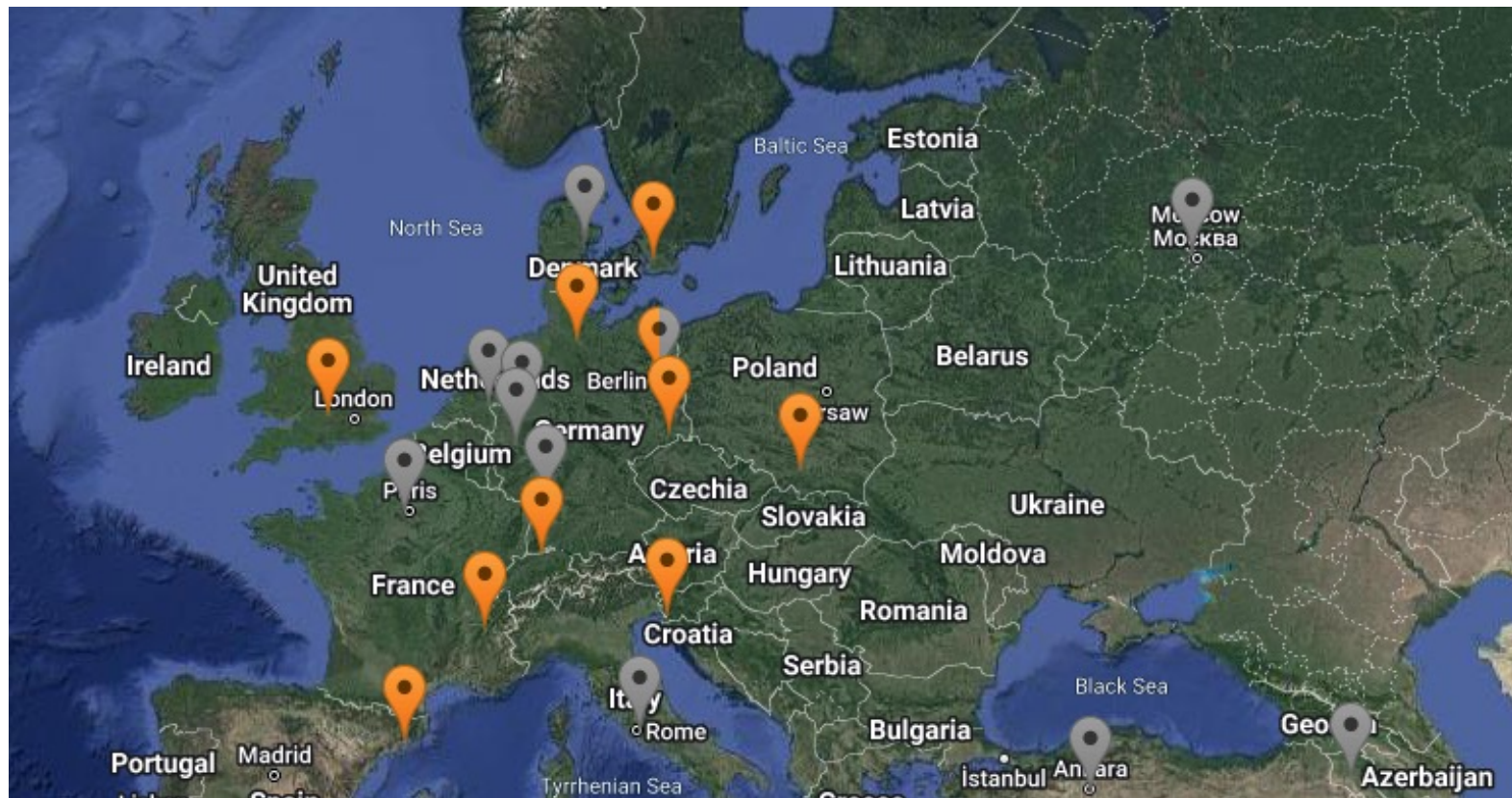


Synchrotrons around the world



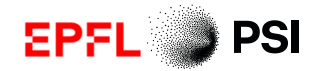
lightsources.org/lightsources-of-the-world

Synchrotrons around the world



lightsources.org/lightsources-of-the-world

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- **X-ray Free Electron Laser**

X-ray free electron lasers



Optical short pulse lasers

Pulse duration: +++ (~fs)

Wavelength: - - - (~800 nm)

- Fastest processes can be analyzed
- Spatial resolution limited

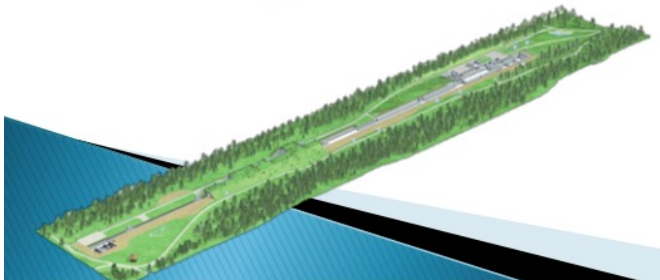


Synchrotrons

Pulse duration: o (~ps)

Wavelength: +++ (~0.1 nm)

- Temporal resolution limited
- Wavelength allows for atomic resolution



X-ray free-electron lasers

Pulse duration: +++ (~fs)

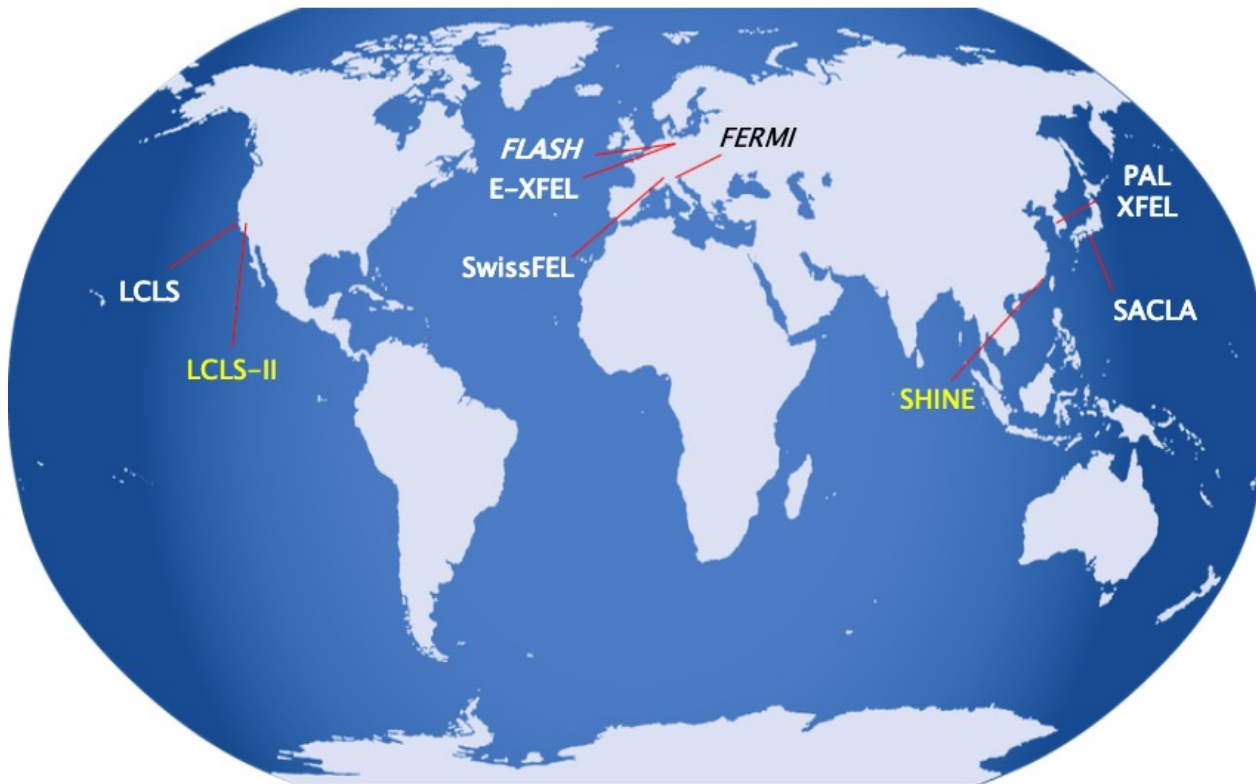
Wavelength: +++ (~0.1 nm)

- Fastest processes can be analyzed
- Wavelength allows for atomic resolution

X-ray free electron lasers

	Optical lasers	Synchrotrons	FELs
Time atomic res.	Yes	No	Yes
Spatial atomic res.	No	Yes	Yes
Pulse energy	few mJ	< nJ	few mJ
Rep rate	< 0.1 MHz	~500 MHz	~0.1 – 1 MHz
Facility cost	Moderate (≥ 10 MEUR)	High (≥ 100 MEUR) (but government funded)	Very High (≥ 300 MEUR) (but government funded)
User accessibility	Poor, normally buy it yourself!	Good: ~50 user facilities worldwide (<50 beamlines each)	Fair: ~5 user facilities worldwide (<10 beamlines each)

XFEL around the world



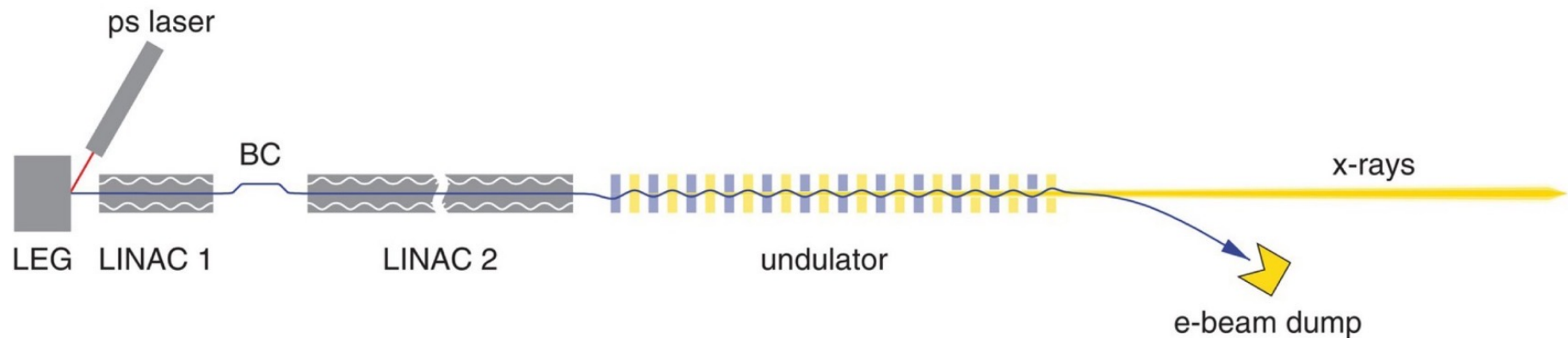
- **FLASH:** first soft X- ray with high repetition rate (MHz), 2007
- **LCLS:** first hard X- ray, 2009
- **SACLA:** compact hard X-ray, 2011
- **FERMI:** first soft X- ray seeded-FEL, 2013
- **PAL-XFEL:** hard X- ray with low (20 fs) timing jitter, 2016
- **E-XFEL:** hard X-ray with high repetition rate, 2017
- **SwissFEL:** compact hard X-ray driven by low emittance beam, 2017
- **Future facilities:** LCLS2 and SHINE (high repetition rate)

XFEL around the world

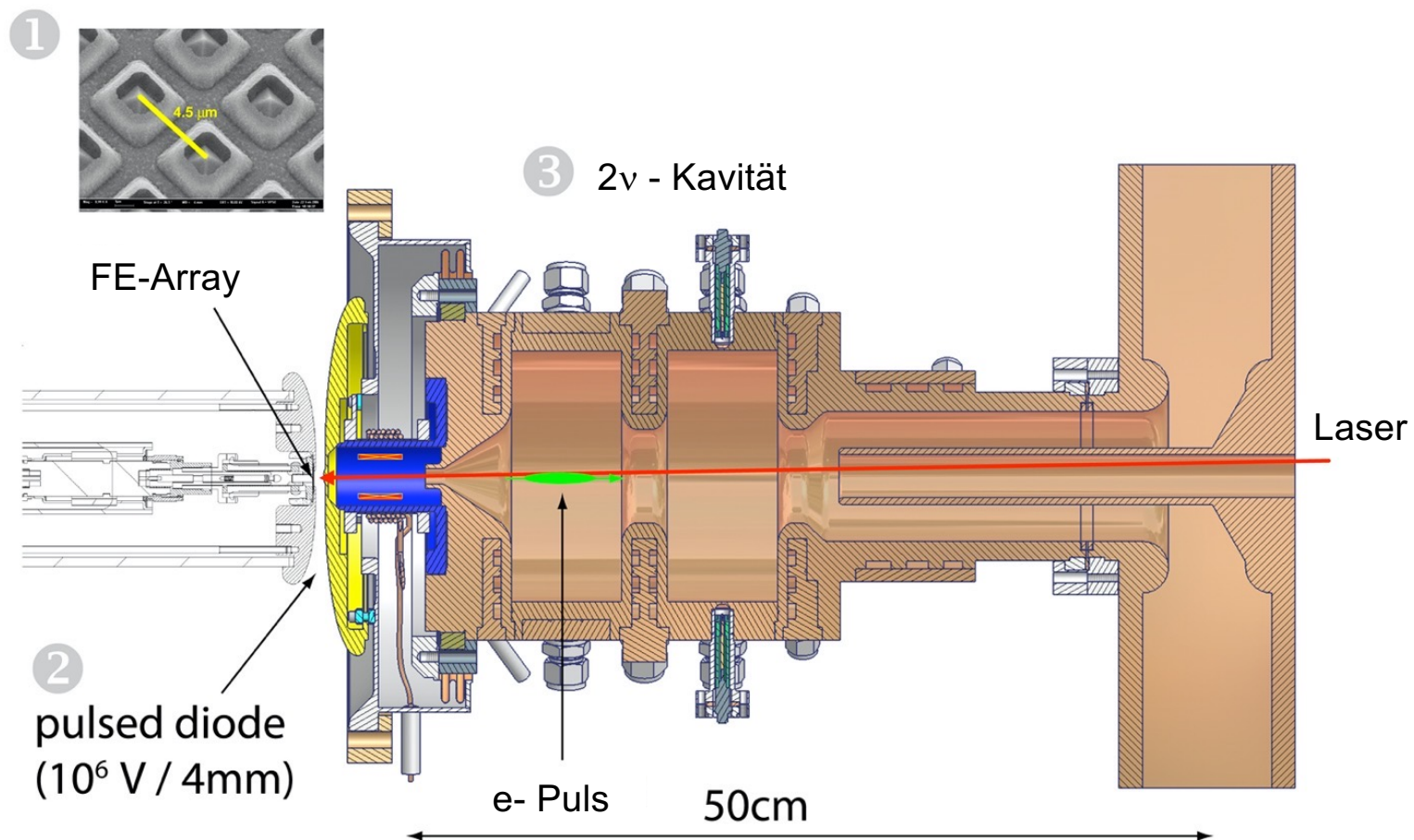
	LCLS	SACLA	PAL-XFEL	E-XFEL	SwissFEL
Country	USA	Japan	Korea	Germany	Switzerland
Starting oper. year	2009	2010	2017	2017	2018
Energy (GeV)	14.3	8.5	10	17.5	5.8
Length (km)	3.0	0.75	1.1	3.4	0.74
Construction cost (M\$)	415	370	400	1600	280
Pulses per second	120	60	60	27000	100
Emittance (nm)	400	1000	550	<600	200

Schematic XFEL

- Electron bunches are emitted from a low-emittance gun (LEG) irradiated by picosecond laser pulses. They are then accelerated in a short LINAC (LINAC 1), compressed longitudinally using one or more bunch-compressor magnet chicanes (BC), then further accelerated using a much longer LINAC (LINAC 2) before entering a long undulator, typically a few hundred metres in length.
- The SASE process along the undulator produces highly intense x-ray pulses with durations of the order of 50 fs.
- The electrons are deflected after the undulator using a bending magnet and subsequently dumped.

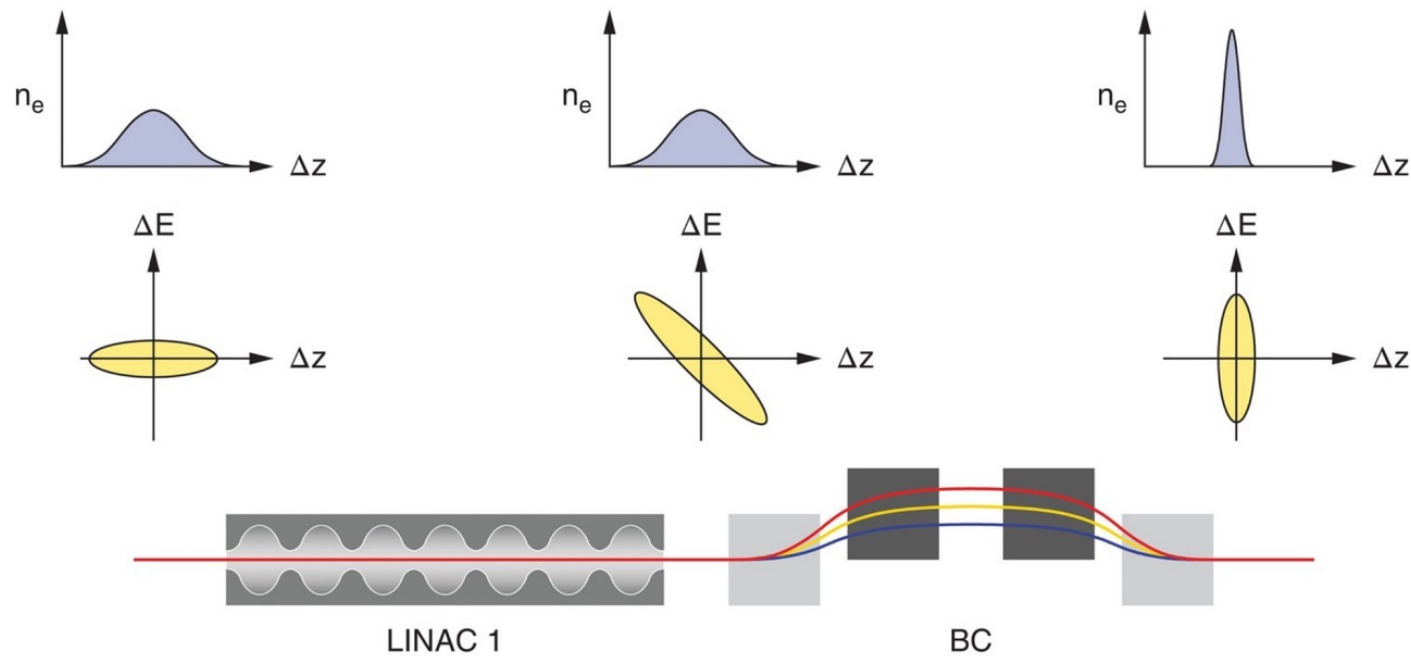


Electron gun

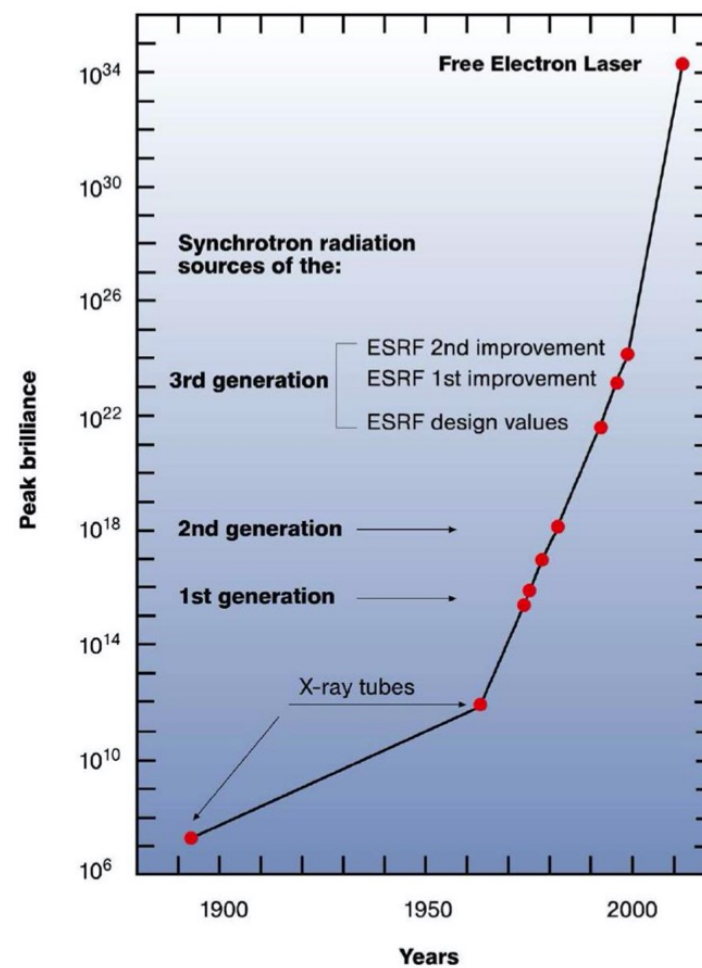


Bunch compression

- Leading electrons are accelerated less than those electrons towards the back of the bunch (this phenomenon is referred to as 'chirping')
- Compression is achieved by allowing an electron bunch to pass through a four-dipole chicane, whereby the faster electrons catch up with the slower electrons at the front, thanks to the shorter path they execute.



Brilliance

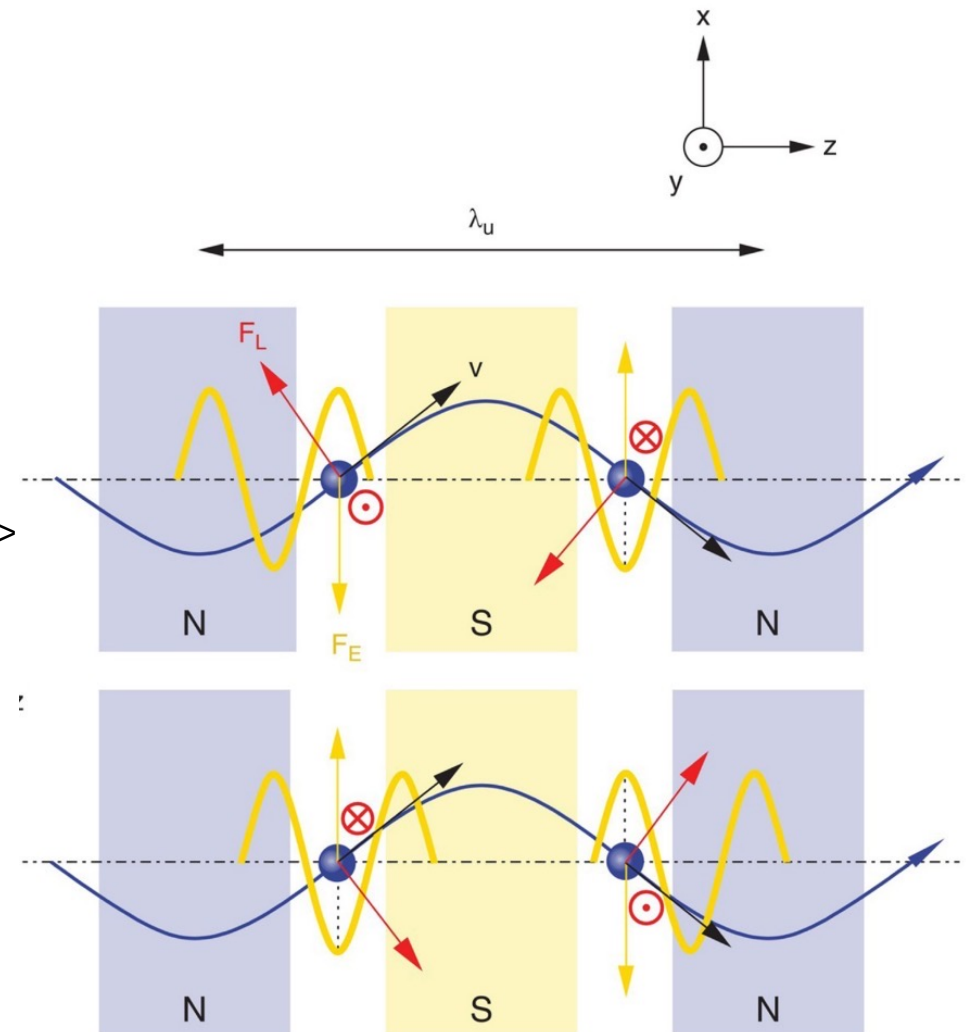


The SASE process

- Self-amplified spontaneous emission
- A typical synchrotron undulator may have a source size of $100 \times 10 \mu\text{m}^2$, and generate bunches of 50 ps duration containing 5×10^6 1 Å photons. This equates to an areal power density of the order of $1.6 \times 10^{11} \text{ W m}^{-2}$, and an electric-field amplitude $E^0 \approx 10^7 \text{ V m}^{-1}$, and **$B_0 \approx 40 \text{ mT}$** , nearly two orders of magnitude smaller than that imposed by the undulator's magnet array. We can thus conclude that, *in a conventional third-generation facility, the impact of the electromagnetic field generated by undulator radiation on the electrons' trajectory is completely negligible.*
- In the case of XFELs, however, the beam is tailored to have as low an emittance and as high an electron density as possible, through the low-emittance gun and bunch compression

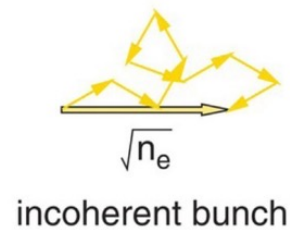
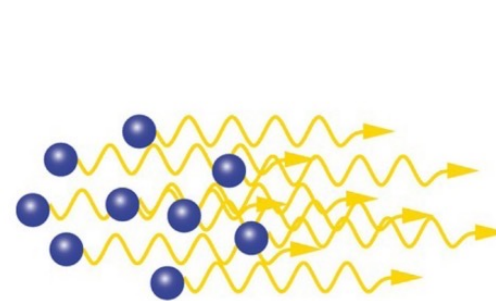
The SASE process

- Magnetic field of the emitted electromagnetic radiation always lies perpendicular to the electrons' trajectory -> no transfer of energy
- Difference in length between the path executed by the electron through a single undulator period and that of the electromagnetic field is a single wavelength
- Velocity along x (v_x) opposite to electric field -> deceleration the electron
- Velocity along x (v_x) in the same direction as electric field -> acceleration the electron
- This causes spatial redistribution of the electrons within the wavelength of the emitted radiation, causing microbunching

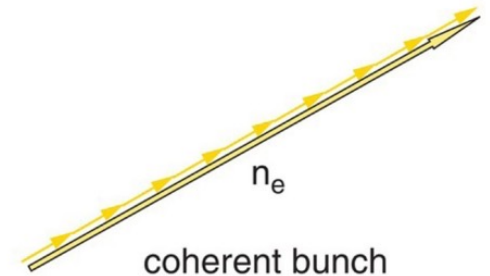
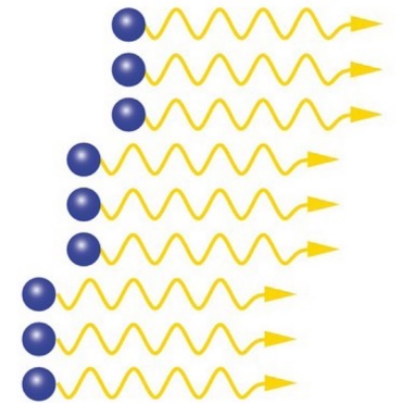


The SASE process

- If the positions of the electrons bear no correlation to the radiation they emit, the phases of that radiation are likewise random. Their summed amplitude is equal to the amplitude of the radiation emitted from one electron multiplied by the square root of the number of electrons in the bunch.
- If emission from all electrons are in phase. The amplitudes of the radiation from each of the n_e electrons add linearly, resulting in an increase in intensity by n_e^2 .



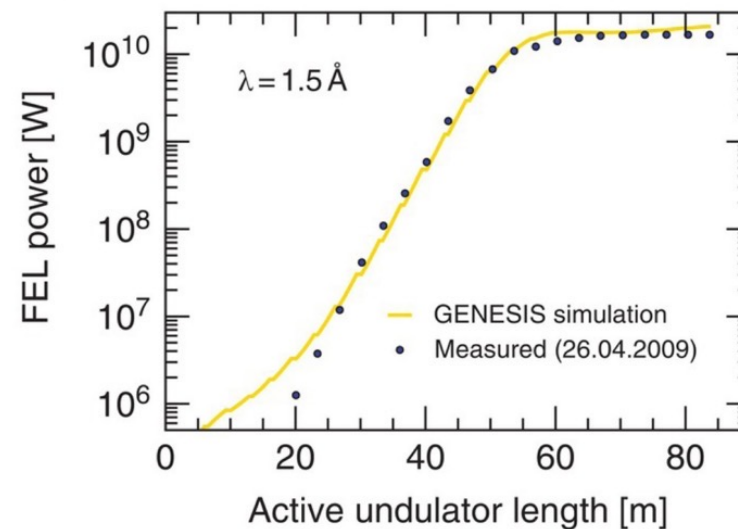
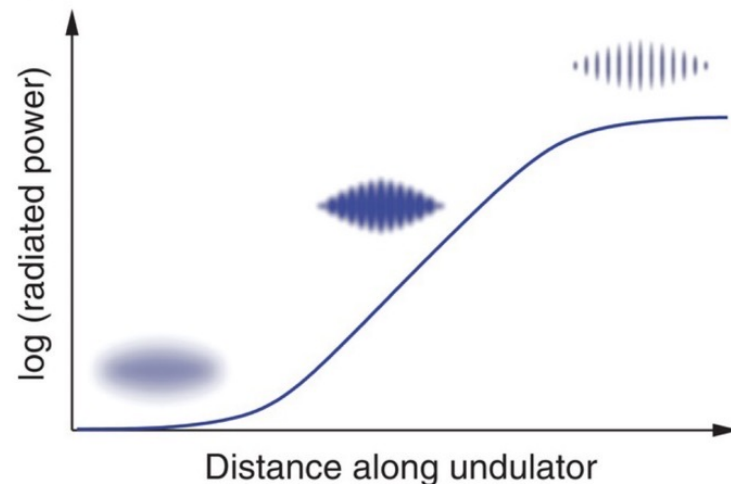
incoherent bunch



coherent bunch

The SASE process

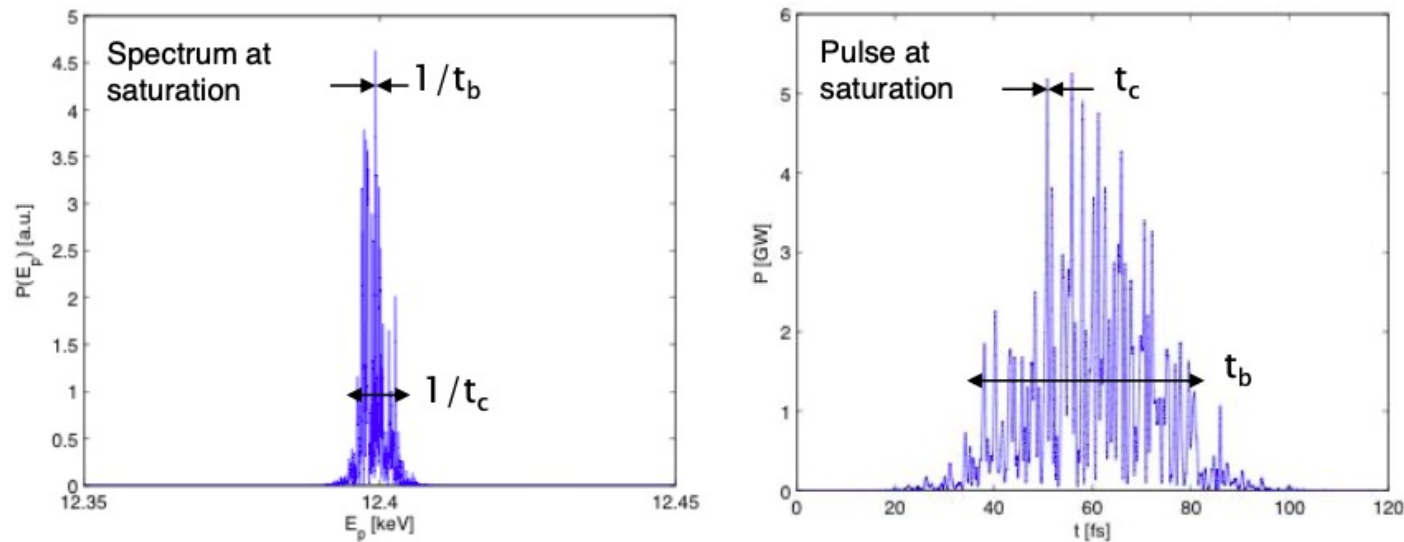
- The more the electrons bunch and radiate in phase (that is, coherently), the stronger is the electromagnetic-field interaction with them, which in turn further enhances the microbunching phenomenon associated with coherent emission, thus increasing the degree of coherent emission ...
- This results in a longitudinal density modulation (microbunching) together with an exponential growth of the radiation power along the undulator until saturation.



The SASE process

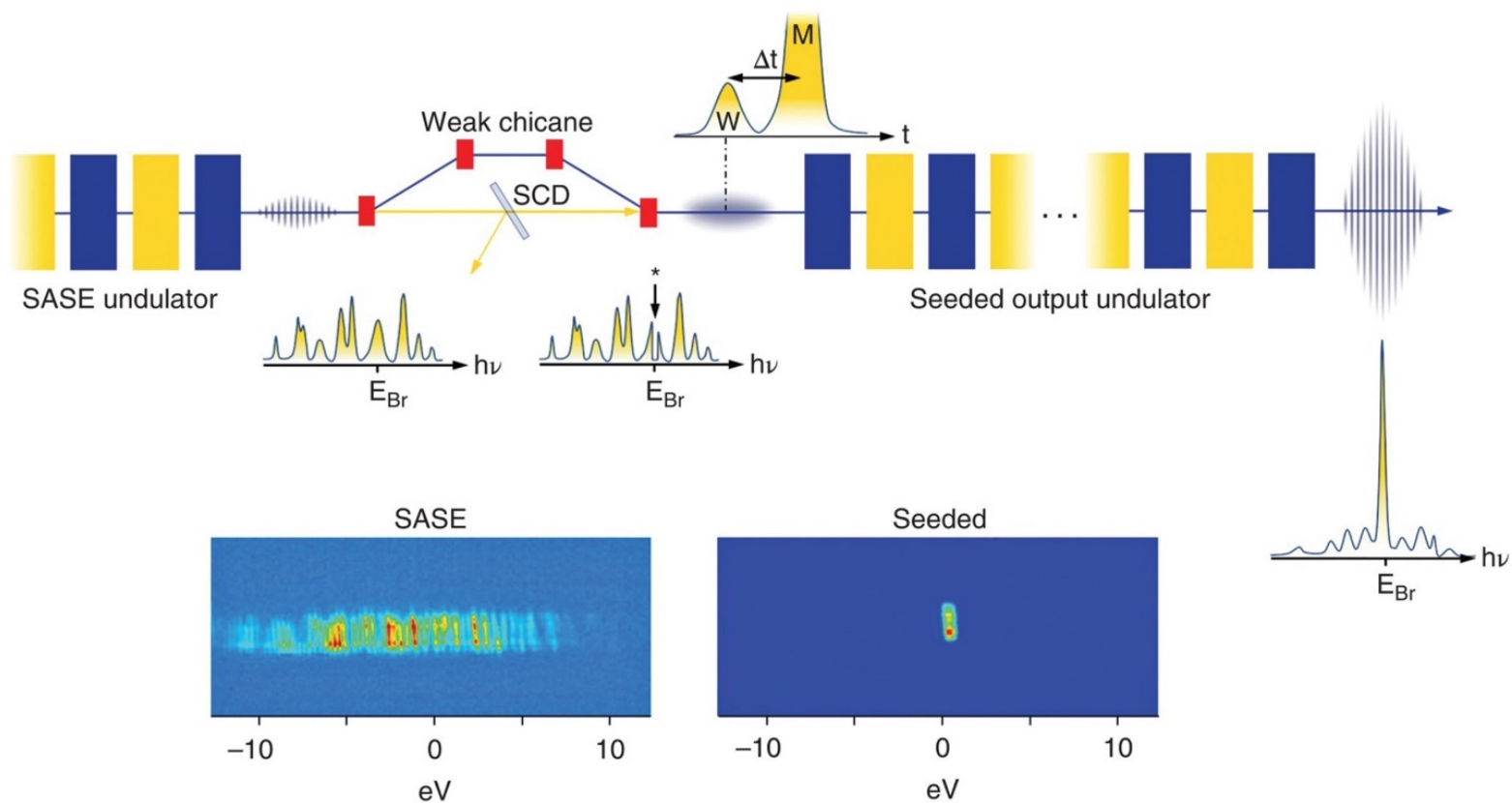
- The microbunches produced by the SASE process each contain approximately $n_e = 10^9$ electrons, resulting in a peak brilliance for XFELs approximately one billion times more intense than that from fourth-generation synchrotron sources.
- X-ray is fully transverse coherent, but not longitudinal

SwissFEL: Simulation for 1 Angstrom radiation



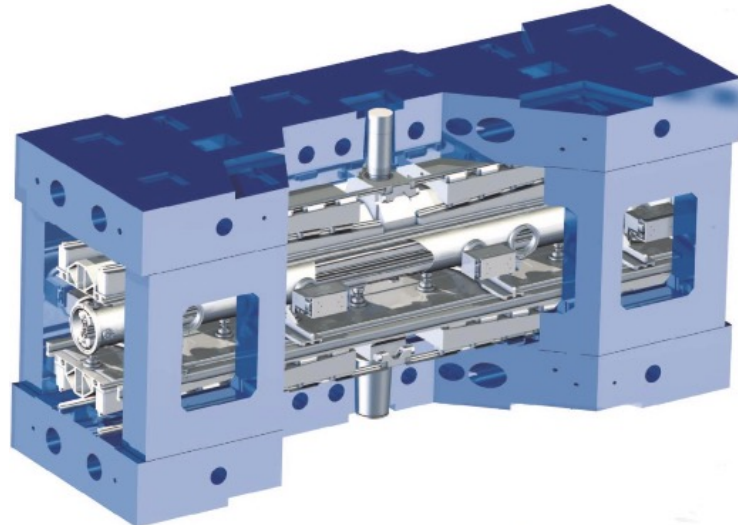
Beyond SASE

Longitudinal coherence can be improved by 'seeding'.





Aramis undulator



- In-vacuum, variable-gap undulators
- Undulator period 15 mm
- Nominal gap 4.5 mm for $K = 1.2$
- Array of 1060 permanent magnets (NdFeB with diffused Dy) per module
- 13 modules of 4 m length, total length 65 m



Coming next: beamlines and detectors

