

Heterophase Grain Boundaries

Lesson 11

MSE 304

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Reading for this Class

Key Topics in the Previous Class

Heterophase Grain Boundaries

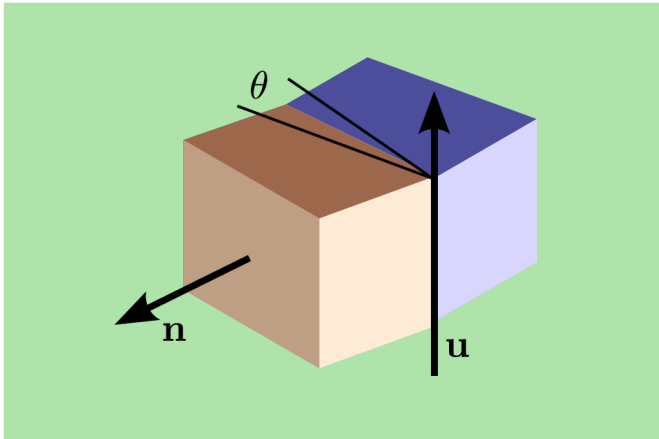
Coherent [5 -200 mJ/m²]

Semi-coherent [200-800 mJ/m²]

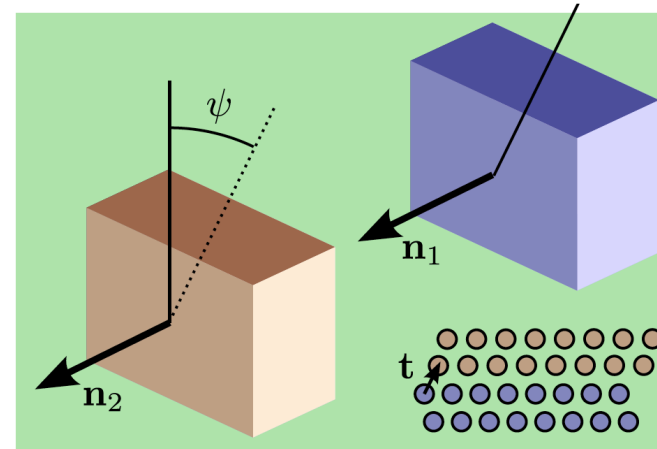
Incoherent [800-2500 mJ/m²]

Heterophase Grain Boundaries: Nomenclature

For simple interfaces
we have used this approach



For complex interfaces
we to use:



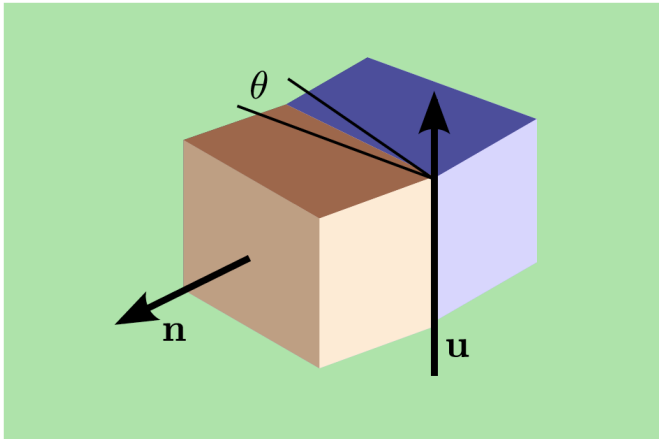
Planes (hkl)

Directions [uvw]

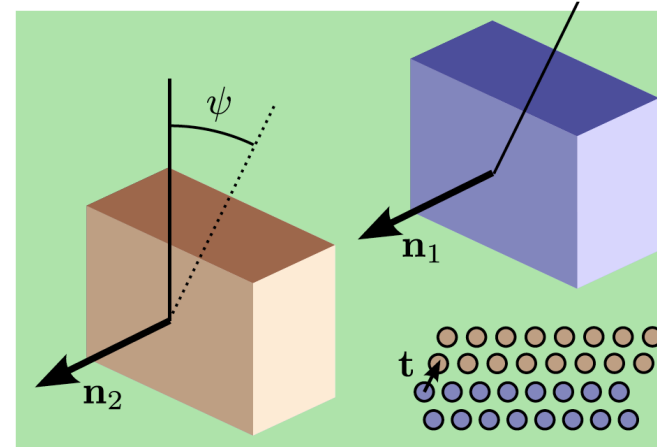
$$(hkl)_\alpha || (hkl)_\beta; [uvw]_\alpha || [uvw]_\beta$$

Heterophase Grain Boundaries: Nomenclature

For simple interfaces
we have used this approach



For complex interfaces
we to use:



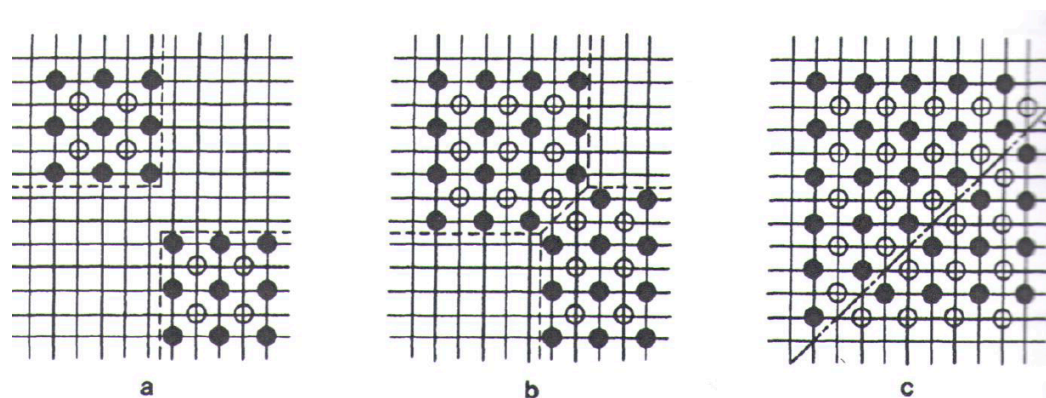
Planes (hkl)

Directions [uvw]

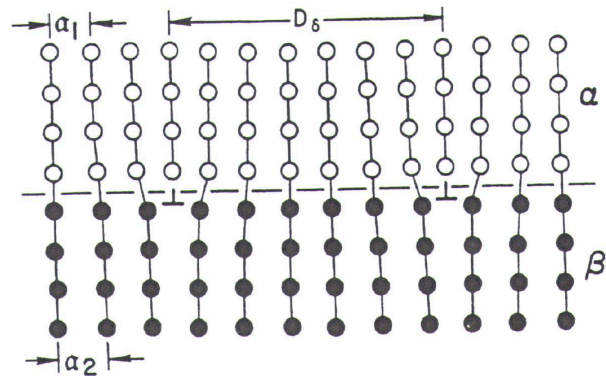
$$(hkl)_\alpha || (hkl)_\beta; [uvw]_\alpha || [uvw]_\beta$$

Special Grain Boundaries

Antiphase Domain Boundaries



Energetics of Semi-coherent Grain Boundaries



$$D_{\delta} = \frac{(a_1 + a_2)^2}{4(a_1 - a_2)}$$

$$\delta = \frac{2(a_1 - a_2)}{a_1 + a_2} = \frac{a_1 - a_2}{a'}$$

$$D_{\delta} = \frac{a'}{\delta} = \frac{b}{\delta}$$

Turnbull $\gamma_{sc}^{SS} = \gamma_c^{SS} + \gamma_s^{SS}$

$$\gamma_s^{SS} = \frac{\mu a'}{4\pi^2} \left\{ 1 + \Lambda - (1 + \Lambda^2)^{1/2} - \Lambda \ln \left[2\Lambda(1 + \Lambda^2)^{1/2} - 2\Lambda^2 \right] \right\}$$

$$\Lambda = 2\pi\delta \left(\frac{C_8}{\mu} \right)$$

$$\frac{1}{C_8} = \frac{1 - \nu_{\alpha}}{\mu_{\alpha}} + \frac{1 - \nu_{\beta}}{\mu_{\beta}}$$

Eterophase Interphases

$$(hkl)_\alpha || (hkl)_\beta; [uvw]_\alpha || [uvw]_\beta$$

$$\mathbf{A} = \begin{bmatrix} e_{11} & e_{12} \\ 0 & 1 \end{bmatrix}.$$

Eterointerphases

$$(hkl)_\alpha \parallel (hkl)_\beta; [uvw]_\alpha \parallel [uvw]_\beta$$

$$A = \begin{bmatrix} e_{11} & e_{12} \\ 0 & 1 \end{bmatrix}$$

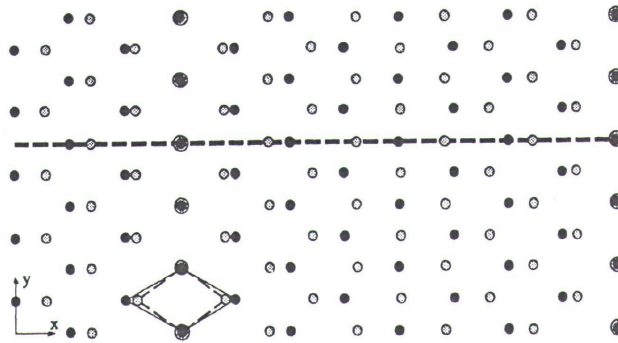


Figure 14.38. Schematic diagram representing the rhombic surface unit cells of a {111} f.c.c. plane (filled circles) and a {110} b.c.c. plane (open circles) when $\tau = a_{f.c.c.}/a_{b.c.c.} = 1.15$. From [143].

High Index Interfaces

$$(hkl)_\alpha \parallel (hkl)_\beta; [uvw]_\alpha \parallel [uvw]_\beta$$

$$\mathbf{A} = \begin{bmatrix} e_{11} & e_{12} \\ 0 & 1 \end{bmatrix}$$

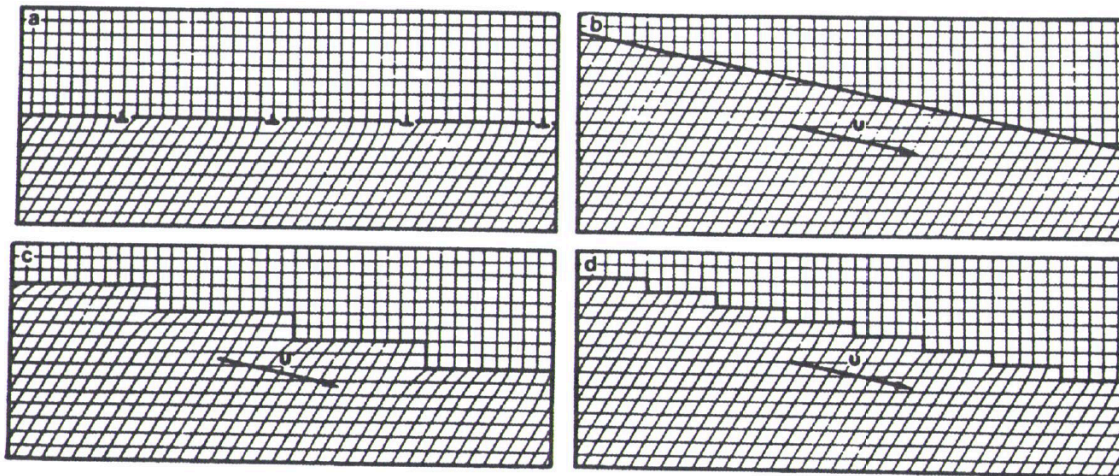


Figure 14.39. Possible interfaces for a transformation described by a shear e_{12} coupled with an expansion e_{11} in the shear plane. (a) The expansion is accommodated by a set of perfect lattice (misfit) dislocations. (b) The interface lies along an invariant line u and no dislocations are necessary. (c) The invariant line is resolved onto every other close-packed plane as in the f.c.c.–h.c.p. transformation where each ledge is an $\alpha/6\langle 112 \rangle$ partial dislocation. (d) The invariant line is resolved onto every close-packed plane as in an f.c.c.–b.c.c. transformation where each ledge contains an $\alpha/12\langle 112 \rangle$ partial dislocation. Reprinted with permission from [153] by Elsevier Science Ltd., Oxford, England.

High Index Interfaces

$(hkl)_\alpha || (hkl)_\beta; [uvw]_\alpha || [uvw]_\beta$

$$\mathbf{A} = \begin{bmatrix} e_{11} & e_{12} \\ 0 & 1 \end{bmatrix}.$$

Coherent Interfaces: The Cahn Hillard Model

Coherent Interfaces: The Cahn Hillard Model

$$f(c, \nabla c, \nabla^2 c, \dots) = f_0 + \sum_i L_i \nabla_i c + \sum_{ij} k_{ij}^{(1)} \frac{\partial^2 c}{\partial x_i \partial x_j} + \frac{1}{2} \sum_{ij} k_{ij}^{(2)} \frac{\partial c}{\partial x_i} \frac{\partial c}{\partial x_j} + \dots$$

$$L_i = \frac{\partial f}{\partial (\partial c / \partial x_i)}$$

$$k_{ij}^{(1)} = \frac{\partial f}{\partial \left(\frac{\partial^2 c}{\partial x_i \partial x_j} \right)}$$

$$k_{ij}^{(2)} = \frac{\partial^2 f}{\partial (\partial c / \partial x_i) \partial (\partial c / \partial x_j)}$$

Coherent Interfaces: The Cahn Hillard Model

$$f(c, \nabla c, \nabla^2 c, \dots) = f_0 + \sum_i L_i \nabla_i c + \sum_{ij} k_{ij}^{(1)} \frac{\partial^2 c}{\partial x_i \partial x_j} + \frac{1}{2} \sum_{ij} k_{ij}^{(2)} \frac{\partial c}{\partial x_i} \frac{\partial c}{\partial x_j} + \dots$$

$$L_i = \frac{\partial f}{\partial (\partial c / \partial x_i)}$$

$$L_i = 0$$

$$k_{ij}^{(1)} = \frac{\partial f}{\partial (\frac{\partial^2 c}{\partial x_i \partial x_j})}$$

$$k_{ij}^{(1)} = k^1 = \frac{\partial f}{\partial (\nabla^2 c)} \text{ for } i = j; \quad k_{ij}^{(1)} = 0 \text{ otherwise}$$

$$k_{ij}^{(2)} = \frac{\partial^2 f}{\partial (\partial c / \partial x_i) \partial (\partial c / \partial x_j)}$$

$$k_{ij}^{(2)} = k^2 = \frac{\partial^2 f}{\partial (|\nabla c|^2)^{1/2}} \text{ for } i = j; \quad k_{ij}^{(2)} = 0 \text{ otherwise}$$

$$f(c, \nabla c, \nabla^2 c, \dots) = f_0 + \sum_i k^1 \frac{\partial^2 c}{\partial^2 x_i} + \sum_i k^2 \left(\frac{\partial c}{\partial x_i} \right)^2 + \dots = f_0 + k^1 \nabla^2 c + k^2 (\nabla c)^2 + \dots$$

Coherent Interfaces: The Cahn Hillard Model

$$F = \int \mu dN = N_V \int_V f dV = N_V \int_V (f_0 + k^1 \nabla^2 c + k^2 (\nabla c)^2) dV$$

Divergence Theorem:

$$\int_V k^1 \nabla^2 c dV = - \int_V \frac{dk^1}{dc} (\nabla c)^2 dV + \int_S k^1 \nabla c \cdot n dS$$

$$F = N_V \int_V (f_0 + k^1 \nabla^2 c + k^2 (\nabla c)^2) dV = N_V \int_V \left(f_0 - \frac{dk^1}{dc} (\nabla c)^2 + k^2 (\nabla c)^2 \right) dV = N_V \int_V (f_0 + k (\nabla c)^2) dV$$

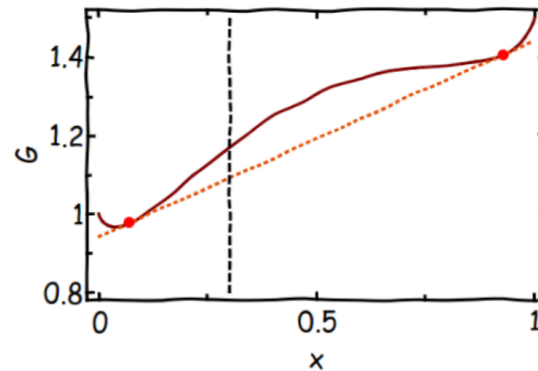
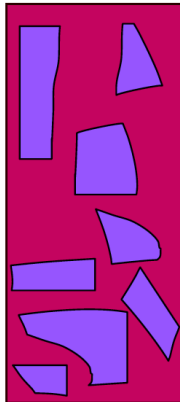
$$k = k^2 - \frac{dk^1}{dc} = \frac{\partial^2 f}{\partial (|\nabla c|^2)^{1/2}} - \frac{d}{dc} \left(\frac{\partial f}{\partial (\nabla^2 c)} \right)$$

Coherent Interfaces: The Cahn Hillard Model

- Take x_A moles of compound A, and $x_B = 1 - x_A$ moles of compound B, and **physically** mix them
- Assume only the entropy of solution matters (**ideal solution**)

$$\mu_{A,B} = \mu_{A,B}^{\ominus} + RT \ln x_{A,B}$$

- Add an enthalpy of mixing term $\Omega x_A x_B$ (**regular solution**)



$$\begin{array}{ll} \mu_A^{\ominus} = 1 & \mu_B^{\ominus} = 1.5 \\ RT = 1 & \Omega = 2.5 \end{array}$$

$$G = x_A (\mu_A^{\ominus} + RT \ln x_A) + (1 - x_A) (\mu_B^{\ominus} + RT \ln (1 - x_A)) + \Omega x_A (1 - x_A)$$

Coherent Interfaces: The Cahn Hillard Model

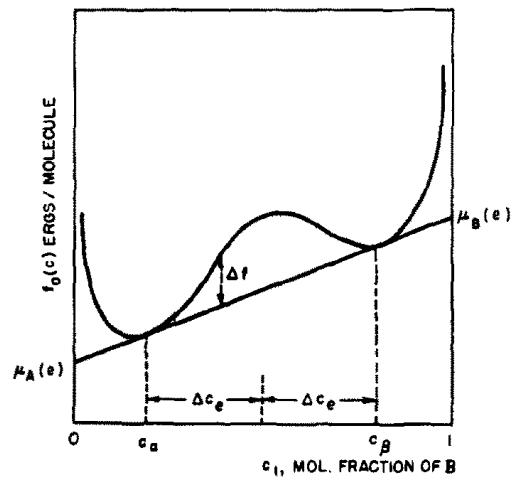


FIG. 1. $f_0(c)$ for $T < T_c$.

$$F = N_V \int_V (f_0 + k(\nabla c)^2) dV = N_V A \int_{-\infty}^{+\infty} \left(f_0 + k \left(\frac{dc}{dx} \right)^2 \right) dx$$

$$\gamma = N_V \int_{-\infty}^{+\infty} \left[f_0 + k \left(\frac{dc}{dx} \right)^2 - c\mu_\beta(e) - (1-c)\mu_\alpha(e) \right] dx$$

$$\gamma = N_V \int_{-\infty}^{+\infty} \left[\Delta f(c) + k \left(\frac{dc}{dx} \right)^2 \right] dx$$

Coherent Interfaces: The Cahn Hillard Model

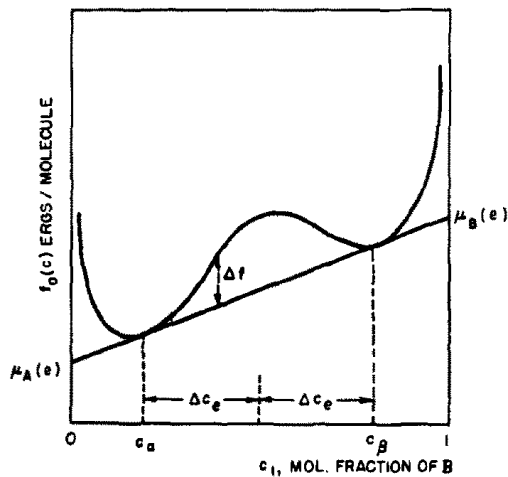


FIG. 1. $f_0(c)$ for $T < T_c$.

Euler equation:

$$\gamma = N_V \int_{-\infty}^{+\infty} \left[\Delta f(c) + k \left(\frac{dc}{dx} \right)^2 \right] dx$$

$$I - \left(\frac{dc}{dx} \right) \left[\frac{\partial I}{\partial \left(\frac{dc}{dx} \right)} \right] = 0$$

$$\Delta f(c) - k \left(\frac{dc}{dx} \right)^2 = 0$$

$$\Delta f(c) = k \left(\frac{dc}{dx} \right)^2$$

$$dx = \left(\frac{k}{\Delta f(c)} \right)^{1/2} dc$$

$$\gamma = 2N_V \int_{-\infty}^{+\infty} \Delta f(c) dx$$

$$\gamma = 2N_V \int_{c_\alpha}^{c_\beta} (k\Delta f(c))^{1/2} dc$$

Coherent Interfaces: The Cahn Hillard Model

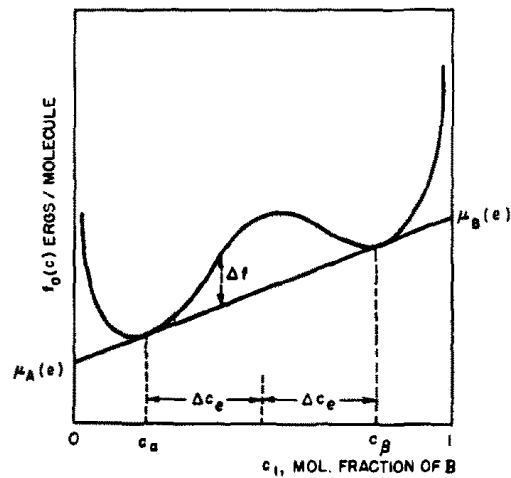


FIG. 1. $f_0(c)$ for $T < T_c$.

$$\gamma = 2N_V \int_{c_\alpha}^{c_\beta} (k\Delta f(c))^{1/2} dc$$

$$dx = \left(\frac{k}{\Delta f(c)} \right)^{1/2} dc$$

$$\frac{dc}{dx} = \left(\frac{k}{\Delta f(c)} \right)^{1/2}$$

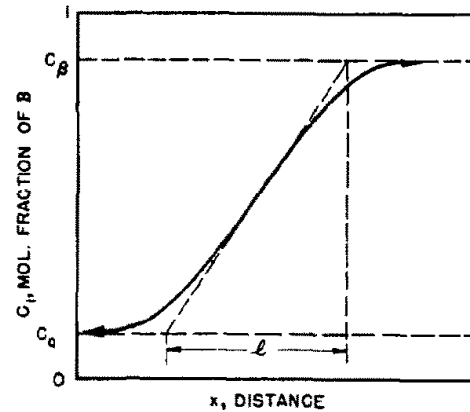


FIG. 2. Interface profile.

Coherent Interfaces: The Cahn Hillard Model

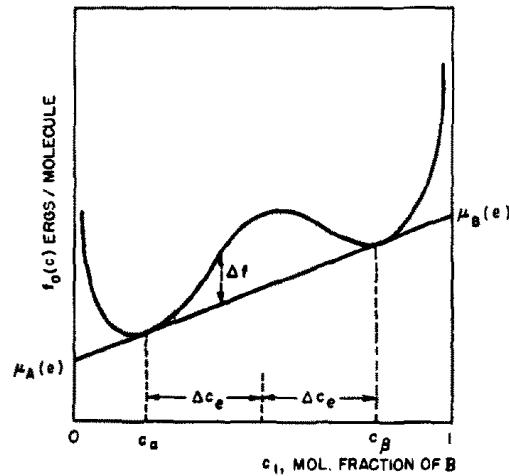


FIG. 1. $f_0(c)$ for $T < T_c$.

$$\gamma = 2N_V \int_{c_\alpha}^{c_\beta} (k\Delta f(c))^{1/2} dc$$

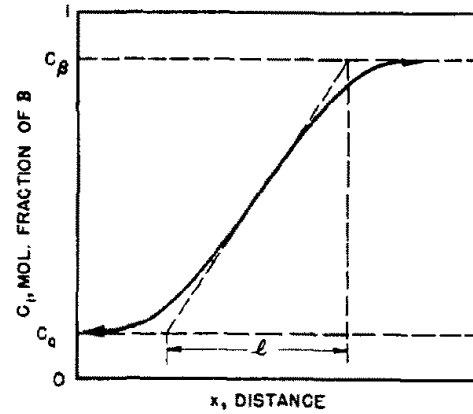
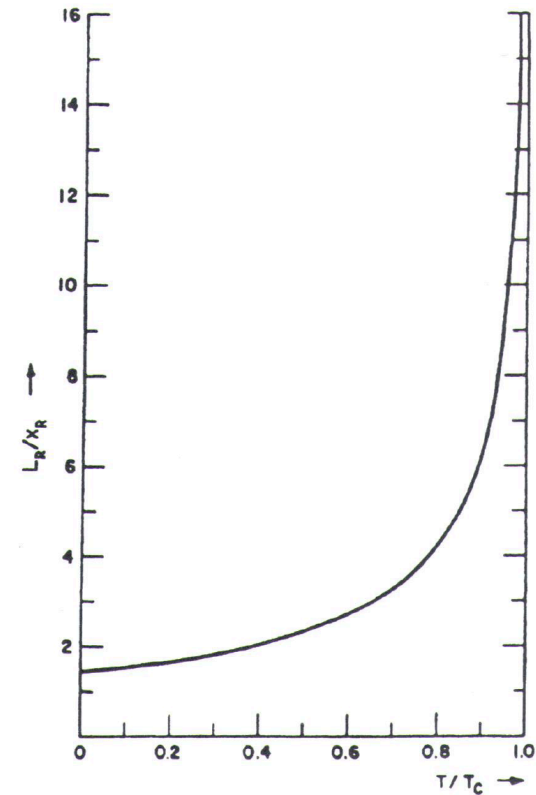
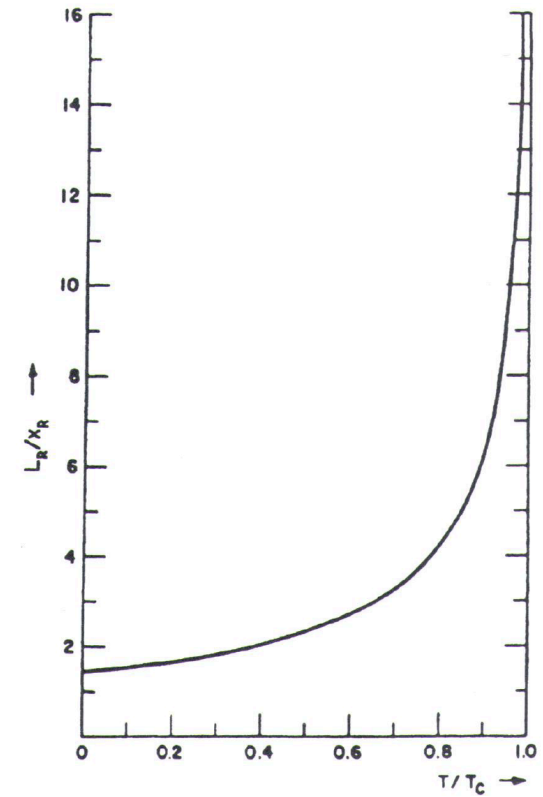


FIG. 2. Interface profile.

$$\frac{dc}{dx} = \left(\frac{k}{\Delta f(c)} \right)^{1/2}$$



Coherent Interfaces: The Cahn Hillard Model



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Conclusions
