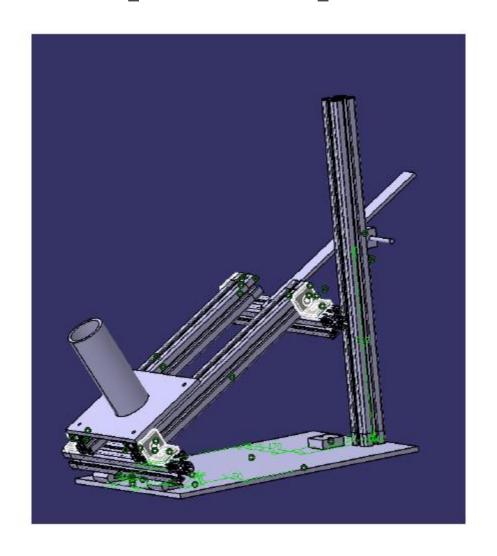
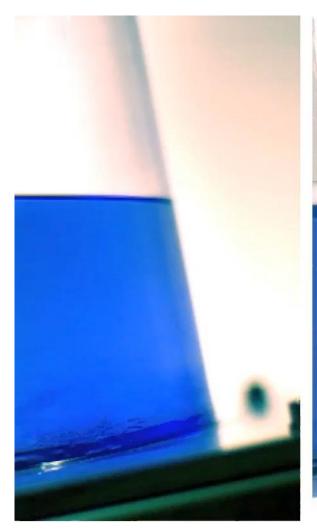
## Impulse response

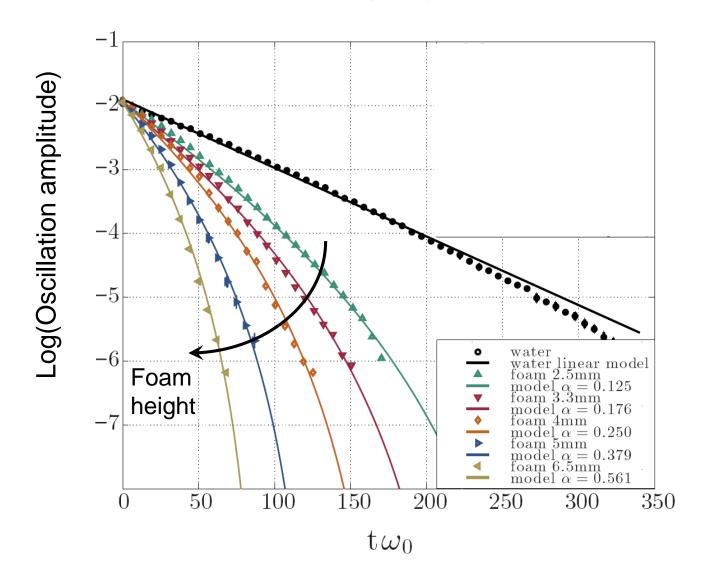








see also Sauret et al. 2015



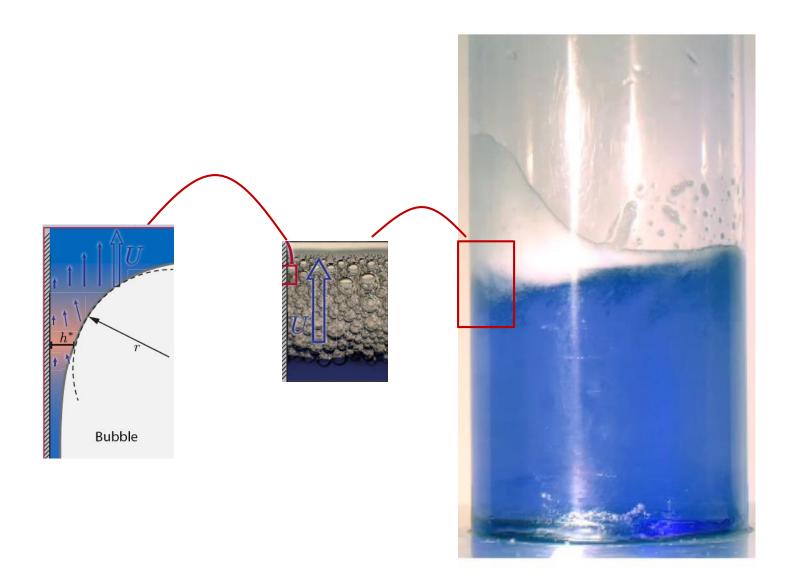
There must be something nonlinear!

# Damping by foam to exacerbate nonlinear effects









### A nonlinearly damped oscillator

$$\pi R^2(\rho \lambda + \rho_f h)\ddot{X} + F_g = F_w + F_b + F_s + F_f$$

Wall friction

$$F_w \sim \dot{X}$$

Bulk dissipation

$$F_b \sim X$$

- Free surface dissipation  $F_s \sim \dot{X}$
- Contact line friction  $F_f \sim \dot{X}|\dot{X}|^{-1/3}$

$$\ddot{X} + \omega_0^2 X = -2\sigma\omega_0 \dot{X} - \alpha \dot{X} |\dot{X}|^{-1/3}$$

$$\ddot{x} + \omega_0^2 x = -\epsilon \dot{x} |\dot{x}|^{-1/3}$$

Asymptotic expansion:

$$x = x_0 + \epsilon x_1 + \epsilon^2 x_2 + \dots \qquad t = t + \epsilon t_1$$

 $\epsilon^0$ 

$$\ddot{x}_0 + \omega_0^2 x_0 = 0$$
$$x_0 = A(t_1)e^{i\omega_0 t} + \overline{A}(t_1)e^{-i\omega_0 t}$$

 $\epsilon^1$ 

$$\ddot{x}_1 + \omega_0^2 x_1 = -2 \frac{\partial^2 x_0}{\partial t \partial t_1} - \left( \frac{\partial x_0}{\partial t} \right) \left| \frac{\partial x_0}{\partial t} \right|^{-1/3}$$

Let examine the terms on the r.h.s: The first as before:

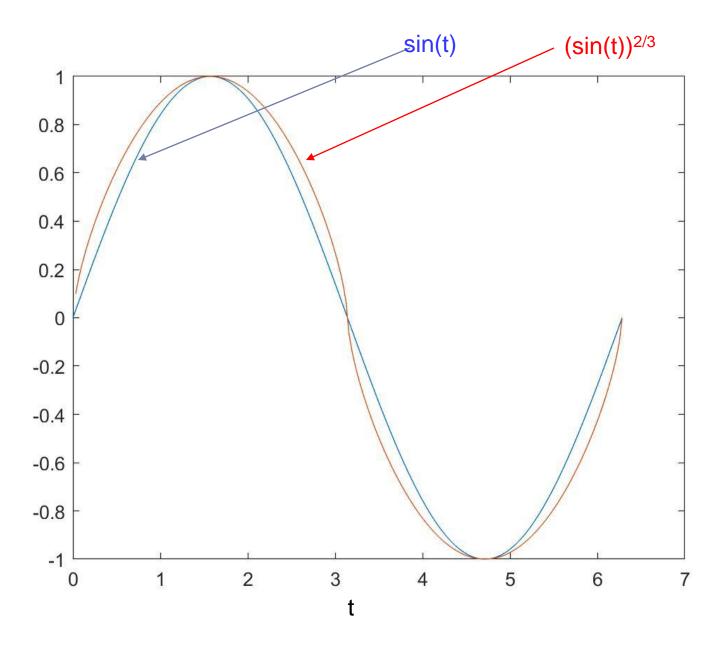
$$2\frac{\partial^2 x_0}{\partial t \partial t_1} = 2i\omega_0 A' e^{i\omega_0 t} + c.c. = 2i\omega_0 (a' + i\beta' a) e^{i\beta} e^{i\omega_0 t} + c.c.$$

The second is non-linear, by using  $A = ae^{i\beta}$ :

$$\dot{x_0}|\dot{x_0}|^{-1/3} = (i\omega_0 A e^{i\omega_0 t} - i\omega_0 \overline{A} e^{-i\omega_0 t})|i\omega_0 A e^{i\omega_0 t} - i\omega_0 \overline{A} e^{-i\omega_0 t}|^{-1/3} = (i\omega_0 a e^{i\omega_0 t + i\beta} - i\omega_0 a e^{-i\omega_0 t - i\beta})|i\omega_0 a e^{i\omega_0 t + i\beta} - i\omega_0 a e^{-i\omega_0 t - i\beta}|^{-1/3} = \omega_0^{2/3} a^{2/3} \underbrace{(i(e^{i\phi} - e^{-i\phi}))|i(e^{i\phi} - e^{-i\phi})|^{-1/3}}_{f(\phi)}$$

where  $\phi = \omega_0 t + \beta$ , the last term can be transformed with Fourier series

$$f(\phi) = (-2\sin(\phi))|-2\sin(\phi)|^{-1/3} =$$



$$f(\phi) = (-2\sin(\phi))|-2\sin(\phi)|^{-1/3} = \sum_{-\infty}^{\infty} c_n e^{in\phi} = c_1 e^{i\phi} + c_{-1} e^{-i\phi} + NST$$

where  $c_{\pm 1}$  are the Fourier series coefficient given by  $c_{\pm 1} = \frac{1}{2\pi} \int_0^{2\pi} f(\phi) e^{\mp i\phi} d\phi$ , with  $c_{-1} = \overline{c}_1$ .  $c_1 \approx 0.85i = c_{1i}i$ .

$$|\dot{x_0}|\dot{x_0}|^{-1/3} = \omega_0^{2/3} a^{2/3} e^{i\beta} c_{1i} i e^{i\omega_0 t} + c.c. + NST$$

The amplitude equation is obtained:

$$2i\omega_0(a'+ia\beta')e^{i\beta}e^{i\omega_0t} + \omega_0^{2/3}a^{2/3}e^{i\beta}c_{1i}ie^{i\omega_0t} + c.c = 0$$

By splitting real and imaginary parts:

$$\begin{cases} 2\omega_0 a' + c_{1i}\omega_0^{2/3}a^{2/3} = 0 \Rightarrow \int_{a_0}^a a^{-2/3} da = -\frac{c_{1i}}{2\omega_0^{2/3}} \int_0^{t_1} dt_1 \Rightarrow a(t_1) = (-\frac{c_{1i}t_1}{6\omega_0^{1/3}} + a_0^{1/3})^3 \\ 2\omega_0 a\beta' = 0 \Rightarrow \beta = \beta_0 \end{cases}$$

With the finite time singularity:

$$t_{inv}^* = \frac{6\omega_0^{1/3}}{\epsilon c_{1i}} a_0^{1/3}$$

