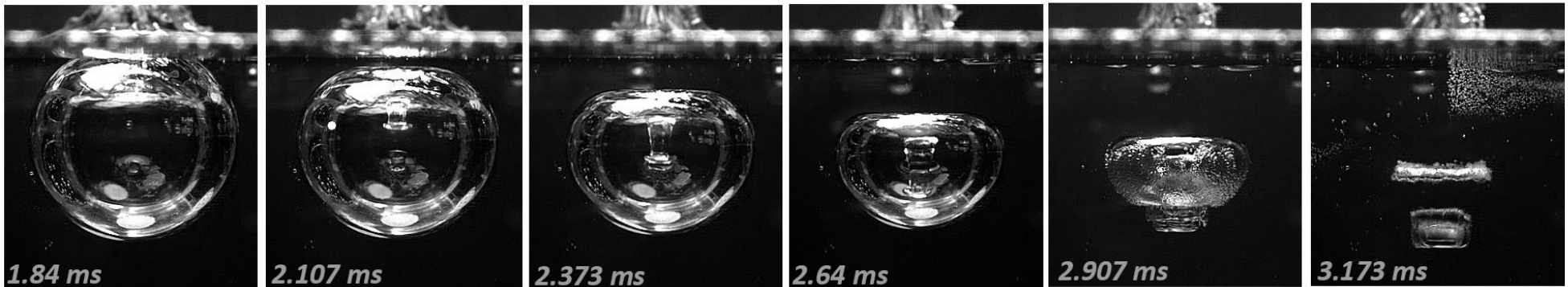


ECOLE POLYTECHNIQUE FEDERALE DE LAUSANNE
SECTION DE GENIE MECANIQUE
6th & 8th Semester, Fall 2024

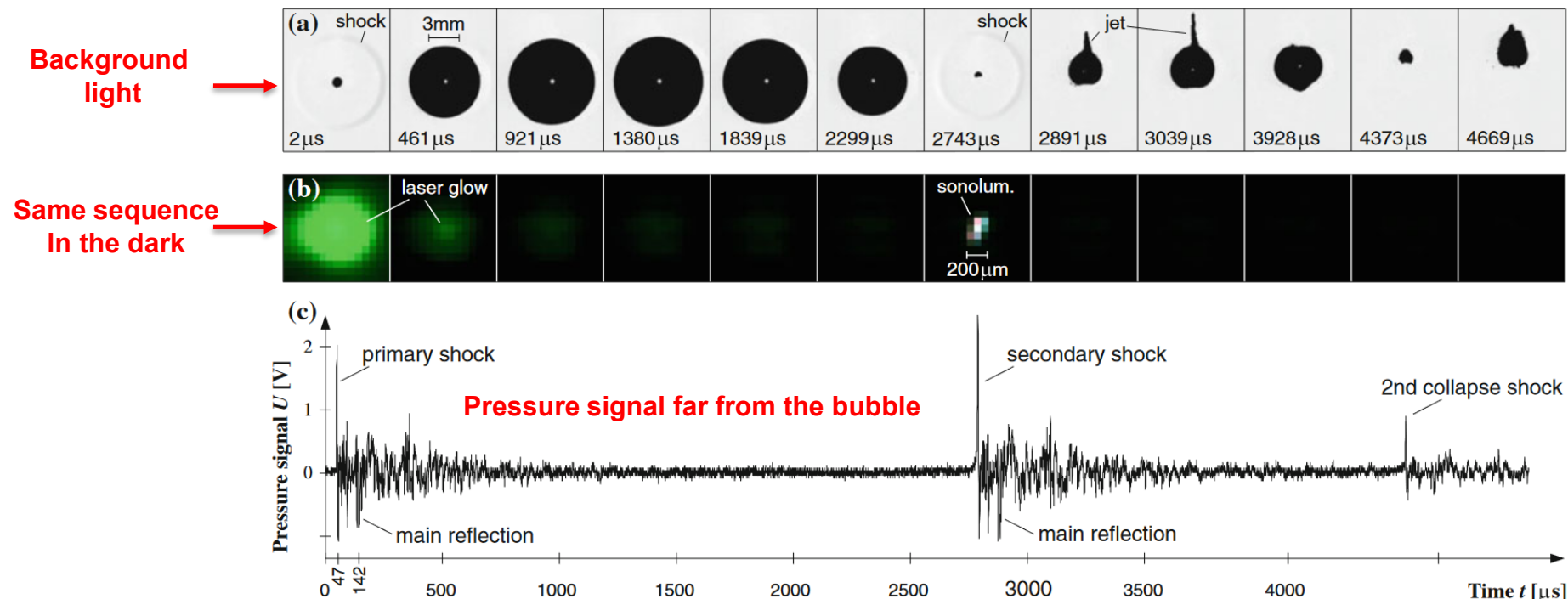
CAVITATION AND INTERFACE PHENOMENA
Chapter 3 : Dynamics of non-spherical cavitation bubbles
3.2 : Micro-Jetting



Dr Mohamed FARHAT Assistants: Th. Berger, A. Sache.
EPFL – Cavitation Research Group
Avenue de Cour 33 bis, 1007 Lausanne

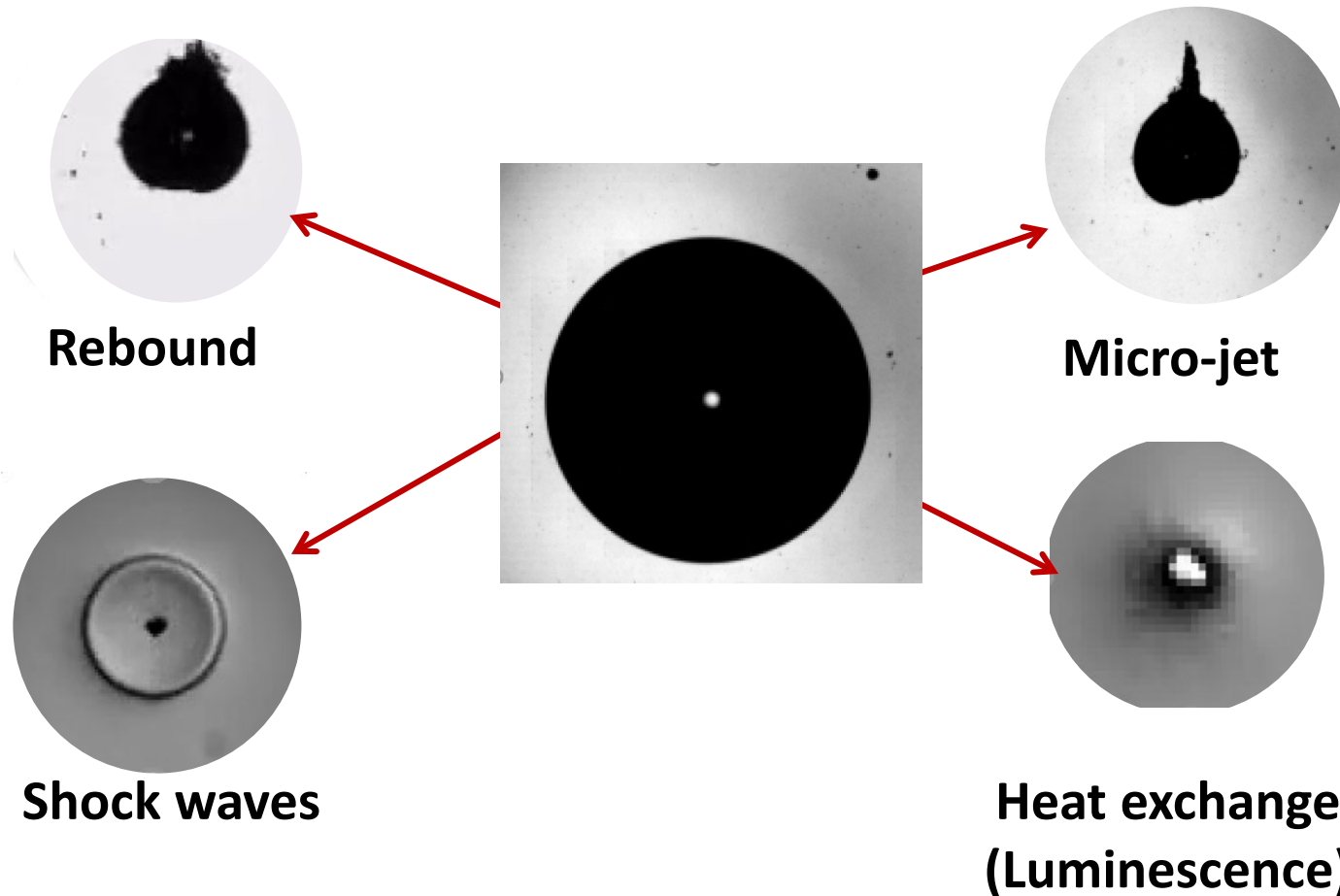
Cavitation Bubble Dynamics

- Key phenomena associated with a laser-induced cavitation bubble:
 - Plasma → Explosive growth of the bubble with emission of a shockwave
 - Deceleration of the expanding interface, up to $R = R_{max}$ ($\dot{R}(R_{max}) = 0$)
 - Bubble collapses violently, leading to:
 1. Shockwave emission (compressibility effects)
 2. Micro-jetting, due to bubble deformation
 3. Bubble rebound (due to the non-condensable gas)
 4. Light (luminescence), due to extreme heating of the non-condensable gas



Cavitation Bubble Dynamics

Energy budget: Where does the energy of a collapsing bubble go ?

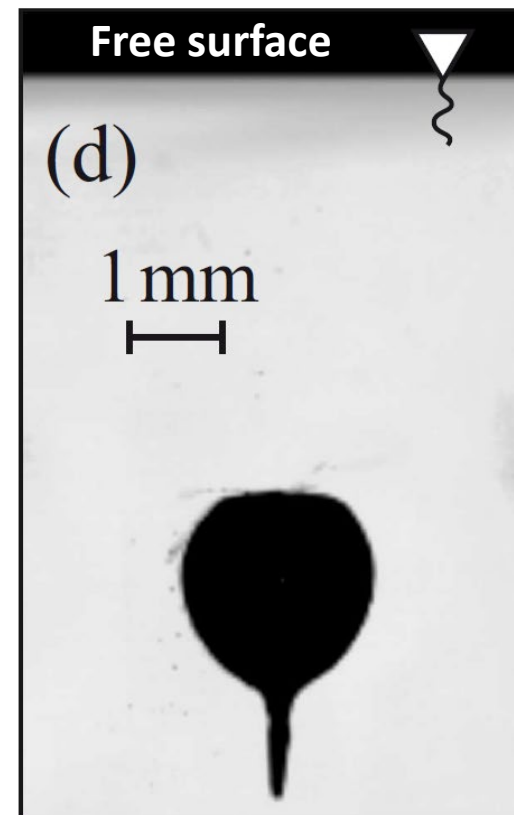
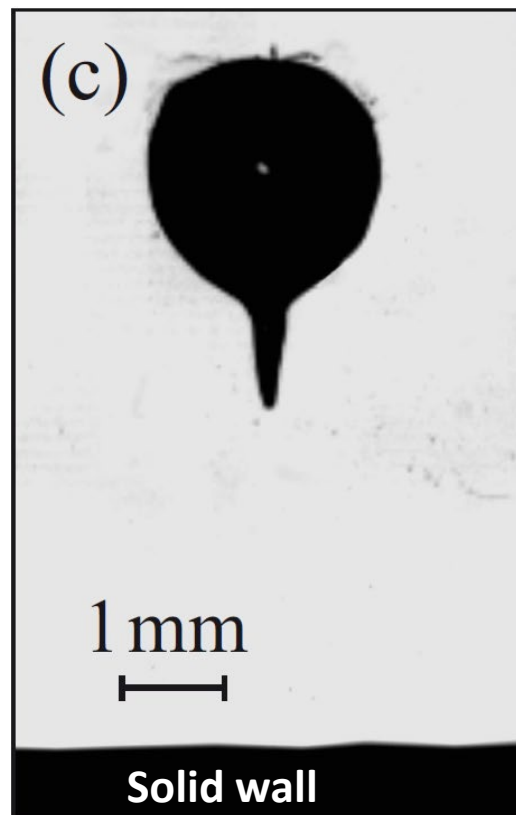


- Microjet
- Shock waves
- Luminescence
- Rebound

Cavitation Bubble Dynamics

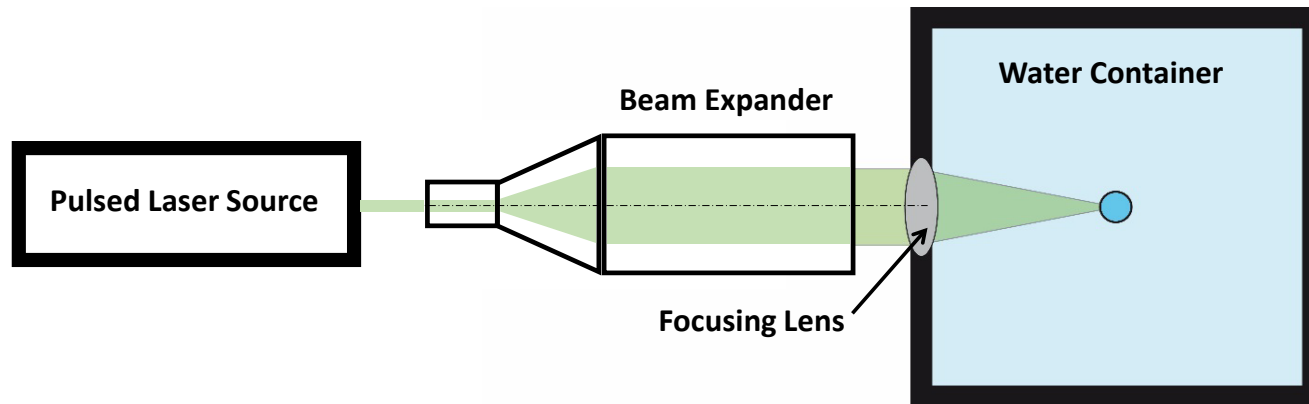
Microjets generated by collapsing bubbles

- Case of solid and free surfaces
 - Micro-jet due to pressure anisotropy
 - Well visible during the bubble rebound (conical shell of vapor)
 - Jet characteristics (size, speed, ...) $\leftrightarrow$ pressure gradient ?

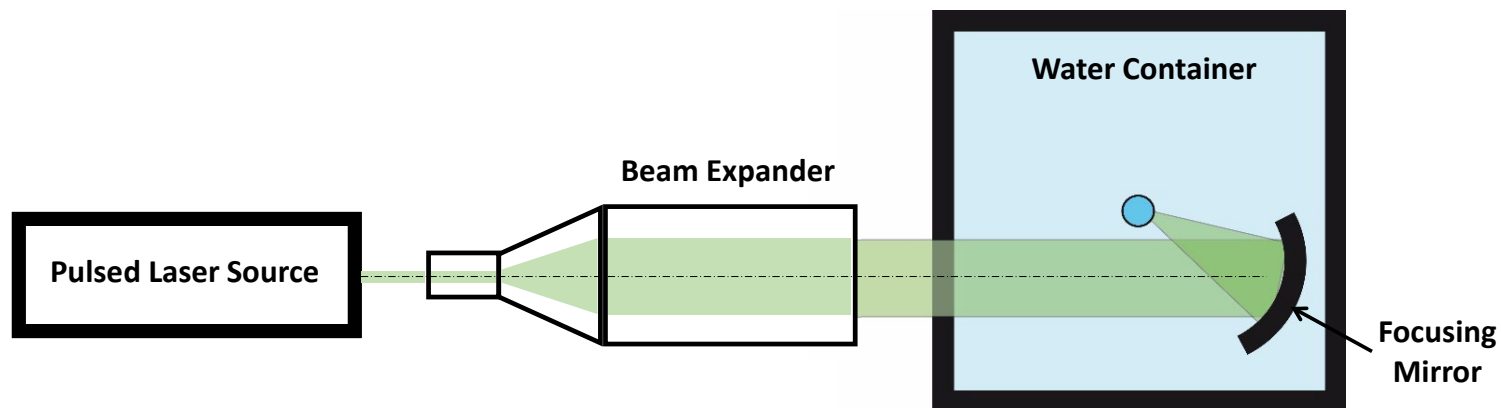


Chasing the most spherical bubble using Pulsed Focused Laser Technique

- Classical Technique:

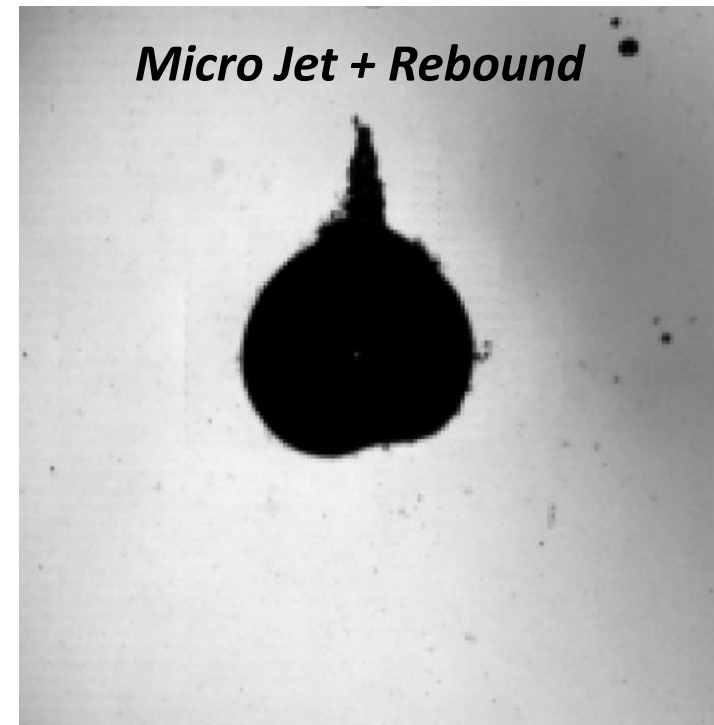
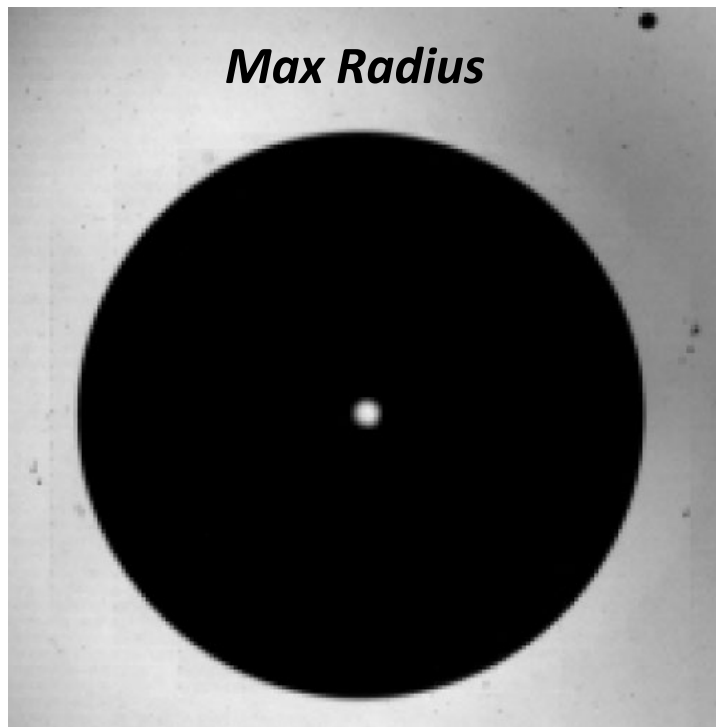


- Improved Technique: Immersed parabolic mirror
→ less aberrations → better focusing → enhanced sphericity of the bubble



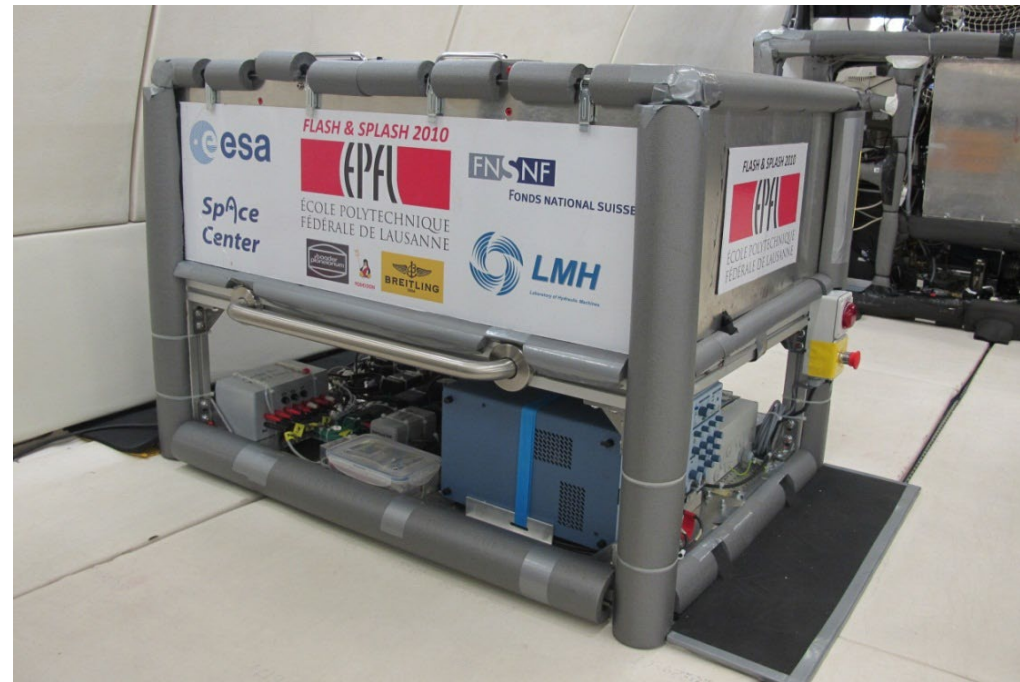
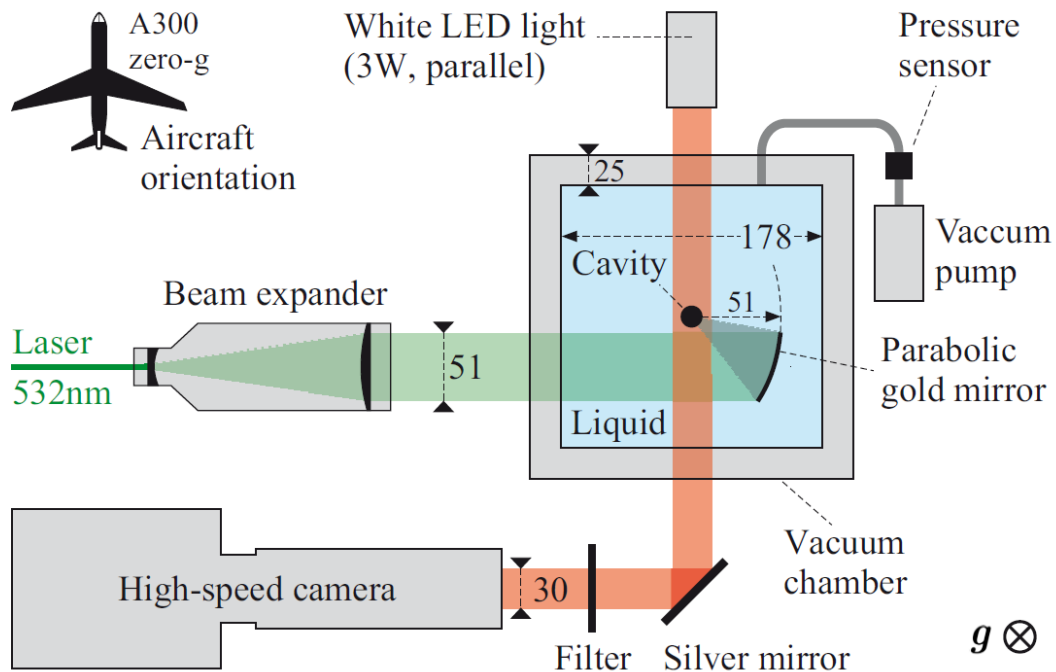
Chasing the most spherical bubble using Pulsed Focused Laser Technique

- Despite all the care given to the experimental setup, we always observe a deformation of the bubble when its maximum radius is in the millimeter range or above (under atmospheric pressure).
- This is due to the pressure anisotropy caused by the gravity field
 - *The bottom of the bubble “feels” more pressure than its top*



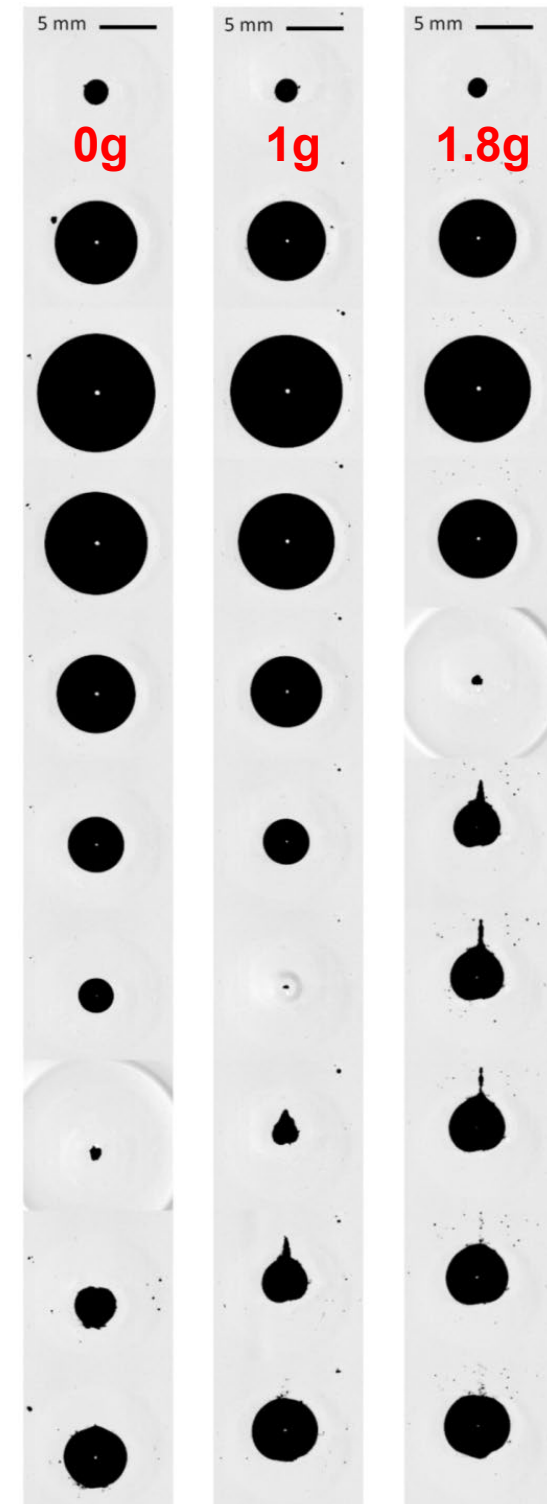
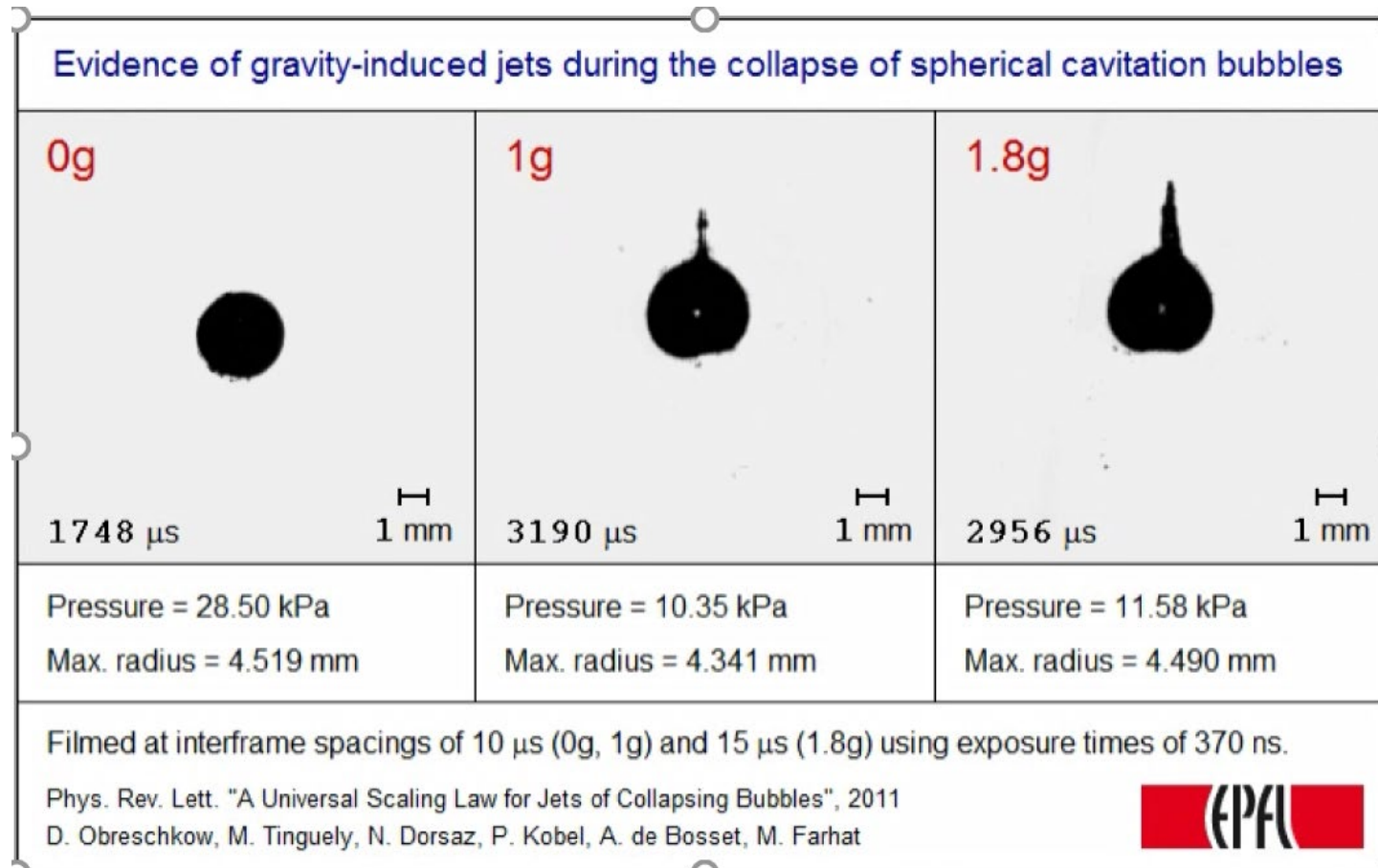
Effects of gravity induced pressure gradient

- Zero gravity tests: Experimental setup
 - Laser induced bubble in a closed water volume
 - Liquid pressure : varied from 8 to 80 Kpa
 - Point-like Plasma with a diameter < 0.1 mm
 - Pressure gradient: 0 to 18 KPa/m (gravity: 0, 1, 1.2, 1.4, 1.6, 1.8 g)



Effects of gravity induced pressure gradient

- Evidence of gravity induced jets (for the 1st time)
 - Upward jet due to gravity induced pressure gradient
 - More pronounced in hypergravity (1.8 g)
 - No jet is formed in microgravity



Scaling laws for collapse induced jets

- **Definition: Dimensionless pressure anisotropy parameter (vector):**

$$\vec{\zeta} = -\frac{\vec{\nabla}p R_{max}}{\Delta p}$$

Where $\vec{\nabla}p$ is the pressure gradient “felt” by the bubble and Δp is the pressure difference between outside and inside the bubble ($p_{\infty} - p_{bubble}$)

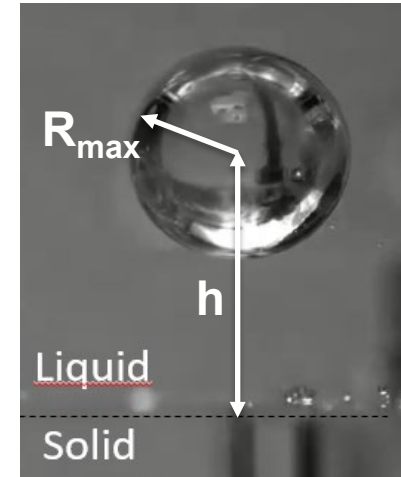
- **Pressure anisotropy induced by gravity:**

$$\zeta = |\vec{\zeta}| = -\frac{|\vec{\nabla}p| R_{max}}{\Delta p} - \frac{\rho g R_{max}}{\Delta p}$$

¹for more details, see Supponen et al. (2016) Journal of Fluid Mechanics

Scaling laws for collapse induced jets

- Quantifying bubble asymmetry:
 - How to characterize the pressure anisotropy in other configurations, such as neighboring rigid or free surfaces ?
 - Kelvin Impulse theory (out of the scope of the present course¹)*, allows to relate the anisotropy parameter to the stand-off parameter $\gamma \left(= \frac{h}{R_{max}} \right)$



The Kelvin Impulse is the momentum gained by the liquid during bubble growth and collapse in a constant pressure gradient:

$$I = \nabla p \int_{-T_c}^{T_c} V dT$$

$$I = \frac{5\pi}{3\sqrt{6}} B \left(\frac{5}{6}, \frac{2}{2} \right) R_0^3 \sqrt{\Delta p \rho} \zeta \approx 4.789 R_0^3 \sqrt{\Delta p \rho} \zeta$$

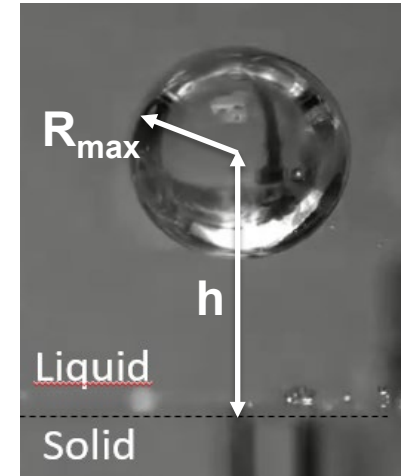
Where V is the volume of the bubble and B is the beta function, defined by:

$$B(x, y) = \int_0^1 t^{x-1} (1-t)t^{y-1} dt = \frac{\Gamma(x)\Gamma(y)}{\Gamma(x+y)}$$

¹for more details, see Supponen et al. (2016) *Journal of Fluid Mechanics*

Scaling laws for collapse induced jets

- Quantifying bubble asymmetry:
 - How to characterize the pressure anisotropy in other configurations, such as neighboring rigid or free surfaces ?
 - Kelvin Impulse theory (out of the scope of the present course¹)*, allows to relate the anisotropy parameter to the stand-off parameter $\gamma \left(= \frac{h}{R_{max}} \right)$



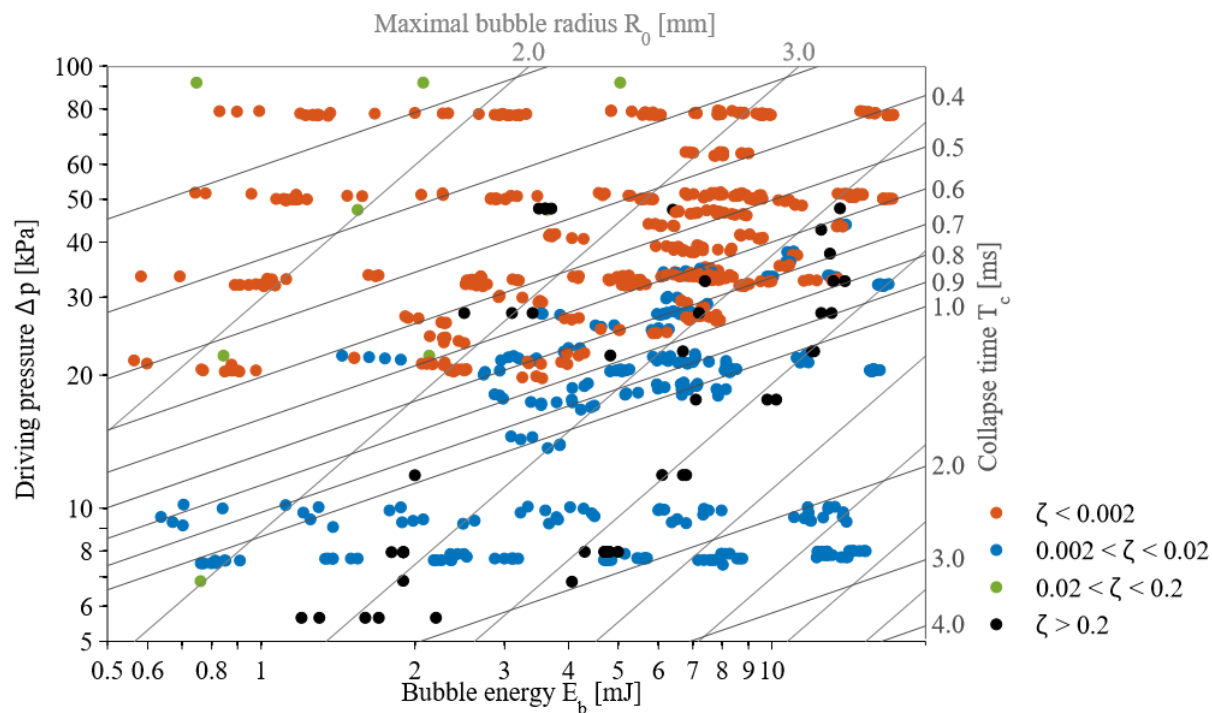
The Kelvin Impulse is the momentum gained by the liquid during bubble growth and collapse in a constant pressure gradient:

$$\zeta = \begin{cases} \begin{array}{l} \text{driving pressure} \\ -\rho g R_0 \Delta p^{-1} \\ -0.195 \gamma^{-2} \mathbf{n} \\ +0.195 \gamma^{-2} \mathbf{n} \end{array} & \begin{array}{l} \text{normal unit vector} \\ \text{on surface} \end{array} & \begin{array}{l} \text{gravitational field} \\ \text{flat rigid surface} \\ \text{flat free surface} \end{array} \\ \begin{array}{l} -\rho (\mathbf{u} \cdot \nabla) \mathbf{u} R_0 \Delta p^{-1} \\ 0.195 \gamma^{-2} (\rho_1 - \rho_2) (\rho_1 + \rho_2)^{-1} \mathbf{n} \\ 0.195 \gamma^{-2} (4\alpha - 1 - 8\alpha^2 e^{2\alpha} E_1(2\alpha)) \mathbf{n} \end{array} & & \begin{array}{l} \text{stationary potential flow} \\ \text{liquid interface} \\ \text{inertial boundary} \end{array} \end{cases} \quad \zeta = |\vec{\zeta}|$$

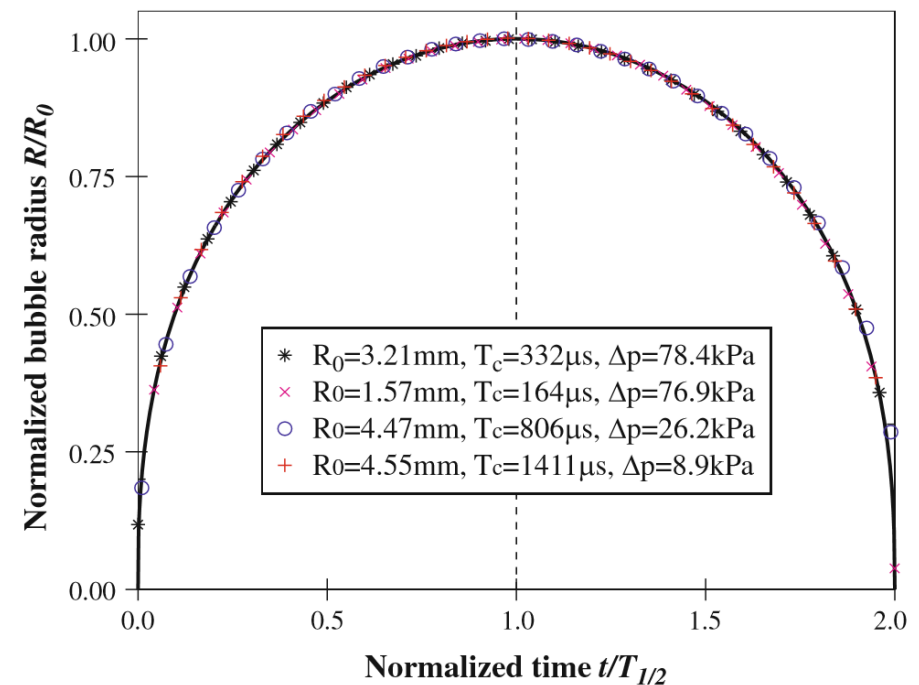
¹for more details, see Supponen et al. (2016) Journal of Fluid Mechanics

Scaling laws for collapse induced jets

- Pressure anisotropy : Relevance ?
 - Experimental validation:
 - Overview of the parameter space covered by the experiment (gravity and a nearby rigid/free surface)



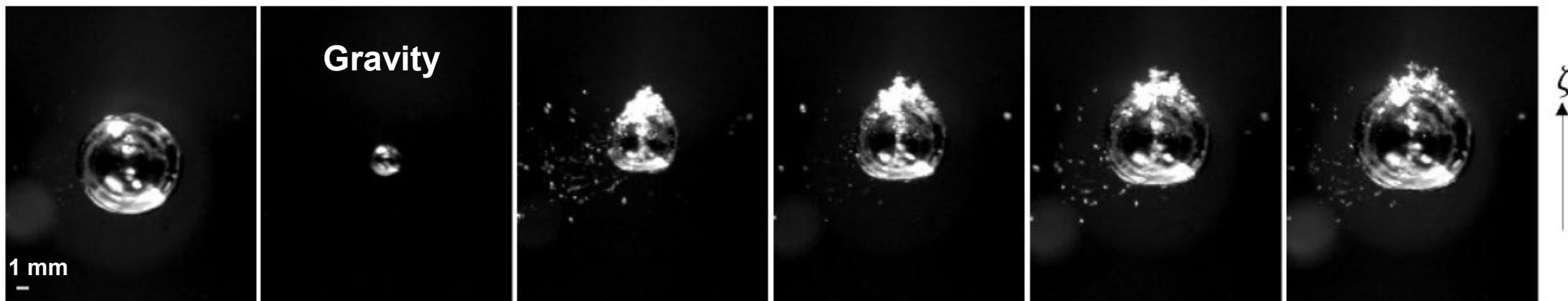
Rayleigh Model: A remarkable agreement



¹for more details, see Supponen et al. (2016) Journal of Fluid Mechanics

Scaling laws for collapse induced jets

- Jet types:
 - 1st type: Weak Jets ($\zeta < 10^{-3}$) :
 - Slight deformation of the bubble
 - The jet hardly pierces the interface and remains within the bubble throughout the collapse and rebound phases.



Scaling laws for collapse induced jets

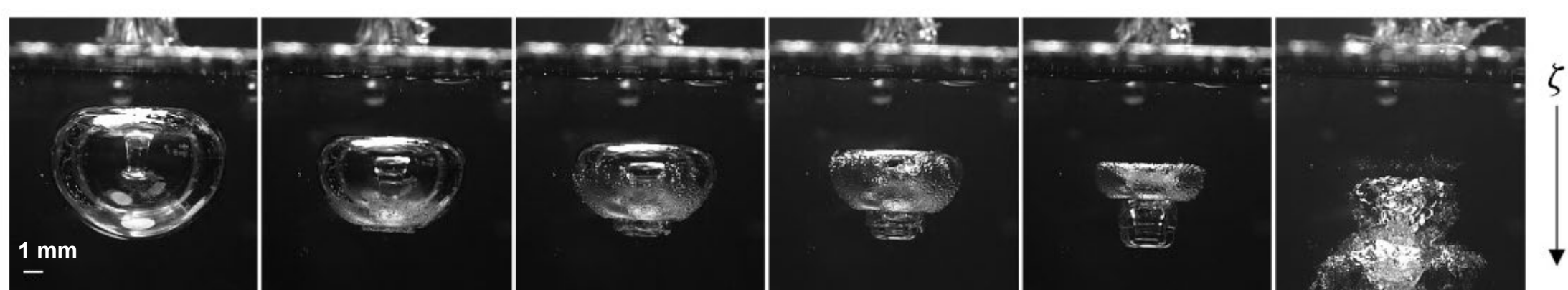
- Jet types:
 - 2nd type: Intermediate jets ($10^{-3} < \zeta < 0.1$) :
 - The bubble deforms significantly at the end of the collapse phase
 - The jet pierces the bubble and clearly emerges during the rebound phase



¹for more details, see Supponen et al. (2016) *Journal of Fluid Mechanics*

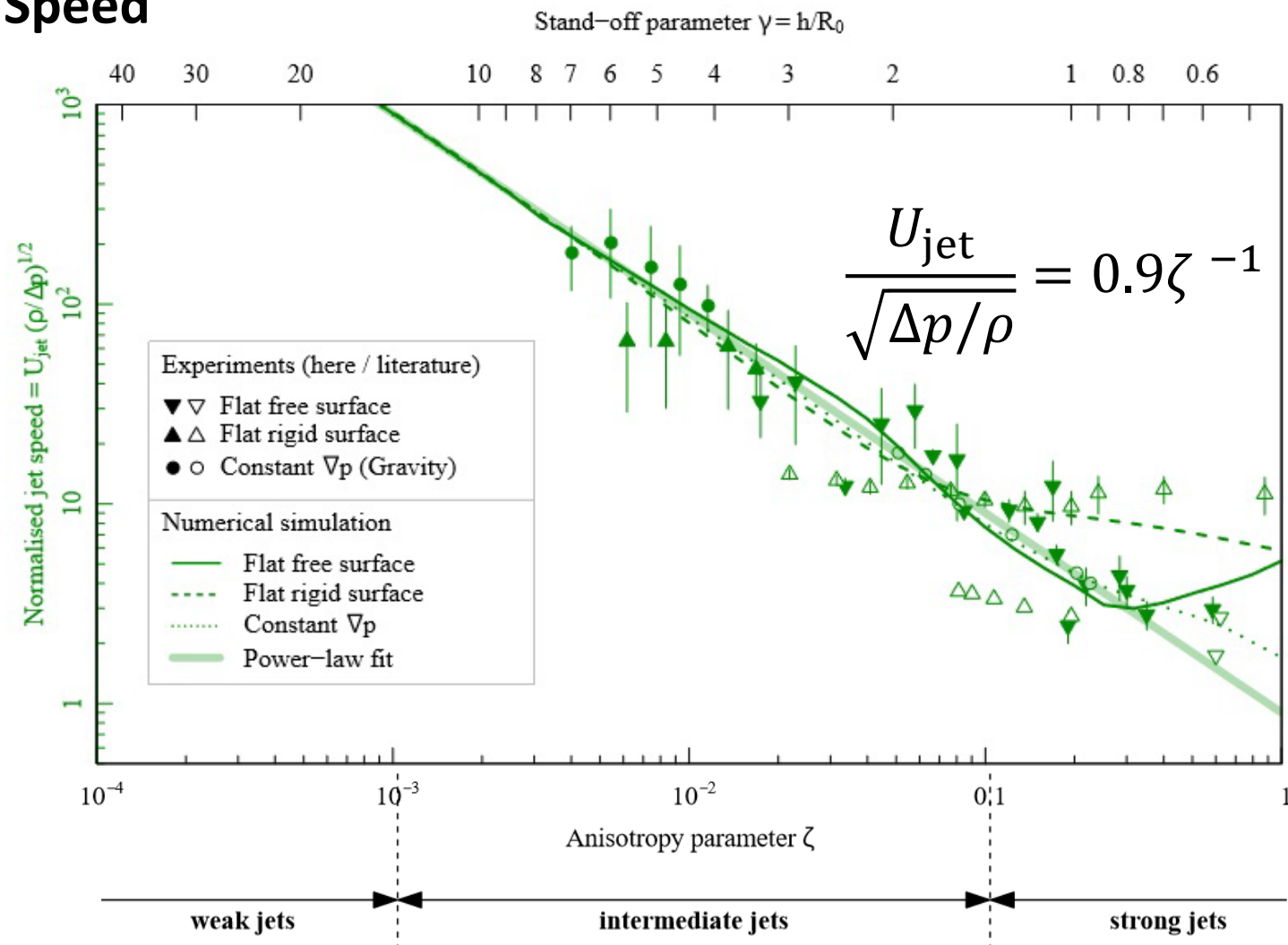
Scaling laws for collapse induced jets

- Jet types:
 - 3rd type: Strong jets ($\zeta > 0.1$) :
 - The bubble deforms significantly early in the collapse phase
 - The jet pierces the bubble “long time” before the collapse



Scaling laws for collapse induced jets

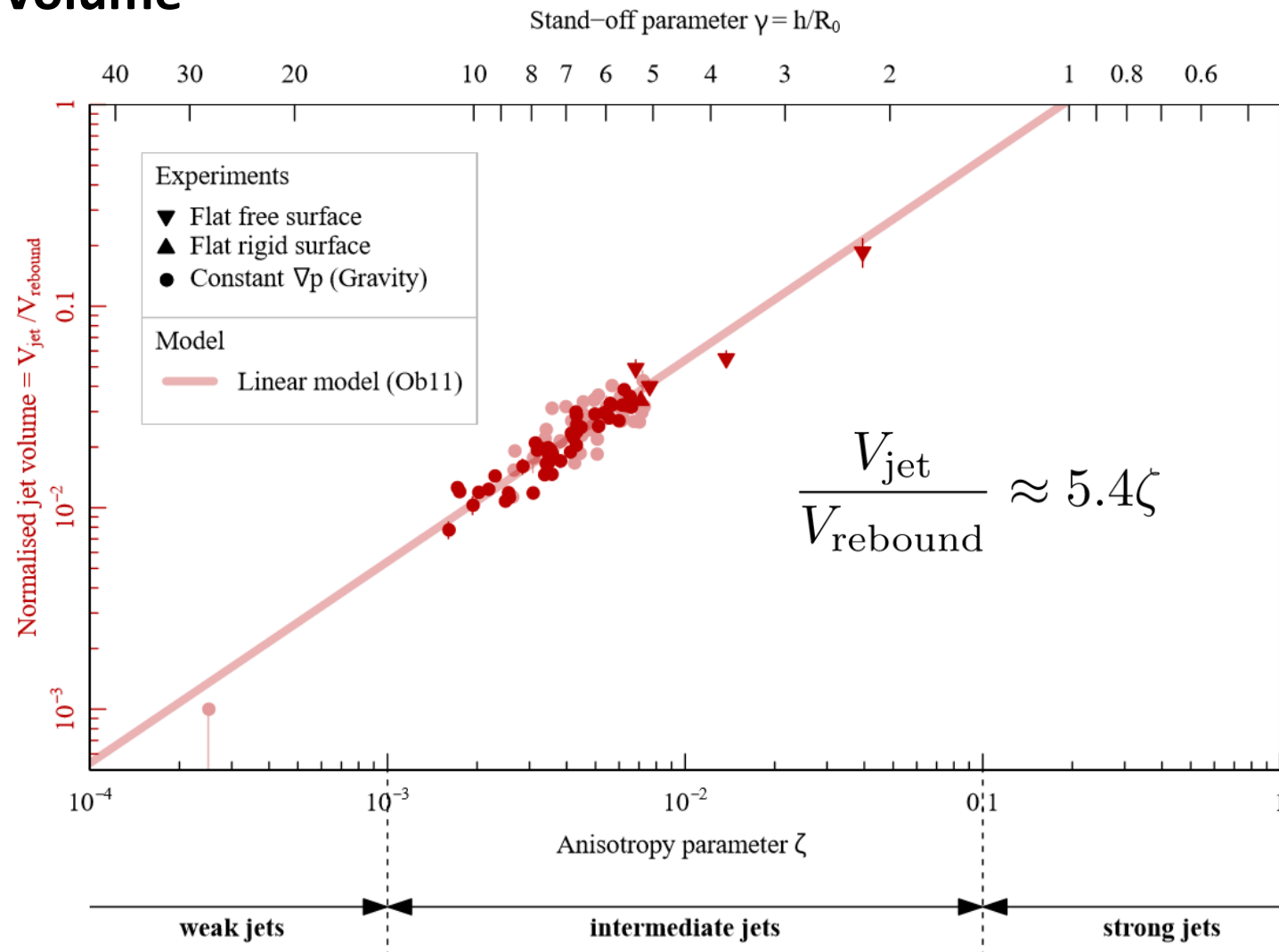
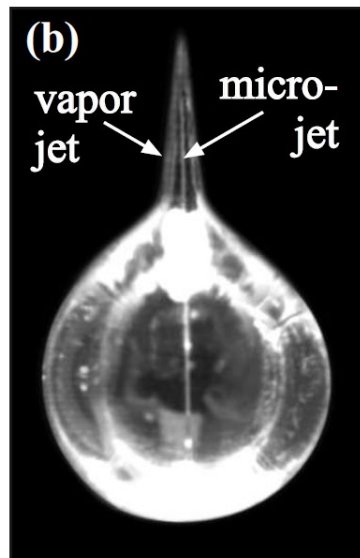
- Jet Speed



¹for more details, see Supponen et al. (2016) Journal of Fluid Mechanics

Scaling laws for collapse induced jets

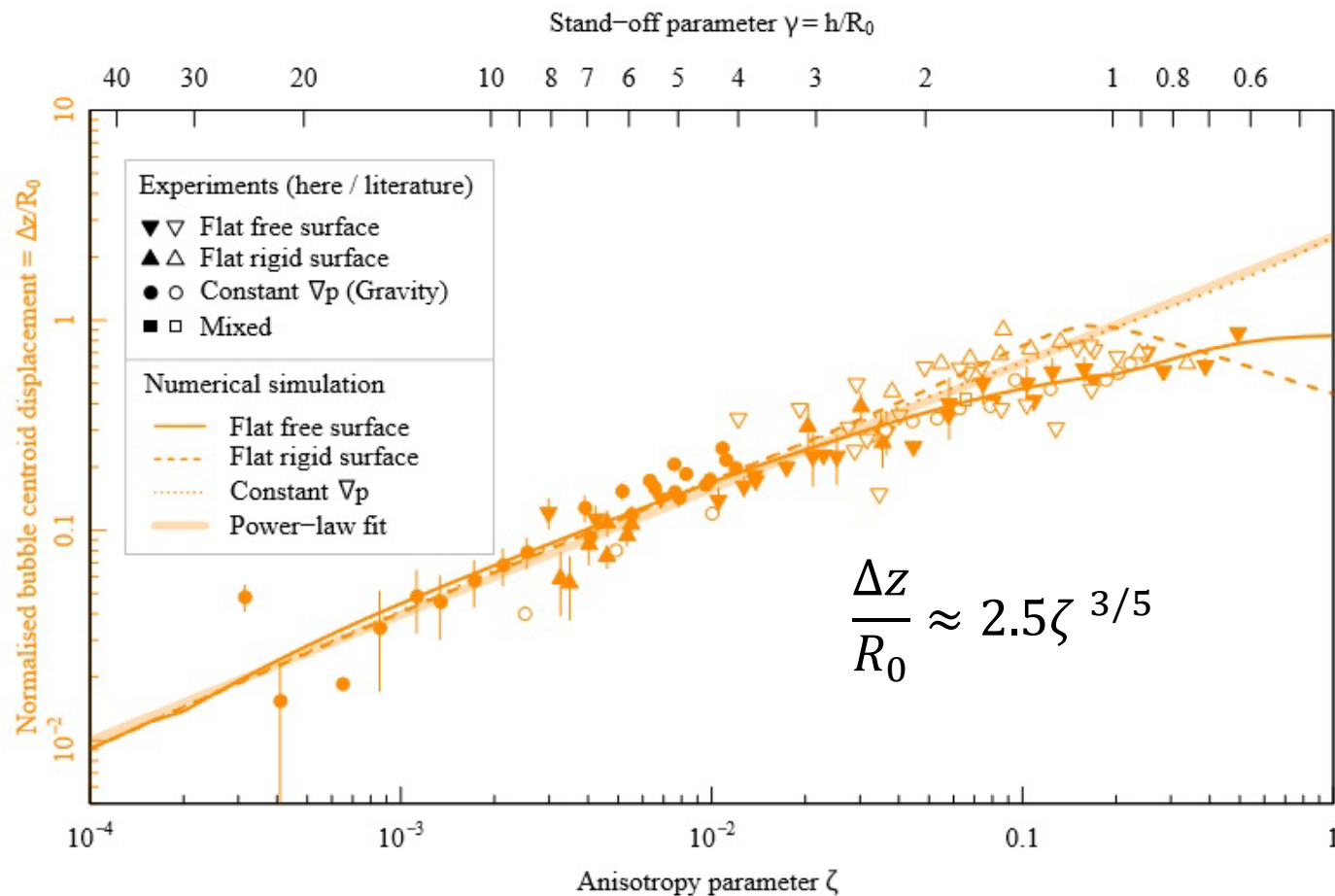
- Vapor Jet Volume



¹for more details, see Supponen et al. (2016) Journal of Fluid Mechanics

Scaling laws for collapse induced jets

- Bubble centroid displacement

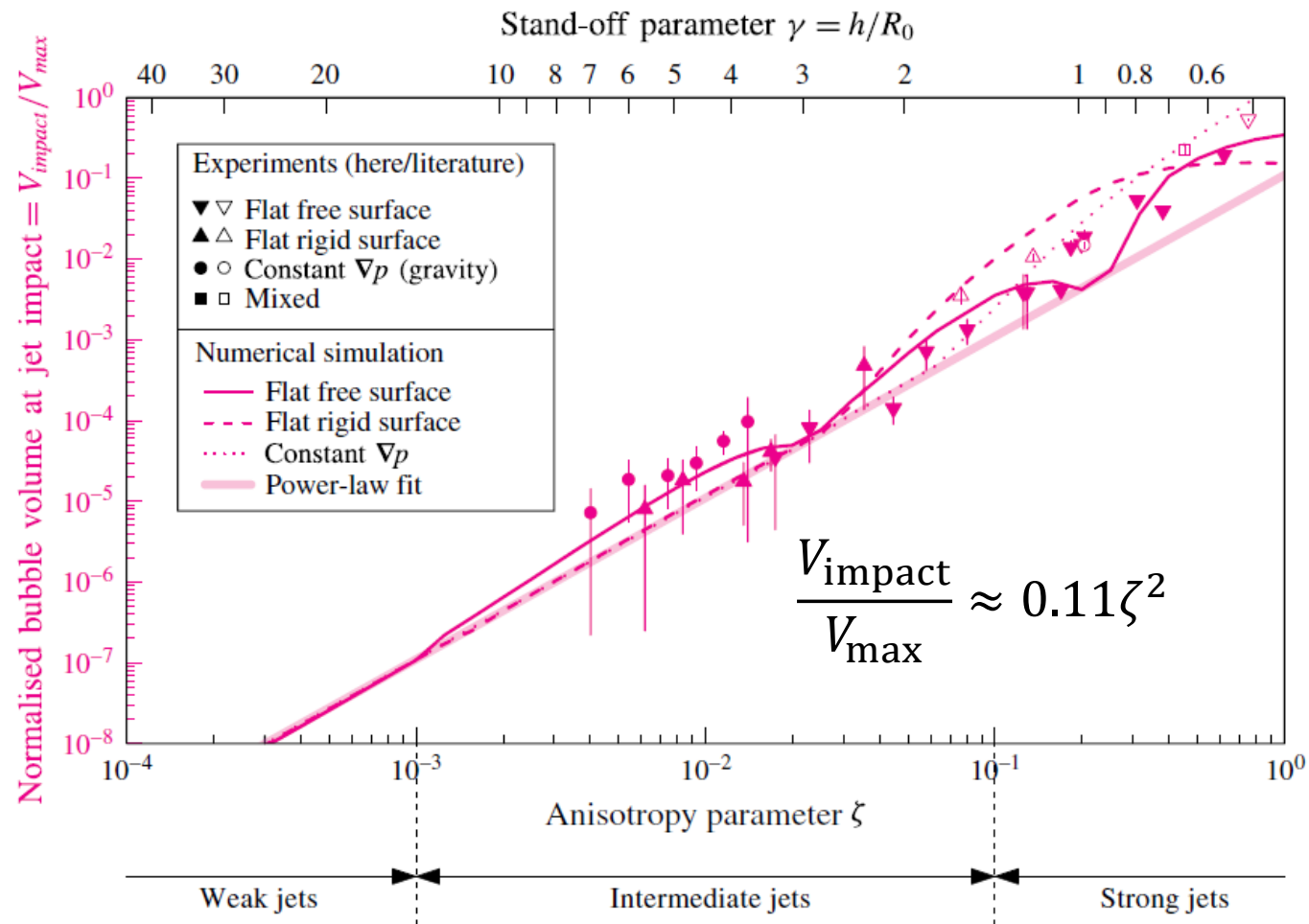


¹for more details, see Supponen et al. (2016) *Journal of Fluid Mechanics*

Scaling laws for collapse induced jets

Bubble volume at jet impact

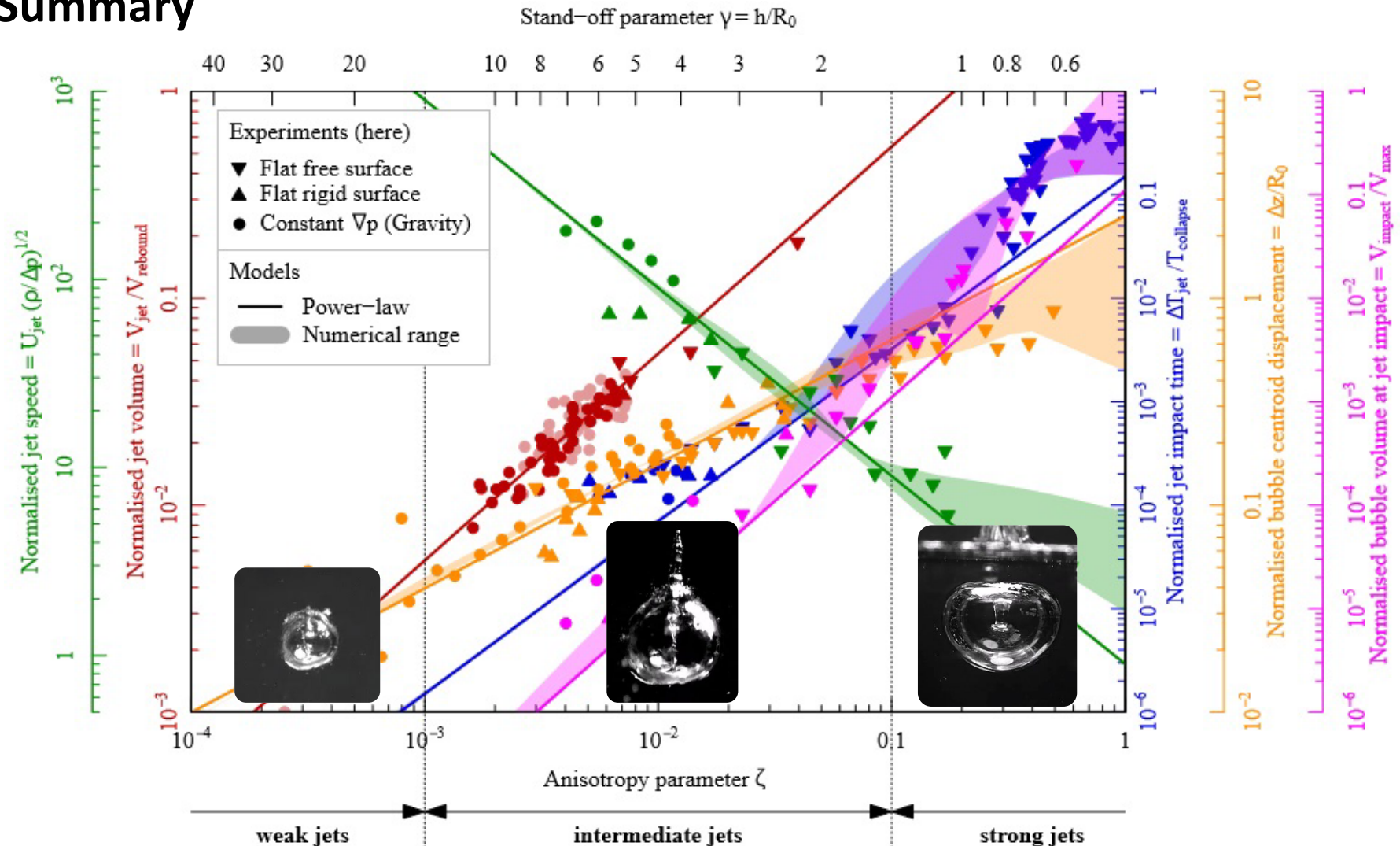
- Bubble volume at jet impact



¹for more details, see Supponen et al. (2016) Journal of Fluid Mechanics

Scaling laws for collapse induced jets

- Summary



¹for more details, see Supponen et al. (2016) Journal of Fluid Mechanics

Bubble Dynamics: Jet shapes by potential flow computation

- Flow simulation of a collapsing bubble near boundaries :
 - Boundary integral method
 - Experimental validation

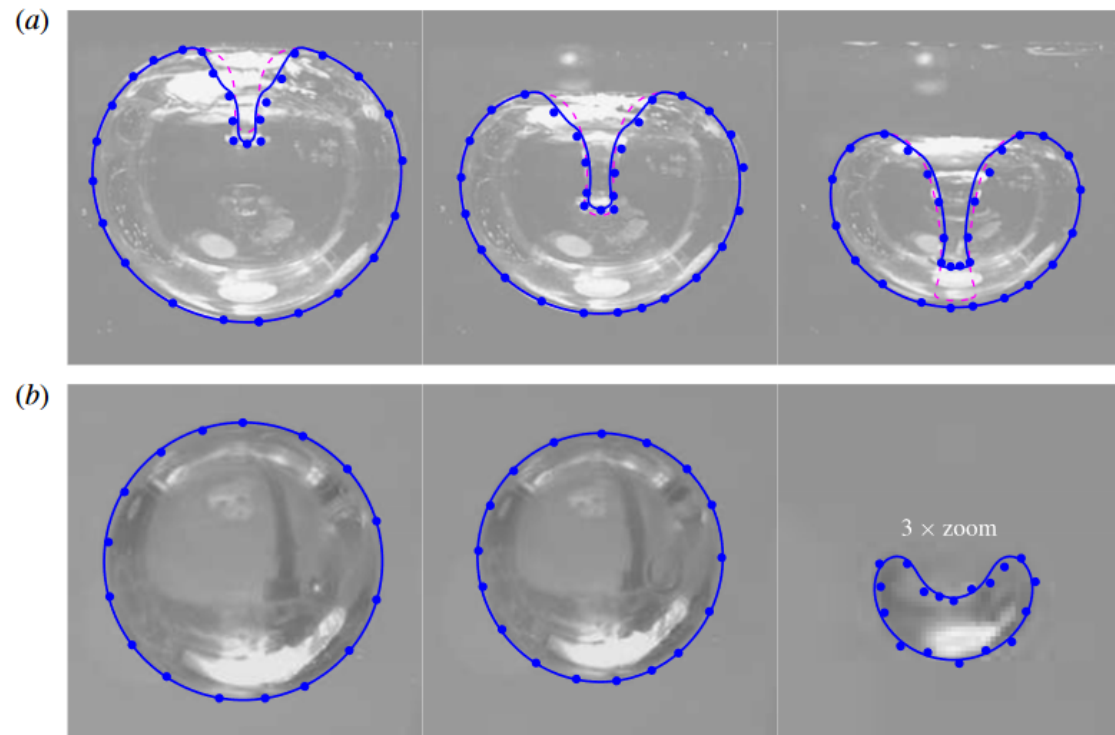


FIGURE 11. (Colour online) The numerical simulations superimposed on the experimental visualisations for a bubble collapsing near a free surface $\gamma = 0.56$ (a) and near a rigid surface $\gamma = 2.32$ (b). The blue points are extracted from the observed bubble shapes and the lines represent simulated data. In (a), the simulated bubble shape (dashed purple line) was corrected for optical refraction (solid blue line) by the outer bubble boundary, assuming a refraction by a sphere with equations analogous to those in Kobel *et al.* (2009) (with water and vacuum inverted).

¹for more details, see Supponen *et al.* (2016) *Journal of Fluid Mechanics*

Bubble Dynamics: Jet shapes by potential flow computation

- Role of the jet driver (rigid boundary, free surface, gravity) on the shape of the jet induced by a collapsing bubble:

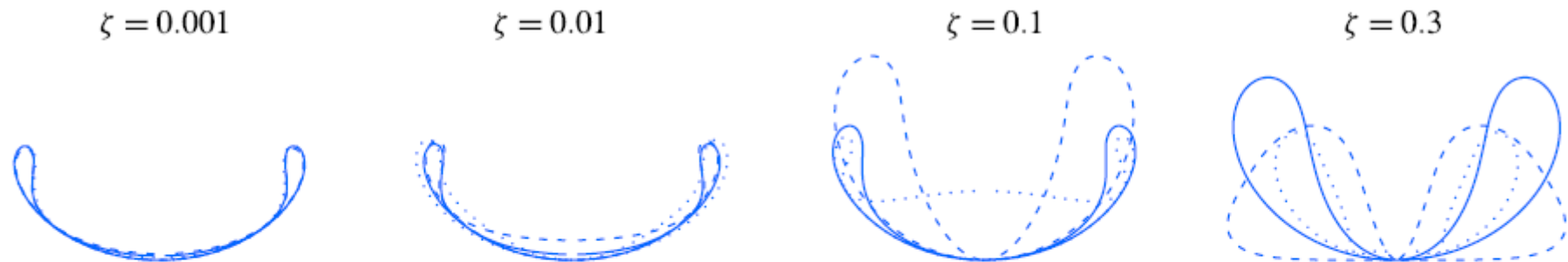


FIGURE 13. (Colour online) Zoomed bubble shapes at the jet impact from figure 12. The different jet drivers are indicated by solid (constant ∇p), dashed (rigid surface) and dotted (free surface) lines.

¹for more details, see Supponen et al. (2016) *Journal of Fluid Mechanics*

Scaling laws for collapse induced jets

$$\left. \begin{aligned}
 \Delta T_{jet}/T_{collapse} &= 0.15 \zeta^{5/3} && \text{(normalised jet impact time),} \\
 U_{jet}/(\Delta p/\rho)^{1/2} &= 0.9 \zeta^{-1} && \text{(normalised jet speed),} \\
 \Delta z/R_0 &= 2.5 \zeta^{3/5} && \text{(normalised bubble displacement),} \\
 V_{impact}/V_{max} &= 0.11 \zeta^2 && \text{(normalised bubble volume at jet impact),} \\
 V_{jet}/V_{rebound} &= 5.4 \zeta && \text{(normalised volume of vapour-jet).}
 \end{aligned} \right\} \text{Power laws}$$

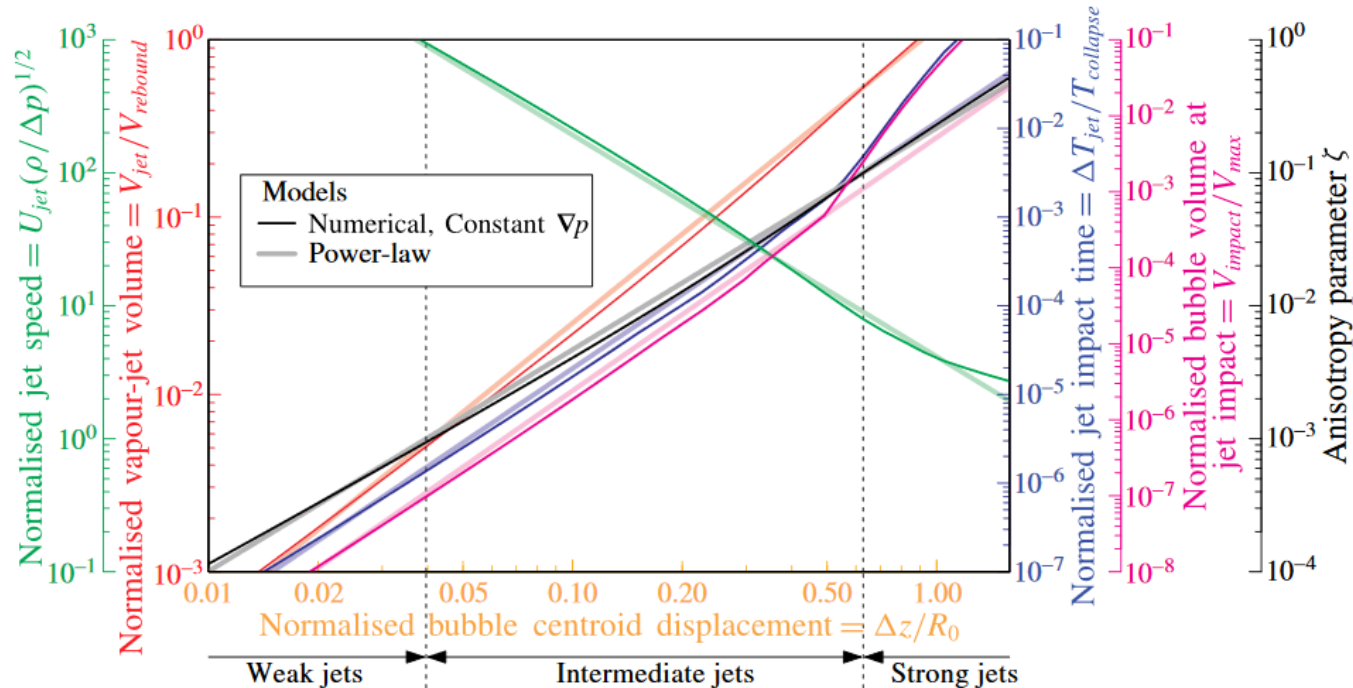


FIGURE 20. The pressure anisotropy parameter ζ and the normalised jet impact time, the normalised jet speed, the normalised bubble volume at jet impact and the normalised jet volume (Obreschkow *et al.* 2011) are plotted as functions of the normalised bubble centroid displacement for jets driven by a constant pressure gradient. The simulated models and the power-law fits are plotted with dark and light lines, respectively.

Bubble dynamics near granular boundaries

PhD Thesis, Armand Sieber, 2023

Motivations:

- Cavitation bubble dynamics $\leftrightarrow$ granular media dynamics.
 - Compelling scientific interest with rich physics.
- Applications in the field of underwater explosions, when the charges are detonated near the seafloor.
 - Information about the bubble acoustic signature (*Krieger et al, 2005*).

Method:

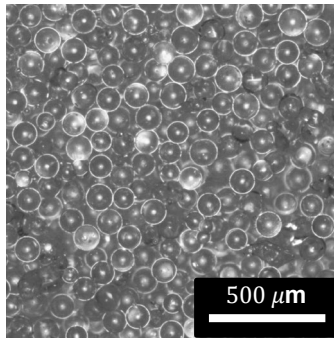
- Experimental investigation laser-induced cavitation bubbles near beds of sand with different grain sizes.

Bubble dynamics near granular boundaries

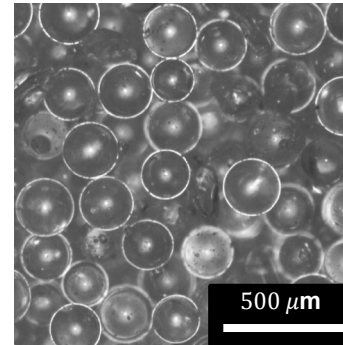
Granular boundaries → experiment:

- Beds of sand made out of soda lime glass beads → 3 granularities investigated:

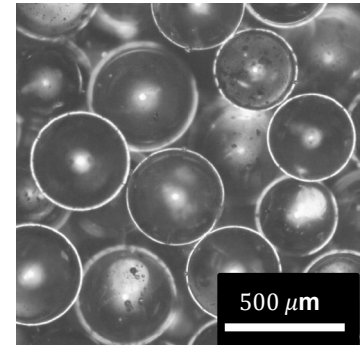
Fine sand:
 $d_g \approx 90\mu\text{m}$



Medium sand:
 $d_g \approx 200\mu\text{m}$



Coarse sand:
 $d_g \approx 350\mu\text{m}$

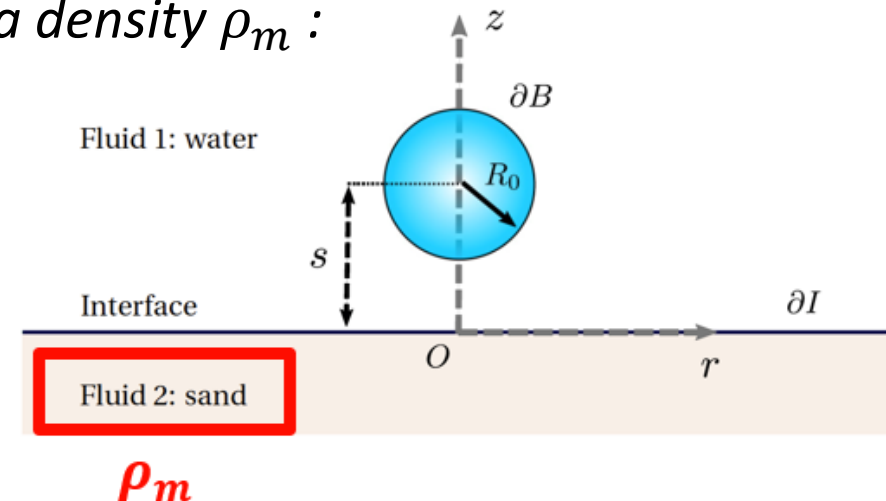


Granular boundaries → simulation (BIM):

- *Water-saturated sand treated as a fluid with a density ρ_m :*

$$\rho_m = (1 - \phi)\rho_s + \phi\rho_w$$

glass beads densitywater density
measured porosity



Bubble dynamics near granular boundaries

Impact of the standoff distance $\rightarrow$ 3 regimes depending on γ



First regime:

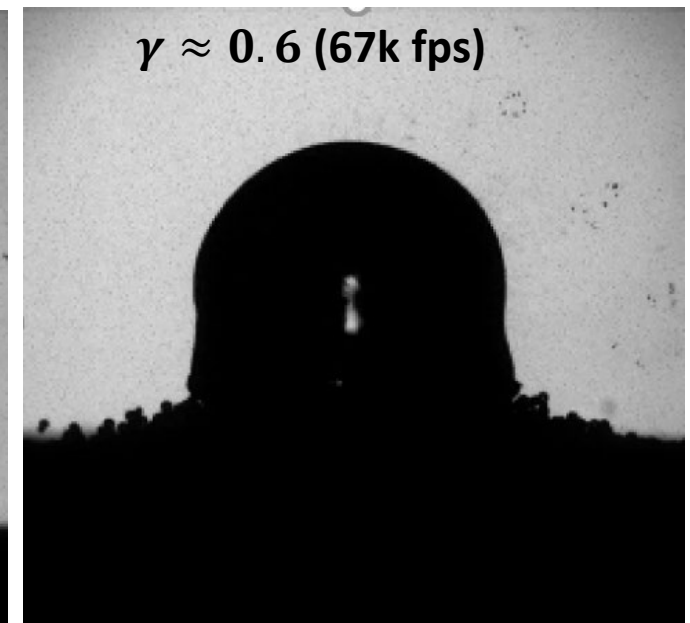
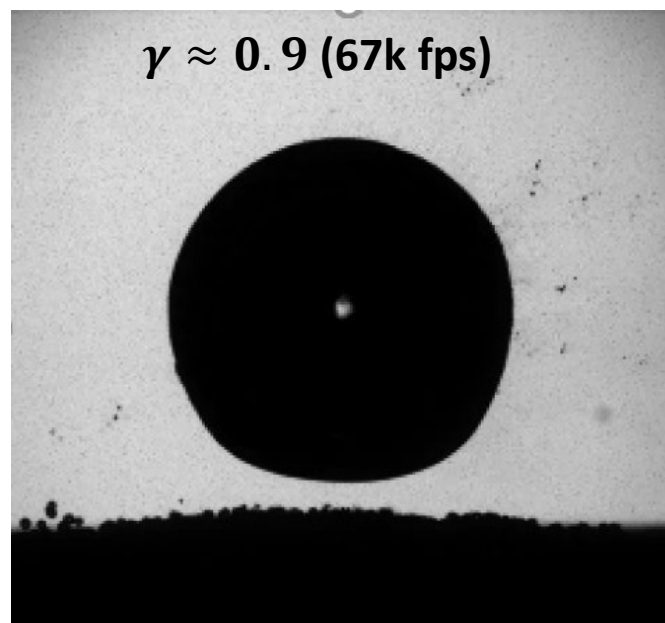
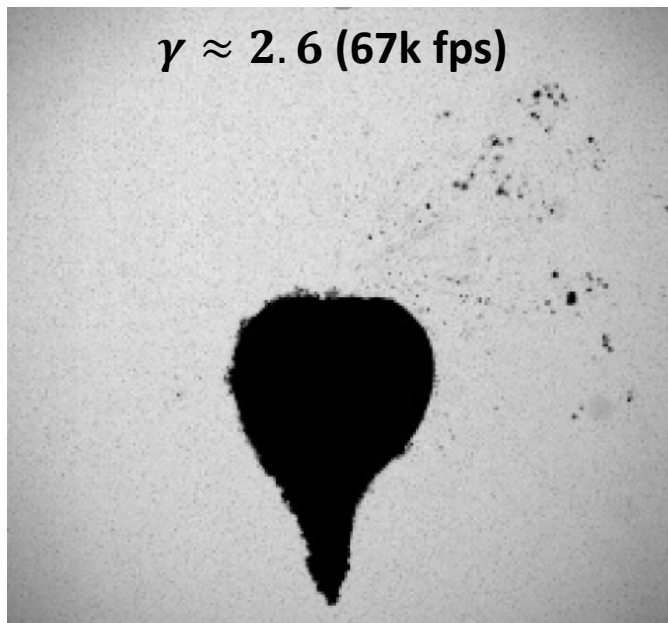
- No influence of grain size.
- Simulations fairly reproduce bubble dynamics.

Second regime:

- Bubble's rebound 'swallowed' by sand.
- Sand grain size matters.

Third regime:

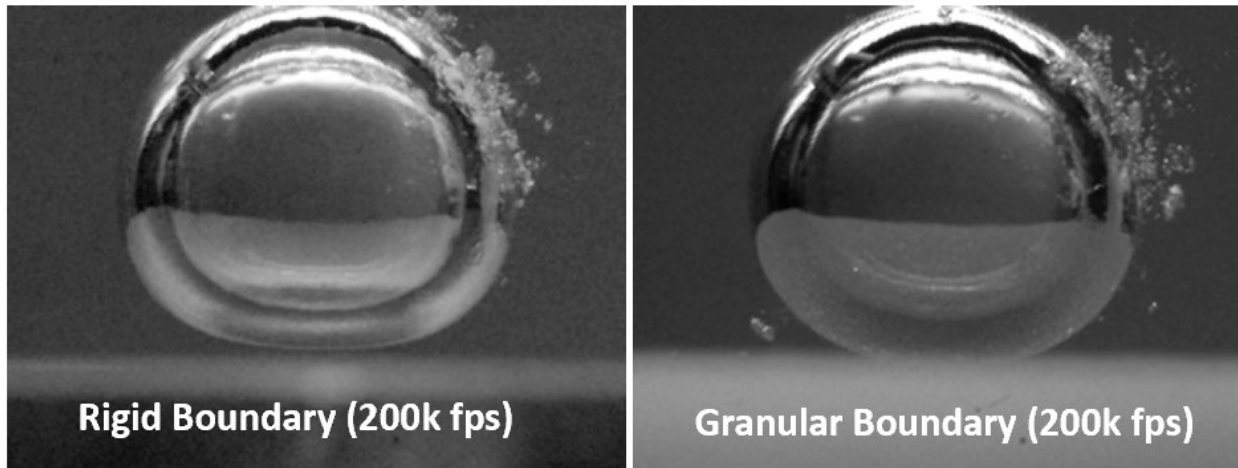
- Bubble 'swallowed' by sand.
- Sand grain size matters.



Bubble dynamics near granular boundaries

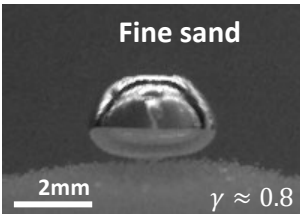
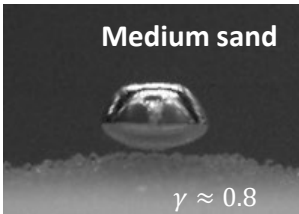
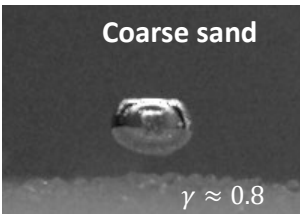
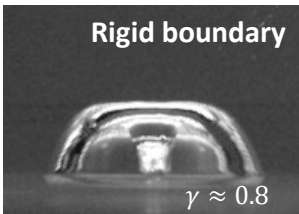
Results: second regime $1.3 > \gamma > 0.6$

Bubbles dynamics at $\gamma \approx 0.8$ ($R_{max} = 4.4$ mm)



- Formation of a granular mound upon bubble collapse.
- Bubble takes on a conical shape.
- Bubble rebound 'swallowed' by sand.

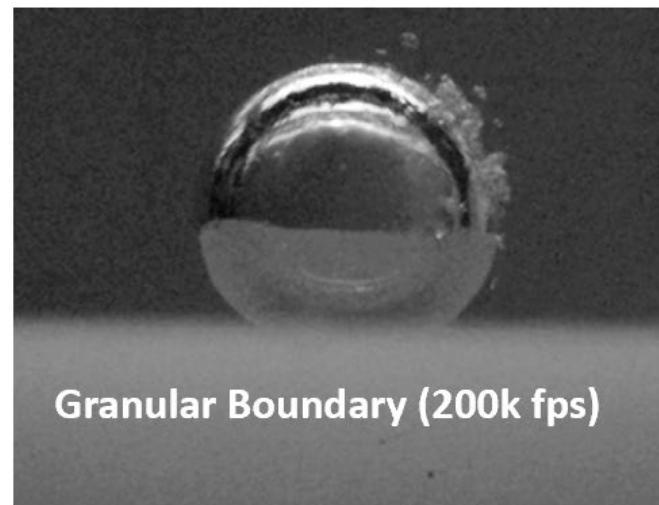
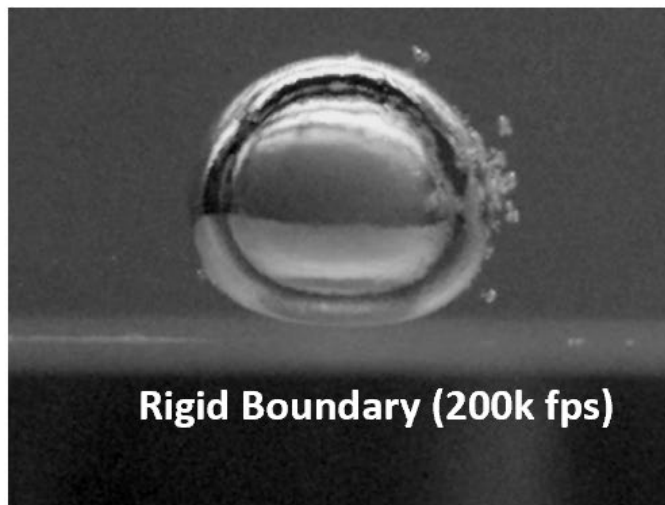
■ Effect of the sand grain size on the bubble shape and microjet speed:

$v_{jet} = 328 \text{ m/s} \leftarrow$			$\rightarrow v_{jet} = 289 \text{ m/s}$	• Finer sand $\rightarrow$ larger bubble volume at microjet impact. • Finer sand $\rightarrow$ faster and thinner microjets.
$v_{jet} = 218 \text{ m/s} \leftarrow$			$\rightarrow v_{jet} = 87 \text{ m/s}$	

Bubble dynamics near granular boundaries

Results: third regime $0.6 > \gamma > 0.3$

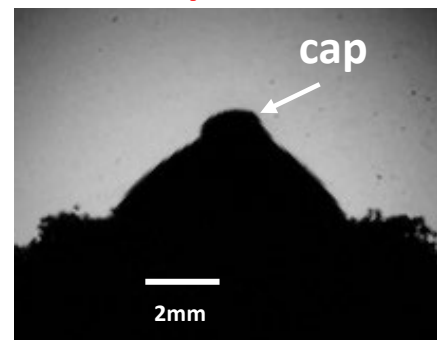
Bubbles dynamics at $\gamma \approx 0.5$ ($R_{max} = 4.4$ mm)



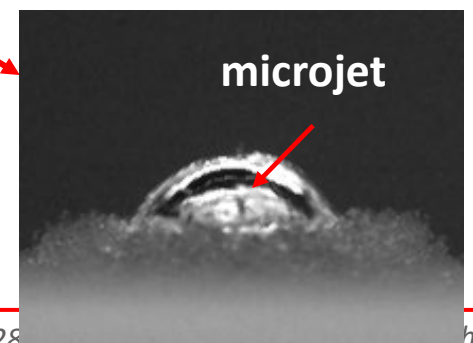
- Bubble takes a bell-shape.
- Bubble 'swallowed' by sand
- Bubble rebounds within the sand
- Formation of a post-collapse crater

- Final instants of the bubble collapse:
 - The bell-shaped bubble cap collapses faster than the rest of the bubble
 - Emission of a shockwave
 - Very fast microjet : ~ 1000 m/s !

Bell-shaped bubble

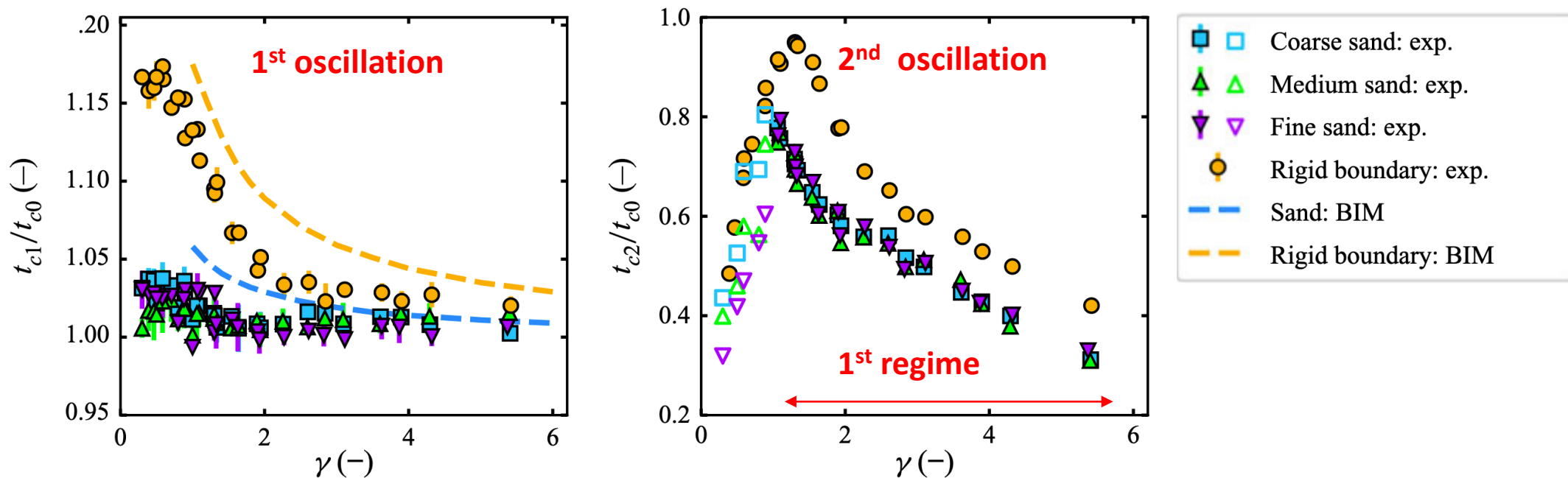


+ 5 μ s



Bubble dynamics near granular boundaries

Results: bubble 1st and 2nd oscillation periods (t_{c1} and t_{c2})



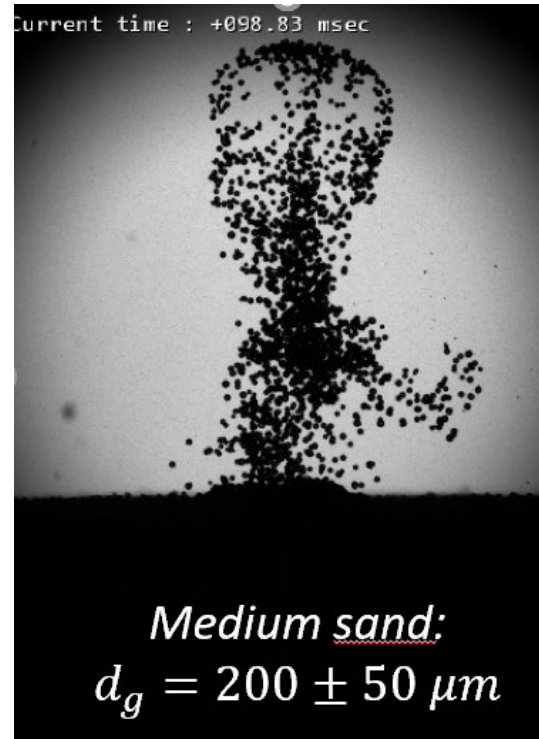
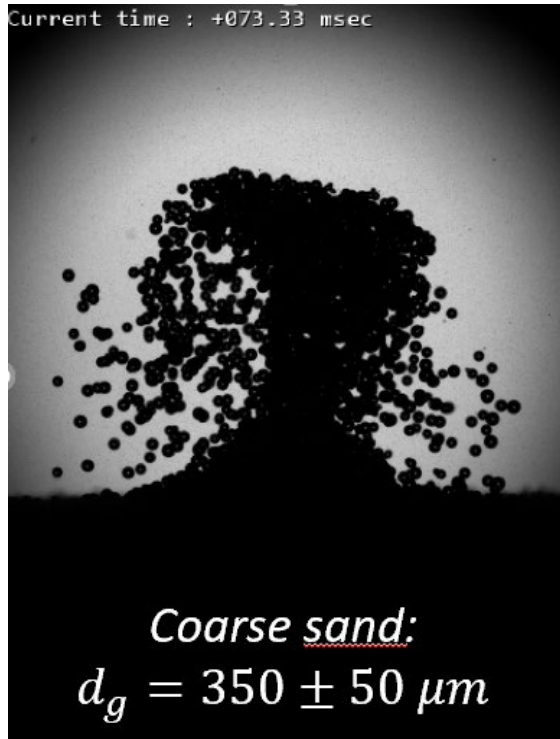
- Weak influence of the sand grain size.
- Bubble oscillations are longer near a rigid boundary.

$$t_{c0} = 1.83R_{max}\sqrt{\rho/\Delta p} \quad (2x \text{ Rayleigh collapse time})$$

Bubble dynamics near granular boundaries

Granular boundaries: jets of sand

$\gamma \approx 0.6$
 $R_{max} \approx 4.44 \text{ mm}$



Two possible mechanisms may explain the sand jetting:

1. Rebound of the bubble below the sand surface creates a cavity in the granular medium $\rightarrow$ collapse of the cavity $\rightarrow$ formation of a jet of sand that shoots upwards
2. Flow field induced by the collapsing bubble $\rightarrow$ ascending vortex ring observed after the toroidal 2nd collapse of bubbles near rigid boundaries (Reuter et al., 2017).

Ref.: Sieber et al. Dynamics of cavitation bubbles near granular boundaries. J. of Fluid Mech., 2022.

Bubble dynamics near tissue mimicking surfaces

PhD Thesis, Armand Sieber, 2023

Motivations:

- Various applications in biomedicine where bubbles may interact with human tissues:
 - Targeted drug delivery (*Stride et Coussios, 2019*).
 - Ultrasound-based tissue ablation (*Bader et al., 2019*).
 - Laser-based medical procedures, e.g. laser lithotripsy (*Ho et al., 2021*) or laser atherectomy (*Golino et. al, 2019*).

Method:

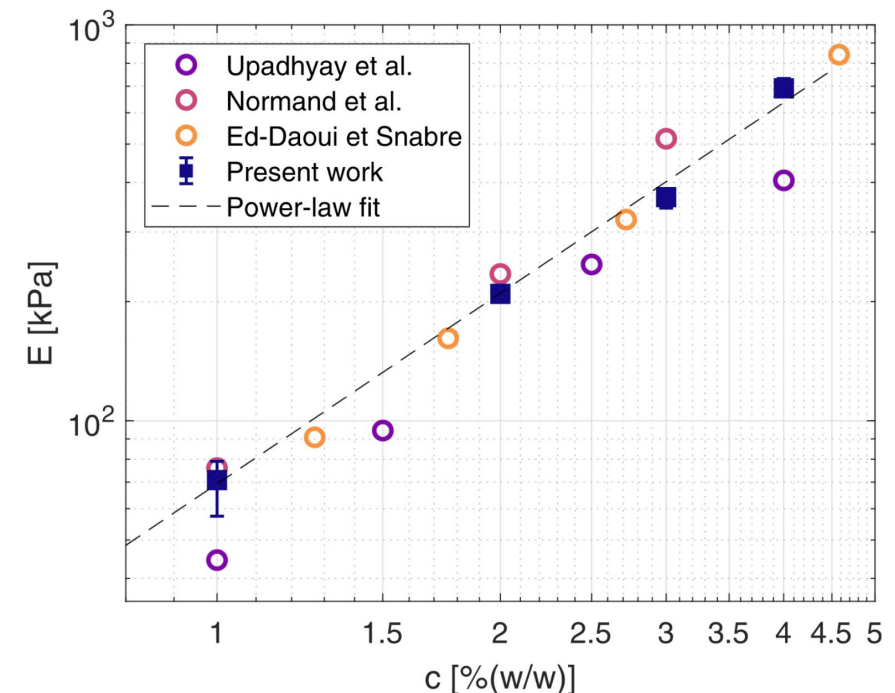
- Simplified test case: study of millimetric cavitation bubbles near initially flat, tissue-mimicking boundaries with different levels of elasticity.

Bubble dynamics near tissue mimicking surface

Interactions of cavitation bubbles with tissue-mimicking materials

→ We used agarose hydrogels

- Agarose hydrogels with different agarose concentration c : 1%, 2%, 3% and 4% w/w.
- Uniaxial compression tests → as c is increased the Young's modulus E increases.
- Power-law scaling: $E = 69.9c^{1.6}$ kPa
- Hydrogels mimic different human tissues, e.g. :
 - Cardiac muscle cells → $E \sim 100$ kPa
 - Cornea → $E \sim 290$ kPa
 - Gastrointestinal tissues → $E \sim 320 - 790$ kPa



Bubble dynamics near tissue mimicking surface

- Numerical method: Boundary Integral Method (BIM)

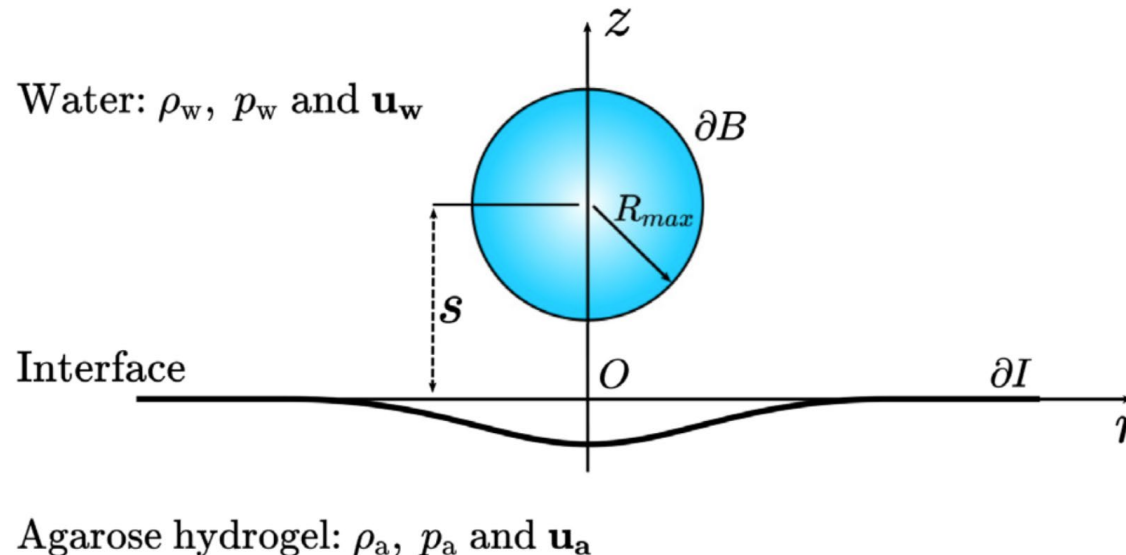
- Agarose hydrogel modelled as a liquid with density ρ_a ,

$$\rho_a = \frac{c\rho_s + (100 - c)\rho_w}{100}$$

- Elasticity of the hydrogel modelled as an interfacial surface tension,

Derived from an energy
balance reasoning

$$\bar{\sigma} = \frac{4}{3\pi} \frac{E}{(1 - \nu^2)\Delta p}$$

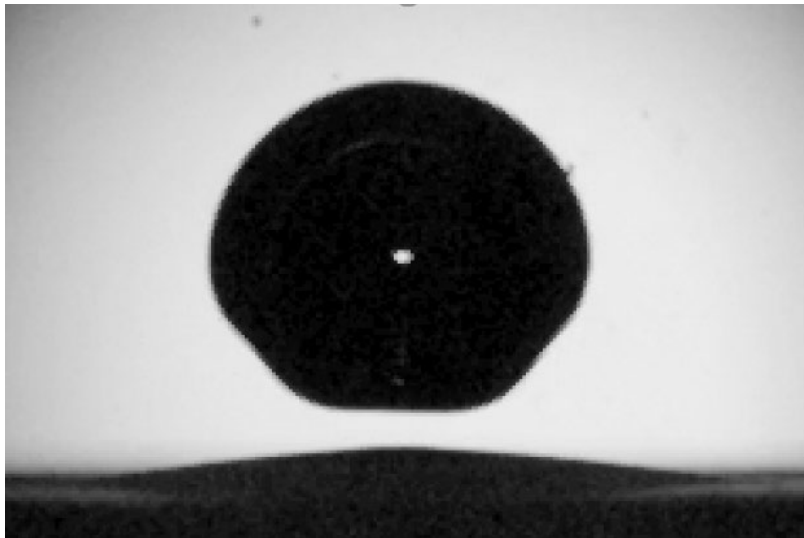


Bubble dynamics near tissue mimicking surface

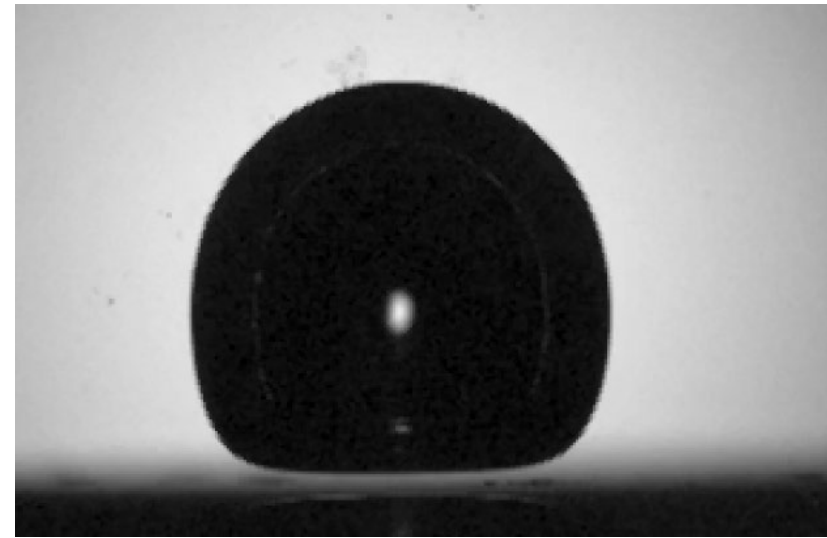
Bubble dynamics near elastic boundary v.s. rigid boundary

Bubbles at $\gamma \approx 0.8$:

- Compression and elastic rebound of the hydrogel $\rightarrow$ different bubble dynamics.
 - Bubble shape: mushroom-like **v.s.** hemispherical.
 - Micro-jet precursor: side walls involution **v.s.** upper hemisphere involution.
 - Rebound: two sub-bubbles **v.s.** one single bubble.



2% agarose hydrogel:
 $\gamma \approx 0.8$, $R_{max} \approx 3.7\text{mm}$, 200k fps

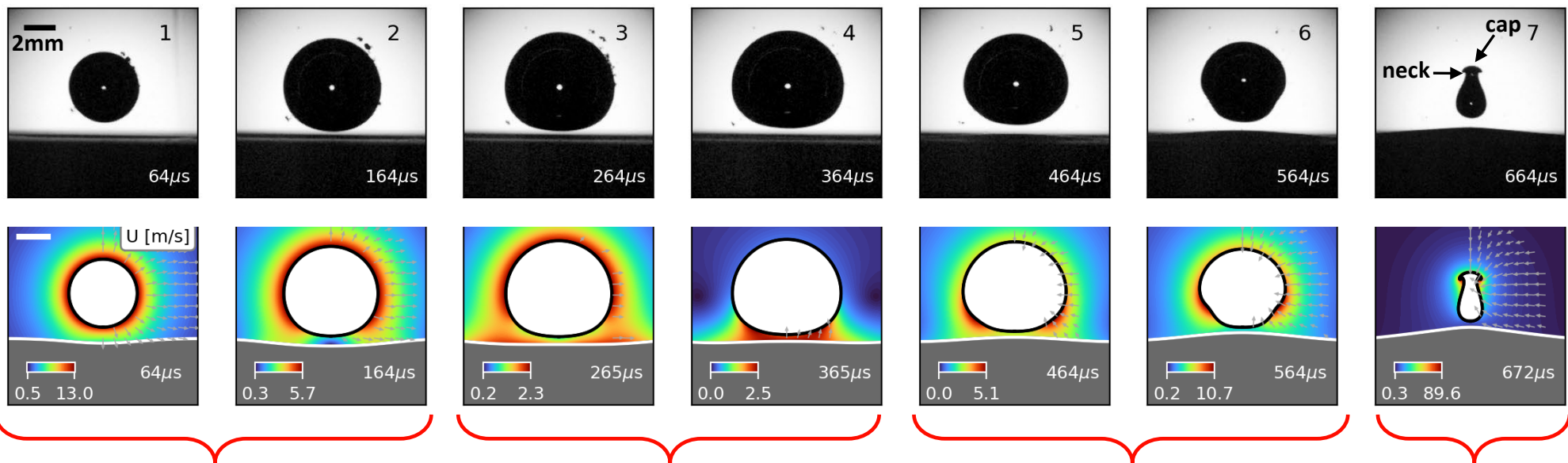


Rigid boundary:
 $\gamma \approx 0.8$, $R_{max} \approx 3.7\text{mm}$, 200k fps

Bubble dynamics near tissue mimicking surface

Towards mushroom-shaped bubbles

3% agarose hydrogel, $\gamma \approx 0.8$, $R_{max} \approx 3.7$ mm: experiment and simulation



Early bubble growth:

- Elastic boundary compressed downwards.
- Flattening of the bubble lower hemisphere.

Late bubble growth:

- Elastic boundary bounces back.
- Flattening of the bubble lower hemisphere

Bubble collapse:

- Elastic boundary forms a mound.
- Indentation of the bubble sidewalls.

Mushroom-shaped bubble:

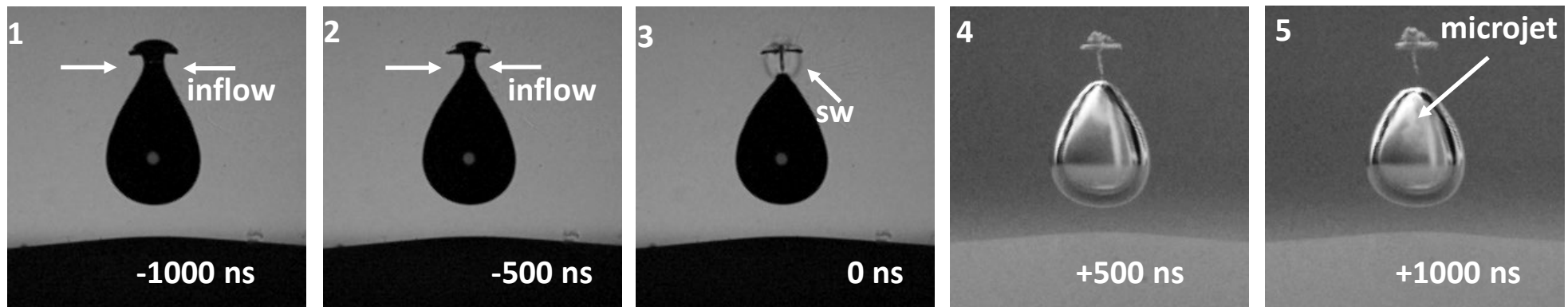
- Hemispherical cap.
- Indented neck.

Bubble dynamics near tissue mimicking surface

Micro-jets (1)

- *Micro-jet initiation mechanism:*
 - *Annular inflow of liquid at the bubble axis of symmetry (1-2).*
 - *Collision of inflow produces a shockwave (3).*
 - *Formation of thin and very fast micro-jet forms (4-5)*

$$\rightarrow u_{jet} \in [200 - 1000] \text{ m/s}$$

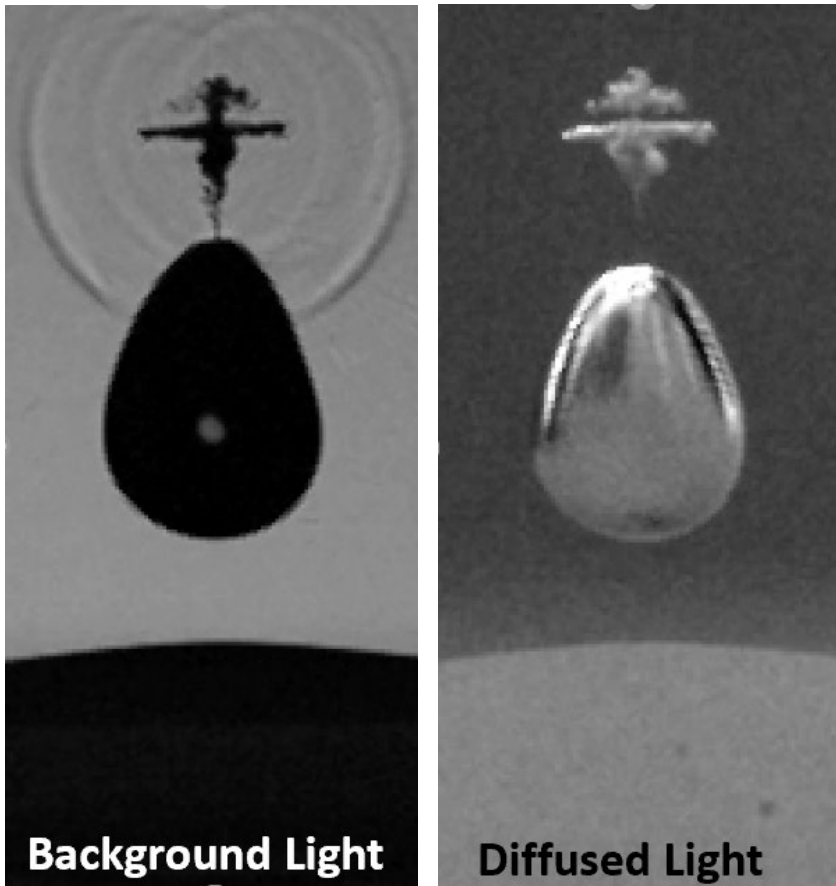


3% agarose , $\gamma \approx 0.8$, $R_{\max} \approx 3.7 \text{ mm}$

Bubble dynamics near tissue mimicking surface

Micro-jets (2)

- Micro-jet atomization preceding the establishment of a fully liquid micro-jet



Possible causes of jet atomization :

- *Instability of the high speed micro-jet:*

➤ *Liquid jet stability is characterized by Weber Nb:*

$$We = \frac{\rho_g u_{jet}^2 d}{\sigma_s}$$

Agarose hydrogel ($\gamma = 0.8$) → $We \approx 35$

Rigid boundary at ($\gamma = 0.8$) → $We \approx 1.1$

- *Interactions with shock-waves:*

➤ *Richtmyer–Meshkov instabilities.*

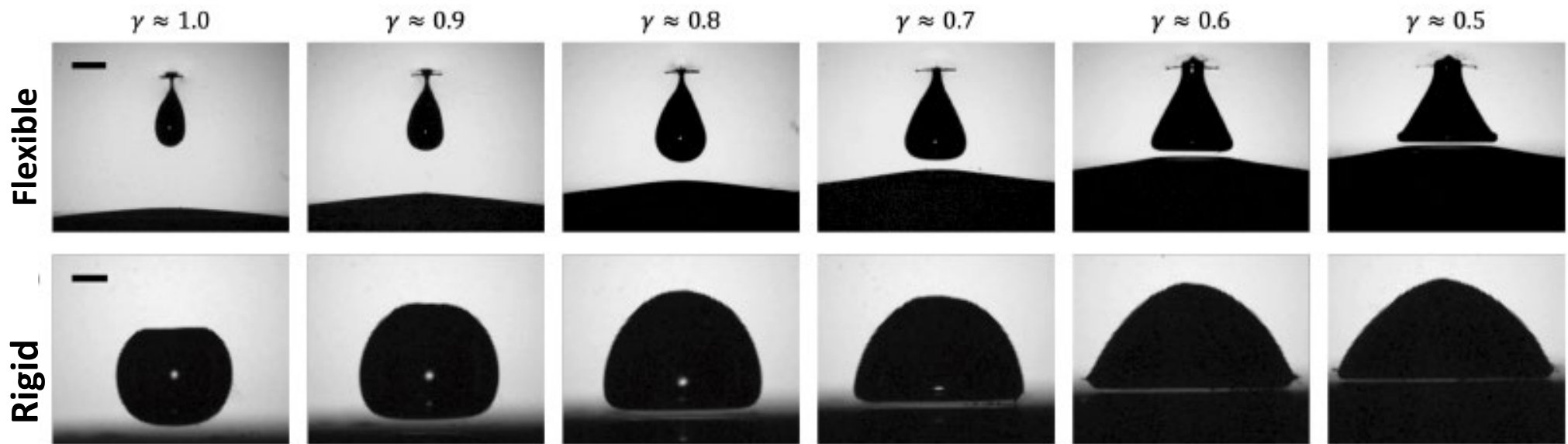
- *Sudden vaporization in the rarefied bubble atmosphere.*

2% agarose , $\gamma \approx 0.8$, $R_{max} \approx 3.7$ mm
10 million fps

Bubble dynamics near tissue mimicking surface

Rigid v.s. elastic boundary (effect of standoff distance):

- *Snapshots of the bubble near the 2% (w/w) agarose hydrogel taken $0.2 \mu\text{s}$ before microjet formation*
- *Snapshots of a similar bubble near a rigid boundary taken at the same time instant after bubble generation.*

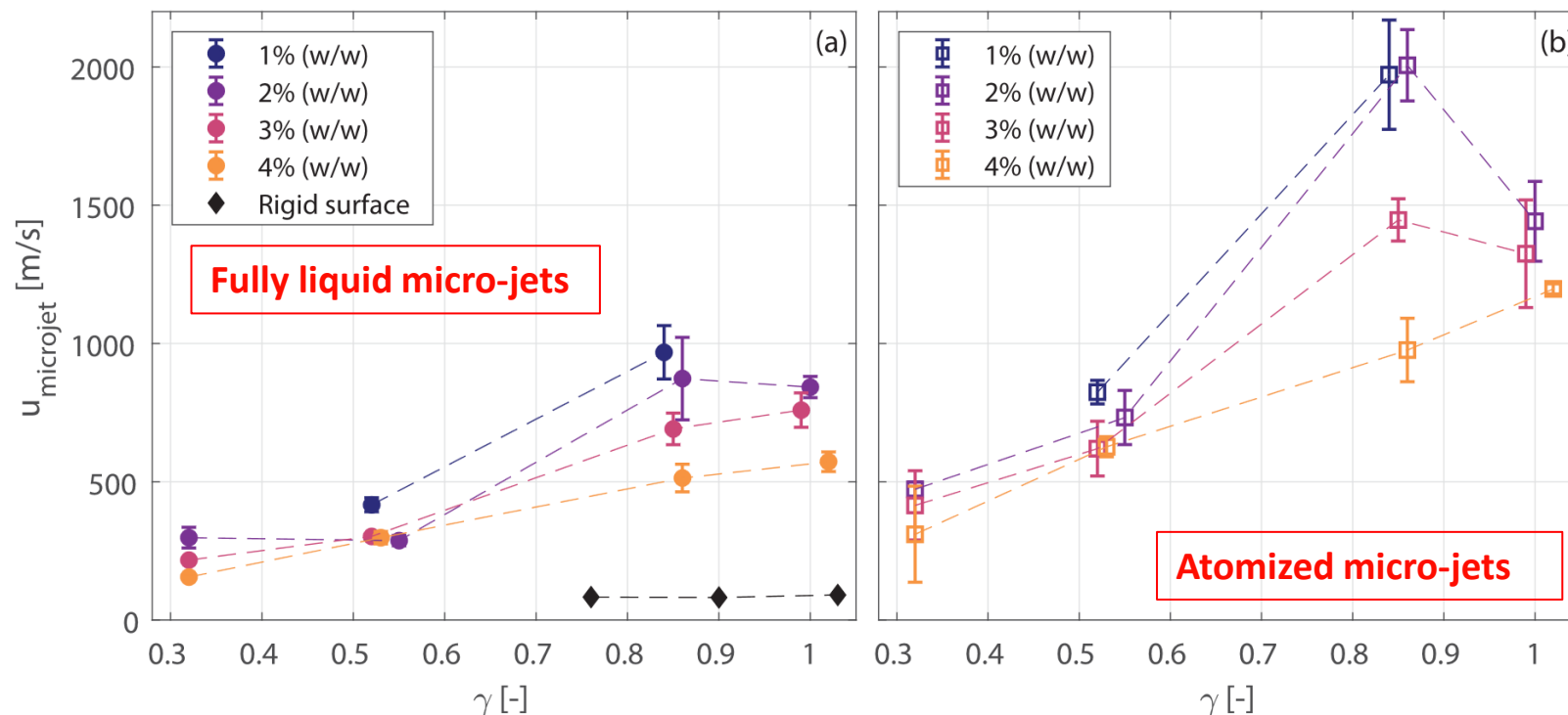


• Bubble dynamics near tissue mimicking surface

Micro-jets (3)

▪ Micro-jet velocities

- Depends on both the stand-off distance and hydrogel elasticity
- Maximum velocity of the **fully liquid micro-jet**: ~ 1000 m/s or $Ma = 2.4$
- Maximum velocity of the **atomized part of the jet**: ~ 2000 m/s or $Ma = 4.8$

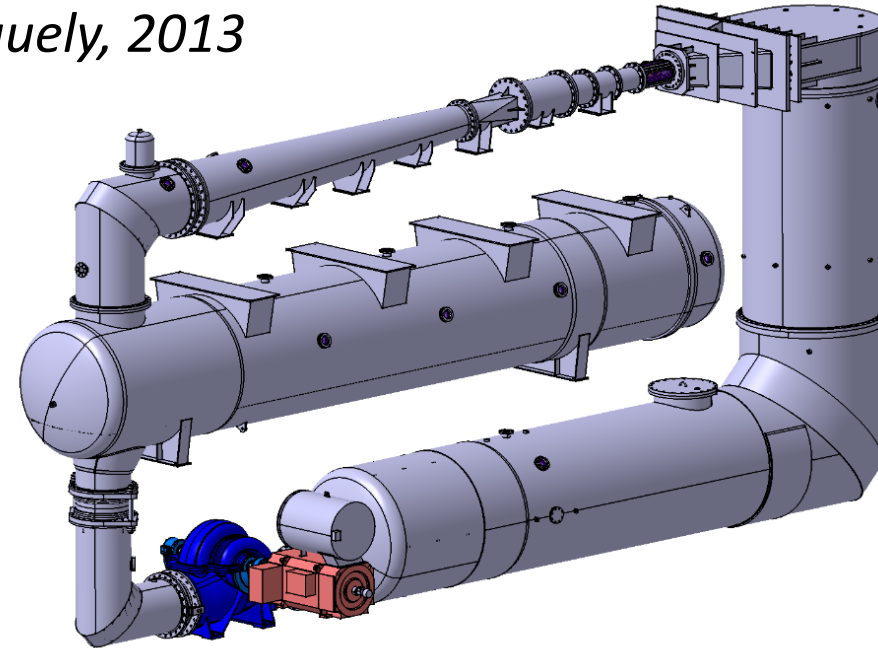


A. Sieber et al., Cavitation bubble dynamics & microjet atomization near tissue-mimicking materials, Phys. of Fluids, 2023.

Collapse of laser induced bubble in a flowing water

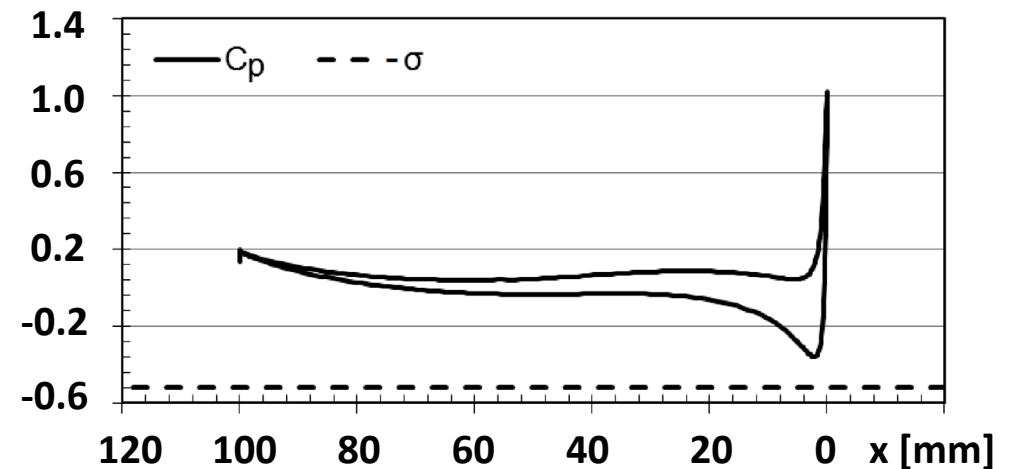
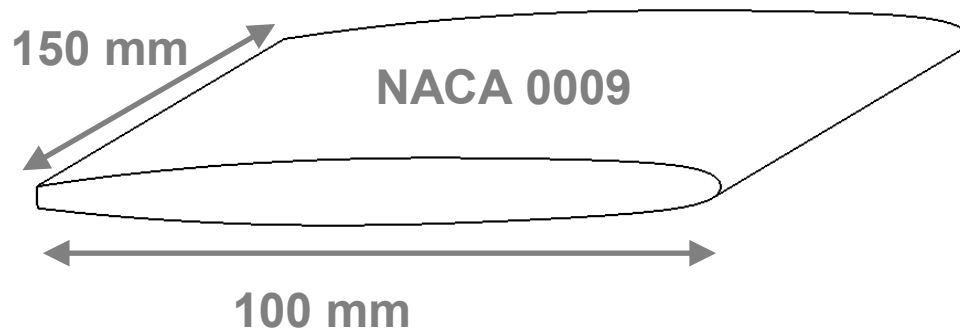
PhD Thesis, Marc Tinguely, 2013

- Experimental Conditions:
 - Hydrofoil : NACA 0009 truncated
 - Inlet velocity : 15 m/s ($Re = 1.5 \cdot 10^6$)
 - Incidence Angle : $\alpha = 1^\circ$
 - Cavitation Coefficient : $\sigma = 0.4$



$$c_p(M) = \frac{p_M - p_{ref}}{\frac{1}{2} \rho C_{ref}^2}$$

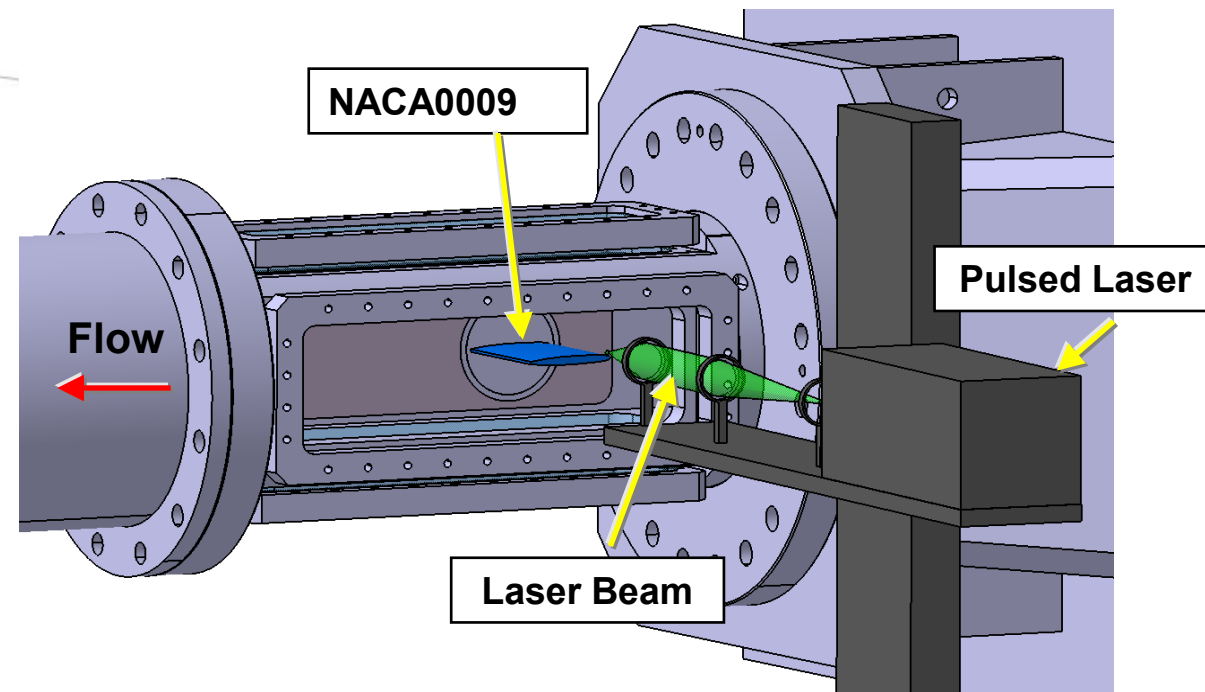
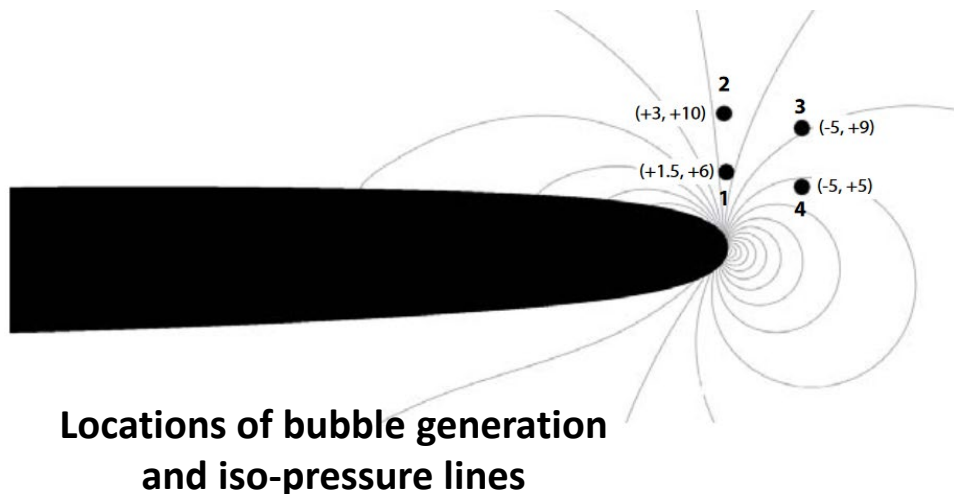
$$\sigma = \frac{p_{ref} - p_v(T)}{\frac{1}{2} \rho C_{ref}^2}$$



Collapse of laser induced bubble in a flowing water

Bubble Generation

- **On demand bubble generation in a flowing liquid**
 - Pulsed Laser focused within the liquid flow
 - Transparent glass window (high transmissibility for the specific wavelength of the laser)
 - Possibility to create a bubble anywhere in the test section
 - investigation of the dynamic of a single bubble in a flowing water
 - Visualization with a high speed camera



Collapse of laser induced bubble in a flowing water

Effect of initial position of the bubble on its collapse

- **Case 1: Jet towards the solid surface**

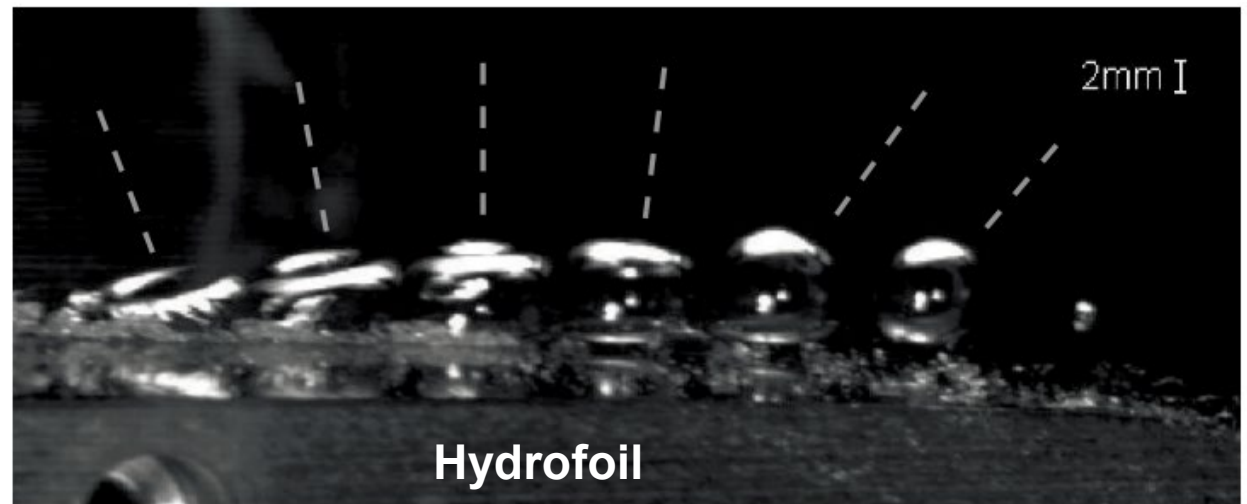
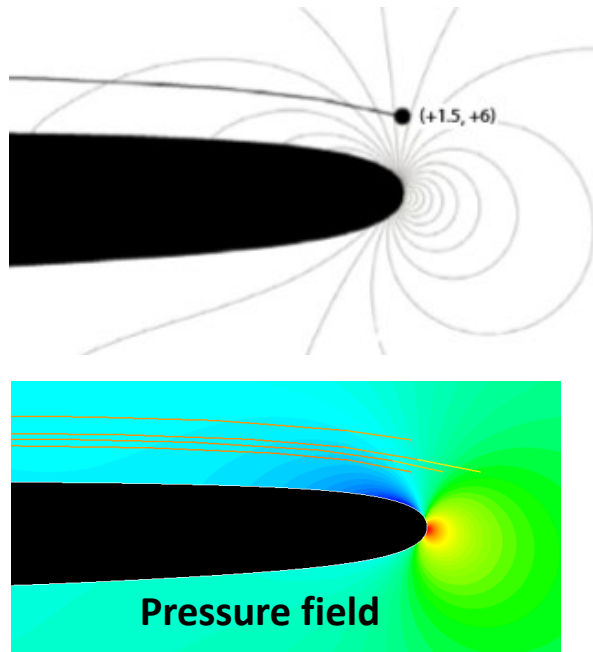
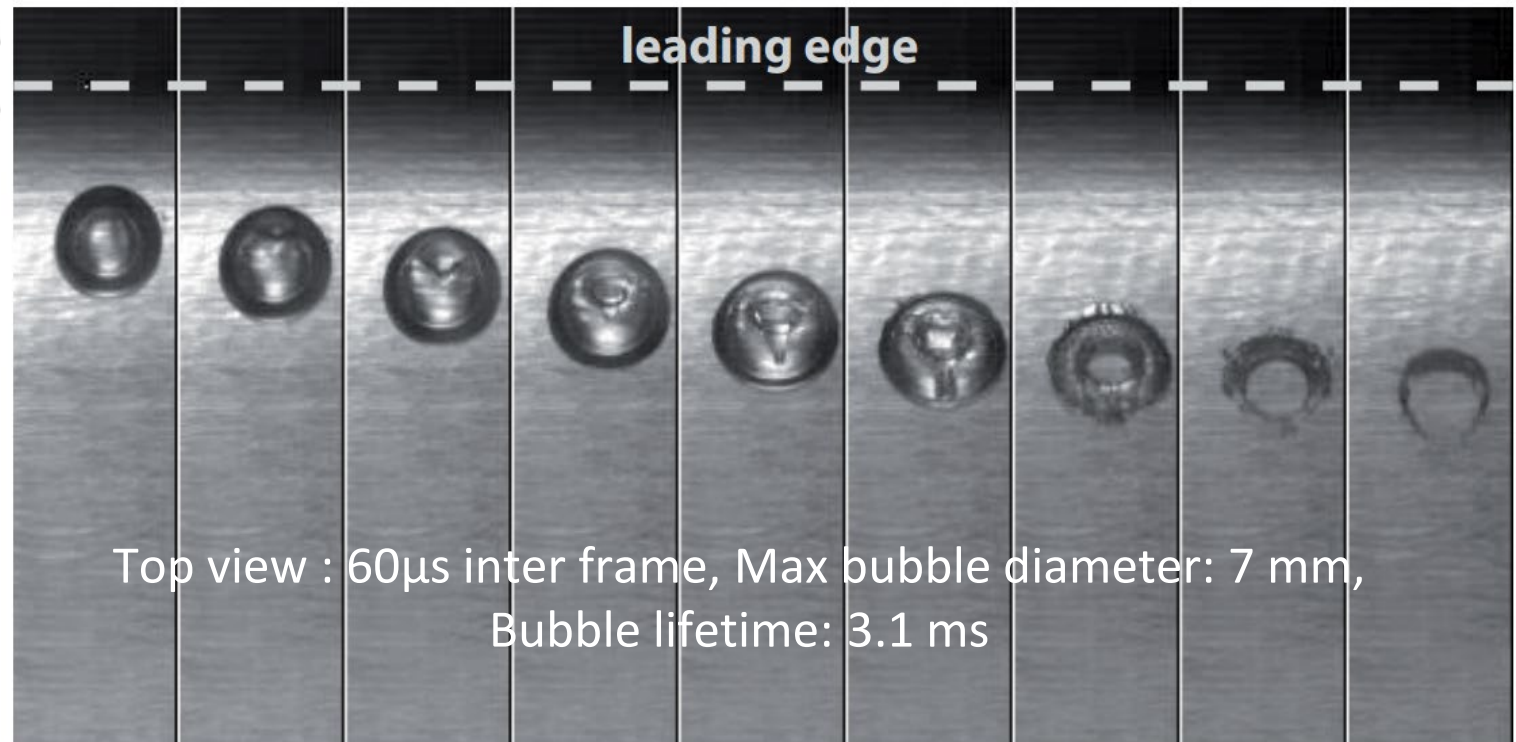
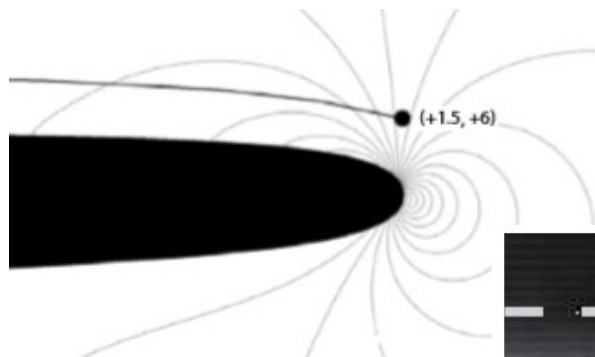


Figure 3.11: Case 1: The bubble traveling on the surface of the hydrofoil. The grey dash lines are the axis perpendicular to the flap top of the bubble and passing through the center. Interframe time: $400 \mu\text{s}$, maximum diameter: 7 mm, bubble lifetime: 3.1 ms.

Collapse of laser induced bubble in a flowing water

Effect of initial position of the bubble on its collapse

- **Case 1: Jet towards the solid surface (top view)**



Collapse of laser induced bubble in a flowing water

Effect of initial position of the bubble on its collapse

- **Case 2: Jet towards the solid surface**

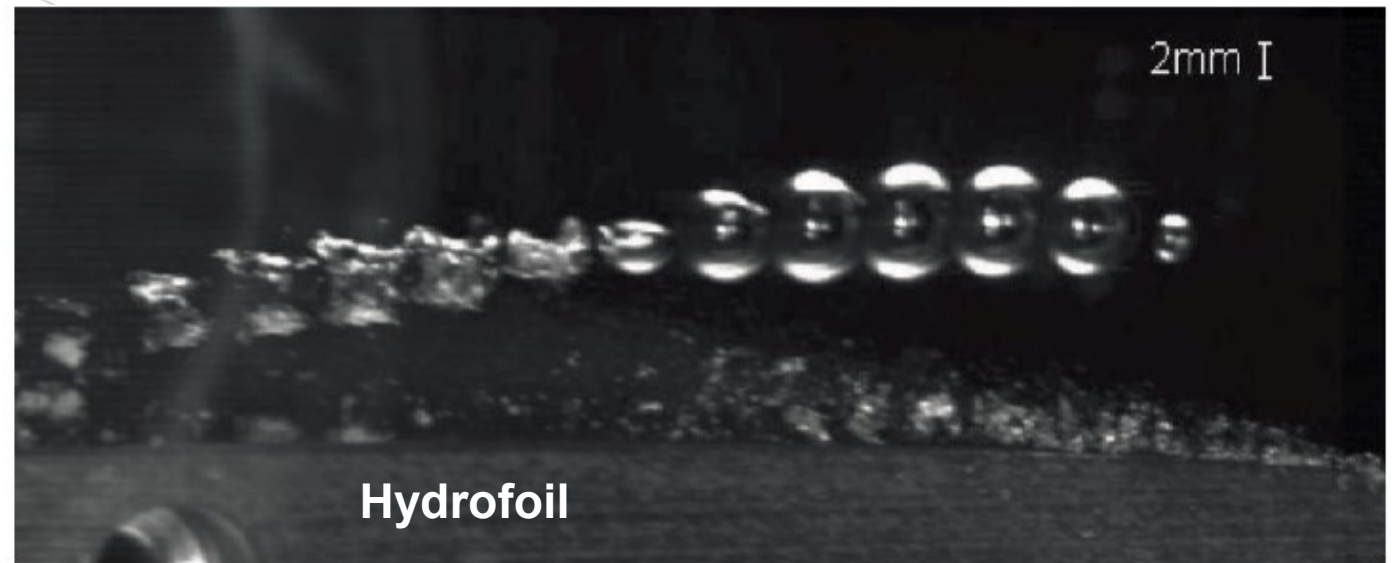
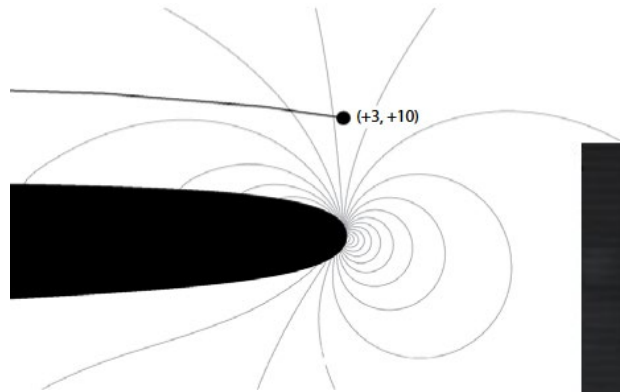


Figure 3.16: Case 2: Collapse of a bubble above the hydrofoil. The collapse is not spherical, it is initiated at the top of the bubble. Interframe time: $200\ \mu\text{s}$, maximum diameter: 7 mm, bubble lifetime: 1.28 ms.

Collapse of laser induced bubble in a flowing water

Effect of initial position of the bubble on its collapse

- **Case 3: Jet towards the pressure gradient**

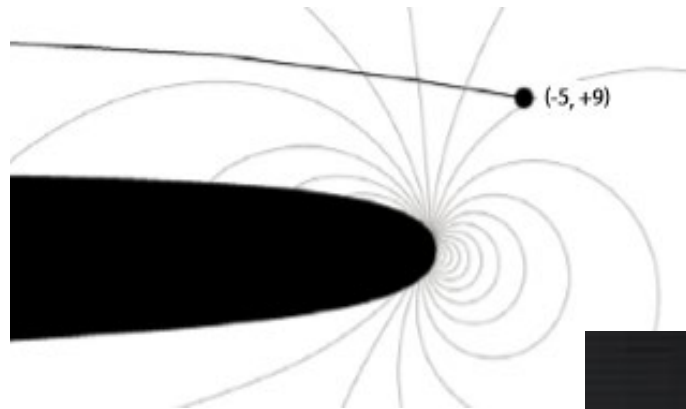


Figure 3.18: Case 3 : Superposition of the high speed images of the collapse of a bubble generated at $(-5, +9)$. The bubble has a microjet not towards the solid surface. Interframe time: $180 \mu\text{s}$, maximum diameter: 5 mm, bubble lifetime: $440 \mu\text{s}$.

Collapse of laser induced bubble in a flowing water

Effect of initial position of the bubble on its collapse

- **Case 4:**
 - **1st collapse** Jet towards the pressure gradient
 - **2nd collapse** towards the solid surface

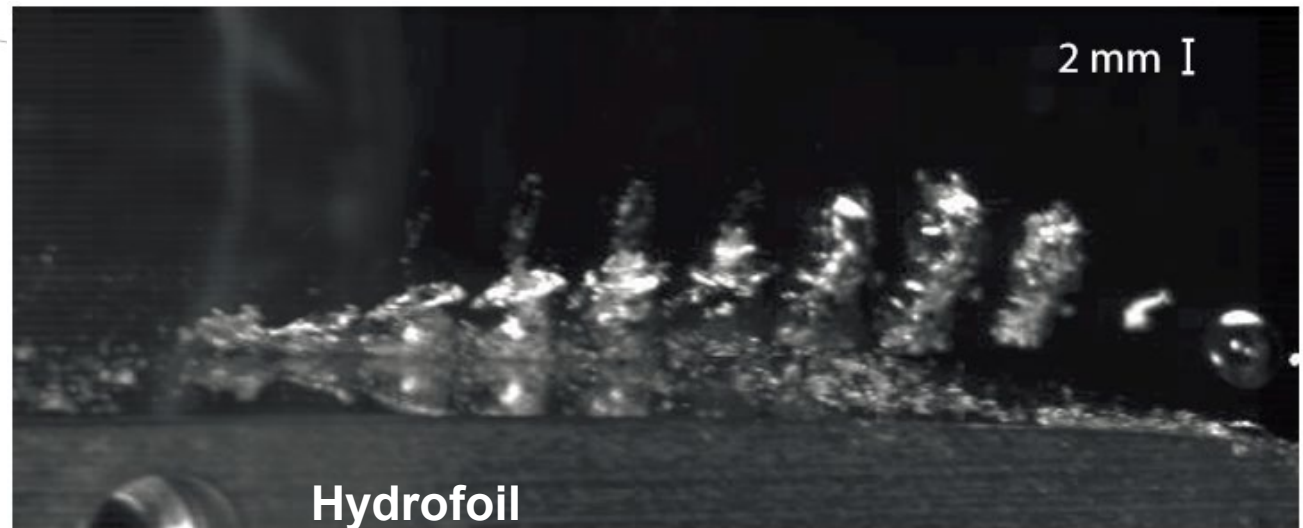
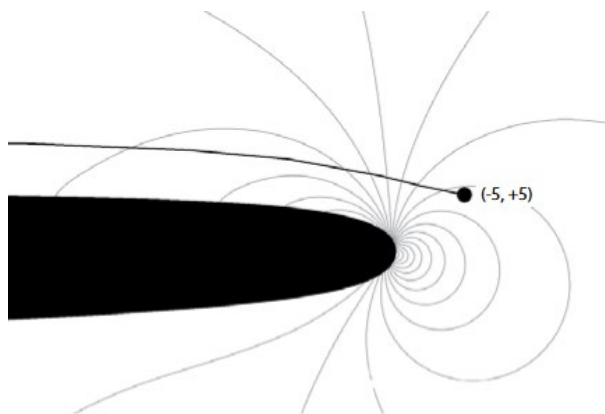


Figure 3.22: Case 4 : Superposition of the high speed images of the collapse of a bubble generated at $(-5, +5)$. The bubble collapses upstream to the leading edge. The rebound, a cloud of microbubbles, travels and collapses on the surface of the hydrofoil. Interframe time: $200 \mu\text{s}$. Initial bubble maximum diameter: 5 mm, lifetime: $380 \mu\text{s}$. Rebound bubble maximum diameter: 6 mm, lifetime: 2 ms.

Collapse of laser induced bubble in a flowing water

Dynamics of rebound cavity

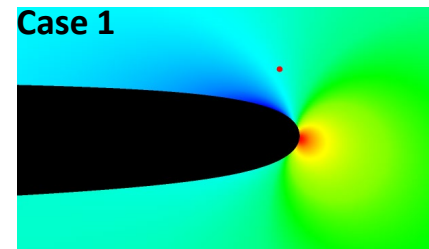
- Rebound cavity is made of a cluster of microbubbles

- The cluster of microbubbles behaves as a single bubble
- Micro jet towards the solid surface
- → Rebound bubble may be more aggressive than the original bubble !

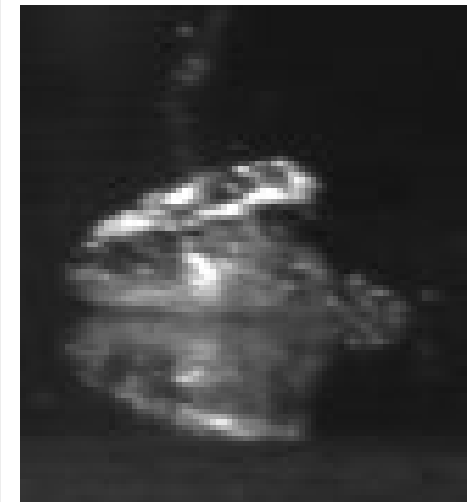
1st collapse



Case 1



2nd collapse



Case 4

