
AEROELASTICITY AND FLUID-STRUCTURE INTERACTION

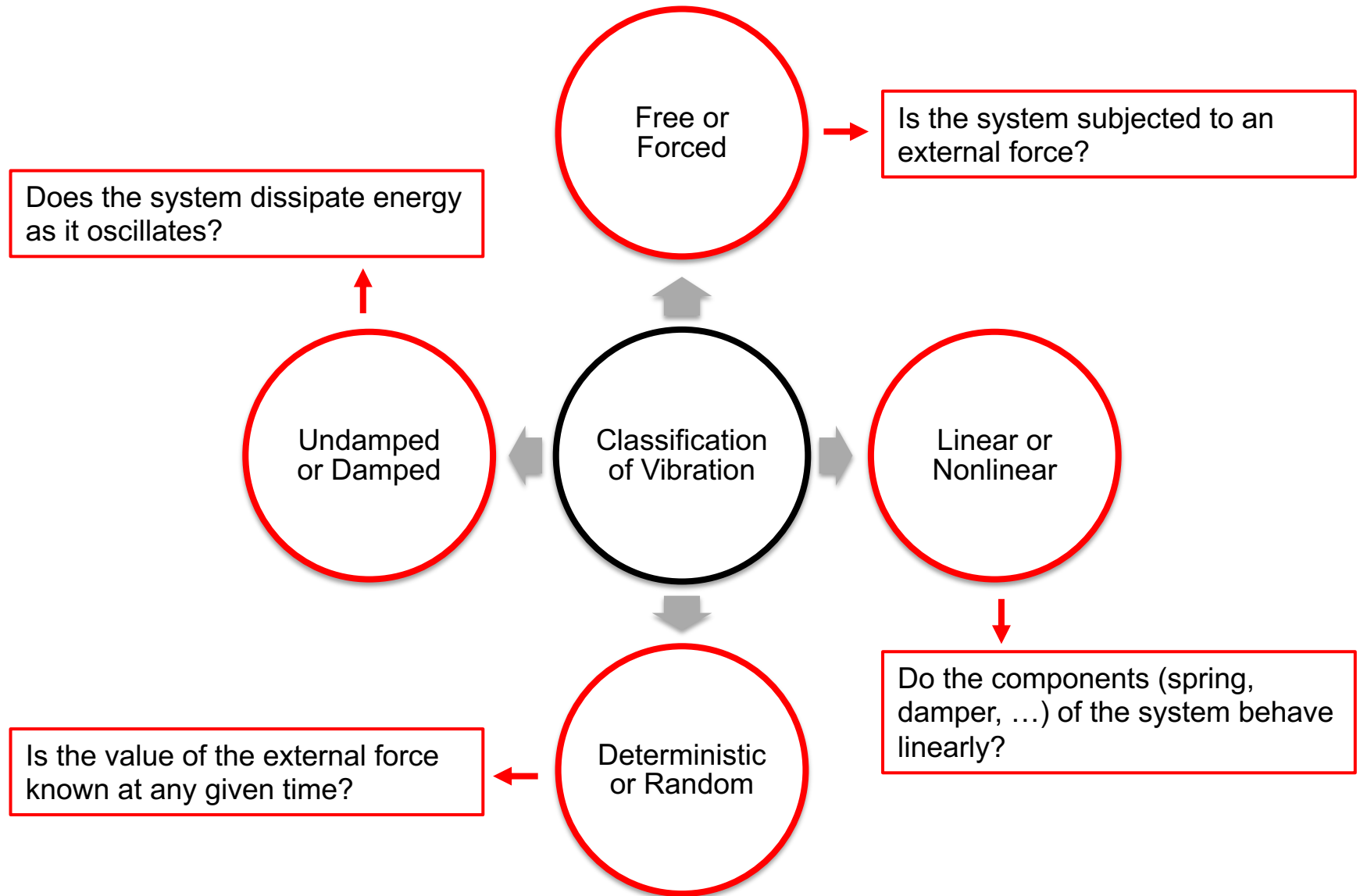
Supplementary Material

Mechanical Vibrations, A Reminder

Basic Concepts of Vibration

- **What are vibrations?**
 - *Any motion that repeats itself after an interval of time (a.k.a. oscillations)*
- **Elementary parts of a vibrating system:**
 - *Means for storing kinetic energy (mass or inertia)*
 - *Means for storing potential energy (spring or elasticity)*
 - *Means for dissipating energy (damper)*
- **Number of degree of freedom:**
 - *The minimum number of coordinates required to completely determine all positions of all parts of the system*
- **Discrete v.s. continuous systems:**
 - *Most systems have an infinite number of degree of freedom → continuous system*
 - *However, continuous systems are often approximated by discrete systems to obtain solutions in a simpler manner → treated as finite lumped masses, springs and dampers*

Basic Concepts of Vibration



Vibrating System Modeling

From a complex vibrating system to a simplified model:

- *Intended to represent all the **important features** of the system in order to derive the equations governing the **system behavior**.*
 - *Elementary models can give a quick insight into the behavior of a model → but the model can be further refined to obtain more accurate results.*
- *Let's start by the basis: the **single-degree-of-freedom** system or SDOF.*

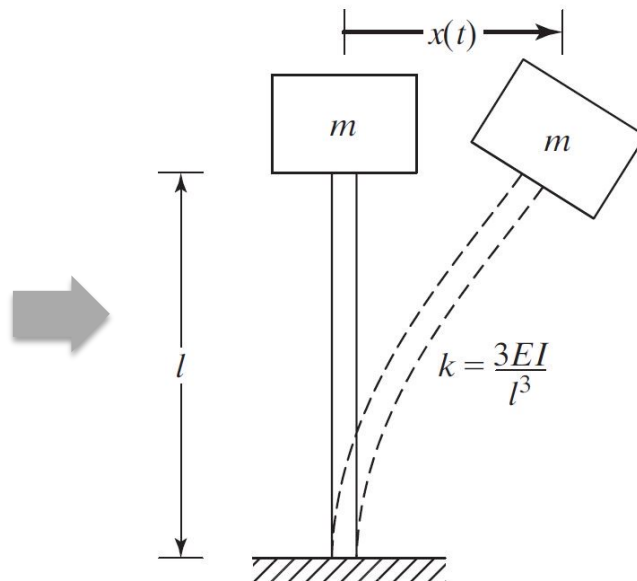
Single Degree of Freedom Systems

Free Vibration of undamped SDOF systems

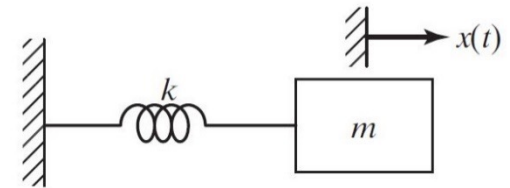
- The system oscillates due to an **initial disturbance** only.
- An example: the transverse vibration of a tall structure:



Actual structure



Simplified model



Equivalent spring-mass system

Single Degree of Freedom Systems

Free Vibration of undamped SDOF systems

- *Deriving the equation of motion (Newton second law of motion, conservation of energy, Lagrange's energy equation, ...):*

$$m\ddot{x} + kx = 0$$

- *The solution can be found assuming: $x(t) = Ce^{st}$, where C and s are complex numbers to be determined. Injecting in the equation of motion:*

$$\begin{aligned} C(ms^2 + k) &= 0 \Leftrightarrow ms^2 + k = 0 \\ \Rightarrow s &= \pm \left(-\frac{k}{m}\right)^{0.5} = \pm i\omega_n \end{aligned}$$

Where $\omega_n = \sqrt{\frac{k}{m}}$ is the natural frequency of the system.

→ Hence the general solution reads: $x(t) = C_1 e^{i\omega_n t} + C_2 e^{-i\omega_n t}$

Single Degree of Freedom Systems

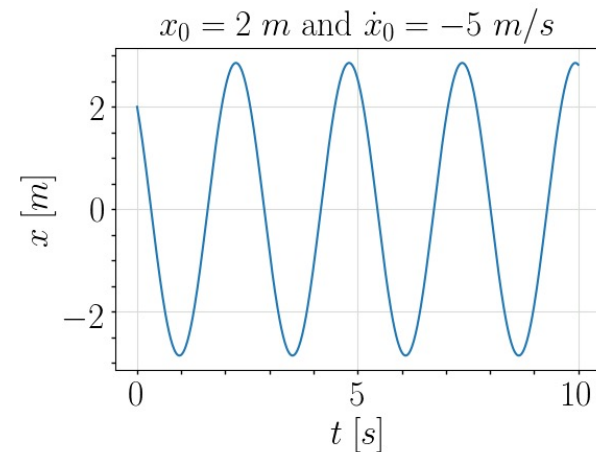
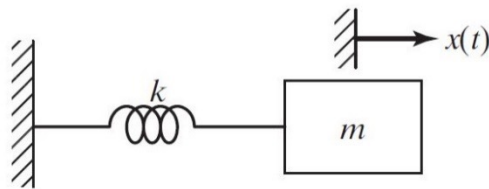
Free Vibration of undamped SDOF systems

- Considering the real part of the response $x(t)$ only, it can be shown that the solution is given by:

$$x(t) = A_1 \cos(\omega_n t) + A_2 \sin(\omega_n t) = C \cos(\omega_n t - \phi)$$

Where A_1 , A_2 , C and ϕ are real constants determined by the initial conditions.

$$\Rightarrow x(t) = x_0 \cos(\omega_n t) + \frac{\dot{x}_0}{\omega_n} \sin(\omega_n t) = \sqrt{x_0^2 + \left(\frac{\dot{x}_0}{\omega_n}\right)^2} \cos\left(\omega_n t - \tan^{-1}\left(\frac{\dot{x}_0}{x_0 \omega_n}\right)\right)$$



Single Degree of Freedom Systems

Adding damping to freely vibrating SDOF systems

- The **viscous damping force** is proportional to the **velocity** of the system and opposes the motion. Hence, the equation of motion reads:

$$m\ddot{x} + c\dot{x} + kx = 0$$

- The solution can be found assuming: $x(t) = Ce^{st}$, where C and s are to be determined

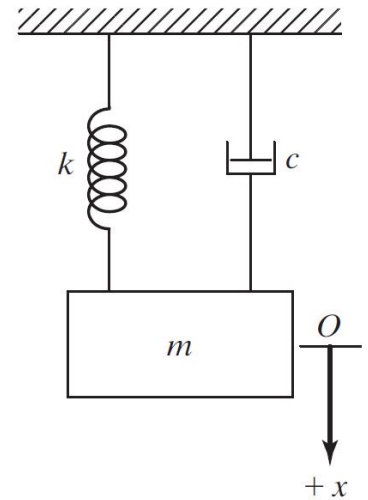
$$C(ms^2 + cs + k) = 0 \Leftrightarrow ms^2 + cs + k = 0$$

$$\Rightarrow s_{1,2} = -\frac{c}{2m} \pm \sqrt{\left(\frac{c}{2m}\right)^2 - \frac{k}{m}}$$

→ Hence the general solution reads: $x(t) = C_1 e^{s_1 t} + C_2 e^{s_2 t}$

- A couple notations: critical damping, c_c , and damping ratio ζ

$$c_c = 2m\omega_n \quad \text{and} \quad \zeta = \frac{c}{c_c}$$



Single Degree of Freedom Systems

Adding damping to freely vibrating SDOF systems

- The roots and the solution can thus be rewritten as*

$$s_{1,2} = \left(-\zeta \pm \sqrt{\zeta^2 - 1} \right) \omega_n$$

$$x(t) = C_1 e^{(-\zeta + \sqrt{\zeta^2 - 1})\omega_n t} + C_2 e^{(-\zeta - \sqrt{\zeta^2 - 1})\omega_n t}$$

→ *The behavior of the system hence depends upon the magnitude of the damping.*

- Three damping conditions exist (excluding undamped vibrations, $\zeta = 0$)*

- Underdamped ($\zeta < 1$ or $c < c_c$) → s_1 and s_2 are complex*

$$s_{1,2} = \left(-\zeta \pm i\sqrt{1 - \zeta^2} \right) \omega_n$$

- Critically damped ($\zeta = 1$ or $c = c_c$) → s_1 and s_2 are equal*

$$s_1 = s_2 = -\omega_n$$

- Overdamped ($\zeta > 1$ or $c > c_c$) → s_1 and s_2 are real*

$$s_{1,2} = \left(-\zeta \pm \sqrt{\zeta^2 - 1} \right) \omega_n < 0$$

Single Degree of Freedom Systems

Adding damping to freely vibrating SDOF systems.

- Three damping conditions exist (excluding undamped vibrations, $\zeta = 0$)

- Underdamped:

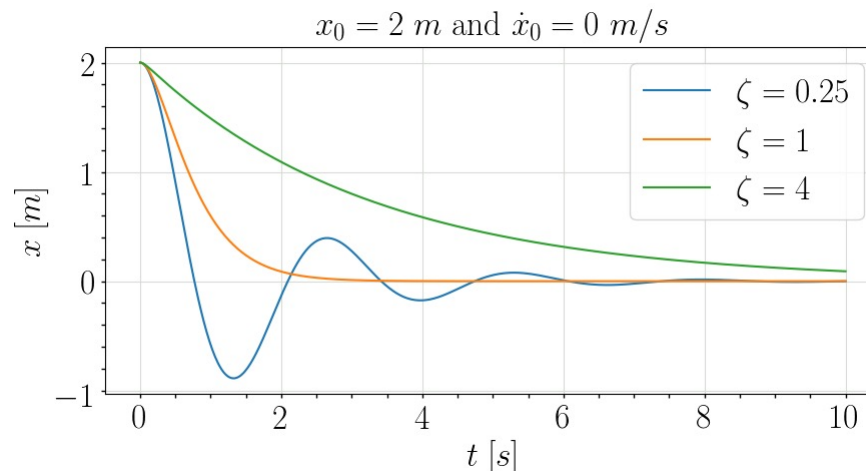
$$x(t) = e^{-\zeta\omega_n t} \left(x_0 \cos(\omega_d t) + \frac{\dot{x}_0 + \zeta\omega_n x_0}{\omega_d} \sin(\omega_d t) \right), \quad \omega_d = \omega_n \sqrt{1 - \zeta^2}$$

- Critically damped:

$$x(t) = e^{-\omega_n t} [x_0 + (\dot{x}_0 + \omega_n x_0)t]$$

- Overdamped:

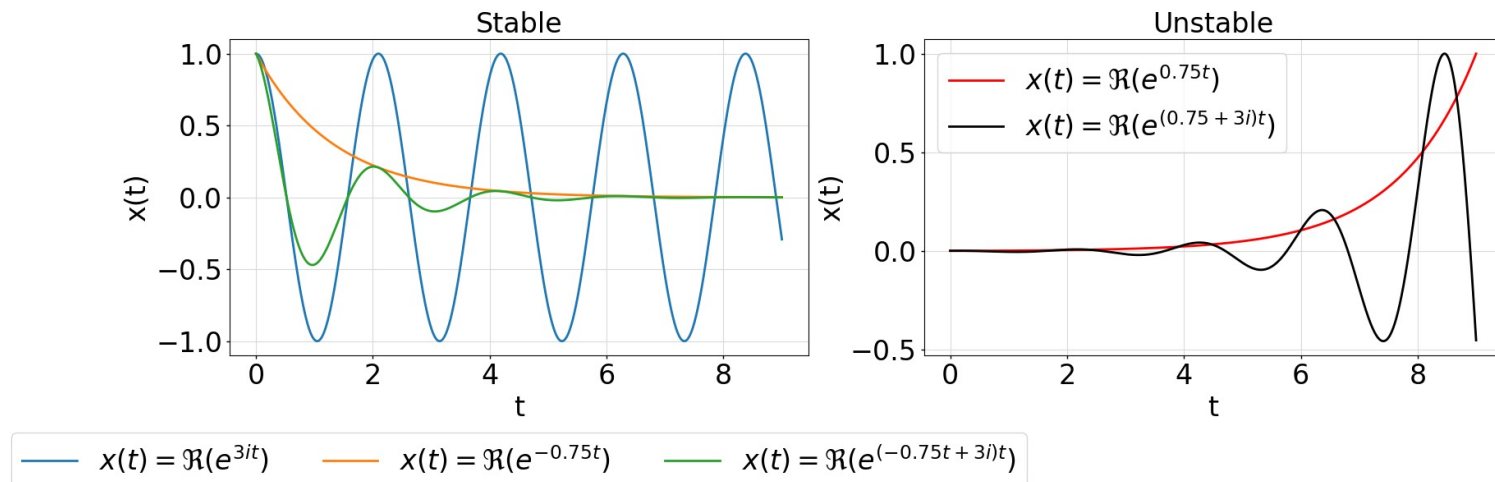
$$x(t) = e^{-\zeta\omega_n t} \left(x_0 \cosh(\omega_* t) + \frac{\dot{x}_0 + \zeta\omega_n x_0}{\omega_*} \sinh(\omega_* t) \right), \quad \omega_* = \omega_n \sqrt{\zeta^2 - 1}$$



Single Degree of Freedom Systems

A note on the nature of the solution

- The solution $x(t) = C_1 e^{s_1 t} + C_2 e^{s_2 t}$ presents the following properties:
 - If $s_{1,2}$ has a **zero or negative real value**, the corresponding response will be **stable**.
 - If $s_{1,2}$ has a **zero imaginary value**, the corresponding response will not oscillate.
 - If $s_{1,2}$ has a **positive real value**, the corresponding response will grow exponentially and will be **unstable**:
 - Divergence instability $\rightarrow s_{1,2}$ has a zero imaginary
 - Flutter instability $\rightarrow s_{1,2}$ has a non-zero imaginary

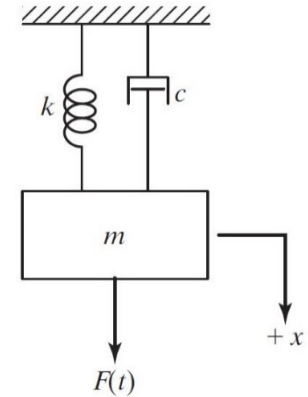


Single Degree of Freedom Systems

SDOF system under harmonic forcing: $F(t) = F_0 e^{i\omega t}$

- The equation of motion reads:

$$m\ddot{x} + c\dot{x} + kx = F_0 e^{i\omega t}$$



→ Non-homogeneous equation whose solution is given by the sum of the homogenous (transient) solution, $x_h(t)$, and the particular (steady-state) solution, $x_p(t)$.

- By assuming a particular solution of the form $x_p(t) = X e^{i(\omega t - \phi)} = X e^{-i\phi} e^{i\omega t}$, we get for the amplitude:

$$X = \frac{F_0}{[(k - m\omega^2)^2 + c^2\omega^2]^{0.5}} \quad \Leftrightarrow \quad \frac{X}{\delta_{st}} = \frac{1}{[(1 - r^2)^2 + (2r\zeta)^2]^{0.5}}$$

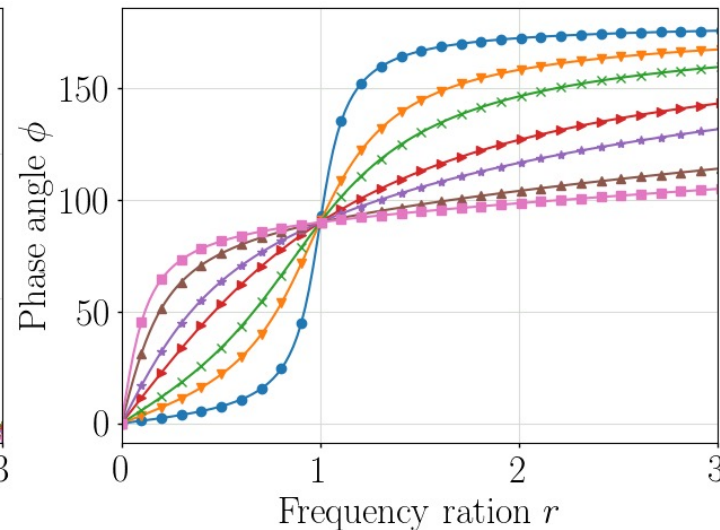
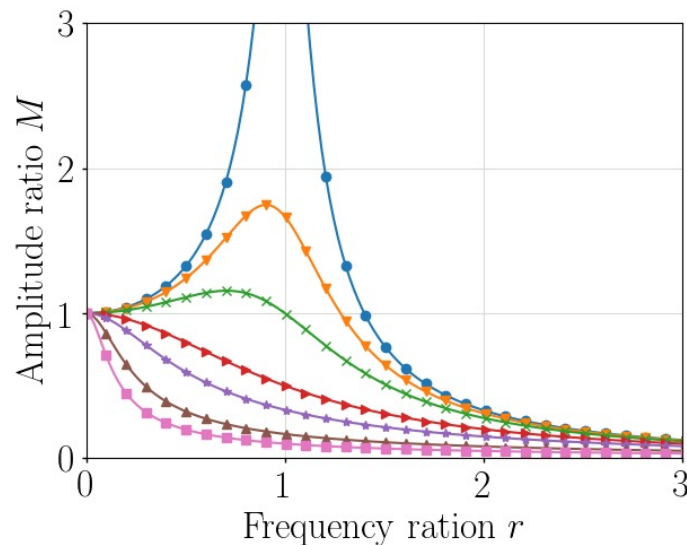
With:

- $\zeta = c/c_c$ is the damping ratio
- $\delta_{st} = F_0/k$ is the deflection under the static force F_0
- $r = \omega/\omega_n$ is the frequency ratio
- $\phi = \tan^{-1} \left(\frac{2\zeta r}{1-r^2} \right)$ is the phase angle

Single Degree of Freedom Systems

SDOF system under harmonic forcing: $F(t) = F_0 e^{i\omega t}$

- Variations of the amplitude ratio, $\frac{X}{\delta_{st}}$, and phase angle as a function of the frequency ratio and damping ratio.*

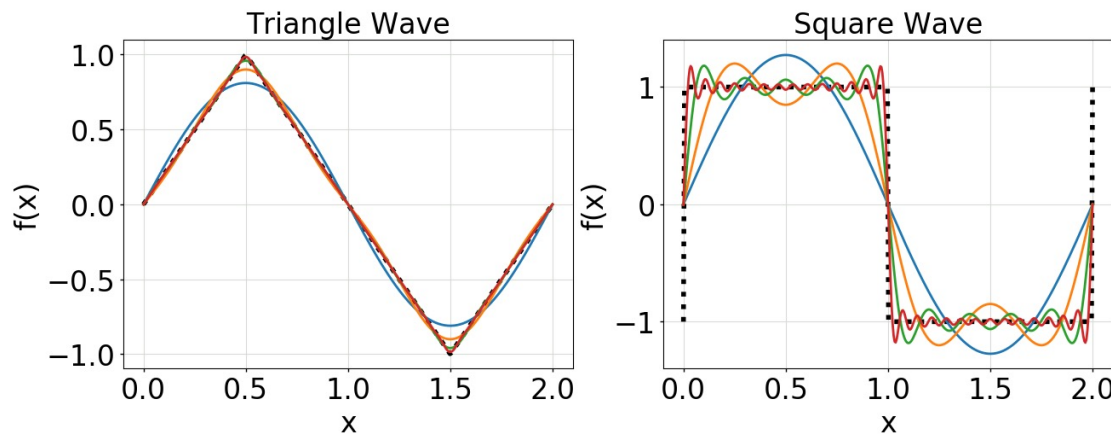


Single Degree of Freedom Systems

SDOF system under general forcing: a very short overview

- **Periodic but non-harmonic forcing:**

- *The periodic forcing can be approximated by a sum of harmonic functions (Fourier series)*
- *The system response is found by superposing the responses of the individual harmonic forcing functions → superposition principle for linear systems*



Notice the Gibbs phenomenon at the jump discontinuity of the periodic function

- **Non-periodic forcing:**

- *Analytical methods: convolution integrals, Laplace transform, ... → may become tedious or even impossible to solve depending on the forcing*
- *Numerical methods: Runge-Kutta methods, ...*

Multi Degree of Freedom Systems

Multi-degree of freedom systems: MDOF

Most practical engineering systems are continuous and have an infinite number of degree of freedom.

However:

The vibration analysis of continuous systems requires the solution of partial differential equations → difficult and/or time consuming.

A solution:

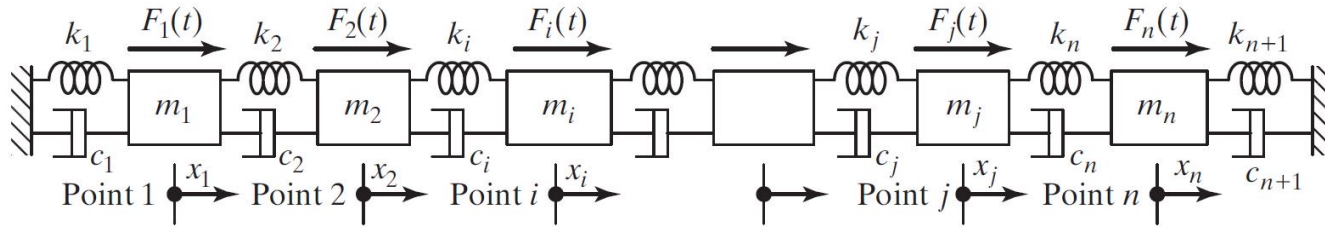
The analysis of multi-degree of freedom system only requires the solution of a set of ordinary differential equations → easier

So:

Continuous systems are often approximated as multi-degree of freedom systems.

Multi Degree of Freedom Systems

Multi-degree of freedom systems: MDOF



- The equation of motion reads for the mass m_i reads:

$$m_i \ddot{x}_i = \sum_j F_{ij}$$

$$m_i \ddot{x}_i = -k_i(x_i - x_{i-1}) + k_{i+1}(x_{i+1} - x_i) - c_i(\dot{x}_i - \dot{x}_{i-1}) + c_{i+1}(\dot{x}_{i+1} - \dot{x}_i) + F_i$$

- Expressing the equation of motion of the complete system in matrix form yields a system of ordinary differential equations

$$[M]\ddot{x} + [C]\dot{x} + [K]x = F$$

$[M]$, $[C]$ and $[K]$ are the mass, damping and stiffness matrices, respectively. They are symmetric.

Multi Degree of Freedom Systems

Multi-degree of freedom systems: MDOF

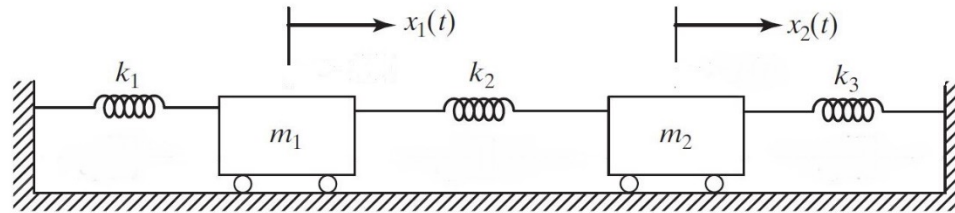
- *Mass, damping and stiffness matrices:*
 - *If the **mass matrix** is not diagonal $\rightarrow$ the system is said to have **mass (or inertia) coupling***
 - *If the **damping matrix** is not diagonal $\rightarrow$ the system is said to have **velocity coupling***
 - *If the **stiffness matrix** is not diagonal $\rightarrow$ the system is said to have **static coupling***
- $\rightarrow$ *If the system of equation is coupled, then the equations have to be solved simultaneously.*

$$[M]\ddot{x} + [C]\dot{x} + [K]x = F$$

Two Degree of Freedom Systems

2 DOF system: undamped free-vibration

- Consider the following system:



- The equation of motion reads:

$$\begin{bmatrix} m_1 & 0 \\ 0 & m_2 \end{bmatrix} \begin{pmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{pmatrix} + \begin{bmatrix} k_1 + k_2 & -k_2 \\ -k_2 & k_2 + k_3 \end{bmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

- As for the SDOF case, we can assume the masses undergo harmonic motion. The solution therefore takes the form: $x_j(t) = X_j e^{i\omega t}$. Substituting in the equation of motion:

$$\begin{bmatrix} -\omega^2 m_1 + k_1 + k_2 & -k_2 \\ -k_2 & -\omega^2 m_2 + k_2 + k_3 \end{bmatrix} \begin{pmatrix} X_1 \\ X_2 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

Two Degree of Freedom Systems

2 DOF system: undamped free-vibration

- For the system to have a non-trivial solution, the determinant of the coefficients of X_1 and X_2 must be zero:

$$\det \begin{bmatrix} -\omega^2 m_1 + k_1 + k_2 & -k_2 \\ -k_2 & -\omega^2 m_2 + k_2 + k_3 \end{bmatrix} = 0$$

Or

$$(m_1 m_2) \omega^4 - ((k_1 + k_2) m_2 + (k_2 + k_3) m_1) \omega^2 + ((k_1 + k_2)(k_2 + k_3) - k_2^2) = 0$$
$$\Leftrightarrow A \omega^4 + B \omega^2 + C = 0$$

- The solution to the above equation yields the natural frequencies of the system: ω_1^2 and ω_2^2 .

$$\omega_{1,2}^2 = \frac{-B \pm \sqrt{B^2 - 4AC}}{2A}$$

- The values of X_1 and X_2 remain to be determined, but it is out of the scope of this course.

Two Degree of Freedom Systems

2 DOF system: the general case → damped with harmonic forcing

- The general equation of motion reads:*

$$\begin{bmatrix} m_{11} & m_{12} \\ m_{12} & m_{22} \end{bmatrix} \begin{pmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{pmatrix} + \begin{bmatrix} c_{11} & c_{12} \\ c_{12} & c_{22} \end{bmatrix} \begin{pmatrix} \dot{x}_1 \\ \dot{x}_2 \end{pmatrix} + \begin{bmatrix} k_{11} & k_{12} \\ k_{12} & k_{22} \end{bmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \begin{pmatrix} F_1 \\ F_2 \end{pmatrix}$$

- The forcing is assumed harmonic: $F_j = F_{0j}e^{i\omega t}$. The steady-state solution therefore takes the form: $x_j(t) = X_j e^{i\omega t}$. Substituting in the equation of motion:*

$$\begin{bmatrix} (-\omega^2 m_{11} + i\omega c_{11} + k_{11}) & (-\omega^2 m_{12} + i\omega c_{12} + k_{12}) \\ (-\omega^2 m_{12} + i\omega c_{12} + k_{12}) & (-\omega^2 m_{22} + i\omega c_{22} + k_{22}) \end{bmatrix} \begin{pmatrix} X_1 \\ X_2 \end{pmatrix} = \begin{pmatrix} F_{01} \\ F_{02} \end{pmatrix}$$

$$\begin{bmatrix} Z_{11}(i\omega) & Z_{12}(i\omega) \\ Z_{12}(i\omega) & Z_{22}(i\omega) \end{bmatrix} \begin{pmatrix} X_1 \\ X_2 \end{pmatrix} = \begin{pmatrix} F_{01} \\ F_{02} \end{pmatrix}$$

- $Z_{mn}(i\omega)$ is the mechanical impedance.*

Two Degree of Freedom Systems

2 DOF system: the general case $\rightarrow$ damped with harmonic forcing

- The equation can be solved as:*

$$\begin{pmatrix} X_1 \\ X_2 \end{pmatrix} = \begin{bmatrix} Z_{11}(i\omega) & Z_{12}(i\omega) \\ Z_{12}(i\omega) & Z_{22}(i\omega) \end{bmatrix}^{-1} \begin{pmatrix} F_{01} \\ F_{02} \end{pmatrix}$$

Which leads to:

$$X_1(i\omega) = \frac{Z_{22}(i\omega)F_{01} - Z_{12}(i\omega)F_{02}}{Z_{11}(i\omega)Z_{22}(i\omega) - Z_{12}^2(i\omega)}$$

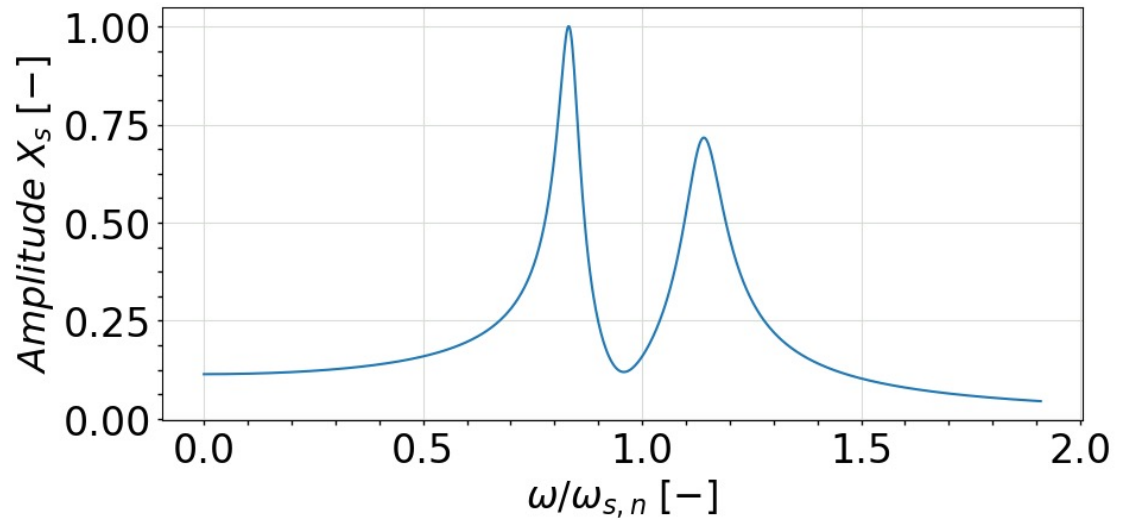
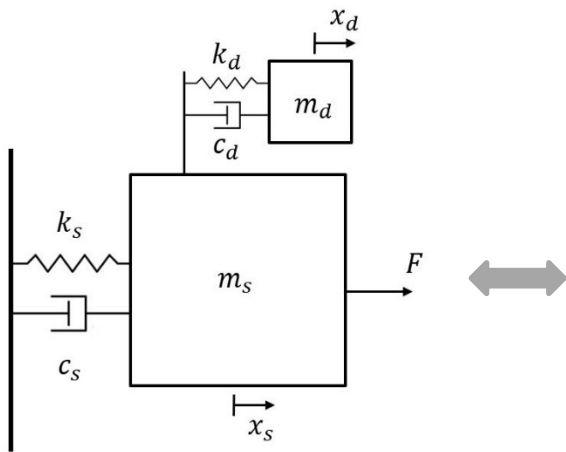
$$X_2(i\omega) = \frac{Z_{11}(i\omega)F_{02} - Z_{12}(i\omega)F_{01}}{Z_{11}(i\omega)Z_{22}(i\omega) - Z_{12}^2(i\omega)}$$

- The complete steady state solution can be found by injecting X_1 and X_2 into $x_j(t) = X_j e^{i\omega t}$*
- The frequency response curves can be found by considering the moduli of $X_1(i\omega)$ and $X_2(i\omega)$.*

Two Degree of Freedom Systems

2 DOF system: the general case $\rightarrow$ damped with harmonic forcing

- Frequency response curves of a 2DOF system: $F = F_0 e^{i\omega t}$



Numerical Approach

MDOF systems: let's simplify and consider a 2 DOF system

An alternative to analytically solving the systems: numerically solving the systems

→ *The state space representation:*

- *Converts the system into a 1st order ODE system*
- *Formulation that allows for an easier numerical solution*

$$\begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \\ \ddot{x}_1 \\ \ddot{x}_2 \end{bmatrix} = \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ [-\mathbf{M}^{-1}\mathbf{K}] & [-\mathbf{M}^{-1}\mathbf{C}] \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ \dot{x}_1 \\ \dot{x}_2 \end{bmatrix} + \begin{bmatrix} 0 & 0 \\ 0 & 0 \\ [\mathbf{M}^{-1}] \end{bmatrix} \begin{bmatrix} F_1 \\ F_2 \end{bmatrix}$$