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# AEROELASTICITY AND FLUID-STRUCTURE INTERACTION

Complement  
Sloshing Dynamics  
*Equivalent Mechanical Model*

# Sloshing Dynamics – Equivalent Mechanical Model

*In certain situations, sloshing can influence the behavior of the structure surrounding the moving liquid  
→ there is a need to incorporate its effects into structural models.*

1.



2.



1. *Coffee mug*
2. *Cruise ship swimming pool.*
3. *Rocket liquid O<sub>2</sub> fuel tank*

- *Find the first asymmetric natural frequency for case 1 and 2.*
- *How to model the fluid-structure interactions?*

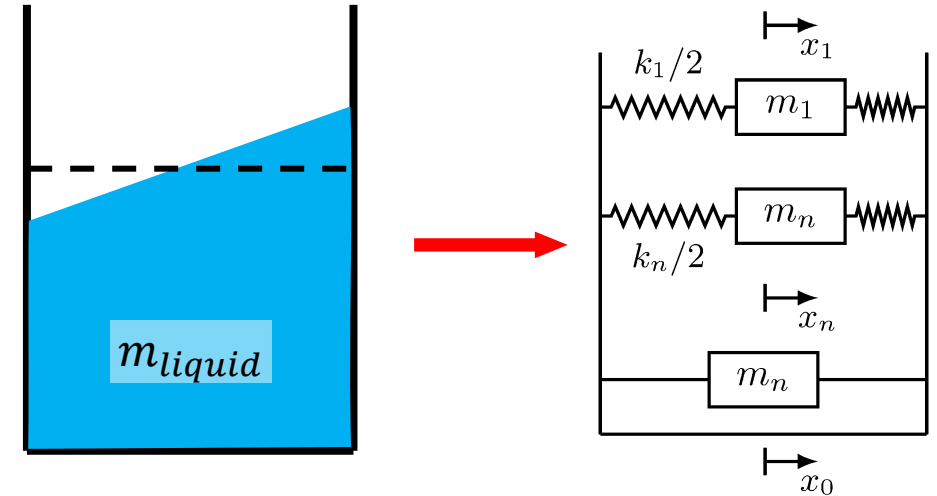
3.



# Sloshing Dynamics – Equivalent Mechanical Model

*To incorporate the effects of sloshing in a dynamical system, it is often convenient to conceptually replace the liquid by an equivalent linear mechanical system:*

- We will use a linear spring-mass model*  
*Each of the  $n$  spring-mass corresponds to one of the infinite sloshing modes.*



*Main benefits:*

- Easier to include in a system*
- Reduced computation cost*
- Fluid damping can be easily incorporated in the model by adding linear dashpots*

- How to define the parameters  $m_n$  and  $k_n$ ?*

# Sloshing Dynamics – Equivalent Mechanical Model

## Defining the model parameters for horizontal motion

- **Static properties:**

- The sum of all the masses must be the same as the liquid mass

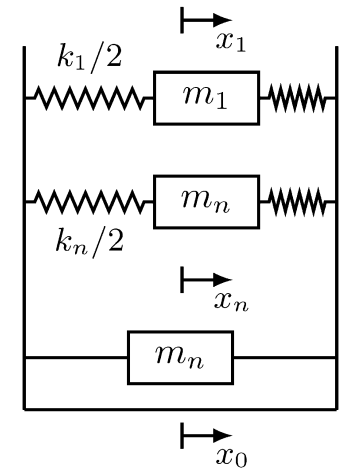
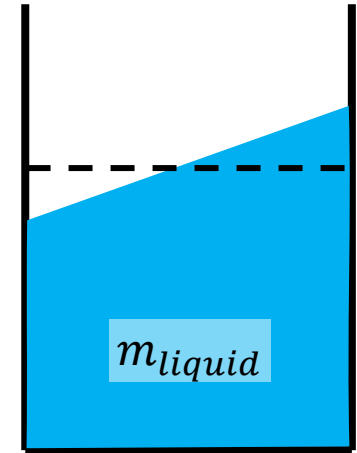
$$m_{liqu} = m_0 + \sum_{n=1}^{\infty} m_n$$

- **Dynamic properties:**

- The natural frequencies must duplicate the ones of the liquid

$$\frac{k_n}{m_n} = \omega_n^2$$

where  $\omega_n$  is the natural frequency of the  $n^{th}$  sloshing mode.  
(known from potential theory)



# Sloshing Dynamics – Equivalent Mechanical Model

## Defining the model parameters for horizontal motion

- **Dynamic properties:**
  - The force components exerted on the tank under certain excitation must be equivalent to the one produced by the actual system (= the ones derived from the potential flow theory)

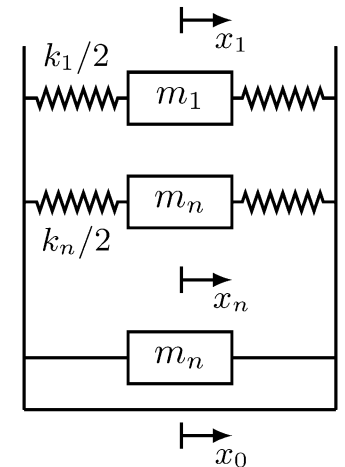
$$-F = m_0 \ddot{x}_0 + \sum_{n=1}^{\infty} m_n (\ddot{x}_n + \ddot{x}_0)$$

- In the case of pure translational excitation of the tank:  $x_0(t) = X_0 \sin(\Omega t)$

From the steady state solution of the undamped equation

$$m_n(\ddot{x}_n + \ddot{x}_0) + k \ddot{x}_n = 0 \text{ we obtain } x_n = \frac{\Omega^2}{\omega_n^2 - \Omega^2} X_0 \sin(\Omega t)$$

$$\begin{aligned} F &= X_0 \Omega^2 \sin(\Omega t) \left[ m_0 + \sum_{n=1}^{\infty} m_n \left( \frac{\Omega^2}{\omega_n^2 - \Omega^2} + 1 \right) \right] \\ &= m_{liqu} X_0 \Omega^2 \sin(\Omega t) \left[ 1 + \sum_{n=1}^{\infty} \frac{m_n}{m_{liqu}} \left( \frac{\Omega^2}{\omega_n^2 - \Omega^2} \right) \right] \end{aligned}$$

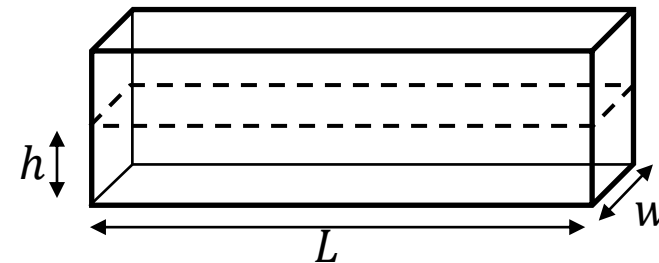


# Sloshing Dynamics – Equivalent Mechanical Model

## Defining the model parameters for horizontal motion

- The model parameters depend on the liquid and on the tank shape (not demonstrated here)  
*Rectangular tank with motion along the  $a$ -direction*

- $m_{liqu} = \rho L b h$
- $\frac{m_n}{m_{liqu}} = 8 \left( \frac{L}{h} \right) \frac{\tanh((2n-1)\pi \frac{h}{L})}{(2n-1)^3 \pi^3}$
- $\frac{k_n}{m_{liqu}} = 8 \left( \frac{g}{h} \right) \frac{\tanh^2((2n-1)\pi \frac{h}{L})}{(2n-1)^2 \pi^2}$



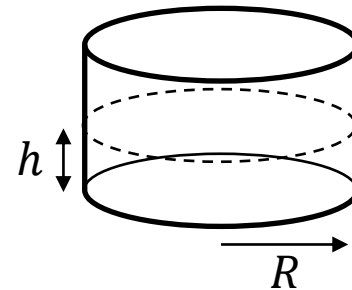
- One may observe that the masses rapidly decrease for all modes exceeding the first one.

# Sloshing Dynamics – Equivalent Mechanical Model

## Defining the model parameters for horizontal motion

- The model parameters depend on the liquid and on the tank shape (not demonstrated here)  
*Cylindrical tank with motion along the R-direction (linear sloshing)*

- $m_{liqu} = \rho \pi R^2 h$
- $\frac{m_n}{m_{liqu}} = \left( \frac{2R}{\varepsilon_{1n} h (\varepsilon_{1n}^2 - 1)} \right) \tanh \left( \frac{\varepsilon_{1n} h}{R} \right)$
- $\frac{k_n}{m_{liqu}} = \left( \frac{2g}{h (\varepsilon_{1n}^2 - 1)} \right) \tanh^2 \left( \frac{\varepsilon_{1n} h}{R} \right)$



- $\varepsilon_{1n}$  corresponds to the roots of the derivatives of the Bessel function of the first kind ( $\varepsilon_{11} = 1.841, \varepsilon_{12} = 5.331, \varepsilon_{13} = 8.536, \dots$ )
- Since  $\varepsilon_{1n}$  increases with  $n$ , the size of the masses also decreases for all modes exceeding the first.

# Sloshing Dynamics – Equivalent Mechanical Model

- *From potential theory we can model the free-response motion.*

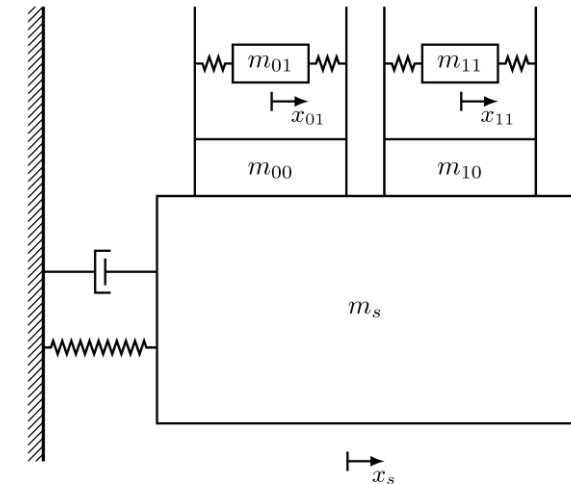
$$\ddot{x} + \omega_1^2 x = 0$$

$$\omega_1 \approx \frac{2\pi g}{L}, 1.841 \frac{g}{R} (h \gg 2L, \text{ respectively Rect. Cyl. })$$

→ *We can now model dynamic interaction with any force as a function of time*

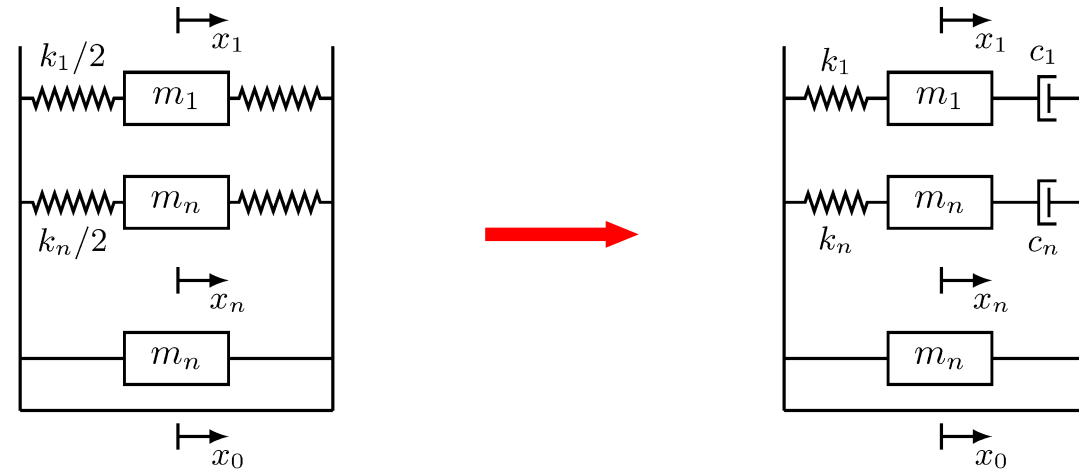
$$\ddot{x} + \frac{k_1}{m_1} x = F(t)$$

→ *Or model coupled dynamical systems*



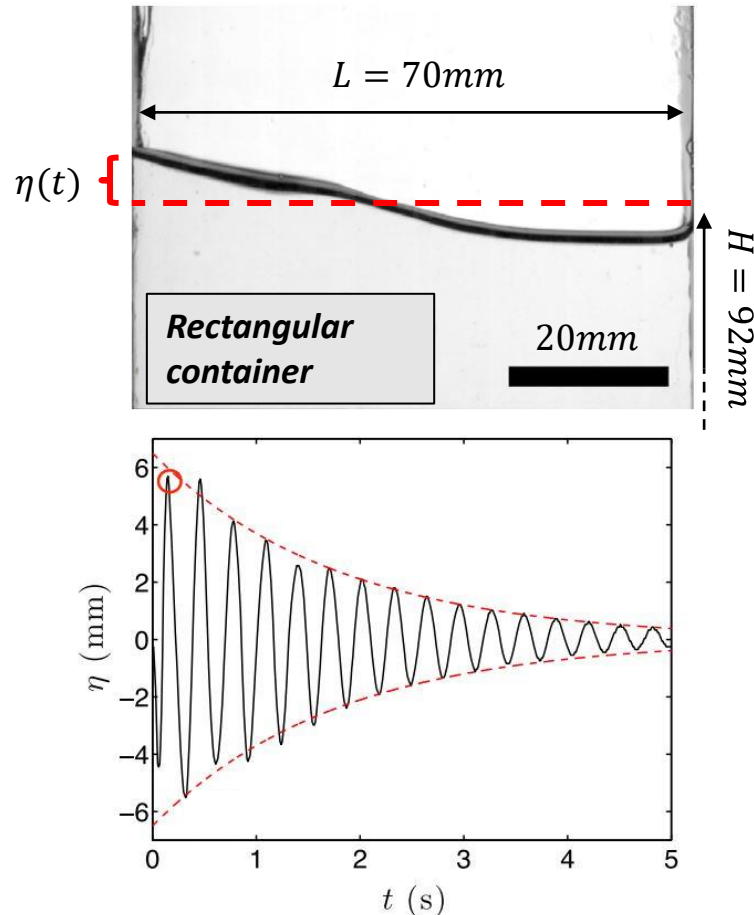
# Sloshing Dynamics – Equivalent Mechanical Model

- *Isn't there something missing to model our experimental observations?*  
→ *Viscous damping !*



# Sloshing Dynamics – Equivalent Mechanical Model

## Modeling the liquid damping – linear dashpots



Damping of liquid sloshing by foams: from everyday observations to liquid transport, Capello et al., 2015

- Experiments show that the free-surface elevation, after an impulse motion, follows a damped harmonic motion:

$$\eta(t) = \eta_0 e^{(-t/\tau)} \cos(2\pi f t)$$

With  $f_1 = \frac{\omega_1}{2\pi} = \frac{1}{2\pi} \sqrt{g k_1 \tanh(k_1 H)} = 3.34 \text{ Hz}$

$\rightarrow f_{exp} = 3.22 \pm 0.11 \text{ Hz}$

- In the linear framework, the effect of damping can thus be modeled by a set of linear dashpots.

# Sloshing Dynamics – Equivalent Mechanical Model

## Modeling the liquid damping – linear dashpots

- *In real liquids, energy dissipation occurs at the tank walls and free surface due to the viscous boundary layer and within the liquid because of viscous stresses.*
- *For small tanks, the boundary layer dissipation dominates, while for large tanks, the dissipation in the liquid interior may be the larger contribution.*
- *Most results for the damping ratio have been obtained experimentally (first sloshing mode):*
  - *In cylindrical tanks (Stephens et al. 1962):*

$$\zeta_1 = 0.83 \sqrt{\frac{\nu}{g^{1/2} R^{3/2}}} \left[ \tanh\left(\frac{\varepsilon_{11} h}{R}\right) \left( 1 + 2 \frac{1 - \frac{h}{R}}{\cosh\left(\frac{\varepsilon_{11} h}{R}\right)} \right) \right]$$

- *In rectangular tanks (Sun, 1991):*

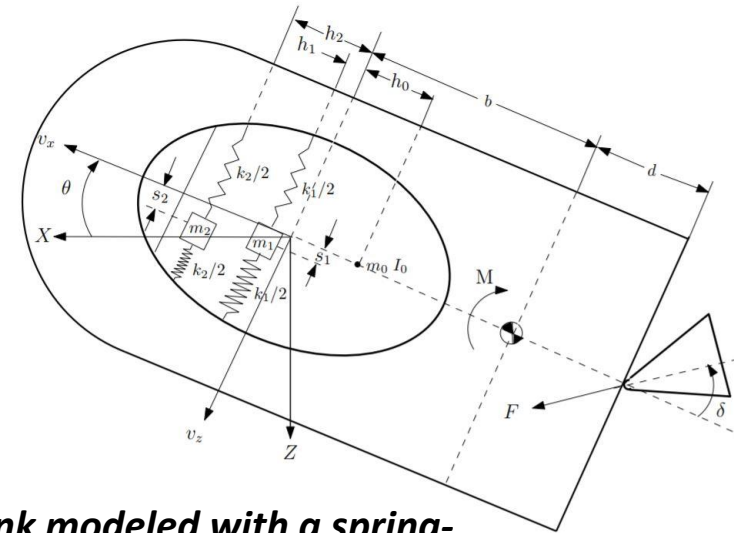
$$\zeta_1 = \frac{1}{2h} \sqrt{\frac{2\nu}{\omega_1}} \left( 1 + \frac{h}{w} \right) \quad \nu \text{ is the liquid kinematic viscosity}$$

# Sloshing Dynamics – Equivalent Mechanical Model

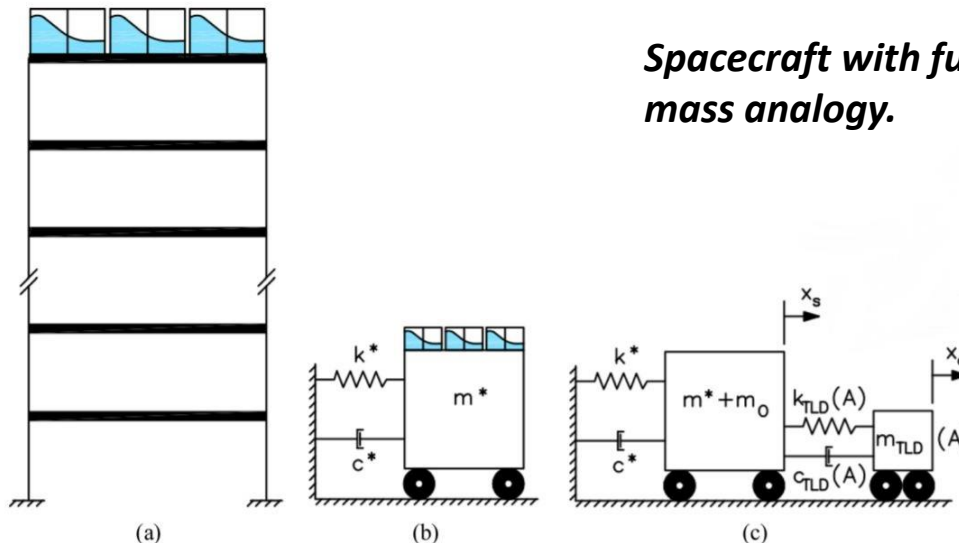
Dynamics and Control of Higher-Order Nonholonomic Systems, Jaime Rubio Hervas, PhD Thesis, 2013

## Examples of application:

- **Control of spacecraft with liquid fuel tanks**
- **Vibration absorbtion (Tuned Liquid Dampers or TLD)**
- **Modeling of fuel tanks in aircraft wings , in oil tankers or in trucks**
- ...



**Spacecraft with fuel tank modeled with a spring-mass analogy.**



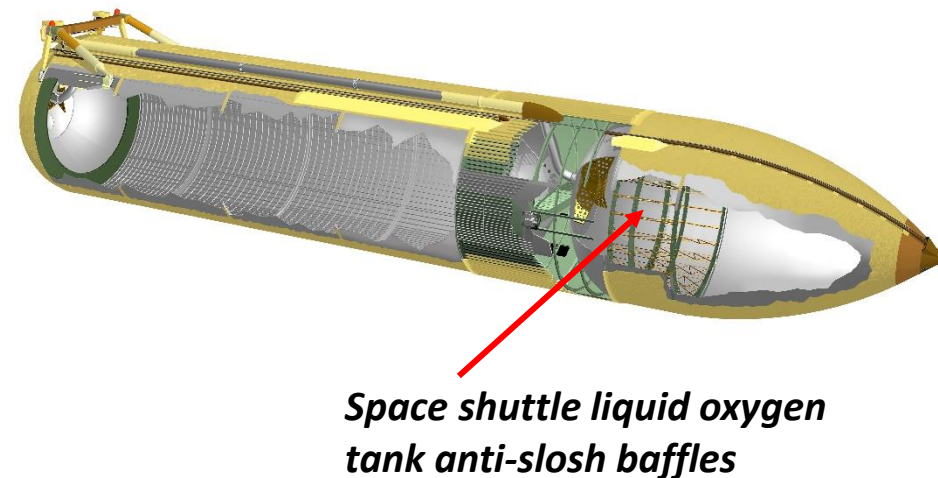
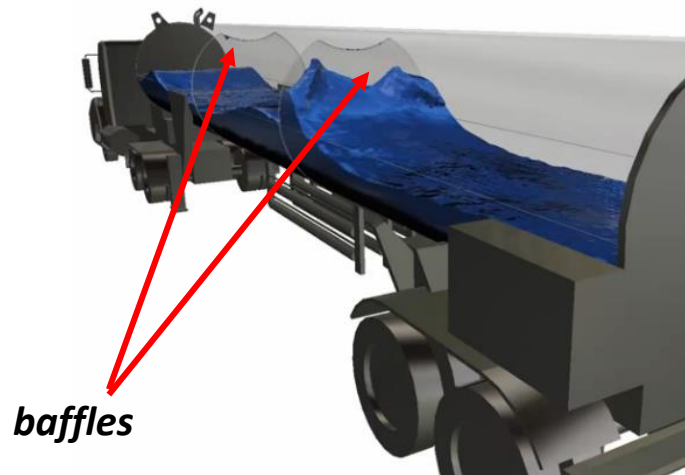
Development and Validation of Finite Element Structure-Tuned Liquid Damper System Models, Soliman et al., 2015

**The evolution of (a) a structure-TLD system into, (b) a generalized structural system with TLDs and then into, (c) a system with equivalent Tuned Mass Damper (TMD) representation**

# Sloshing Dynamics – Equivalent Mechanical Model

## *Sloshing mitigation techniques*

- *Liquid containers: Sub-division with baffles or bulkheads are widely used*
  - *Reduces wave amplitude and wall pressure, increases the liquid damping*
  - *Ongoing research to define the optimal shape, number and locations of the baffles*



# Sloshing Dynamics – Equivalent Mechanical Model

## *Sloshing mitigation techniques – the effect of a foam layer*

- *The addition of foam of the free surface damps the sloshing of the liquid.*

Entry #: V0052

Why beer does not spill:  
foam damps sloshing

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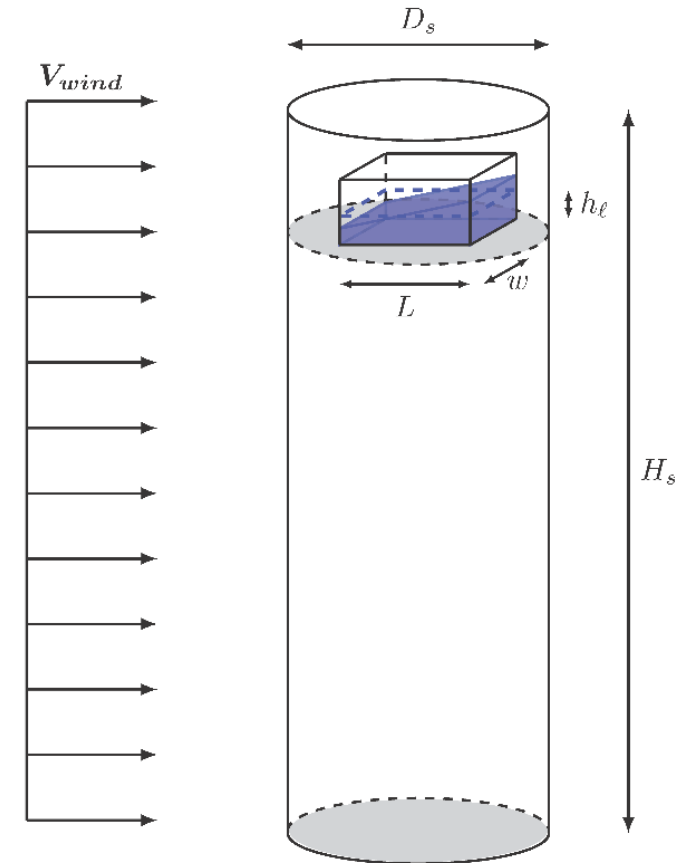
- *Foam generates an additional friction force that adds damping to the fluid oscillations.*
- *This additional force scales as:*

$$f_{\text{foam}} \sim K \gamma Ca^{2/3}$$

*where  $\gamma$  is the foam surface tension,  $K$  represents geometrical properties of the foam layer and  $Ca = \frac{\mu \dot{x}}{\gamma}$  is the capillary number  $\rightarrow$  the foam friction force in non-linear.*

## Exercise 1 – Tuned Liquid damper.

Consider a tall building represented by a cylinder with an equivalent diameter,  $D_s = 50$  m and a total height  $H_s = 300$  m as depicted on figure 1. The structure is subjected to strong winds, vortices are periodically shed from its surface thus forcing the structure to oscillate. To dampen those oscillations, a tuned liquid damper (TLD), in the form of rectangular tank partially filled with water, is installed within the building. The tank measures  $L = 10$  m by  $w = 6$  m.



# Sloshing Dynamics – Equivalent Mechanical Model

- a) *What is the frequency of the vortex shedding if we expect winds ranging up to 50 m/s. Could the associated excitation match the building eigen-frequency measured at  $f_s = 0.15$  Hz.*

*Taking a Strouhal number of  $St = 0.2$  we obtain frequencies up to  $f = \frac{St U}{D} = 0.2$  Hz. The vortex shedding frequency matches the building eigen-frequency when the wind speed is  $U = \frac{f_s D}{St} = 37.5$  m/s.*

# Sloshing Dynamics – Equivalent Mechanical Model

- b) *As said earlier, a rectangular container will be used in the building as a tuned liquid damper. At what height the container should be filled to obtain the same natural frequency as the building.*

*Using the formula of the frequency of the first sloshing mode, we find*

$$\omega_1^2 = \frac{g\pi}{L} \tanh\left(\frac{\pi h}{L}\right)$$
$$h_\ell = \frac{L}{\pi} \tanh^{-1}\left[(2\pi f_s)^2 \frac{L}{\pi g}\right] = 0.944m$$

# Sloshing Dynamics – Equivalent Mechanical Model

- c) As we are only interested in the first asymmetric mode, we will replace the liquid inside the container with two masses:  $m_0$  a rigid mass fixed to the container and  $m_1$  moving mass restrained with a spring of rigidity  $k_1$  and a dashpot with a damping  $c_1$ . Compute the value of all the said parameters.

$$m_{liqu} = \rho Lbh = 56.640 \times 10^3 \text{ kg}$$

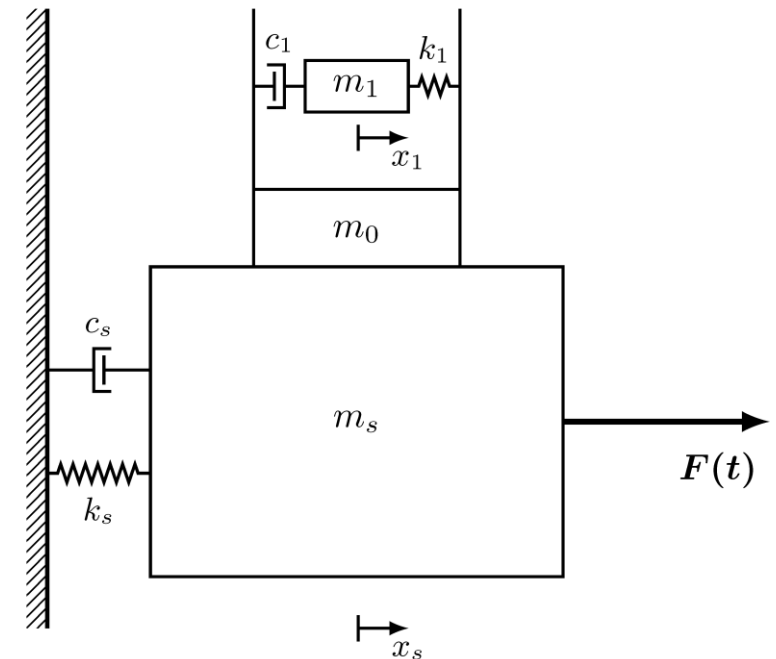
$$m_1 = 8 m_{liqu} \left(\frac{L}{h}\right) \frac{\tanh(\pi \frac{h}{L})}{1^3 \pi^3} = 44.610 \times 10^3 \text{ kg}$$

$$m_0 = m_{liqu} - m_1 = 12.029 \times 10^3 \text{ kg}$$

$$k_1 = 8 m_{liqu} \left(\frac{g}{h}\right) \frac{\tanh^2((2n-1)\pi \frac{h}{L})}{(2n-1)^2 \pi^2} = 39.618 \times 10^3 \text{ N}$$

$$\zeta_1 = \frac{1}{2h} \sqrt{\frac{2\nu}{\omega_1}} \left(1 + \frac{h}{b}\right) = 3.563 \times 10^{-4}$$

$$c_1 = 2\zeta_1 \sqrt{(m_1 k_1)} = 29.955 \text{ Ns/m}$$

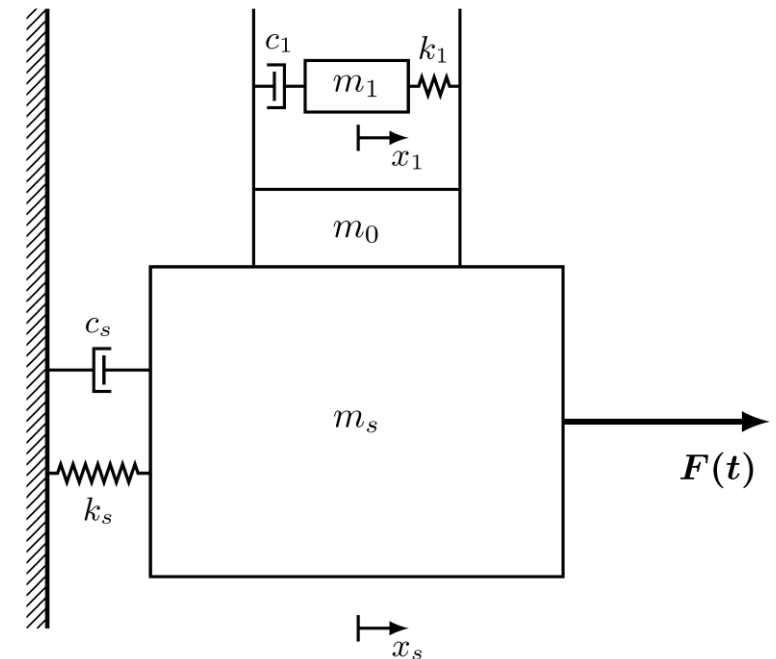


# Sloshing Dynamics – Equivalent Mechanical Model

- d) Assuming the mechanical parameters of the building are  $m_s = 250 \times 6 \text{ kg}$ , and  $c_s = 500 \text{ Ns/m}$ , and the aerodynamic parameter is  $C_{L,osci} = 0.684$ . We will model our system using the mechanical system shown. Find the equation of motion.

$$(m_s + m_0)\ddot{x}_s + c_s\dot{x}_s + k_s x_s + c_l(\dot{x}_s - \dot{x}_l) + k_l(x_s - x_l) = F(t)$$
$$m_l\ddot{x}_l + c_l(\dot{x}_l - \dot{x}_s) + k_l(x_l - x_s) = 0$$

With  $F(t) = \frac{1}{2} \rho V_{wind}^2 D_s H_s C_{L,osci} \sin(2\pi f_{st} t)$



# Sloshing Dynamics – Equivalent Mechanical Model

- e) Numerically compute the response of the system to wind blowing at 30 m/s, 37.25 and 37.5 m/s. Also compute the response of the system if the TLD was not here.

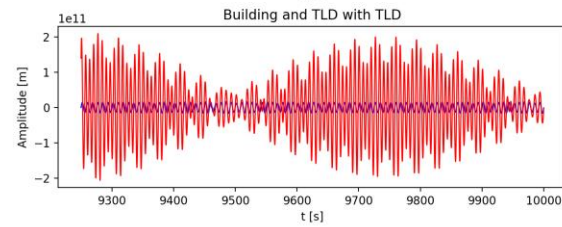
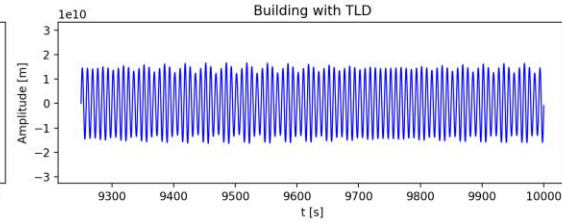
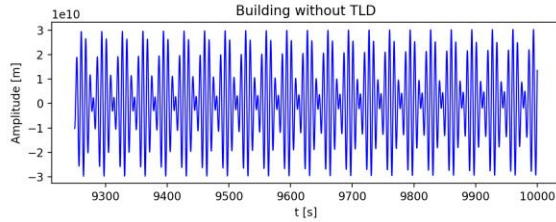
To numerically compute the response, we will need to transform our second order system to a system of 4 first order equations.

We find the following system:

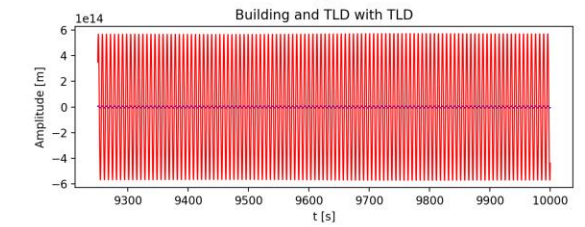
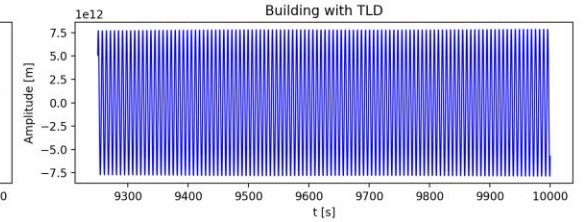
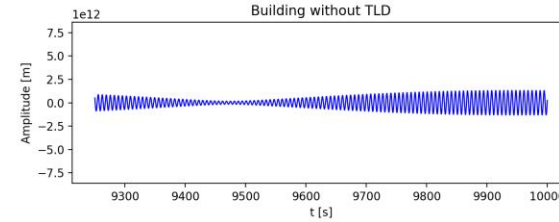
$$\frac{d}{dt} \begin{pmatrix} x_s \\ \dot{x}_s \\ x_\ell \\ \dot{x}_\ell \end{pmatrix} = \begin{pmatrix} -\frac{c_s}{(m_s + m_0)} \dot{x}_s - \frac{k_s}{(m_s + m_0)} x_s - \frac{c_\ell}{(m_s + m_0)} (\dot{x}_s - \dot{x}_\ell) - \frac{k_\ell}{(m_s + m_0)} (x_s - x_\ell) + F \\ \dot{x}_s \\ \dot{x}_\ell \\ -\frac{c_\ell}{m_\ell} (\dot{x}_\ell - \dot{x}_s) - \frac{k_\ell}{m_\ell} (x_\ell - x_s) \end{pmatrix}$$

# Sloshing Dynamics – Equivalent Mechanical Model

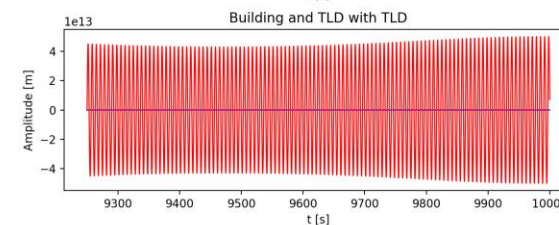
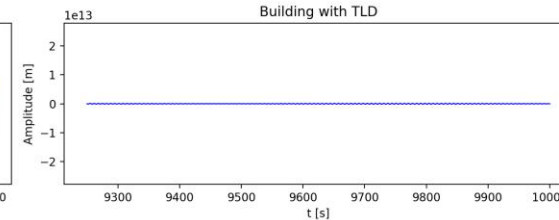
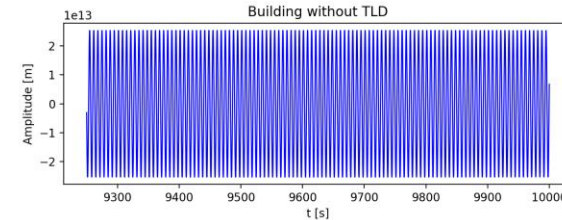
Wind velocity  $U = 30.00$



Wind velocity  $U = 37.25$



Wind velocity  $U = 37.50$

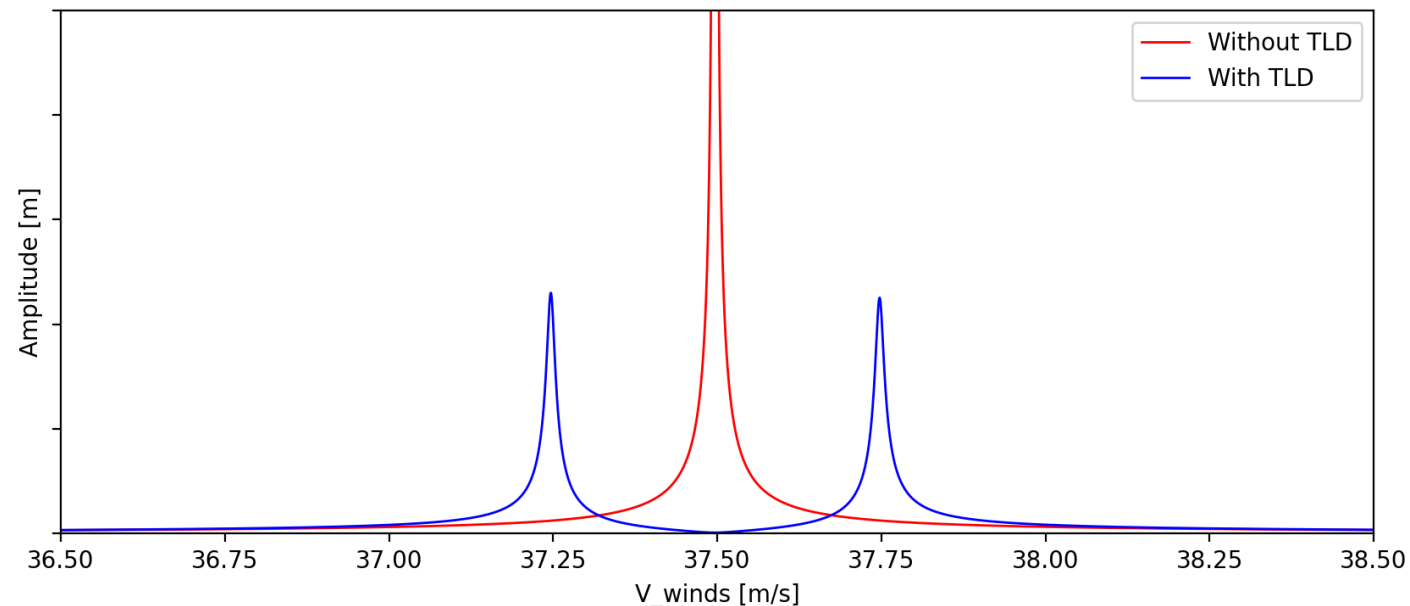


# Sloshing Dynamics – Equivalent Mechanical Model

$$x_s(t) = X_s e^{i\omega t}, \quad x_l(t) = X_l e^{i\omega t}, \quad x_l(t) = X_l e^{i\omega t}$$

$$-\omega^2 M \begin{bmatrix} X_s \\ X_l \end{bmatrix} + i\omega C \begin{bmatrix} X_s \\ X_l \end{bmatrix} + K \begin{bmatrix} X_s \\ X_l \end{bmatrix} = \begin{bmatrix} A \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} X_s \\ X_l \end{bmatrix} = \mathbf{B}^{-1} \begin{bmatrix} A \\ 0 \end{bmatrix}, \quad \mathbf{B} = -\omega^2 M + i\omega C + K$$



# Sloshing Dynamics – Equivalent Mechanical Model

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f) *What advantages and disadvantages can you think of for a tuned liquid damper compared to a classic one.*

## *Advantages:*

- Easy to incorporate in existing buildings or structures. Easy to incorporate on various building shapes (chimneys, ...)*
- The sloshing frequency can easily be adjusted by increasing or decreasing the height of water for a given tank.*
- There are no moving parts.*

## *Disadvantages:*

- Often poor damping capabilities when compared to TMDs (apparatus such as baffles, nets or contaminants are often added to increase the damping properties of the TLDs).*
- Heavy when compared to TMDs.*