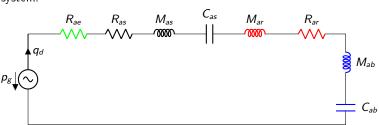
Addendum Closed-box and vented-box formulations

H. Lissek

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Analogue acoustical circuit of a loudspeaker + closed cabinet (Volume V_b) system:



Closed-box: definitions

The components of the analogue scheme are:

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Electrical parts	
$p_{g} = \frac{B\ell U_{g}}{S_{d}R_{e}}$	source pressure analogue to the source electrical voltage $U_{\!g}$
$P_{g} = \frac{1}{S_{d}R_{e}}$ $R_{ae} = \frac{(B\ell)^{2}}{S_{d}^{2}R_{e}}$	acoustic resistances analogue to $R_{ m e_{\perp}}$ DC resistance of the coil
_	(we neglect here the output resistance of voltage source U_S , as well as the electrical inductance of the coil $L_{\mathcal{C}}$)
Mechanical parts	
$R_{as} = \frac{R_{ms}}{s_d^2}$	acoustic resistance corresponding to the mechanical resistance of the loudspeaker
$R_{as} = \frac{R_{ms}}{S_a^2}$ $M_{as} = \frac{M_{ms}}{S_d^2}$ $C_{as} = S_d^2 C_{ms}$	acoustic mass corresponding to the mechanical mass of the loudspeaker
$C_{as} = S_d^2 C_{ms}^d$	acoustic compliance corresponding to the mechanical compliance of the loudspeaker
Acoustic parts	
$M_{ar} = \frac{8}{3\pi} \frac{\rho a}{S_d}$ $R_{ar} \approx \frac{\rho c}{S_d} \frac{(ka)^2}{4}$ $M_{ab} \approx M_{ar} \approx \frac{8}{3\pi} \frac{\rho a}{S_d}$	radiation mass of a piston of radius $a=\sqrt{S_d/\pi}$ mounted on a screen
$R_{ar} \approx \frac{\rho c}{S_{J}} \frac{(ka)^2}{4}$	radiation resistance of a "pulsating" source of radius z (solid angle 4π)
$M_{ab} \approx M_{ar} \approx \frac{8}{3\pi} \frac{\rho_a}{S_d}$	radiation mass inside the cabinet
	(deformation of particle velocity in the cabinet)
$C_{ab} = S_d^2 C_{ms}$	acoustic compliance equivalent to the cabinet volume $V_{oldsymbol{b}}$

Closed-box: response in flow velocity

The closed-box loudspeaker is assimilated to a small semi-monopolar (in 2π steradians) source of radius a (radiation resistance $R_{ar}=2\pi\rho\frac{f^2}{c}$) and flow velocity q_d .

The radiated power is then: $P_a = R_{ar}|q_d|^2$

The power can be derived by expressing q_d as a function of voltage U_g .

Since:
$$q_d = \frac{p_g}{Z_{ac}}$$
, where $Z_{ac} = R_{ac} + j\omega M_{ac} + \frac{1}{j\omega C_{ac}}$ with: $R_{ac} = R_{ae} + R_{as} + R_{ar}$: total losses of the closed-box loudspeaker $M_{ac} = M_{as} + M_{ar} + M_{ab}$: total acoustic mass $C_{ac} = \frac{C_{as}C_{ab}}{C_{as} + C_{ab}}$: total acoustic compliance Then $q_d = \left(\frac{B\ell U_g}{S_d R_e}\right) \frac{j\omega C_{ac}}{(j\omega)^2 M_{ac}C_{ac} + j\omega R_{ac}C_{ac} + 1}$

Closed-box: response in flow velocity

$$q_d = \left(\frac{B\ell U_g}{S_d R_e}\right) \frac{j\omega C_{ac}}{(j\omega)_1^2 M_{ac} C_{ac} + j\omega R_{ac} C_{ac} + 1}$$

$$\begin{split} q_d &= \left(\frac{B\ell U_g}{S_d R_e}\right) \frac{j\omega \, C_{ac}}{(j\omega)^2 M_{ac} \, C_{ac} + j\omega R_{ac} \, C_{ac} + 1} \\ \text{Defining: } \omega_c &= \frac{1}{\sqrt{M_{ac} \, C_{ac}}} \text{ and } Q_{tc} = \frac{1}{\omega_s R_{ac} \, C_{ac}} \text{ the acoustic resonator's} \end{split}$$

resonance frequency and quality factor, we show that:

$$C_{ac} = \frac{1}{\omega_c R_{ac} Q_{tc}}$$

$$M_{ac} C_{ac} = \frac{1}{\omega_c^2}$$

$$R_{ac} C_{ac} = \frac{1}{\omega_c Q_{tc}}$$

$$R_{ac}C_{ac}=\frac{1}{\omega_cQ_{tc}}$$

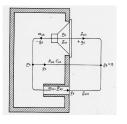
And then
$$q_d = \left(\frac{B\ell U_g}{S_d R_e}\right) \frac{1}{Q_{tc} R_{ac}} \frac{\left(\frac{j\omega}{\omega_c}\right)}{\left(\frac{j\omega}{\omega_c}\right)^2 + \frac{1}{Q_{tc}} \left(\frac{j\omega}{\omega_c}\right) + 1}$$

Since
$$\frac{B\ell}{S_dR_e} = \frac{S_dR_{ae}}{B\ell}$$
, then we can define $q_c = \frac{B\ell U_g}{S_dR_eQ_{tc}R_{ac}} = \frac{S_d}{B\ell Q_{ec}}U_g$

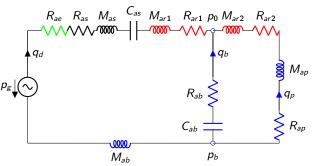
and finally
$$q_d=q_crac{\left(rac{j\omega}{\omega_c}
ight)}{\left(rac{j\omega}{\omega_c}
ight)^2+rac{1}{Q_{tc}}\left(rac{j\omega}{\omega_c}
ight)+1}$$

$$\begin{split} P_{a} &= R_{ar} |q_{d}|^{2} = 2\pi \rho \frac{f^{2}}{c} q_{c}^{2} \left| \frac{\left(\frac{j\omega}{\omega_{c}}\right)}{\left(\frac{j\omega}{\omega_{c}}\right)^{2} + \frac{1}{Q_{tc}} \left(\frac{j\omega}{\omega_{c}}\right) + 1} \right|^{2} \\ &\text{Then } P_{a} = P_{ac} \left| \frac{\left(\frac{j\omega}{\omega_{c}}\right)^{2}}{\left(\frac{j\omega}{\omega_{c}}\right)^{2} + \frac{1}{Q_{tc}} \left(\frac{j\omega}{\omega_{c}}\right) + 1} \right|^{2} \\ &\text{where } P_{ac} = \frac{2\pi \rho}{c} f_{c}^{2} q_{c}^{2} \end{split}$$

Vented-box loudspeaker: definitions



Analogue acoustic circuit of a vented-box + loudspeaker system (bass-reflex):



Vented-box loudspeaker: definitions

In the analogue acoustic representation of the vented-box loudspeaker:

- the electrical voltage U_g applied to the loudspeaker terminals is replaced by the equivalent pressure $p_g=rac{B\ell}{S_dR_e}U_g$.
- the total acoustic impedance of the loudspeaker (excluding the Helmholtz resonator of the vented box) becomes:

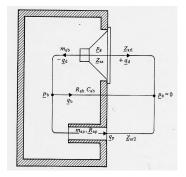
$$\begin{split} Z_{as}' &= \left(R_{ae} + R_{as} + R_{ar}\right) + j\omega\left(M_{as} + M_{ab} + M_{ar}\right) + \frac{1}{j\omega\,\mathcal{C}_{as}} \\ &\approx R_{ae} + R_{as}' + j\omega\,M_{as}' + \frac{1}{j\omega\,\mathcal{C}_{as}} \end{split}$$

• the equivalent acoustic impedance of the Helmholtz resonator constituted with volume V_b (acoustic compliance $C_{ab} = \frac{V_b}{\rho c^2}$, we neglect the resistance R_{ab}) and the port of length I_p and cross-section s_p (acoustic mass $M'_{ap} = M_{ap} + M_{ar2} = \rho \frac{a_p + 2\Delta L}{s_p}$, with ΔL the termination correction accounting for the acoustic radiation at each sides of the port $\Delta L \approx \frac{8a}{3\pi}$) reads:

$$Z_{ab} = rac{j\omega M_{ap}'}{\left(j\omega
ight)^2 M_{ap}' C_{ab} + 1} = rac{j\omega M_{ap}'}{\left(rac{j\omega}{\omega_p}
ight)^2 + 1}, ext{ with } \omega_p = rac{1}{\sqrt{M_{ap}' C_{ab}}}$$

Vented-box loudspeaker: acoustic radiation

The radiated sound power is due to the combination of both the loudspeaker diaphragm and the port radiations. Then, we can consider the flow velocity $q_s+q_p=-q_b$ to process the radiated power, and consider the loudspeaker as a semi-monopole $R_{ar}=\frac{2\pi\rho}{c}f^2$.



The radiated power then writes: $P_a = R_{ar} |q_b|^2$, and can be derived by analyzing the analogue acoustic circuit of the system.

Vented-box loudspeaker: acoustic radiation

We can derive q_b in the analogue acoustic circuit as $q_b = i\omega C_{ab}p_b$, where p_b is the pressure inside the cabinet of volume V_b .

Pressure p_b can be simply deduced by analyzing the "pressure divider":

$$p_b = rac{Z_{ab}}{Z_{ab} + Z_{as}'} p_g$$
, then:
 $q_b = rac{B\ell U_g}{C_{ab}} rac{j\omega C_{ab} Z_{ab}}{Z_{ab}} = rac{B}{C_{ab}}$

$$q_b = \frac{\frac{B\ell U_g}{S_d R_e}}{\frac{j\omega}{Z_{ab} + Z'_{as}}} \frac{j\omega C_{ab} Z_{ab}}{Z_{ab} + Z'_{as}} = \frac{\frac{B\ell U_g}{S_d R_e}}{\left[\left(\frac{j\omega}{\omega_p}\right)^2 + 1\right] \left[j\omega M'_{as} + R'_{as} + \frac{1}{j\omega C_{as}}\right] + j\omega M'_{ap}}$$

 $(j\omega)^3 C_{ab} M'_{ap} C_{as} M'_{as}$

 $\frac{{}^{B\ell U_g \omega_s}}{{}^{S}_d R_e \underline{M_{as}' \omega_s}} \frac{(j\omega)^{\bullet} \, C_{ab} M_{ap}' \, C_{as} M_{as}'}{(j\omega)^{4} \, M_{as}' \, C_{as} M_{ap}' \, C_{ab} + (j\omega)^{3} \, M_{ap}' \, C_{ab} (R_{ae} + R_{as}) \, C_{as} + (j\omega)^{2} [M_{ap}' \, C_{ab} + M_{as}' \, C_{as} + M_{ap}' \, C_{as}] + (j\omega) \, C_{as} (R_{ae} + R_{as}') + 1} \\ \text{Reminding } \omega_s = \frac{1}{\sqrt{M_{as}' \, C_{as}}} \text{ and } Q_{ts} = \frac{1}{\omega_s \, R_{as}' \, C_{as}}, \text{ as well as } \omega_p = \frac{1}{\sqrt{M_{ap}' \, C_{ab}}} \text{ it}$

comes:

$$q_{b} = \frac{B\ell Ug}{S_{d}R_{e}M_{as}^{\prime}\omega_{s}} \frac{\left(\frac{j\omega}{\omega_{s}}\right)^{3}\frac{\omega_{s}^{2}}{\omega_{p}^{2}}}{\left(\frac{j\omega}{\omega_{s}}\right)^{4}\frac{\omega_{s}^{2}}{\omega_{p}^{2}} + \left(\frac{j\omega}{\omega_{s}}\right)^{3}\frac{1}{Q_{ts}}\frac{\omega_{s}^{2}}{\omega_{p}^{2}} + \left(\frac{j\omega}{\omega_{s}}\right)^{2}\left[\frac{\omega_{s}^{2}}{\omega_{p}^{2}} + 1 + \frac{C_{as}}{C_{ab}}\frac{\omega_{s}^{2}}{\omega_{p}^{2}}\right] + \left(\frac{j\omega}{\omega_{s}}\right)\frac{1}{Q_{ts}} + 1}$$

Vented-box loudspeaker: acoustic radiation

$$\begin{split} q_b &= \frac{B\ell Ug}{S_d R_e M_{as} \omega_s} \frac{\left(\frac{j\omega}{\omega_s}\right)^3 \frac{\omega_s^2}{\omega_\rho^2}}{\left(\frac{j\omega}{\omega_s}\right)^4 \frac{\omega_s^2}{\omega_\rho^2} + \left(\frac{j\omega}{\omega_s}\right)^3 \frac{1}{Q_{ts}} \frac{\omega_s^2}{\omega_\rho^2} + \left(\frac{j\omega}{\omega_s}\right)^2 \left[\frac{\omega_s^2}{\omega_\rho^2} + 1 + \frac{C_{as}}{C_{ab}} \frac{\omega_s^2}{\omega_\rho^2}\right] + \left(\frac{j\omega}{\omega_s}\right) \frac{1}{Q_{ts}} + 1}{1}. \end{split}$$
 Introducing $\alpha = \frac{C_{as}}{C_{ab}}$ et $h = \frac{\omega_p}{\omega_s}$, it comes:
$$q_b = \frac{B\ell Ug}{S_d R_e M_{as}' \omega_s} \frac{\left(\frac{j\omega}{\omega_s}\right)^4 h^{-2}}{\left(\frac{j\omega}{\omega_s}\right)^4 h^{-2} + \left(\frac{j\omega}{\omega_s}\right)^3 \frac{h^{-2}}{Q_{ts}} + \left(\frac{j\omega}{\omega_s}\right)^2 \left[1 + h^{-2}(1+\alpha)\right] + \left(\frac{j\omega}{\omega_s}\right) \frac{1}{Q_{ts}} + 1}. \end{split}$$

Vented-box loudspeaker: response in power

$$\begin{split} P_{a} &= \frac{2\pi\rho}{c} f^{2} \left(\frac{B\ell U_{g}}{S_{d}R_{e}M_{as}^{\prime}\omega_{s}} \right)^{2} \left| \frac{\left(\frac{j\omega}{\omega_{s}} \right)^{3} h^{-2}}{\left(\frac{j\omega}{\omega_{s}} \right)^{3} h^{-2}} \left| \frac{\left(\frac{j\omega}{\omega_{s}} \right)^{3} h^{-2}}{\left(\frac{j\omega}{\omega_{s}} \right)^{4} h^{-2} + \left(\frac{j\omega}{\omega_{s}} \right)^{3} \frac{h^{-2}}{Q_{ts}} + \left(\frac{j\omega}{\omega_{s}} \right)^{2} \left[1 + h^{-2}(1+\alpha) \right] + \left(\frac{j\omega}{\omega_{s}} \right) \frac{1}{Q_{ts}} + 1} \right|^{2} \\ &= \frac{2\pi\rho}{c} f_{s}^{2} \left(\frac{B\ell U_{g}}{S_{d}R_{e}M_{as}^{\prime}\omega_{s}} \right)^{2} \left| \frac{\left(\frac{j\omega}{\omega_{s}} \right)^{4} h^{-2} + \left(\frac{j\omega}{\omega_{s}} \right)^{3} \frac{h^{-2}}{Q_{ts}} + \left(\frac{j\omega}{\omega_{s}} \right)^{2} \left[1 + h^{-2}(1+\alpha) \right] + \left(\frac{j\omega}{\omega_{s}} \right) \frac{1}{Q_{ts}} + 1} \right|^{2} \\ &= P_{ao} \left| \frac{\left(\frac{j\omega}{\omega_{s}} \right)^{4} h^{-2} + \left(\frac{j\omega}{\omega_{s}} \right)^{3} \frac{h^{-2}}{Q_{ts}} + \left(\frac{j\omega}{\omega_{s}} \right)^{2} \left[1 + h^{-2}(1+\alpha) \right] + \left(\frac{j\omega}{\omega_{s}} \right) \frac{1}{Q_{ts}} + 1} \right|^{2} \\ &= P_{ao} \left| \frac{\left(\frac{j\omega}{\omega_{s}} \right)^{4} h^{-2} + \left(\frac{j\omega}{\omega_{s}} \right)^{3} \frac{h^{-2}}{Q_{ts}} + \left(\frac{j\omega}{\omega_{s}} \right)^{2} \left[1 + h^{-2}(1+\alpha) \right] + \left(\frac{j\omega}{\omega_{s}} \right) \frac{1}{Q_{ts}} + 1} \right|^{2} \\ &= P_{ao} \left| \frac{\left(\frac{j\omega}{\omega_{s}} \right)^{4} h^{-2} + \left(\frac{j\omega}{\omega_{s}} \right)^{3} \frac{h^{-2}}{Q_{ts}} + \left(\frac{j\omega}{\omega_{s}} \right)^{2} \left[1 + h^{-2}(1+\alpha) \right] + \left(\frac{j\omega}{\omega_{s}} \right) \frac{1}{Q_{ts}} + 1} \right|^{2} \\ &= P_{ao} \left| \frac{\left(\frac{j\omega}{\omega_{s}} \right)^{4} h^{-2} + \left(\frac{j\omega}{\omega_{s}} \right)^{3} \frac{h^{-2}}{Q_{ts}} + \left(\frac{j\omega}{\omega_{s}} \right)^{2} \left[1 + h^{-2}(1+\alpha) \right] + \left(\frac{j\omega}{\omega_{s}} \right) \frac{1}{Q_{ts}} + 1} \right|^{2} \\ &= P_{ao} \left| \frac{\left(\frac{j\omega}{\omega_{s}} \right)^{4} h^{-2} + \left(\frac{j\omega}{\omega_{s}} \right)^{3} \frac{h^{-2}}{Q_{ts}} + \left(\frac{j\omega}{\omega_{s}} \right)^{2} \left[1 + h^{-2}(1+\alpha) \right] + \left(\frac{j\omega}{\omega_{s}} \right) \frac{1}{Q_{ts}} + 1} \right|^{2} \\ &= P_{ao} \left| \frac{\left(\frac{j\omega}{\omega_{s}} \right)^{4} h^{-2} + \left(\frac{j\omega}{\omega_{s}} \right)^{3} \frac{h^{-2}}{Q_{ts}} + \left(\frac{j\omega}{\omega_{s}} \right)^{2} \left[\frac{j\omega}{\omega_{s}} \right] \right|^{2} \\ &= P_{ao} \left| \frac{j\omega}{\omega_{s}} \right|^{2} \left(\frac{j\omega}{\omega_{s}} \right)^{2} \left(\frac{j\omega}{\omega_{s}} \right)$$