Optical methods in chemistry
or
Photon tools for chemical sciences

Session 11

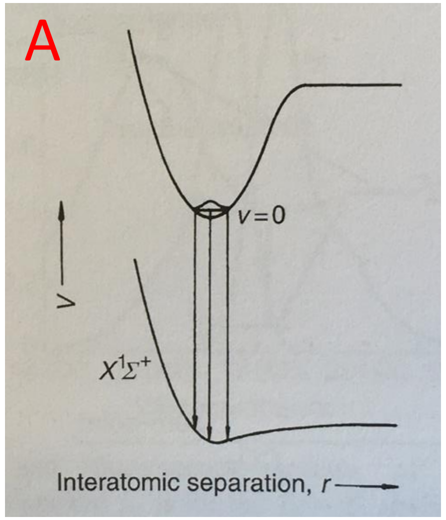
Course layout – contents overview and general structure

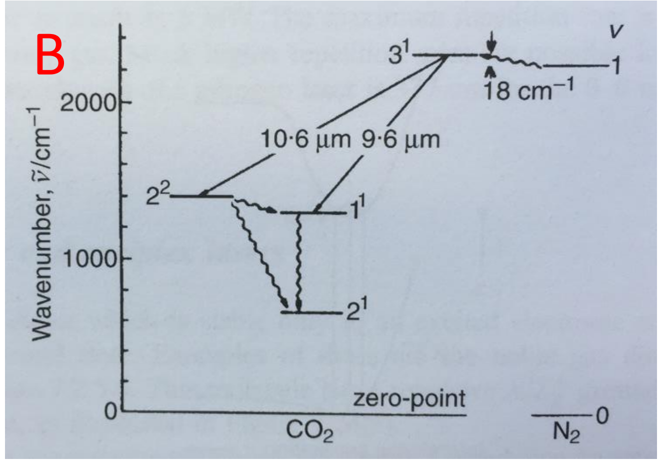
- Introduction and ray optics
- Wave optics
- Beams
- From cavities to lasers
- More lasers and optical tweezers
- From diffraction and Fourier optics
- Microscopy
- Spectroscopy
- Electromagnetic optics
- Absorption, dispersion, and non-linear optics
- Ultrafast lasers
- Introduction to x-rays
- Summary

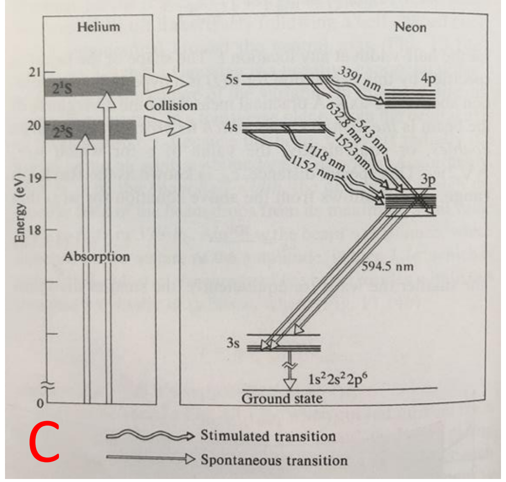
Today:

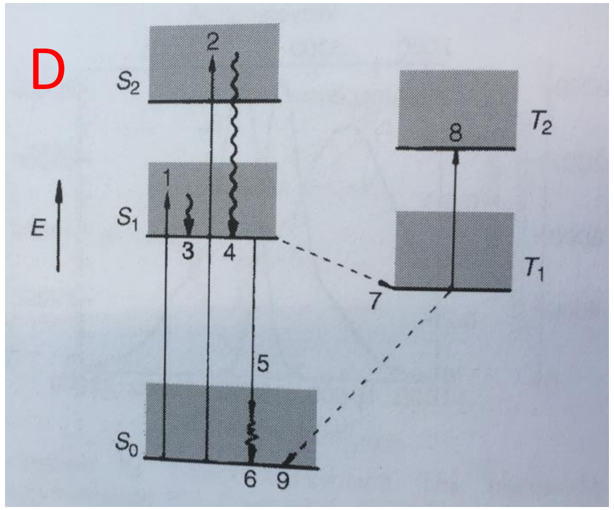
Ultrafast lasers with a bit of non-linear optics

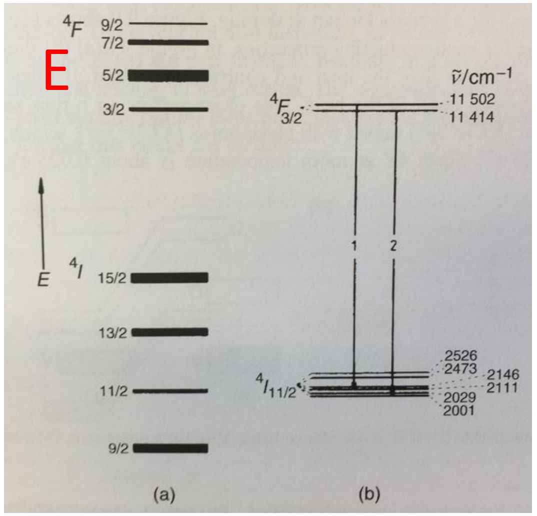
Recall: Many different ideas and version





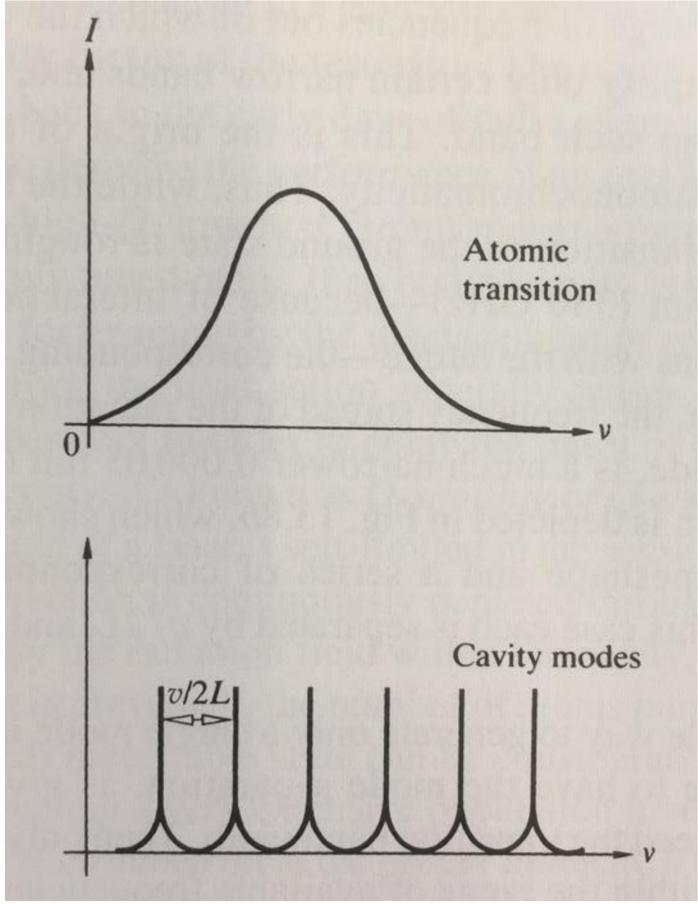




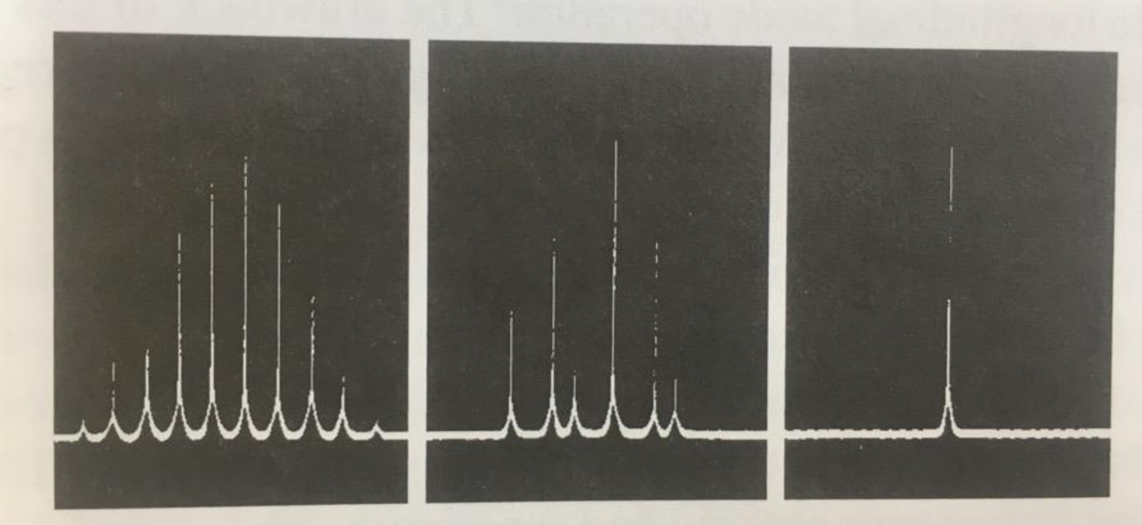


Recall: Lasers and laser cavities

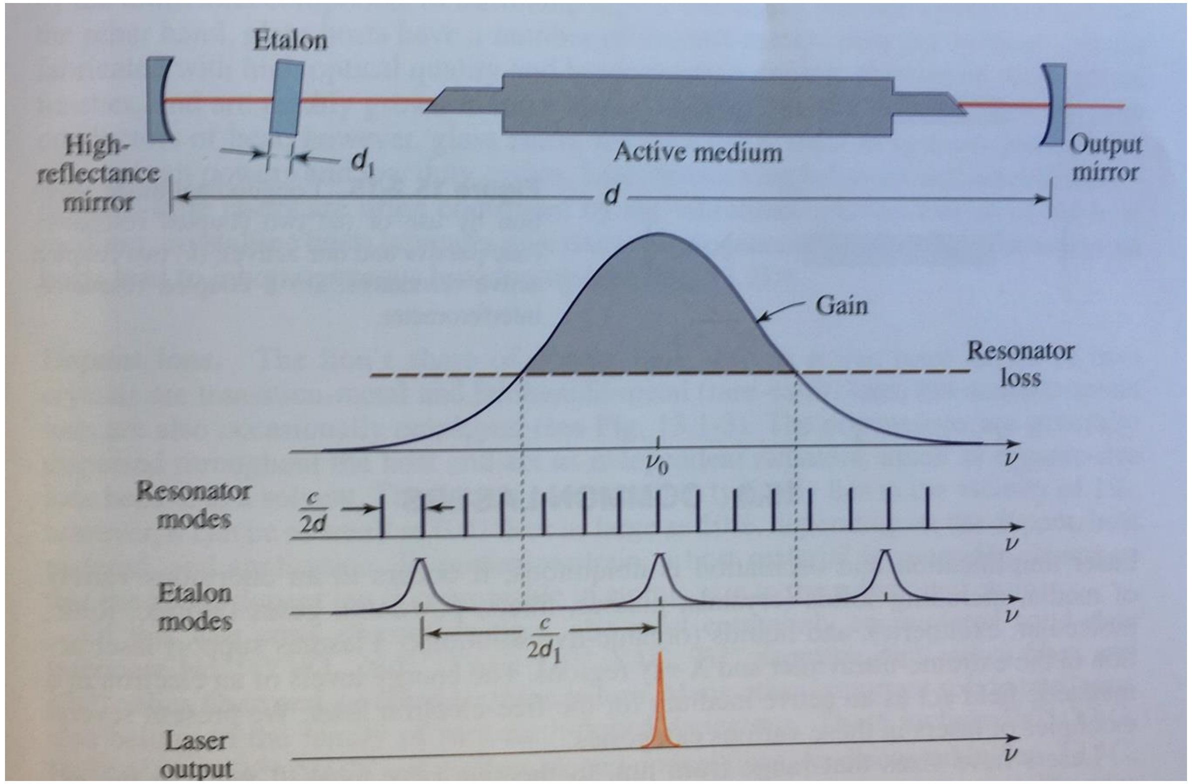
Atomic lines vs cavity modes



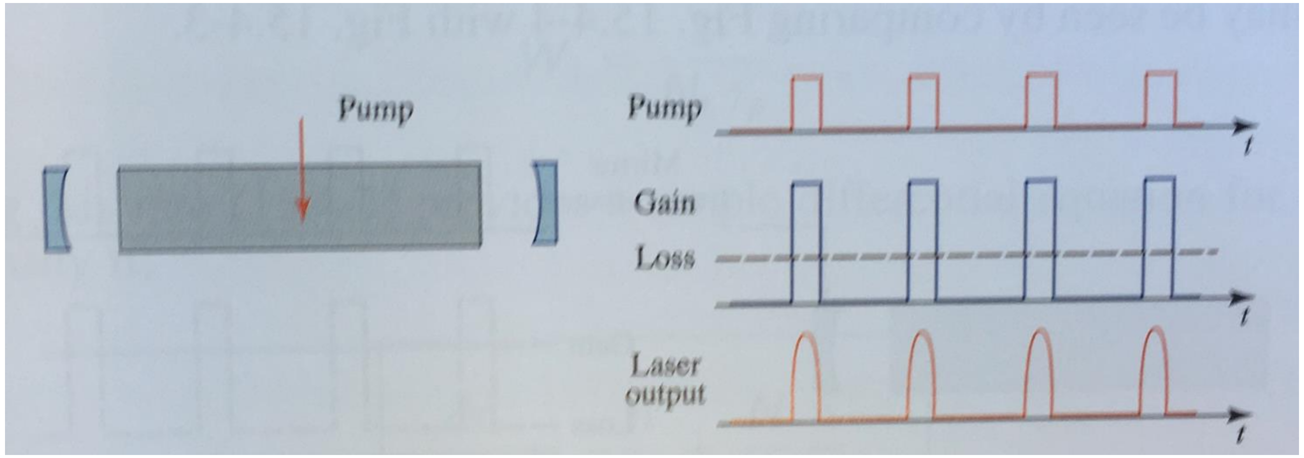
Mode selection



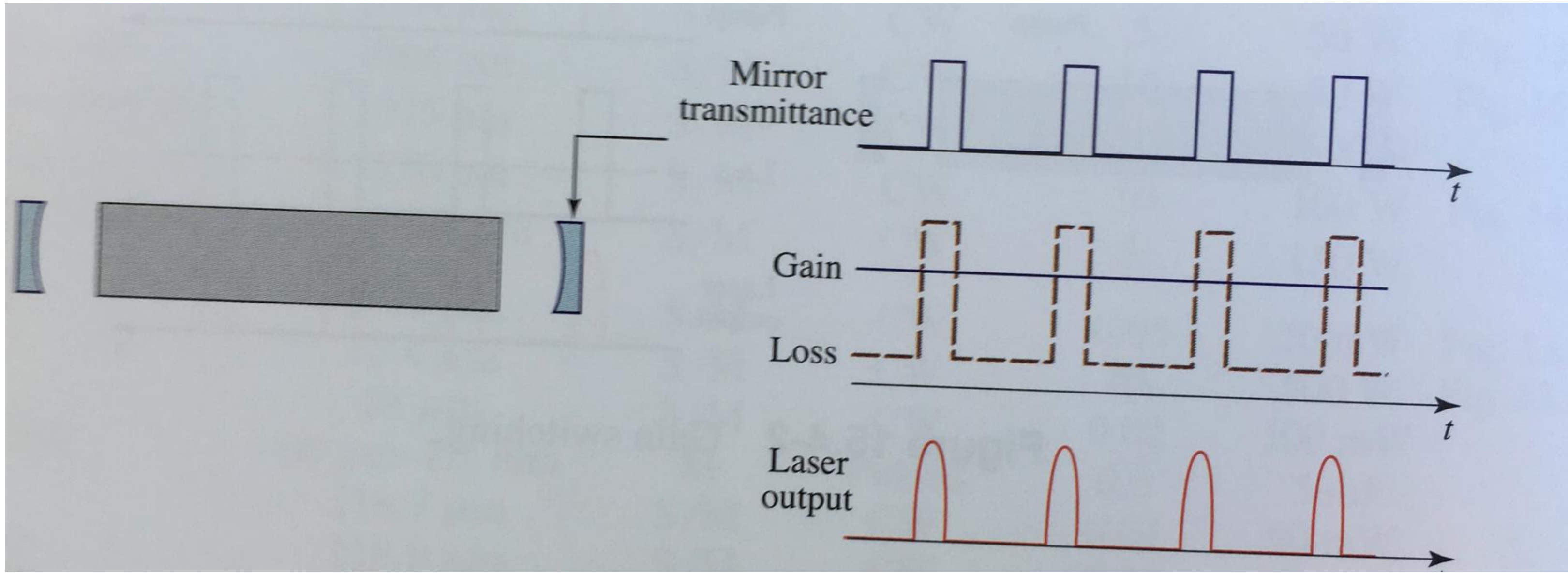
Recall: Mode selection with intracavity wavelength selection (here etalon)



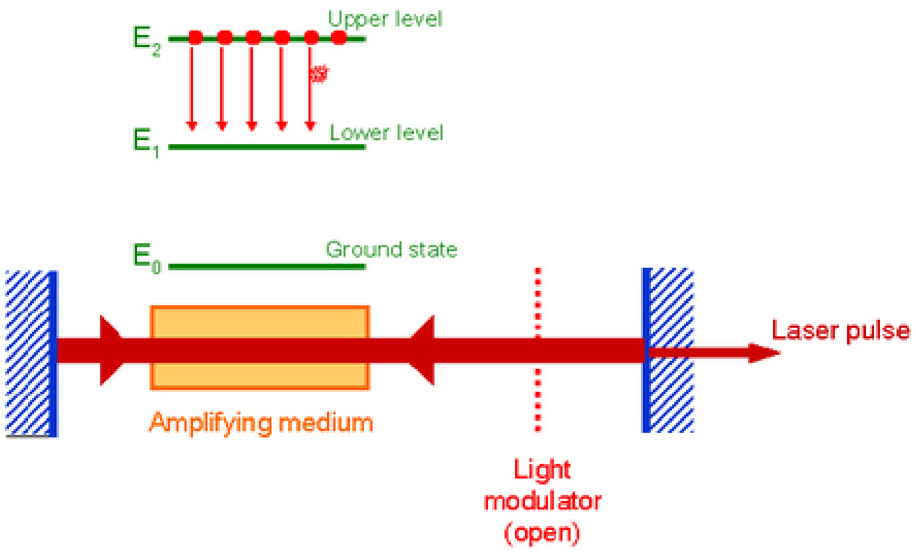
Pulsed lasers: Gain switching

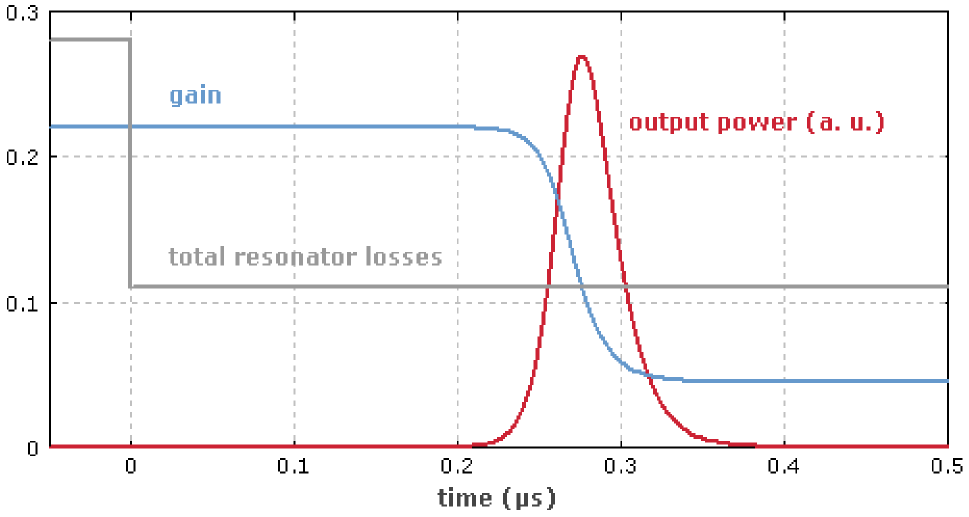


Pulsed lasers: Cavity dumping



Q Switching

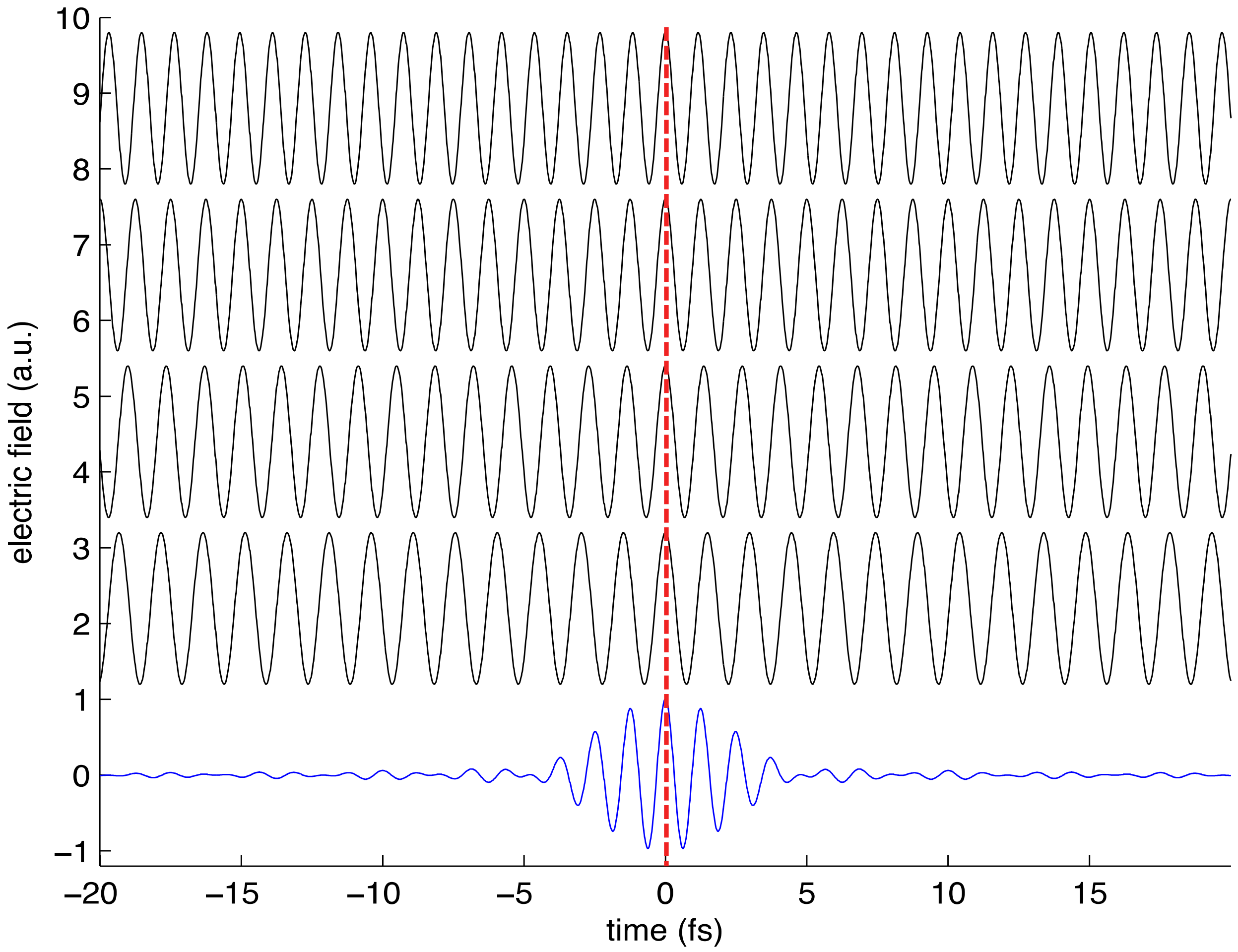




[See video at: https://www.rp-photonics.com/q_switching.html](https://www.rp-photonics.com/q_switching.html)

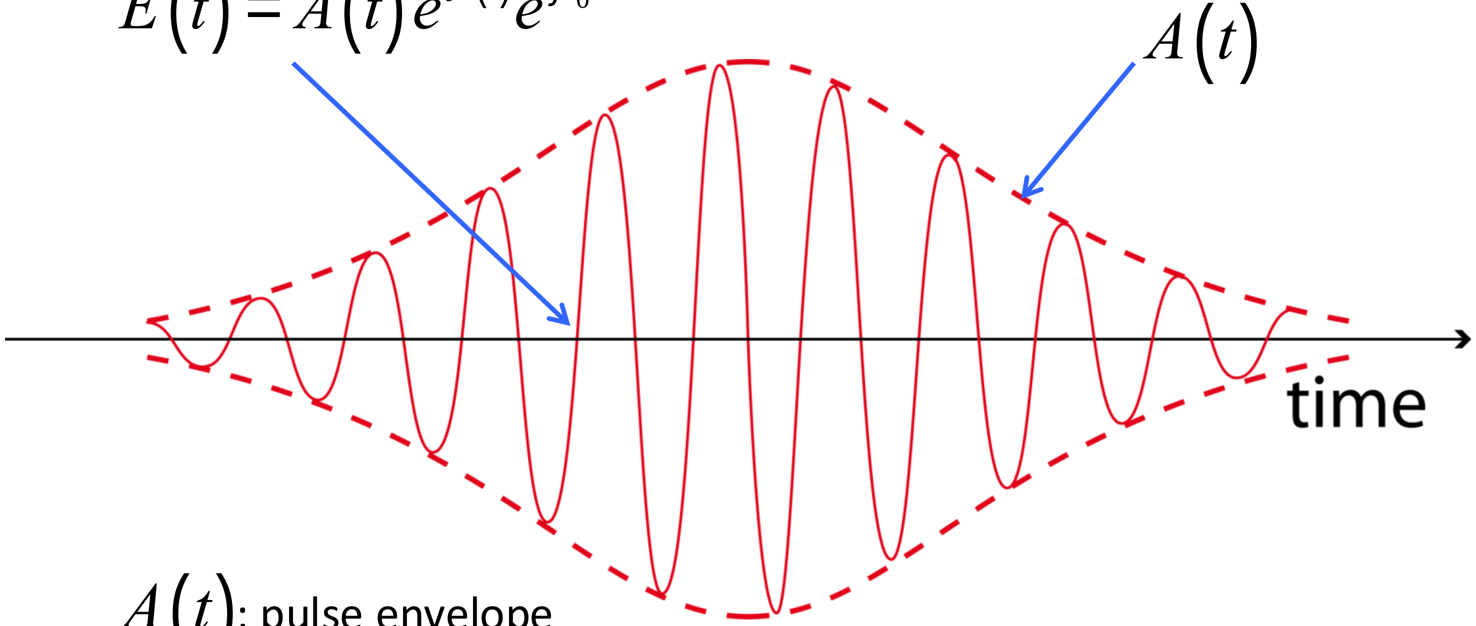
Now going to really short pulses:

Short pulse needs many frequencies: Shortest pulse is Fourier transform limited pulse (aka bandwidth limited pulse)



Note: Mathematical description of optical pulse

$$E(t) = A(t)e^{i\varphi(t)}e^{i\varphi_0}$$



$A(t)$; pulse envelope

$\varphi(t)$; time dependent phase of electric field

$\omega(t) = -\dfrac{d\varphi(t)}{dt}$; instantaneous frequency

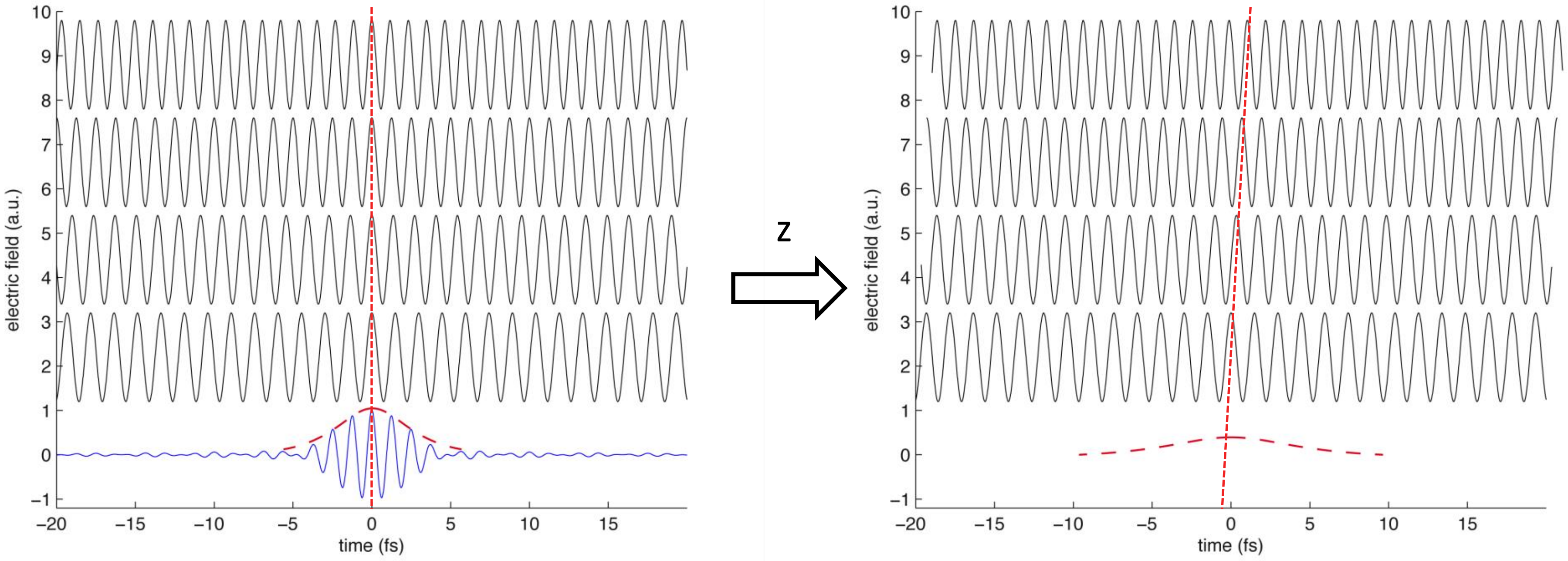
φ_0 ; absolute phase (typically neglected for multi-cycle pulses)

A note on “chirp”

Initially transform limited pulse becomes chirped upon propagation, ($k_2 \neq 0$, $\gamma = 0$)

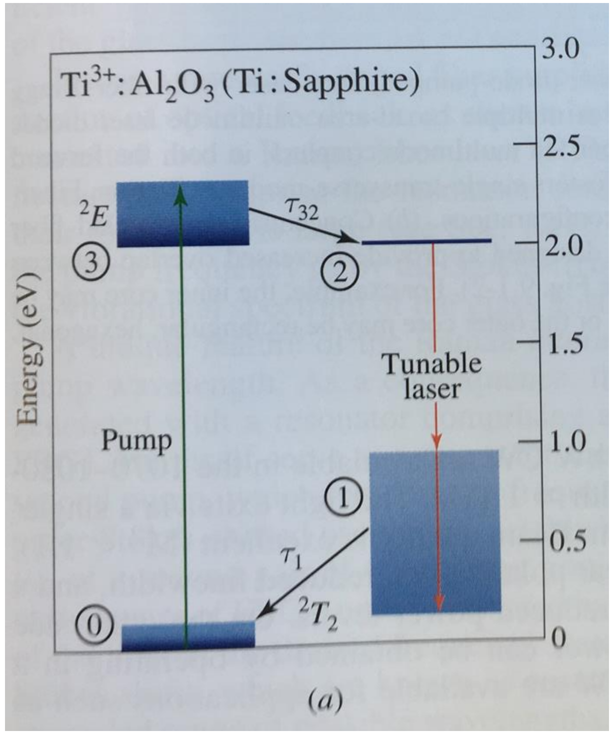
transform-limited

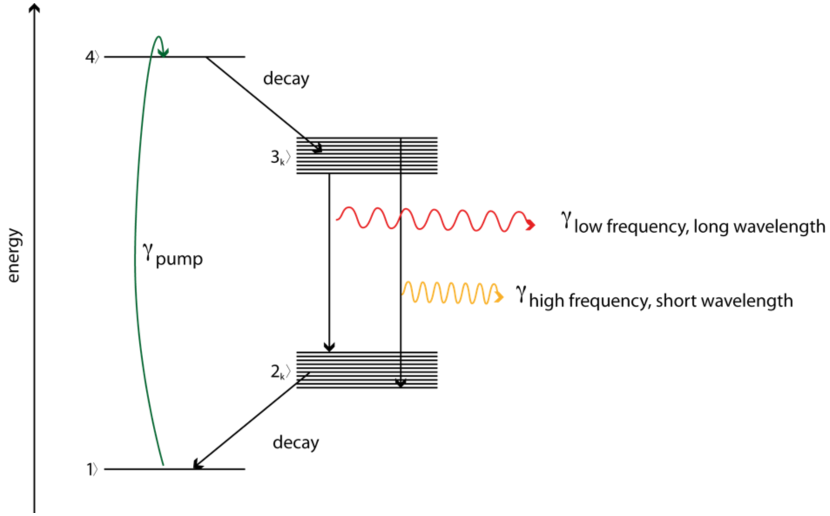
linearly chirped

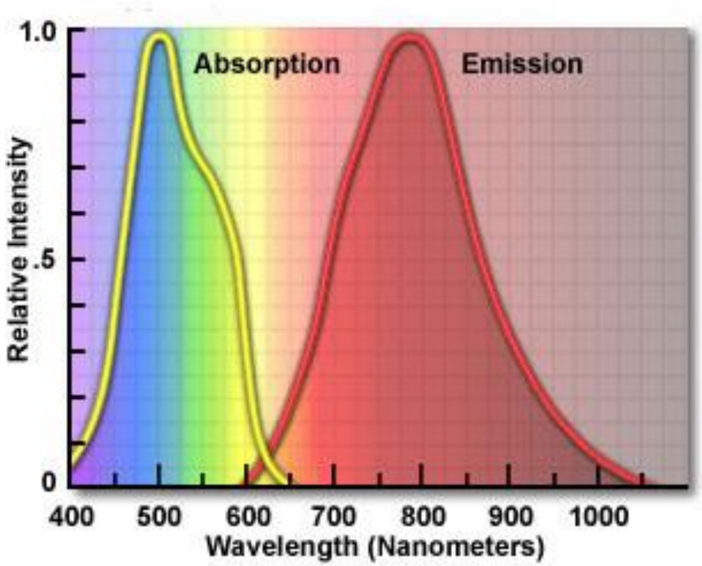


Important laser: TiSa

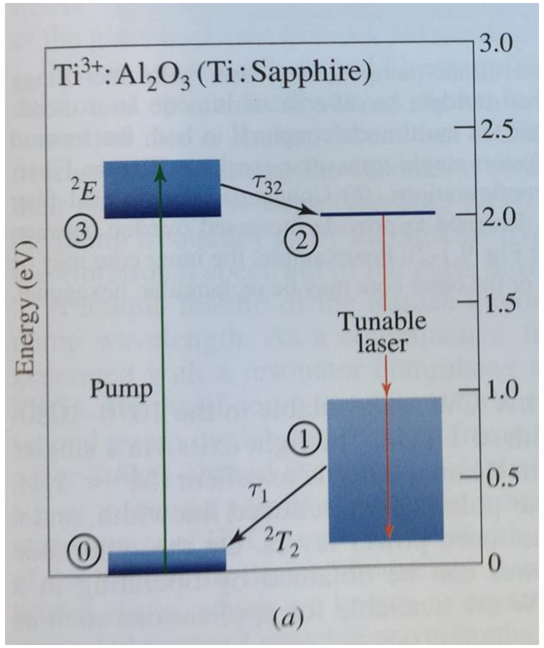
- Ti^{3+} ions in Al_2O_3 (Saphire)
- Tunable, broad bandwidth
- Can provide ultrashort pulses
- Workhorse laser for ultrafast sciences

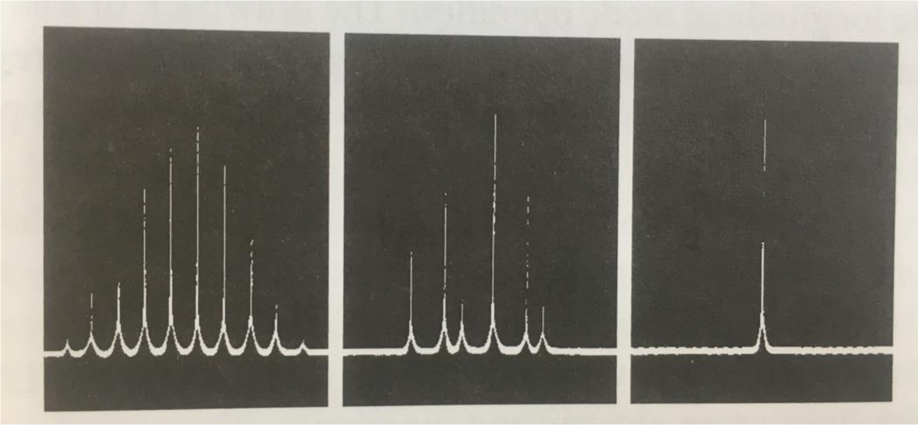






Short pulses and cavities – how do they go together?



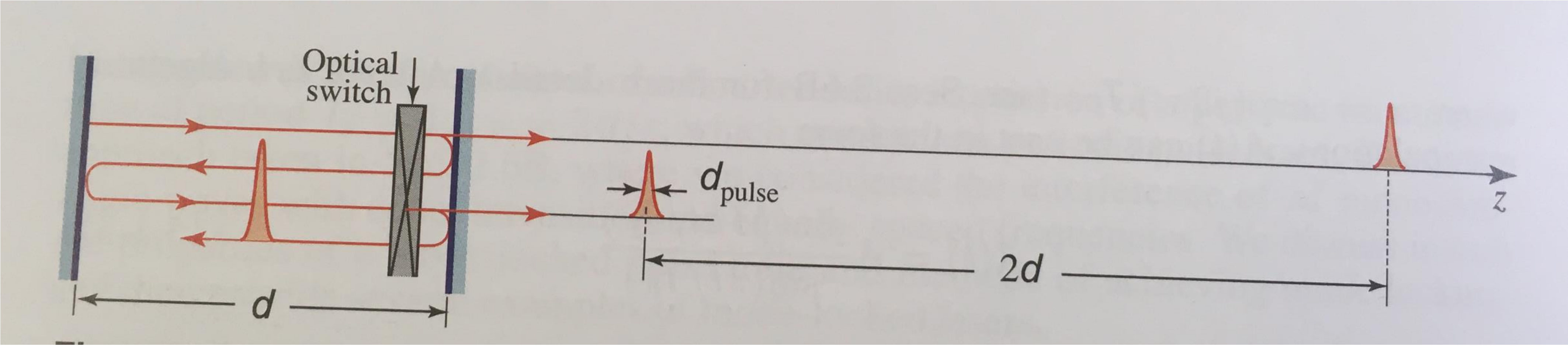


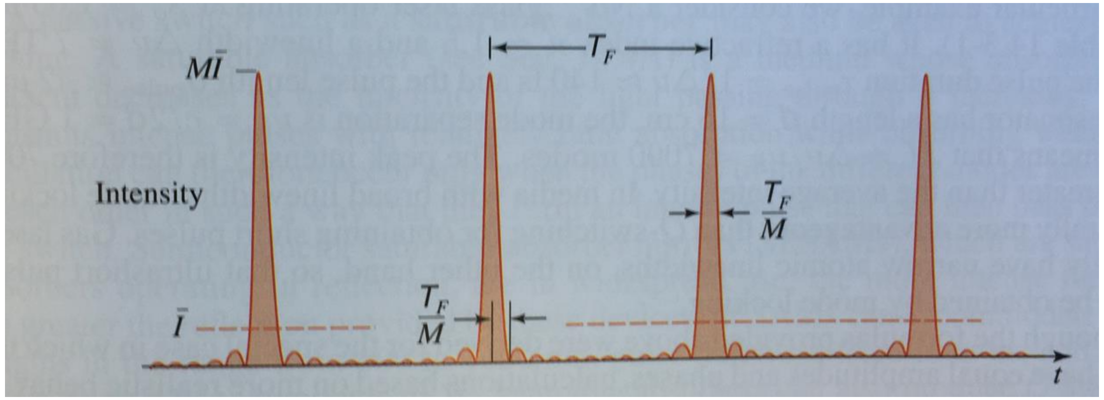
We introduced switching methods

- Gain switching
- Cavity dumping
- Q-switching

but we need to do better…

Mode locking





Expanded description of relationship between Polarization density P and Electric field E

Generally $\quad P = \varepsilon_0 \chi E \qquad \chi \rightarrow$ lets say details unknown

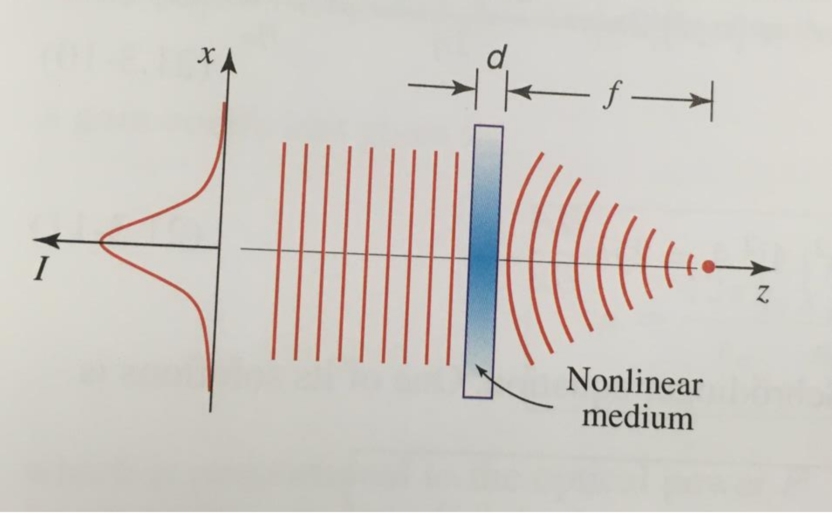
Solution $\quad$ Taylor expansion

$$P = \varepsilon_0 \left(\chi E + \chi^2 E^2 + \chi^3 E^3 + \ldots \right)$$

linear $\qquad$ 2nd order non-linearity $\qquad$ 3rd order non-linearity

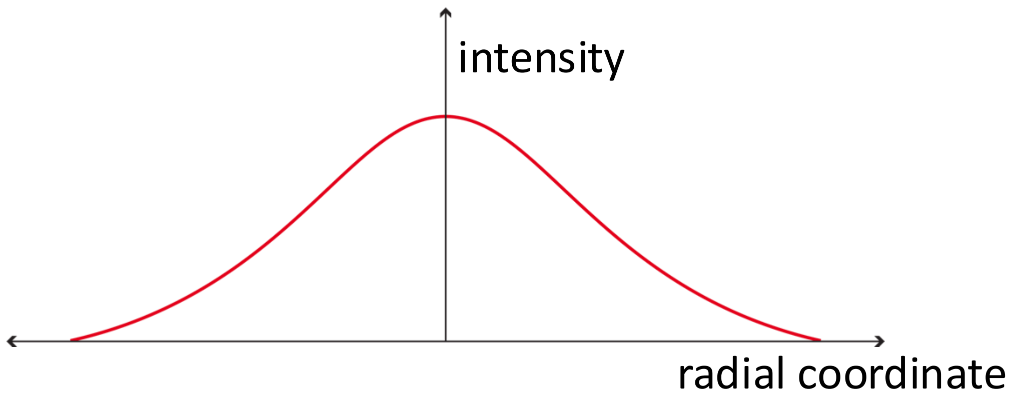
Convention $\quad P = \varepsilon_0 \chi E + 2d E^2 + 4\chi^{(3)} E^3 + \ldots$

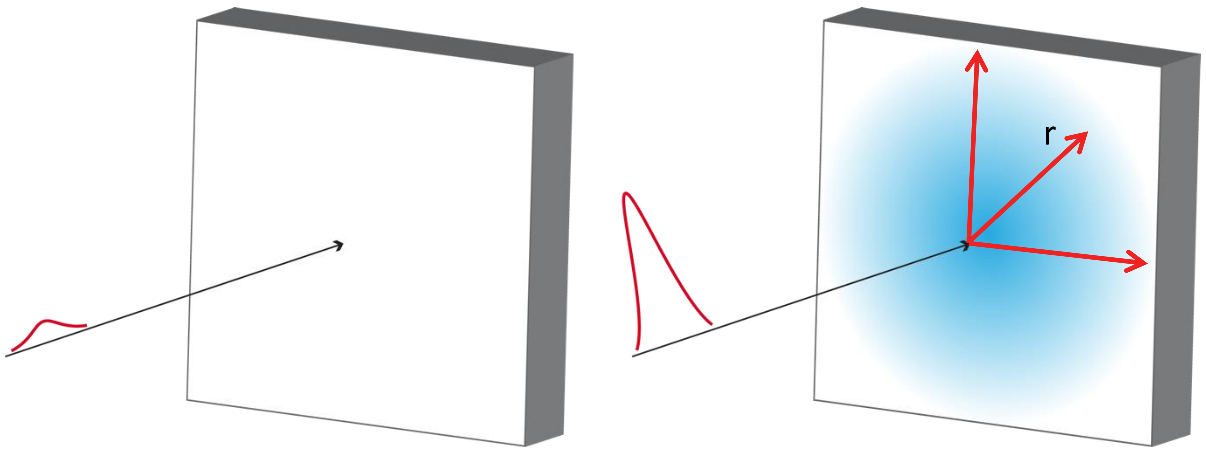
Third order non-linear optics example: Optical Kerr Effect and Self-Focusing



Kerr lens from non-linear refractive index

spatial profile:

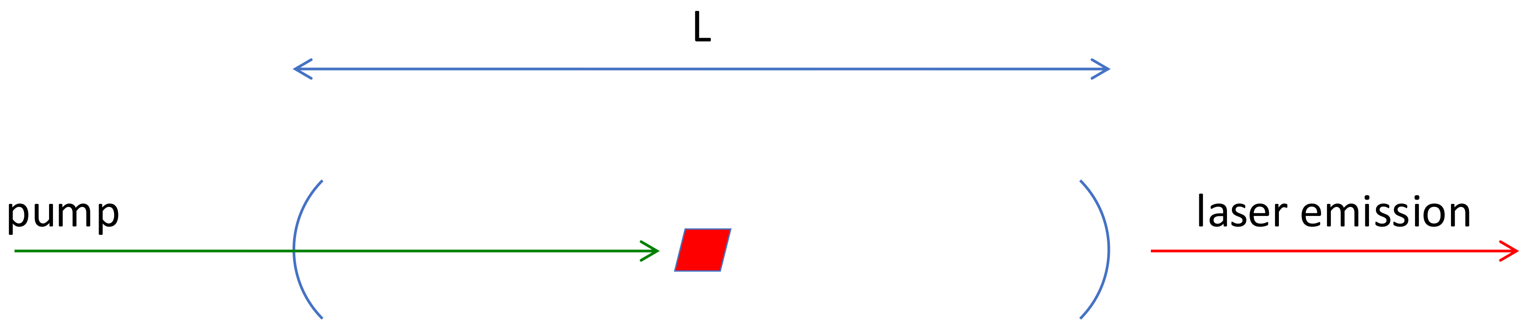




material behaves transparent

material behaves as a lens

Resonant laser cavity

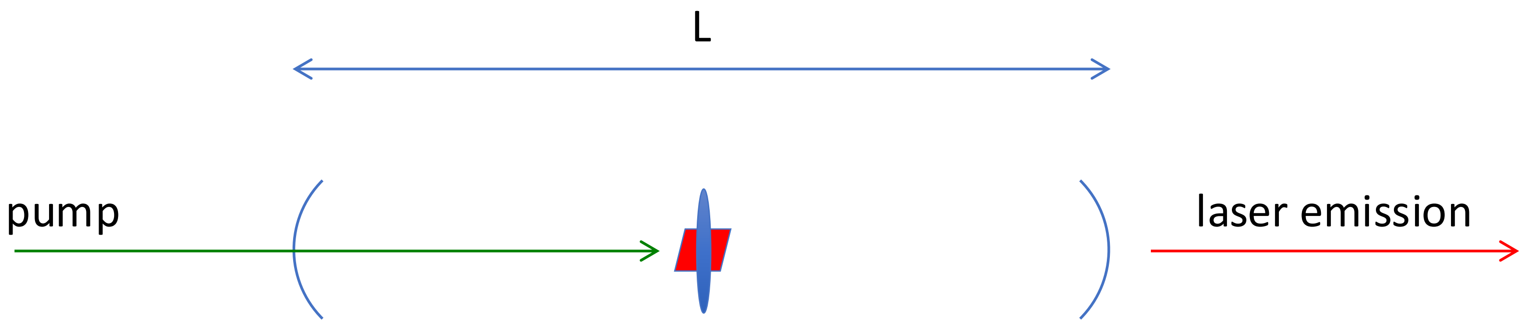


Resonant modes have nodes at cavity end mirrors

Resonant wavelengths and possible frequency modes given by:

$$\lambda_n = \frac{2L}{n} \qquad \nu_n = n\left(\frac{c}{2L}\right)$$

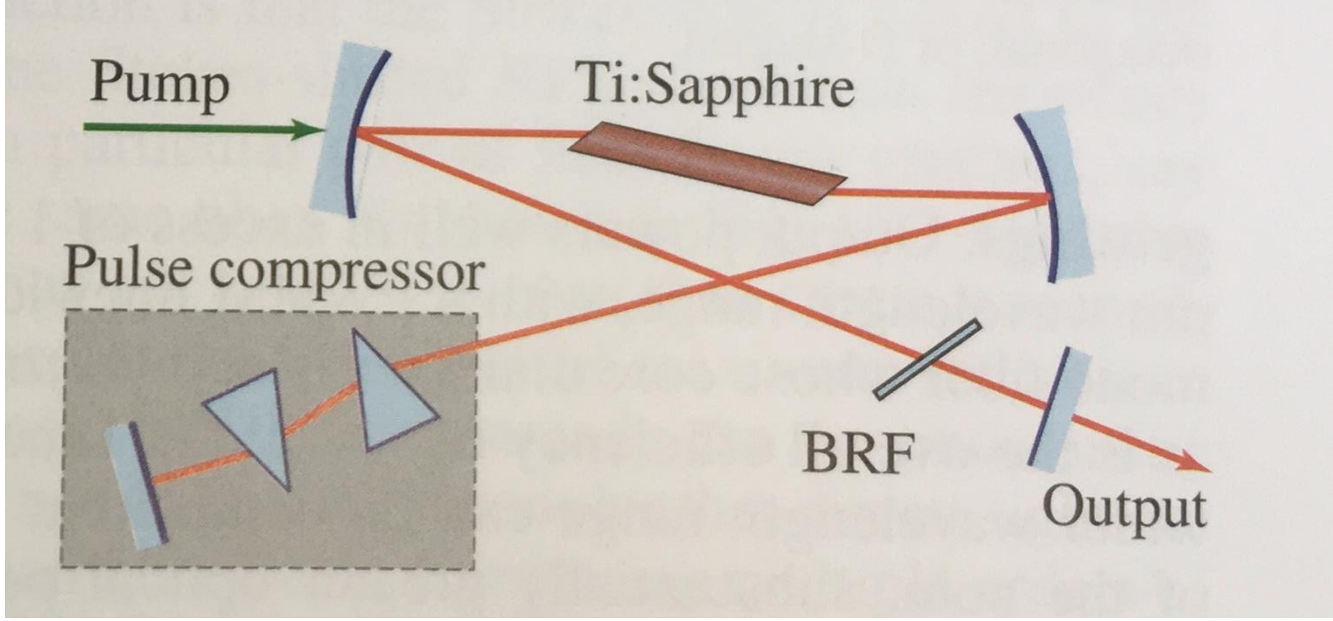
Self mode locking



- Kerr Lens Effect, due to nonlinear index of refraction
- At high intensities, the gain crystal acts like a lens
- Cavity tuned so that is most efficient with the crystal behaving as a lens
- Many modes lase and automatically arrange phases for pulsed high-intensity operation
- Intra-cavity dispersion tuned to support pulsed operation
- nJ pulse energies as short as 4 fs FWHM
- Output ~100MHz repetition rate pulse train

$$\nu_{rep-rate} = \left(\frac{c}{2L}\right)$$

TiSaphire oscillator setup



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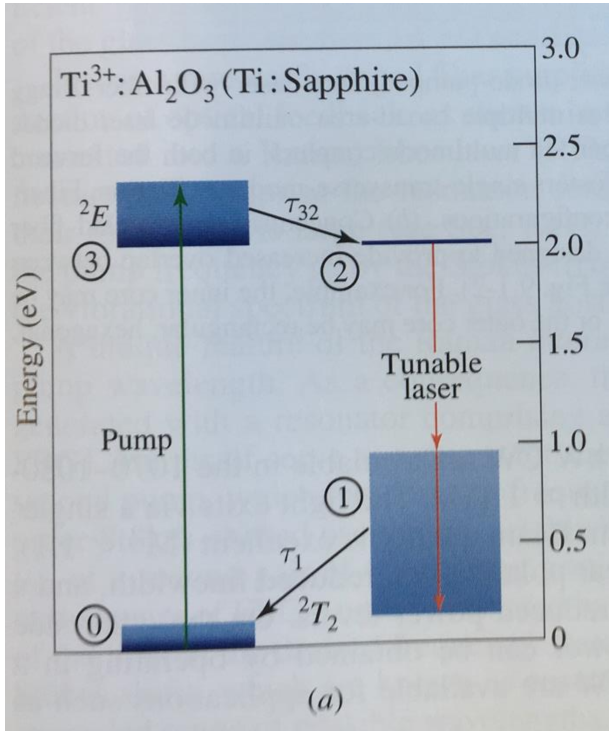
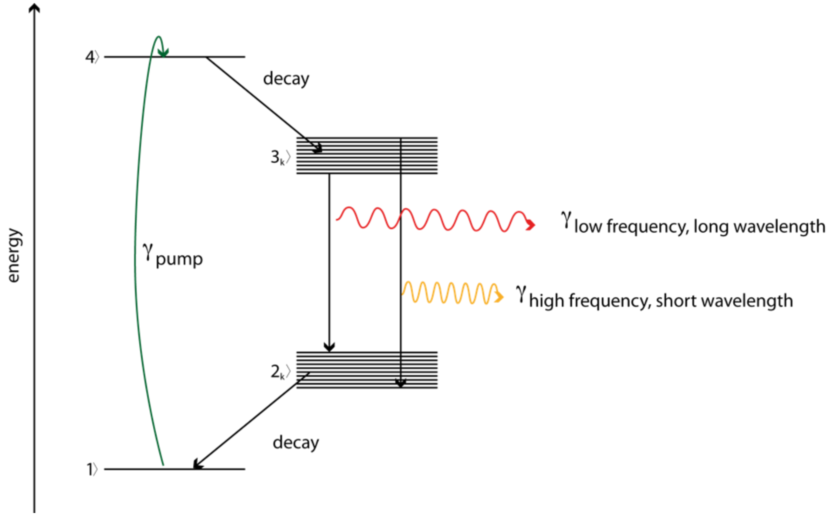
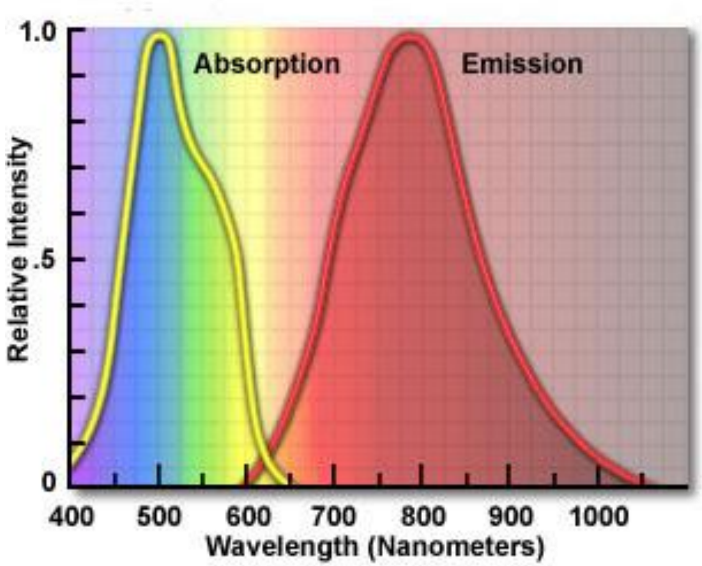
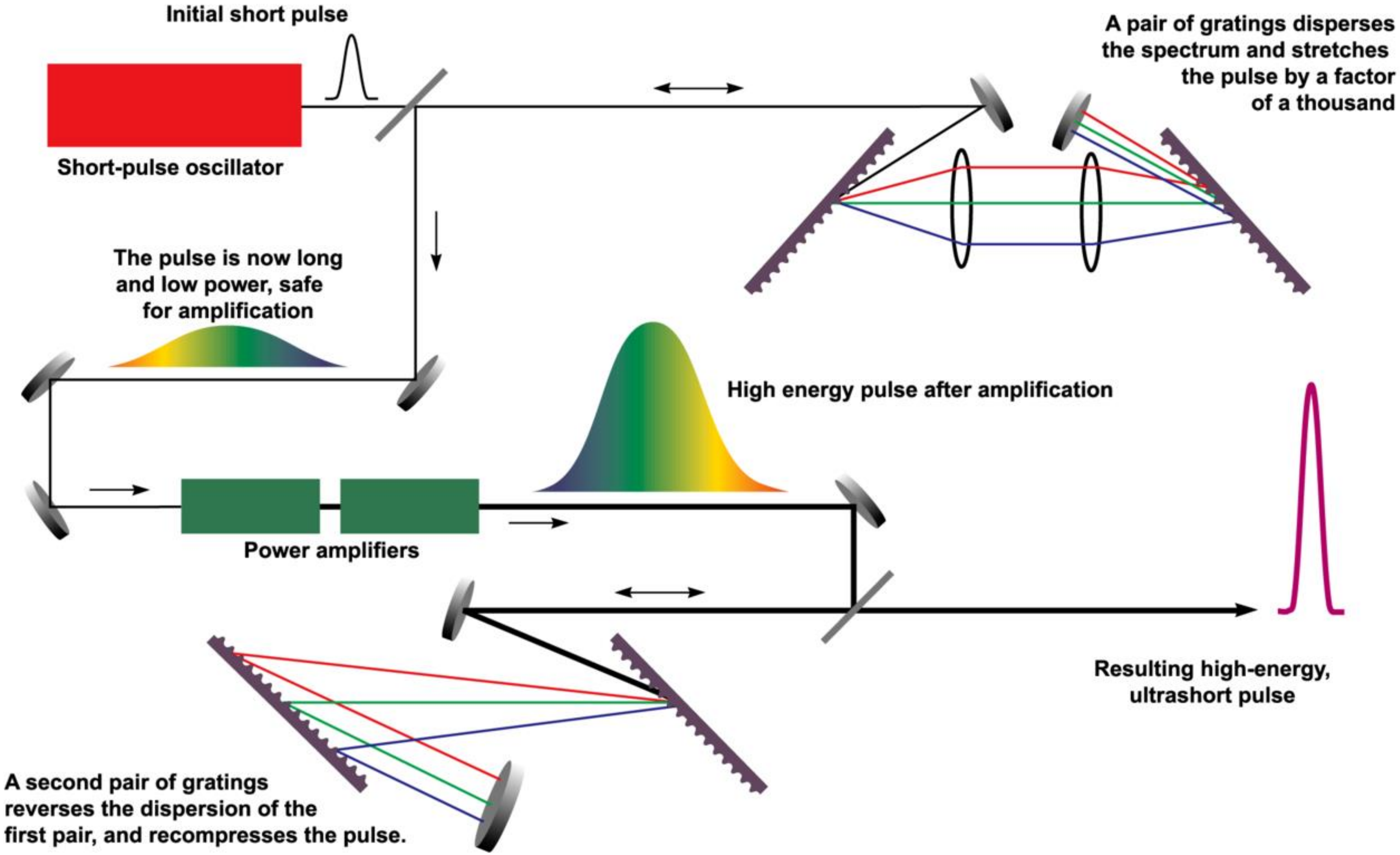


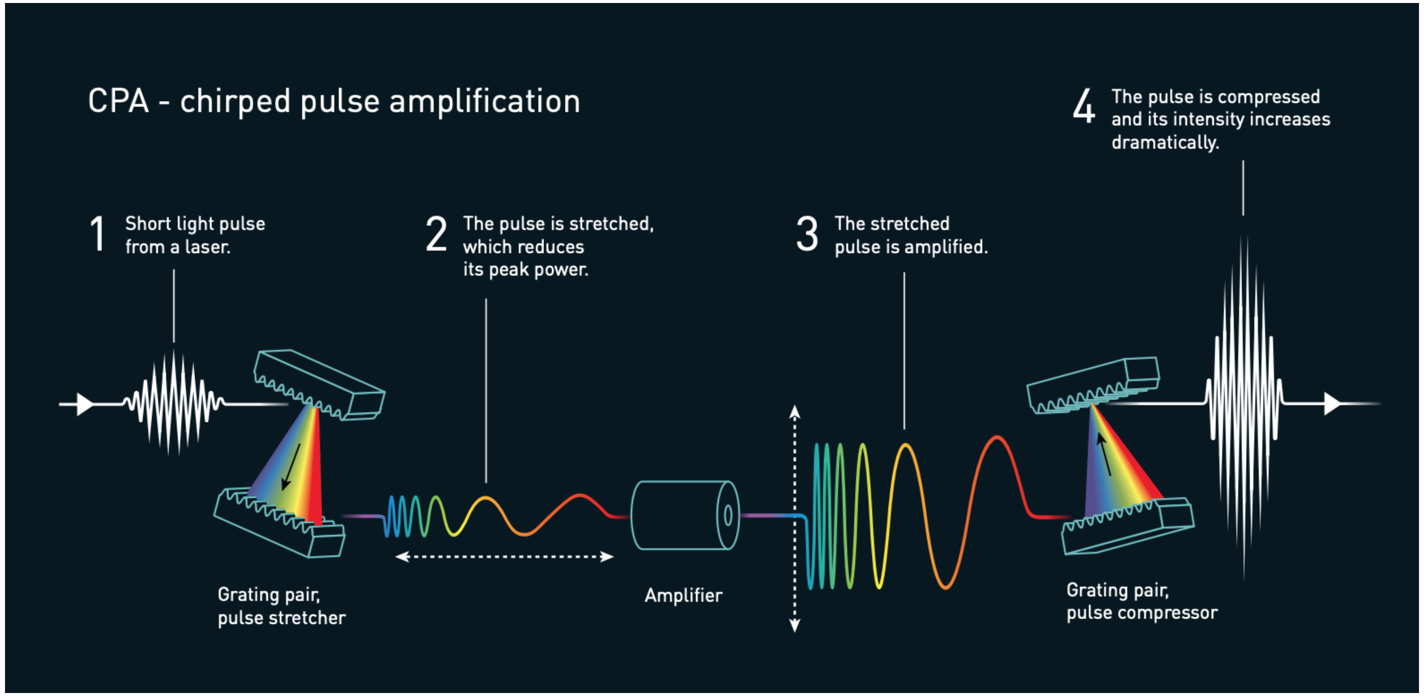




Chirped pulse amplification



Back to start

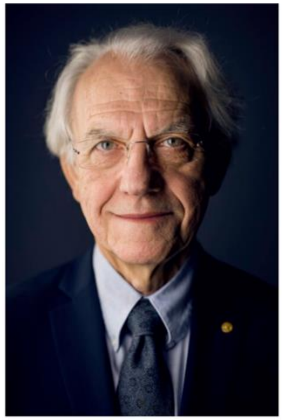


The Nobel Prize in Physics 2018

Ill. Niklas Elmehed. © Nobel Media

Arthur Ashkin

Prize share: 1/2



© Nobel Media AB. Photo: A. Mahmoud

Gérard Mourou

Prize share: 1/4

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Donna Strickland

Prize share: 1/4

The inventions being honoured this year have revolutionised laser physics. Extremely small objects and incredibly fast processes now appear in a new light. Not only physics, but also chemistry, biology and medicine have gained precision instruments for use in basic research and practical applications. Arthur Ashkin invented optical tweezers that grab particles, atoms and molecules with their laser beam fingers. Viruses, bacteria and other living cells can be held too, and examined and manipulated without being damaged. Ashkin's optical tweezers have created entirely new opportunities for observing and controlling the machinery of life. Gérard Mourou and Donna Strickland paved the way towards the shortest and most intense laser pulses created by mankind. The technique they developed has opened up new areas of research and led to broad industrial and medical applications; for example, millions of eye operations are performed every year with the sharpest of laser beams.

The Nobel Prize in Chemistry 1999

Ultrafast lasers in chemistry

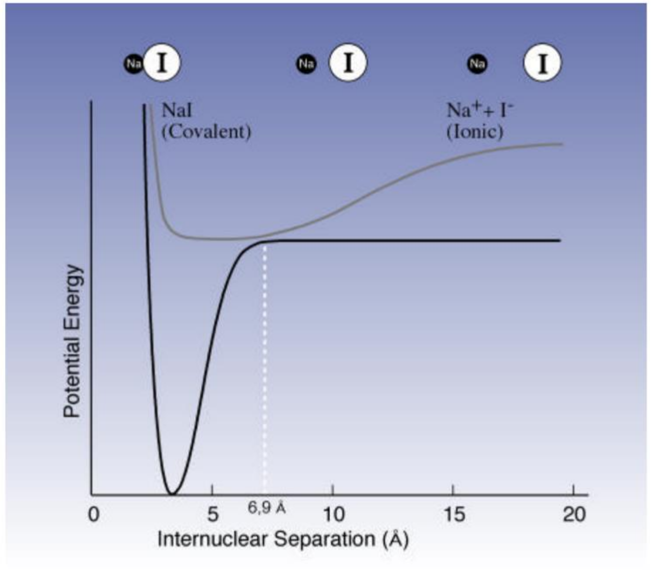


Fig. 1a) Potential energy curves showing the energies of ground state (bottom curve with deep minimum) and excited state (top curve) for NaI as function of the distance between the nuclei.

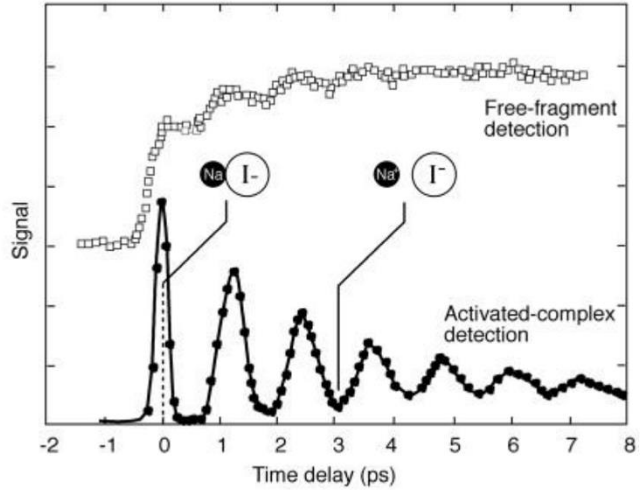


Fig. 1b) Experimental observations of coherent vibrations (so-called wave-packet motion) in femtosecond-excited NaI, on one hand manifested in terms of amount of activated complex [Na-I] at covalent (short) distance, on the other*

Photo from the Nobel Foundation archive.

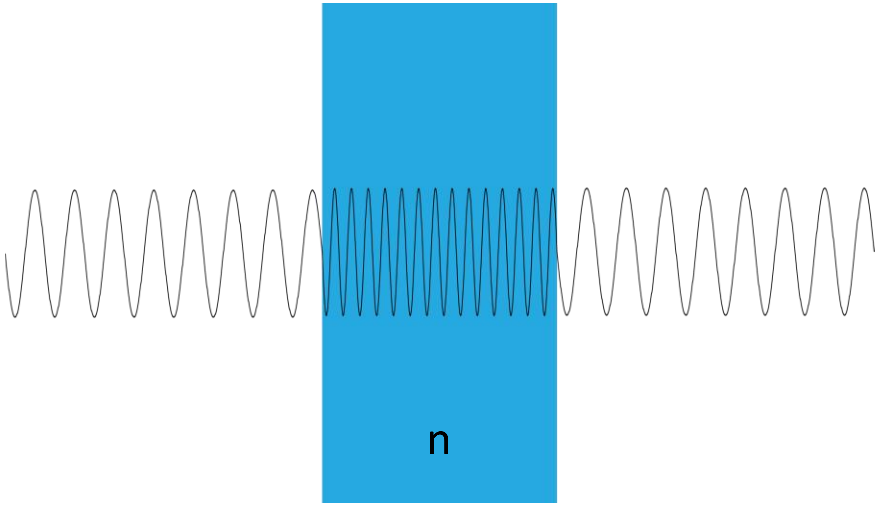
Ahmed H. Zewail

Prize share: 1/1

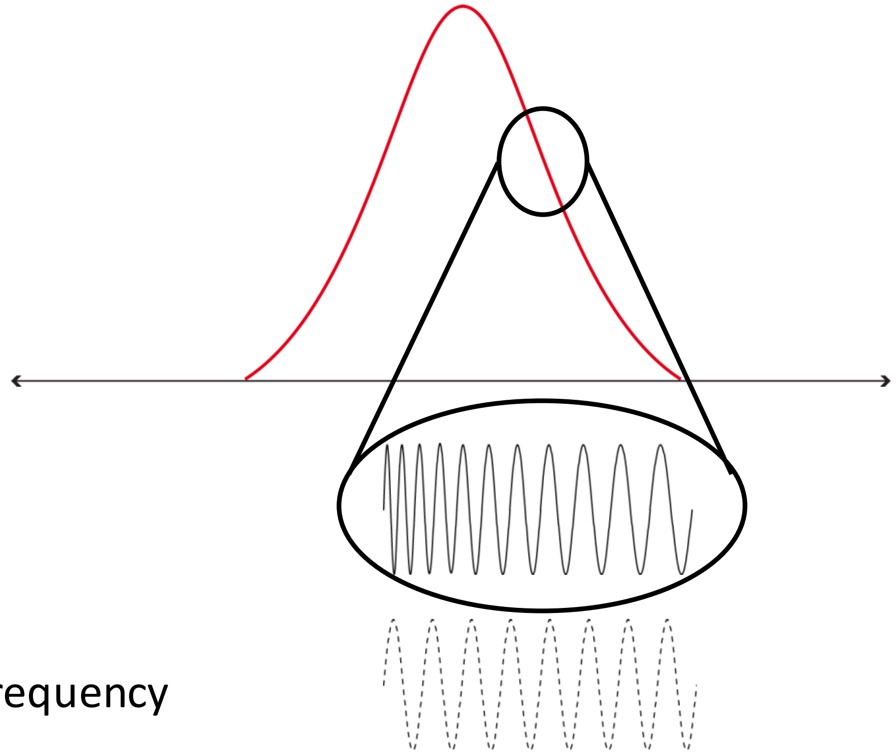
This year's Nobel Laureate, Professor Ahmed Zewail, is rewarded for his pioneering investigations of chemical reactions on the time-scale they really occur. This is the same timescale on which the atoms in the molecules vibrate, namely femtoseconds (1 fs = 10-15 seconds). Only recently have developments in laser technology enabled us to study such rapid processes, using ultra-short laser flashes. Professor Zewail's contributions have brought about a revolution in chemistry, with consequences for many other fields of science, since this type of investigation allows us to understand and predict important processes.

Third order non-linear optics example: Self-phase modulation

normal refractive index



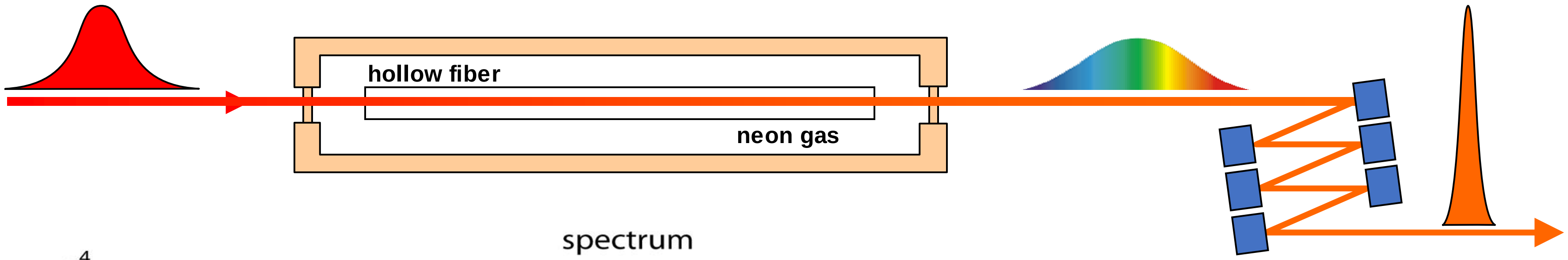
time-dependent refractive index

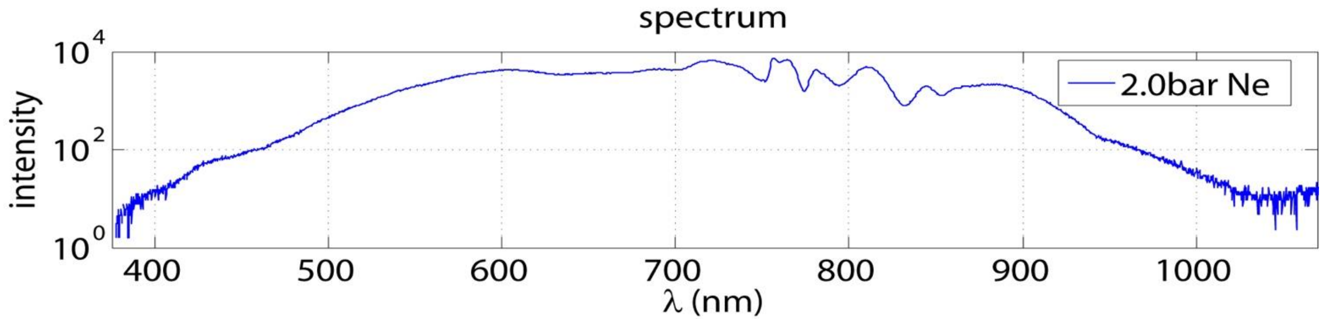


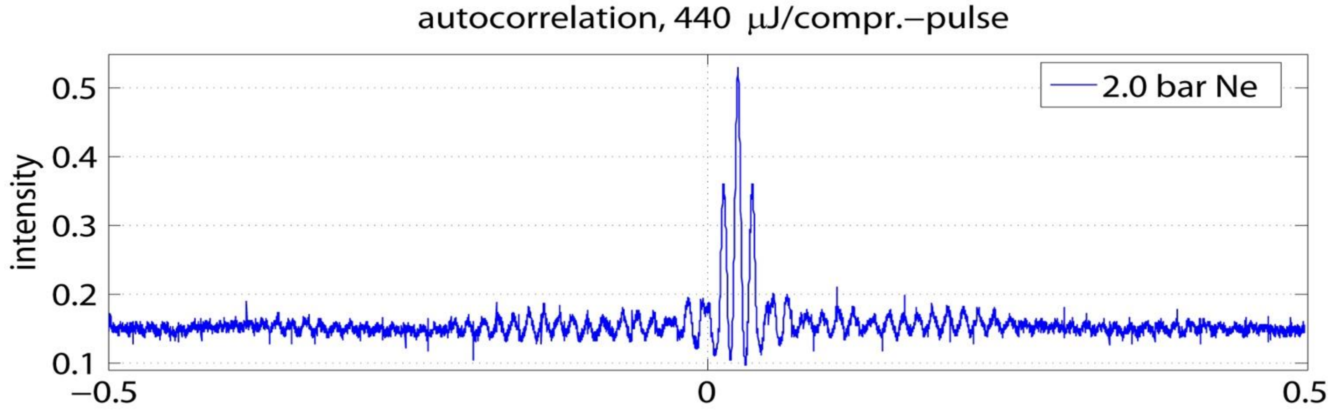
$$\omega(t) = -\frac{d\varphi(t)}{dt}$$; instantaneous frequency

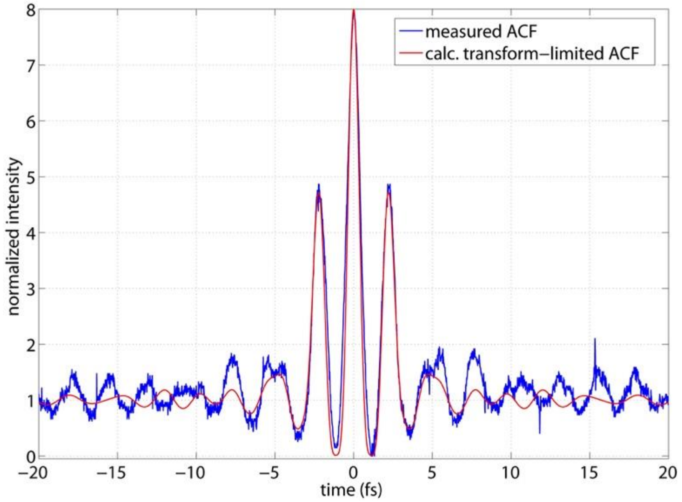
- Conceptually – consider a plane monochromatic wave
- Time-dependent index leads to time-dependent frequency
- New frequency components are generated

Post Compression by SPM & Chirped Mirror Compressor









(Adrian Cavalieri *et. al.*, *NJP* **9**, 242 (2007))

The end.

Kerr lensing needs to be analyzed in wave / beam picture! It is about accumulated phase.

From https://ocw.mit.edu/courses/electrical-engineering-and-computer-science/6-977-ultrafast-optics-spring-2005/lecture-notes/chapter7.pdf

7.1.4 The Kerr Lensing Effects

At high intensities, the refractive index in the gain medium becomes intensity dependent

$$n = n_0 + n_2 I. \tag{7.59}$$

The Gaussian intensity profile of the beam creates an intensity dependent index profile

$$I(r) = \frac{2P}{\pi w^2} \exp\left[-2\left(\frac{r}{w}\right)^2\right]. \tag{7.60}$$

In the center of the beam the index can be appoximated by a parabola

$$n(r) = n_0' \left(1 - \frac{1}{2}\gamma^2 r^2\right), \text{ where} \tag{7.61}$$

$$n_0' = n_0 + n_2 \frac{2P}{\pi w^2}, \quad \gamma = \frac{1}{w^2}\sqrt{\frac{8 n_2 P}{n_0' \pi}}. \tag{7.62}$$

A thin slice of a parabolic index medium is equivalent to a thin lens. If the parabolic index medium has a thickness t, then the ABCD matrix describing the ray propagation through the medium at normal incidence is [16]

$$M_K = \begin{pmatrix} \cos\gamma t & \frac{1}{n_0'\gamma}\sin\gamma t \\ -n_0'\gamma\sin\gamma t & \cos\gamma t \end{pmatrix}. \tag{7.63}$$

Note that, for small t, we recover the thin lens formula ($t \to 0$, but $n_0'\gamma^2 t = 1/f$ =const.). If the Kerr medium is placed under Brewster's angle, we again have to differentiate between the sagittal and tangential planes. For the sagittal plane, the beam size entering the medium remains the same, but for the tangential plane, it opens up by a factor n_0'

$$\begin{aligned} w_s &= w \\ w_t &= w \cdot n_0' \end{aligned} \tag{7.64}$$

The spotsize propotional to w^2 has to be replaced by $w^2 = w_s w_t$.Therefore, under Brewster angle incidence, the two planes start to interact during propagation as the gamma parameters are coupled together by

$$\gamma_s = \frac{1}{w_s w_t}\sqrt{\frac{8 n_2 P}{n_0' \pi}} \tag{7.65}$$

$$\gamma_t = \frac{1}{w_s w_t}\sqrt{\frac{8 n_2 P}{n_0' \pi}} \tag{7.66}$$

Without proof (see [12]), we obtain the matrices listed in Table 7.2. For low

Optical Element	ABCD-Matrix
Kerr Medium Normal Incidence	$M_K = \begin{pmatrix} \cos\gamma t & \frac{1}{n_0'\gamma}\sin\gamma t \\ -n_0'\gamma\sin\gamma t & \cos\gamma t \end{pmatrix}$
Kerr Medium Sagittal Plane	$M_{Ks} = \begin{pmatrix} \cos\gamma_s t & \frac{1}{n_0'\gamma_s}\sin\gamma_s t \\ -n_0'\gamma_s\sin\gamma_s t & \cos\gamma_s t \end{pmatrix}$
Kerr Medium Tangential Plane	$M_{Kt} = \begin{pmatrix} \cos\gamma_t t & \frac{1}{n_0'^3\gamma_t}\sin\gamma_t t \\ -n_0'^3\gamma_t\sin\gamma_t t & \cos\gamma_t t \end{pmatrix}$

Table 7.2: ABCD matrices for Kerr media, modelled with a parabolic index profile $n(r) = n_0'\left(1 - \frac{1}{2}\gamma^2 r^2\right)$.

peak power P, the Kerr lensing effect can be neglected and the matrices in Table 7.2 converge towards those for linear propagation. When the laser is mode-locked, the peak power P rises by many orders of magnitude, roughly the ratio of cavity round-trip time to the final pulse width, assuming a constant pulse energy. For a 100 MHz, 10 fs laser, this is a factor of 10^6. With the help of the matrix formulation of the Kerr effect, one can iteratively find the steady state beam waists in the laser. Starting with the values for the linear cavity, one can obtain a new resonator mode, which gives improved values for the beam waists by calculating a new cavity round-trip propagation matrix based on a given peak power P. This scheme can be iterated until there is only a negligible change from iteration to iteration. Using such a simulation, one can find the change in beam waist at a certain position in the resonator between cw-operation and mode-locked operation, which can be expressed in terms of the delta parameter

$$\delta_{s,t} = \frac{1}{p}\frac{w_{s,t}(P,z) - w_{s,t}(P=0,z)}{w_{s,t}(P=0,z)} \tag{7.67}$$

where p is the ratio between the peak power and the critical power for self-focusing

$$p = P/P_{crit}, \text{ with } P_{crit} = \lambda_L^2 / \left(2\pi n_2 n_0^2\right). \tag{7.68}$$

To gain insight into the sensitivity of a certain cavity configuration for KLM, it is interesting to compute the normalized beam size variations $\delta_{s,t}$ as a function of the most critical cavity parameters. For the four-mirror cavity, the natural parameters to choose are the distance between the crystal and the pump mirror position, x, and the mirror distance L, see Figure 7.12. Figure 7.15 shows such a plot for the following cavity parameters $R_1 = R_2 = 10$ cm, $L_1 = 104$ cm, $L_2 = 86$ cm, $t = 2$ mm, $n = 1.76$ and $P = 200$ kW.

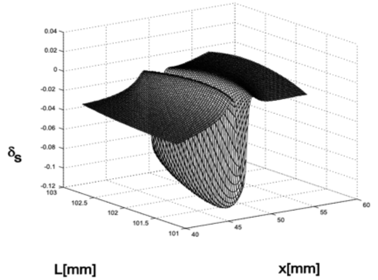


Figure 7.15: Beam narrowing ratio δ_s, for cavity parameters $R_1 = R_2 = 10$ cm, $L_1 = 104$ cm, $L_2 = 86$ cm, $t = 2$ mm, $n = 1.76$ and $P = 200$ kW