

Optical methods in chemistry
or
Photon tools for chemical sciences

Session 10:

Course layout – contents overview and general structure

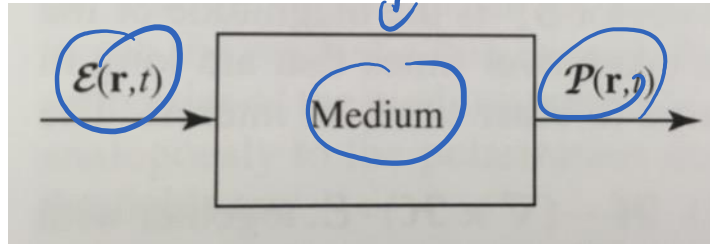
- Introduction and ray optics
- Wave optics
- Beams
- From cavities to lasers
- More lasers and optical tweezers
- From diffraction and Fourier optics
- Microscopy
- Spectroscopy
- Electromagnetic optics
- **Absorption, dispersion, and non-linear optics**
- Ultrafast lasers
- Introduction to x-rays
- X-ray diffraction and spectroscopy
- Summary

Today:

More materials properties, linear and non-linear

Electromagnetic waves in dielectric media

General



interpretation: output response

in response to an electrical field $\mathcal{E}$
medium creates a polarization

But stick with linear, instantaneous, nondispersive, homogenous, and isotropic media right now:

$$\mathcal{E} \rightarrow [\eta] \rightarrow \mathcal{P}$$

Simplifying

The equation $\mathcal{P} = \epsilon_0 \chi \mathcal{E}$ is shown inside a box. The symbol χ is circled in blue.

χ scalar

$$\mathcal{E} \rightarrow [\chi] \rightarrow \mathcal{P}$$

$$\mathcal{D} \approx \epsilon \mathcal{E} \quad | \quad \epsilon = \epsilon_0 (1 + \chi)$$

electric susceptibility

$$\epsilon = \epsilon_0 (1 + \chi)$$

dielectric constant³

This leads to the following Maxwell and wave equations

$$\begin{aligned}\nabla \times \mathcal{H} &= \epsilon \frac{\partial \mathcal{E}}{\partial t} \\ \nabla \times \mathcal{E} &= -\mu \frac{\partial \mathcal{H}}{\partial t} \\ \nabla \cdot \mathcal{E} &= 0 \\ \nabla \cdot \mathcal{H} &= 0.\end{aligned}$$

ϵ - dielectric constant
 μ - magnetic permeability

back to wave equation

$$\nabla^2 u - \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2} = 0 \quad \left. \vphantom{\nabla^2 u} \right\} \text{in media}$$

Homework

familiarize
yourself with
this level of
e-dynamics

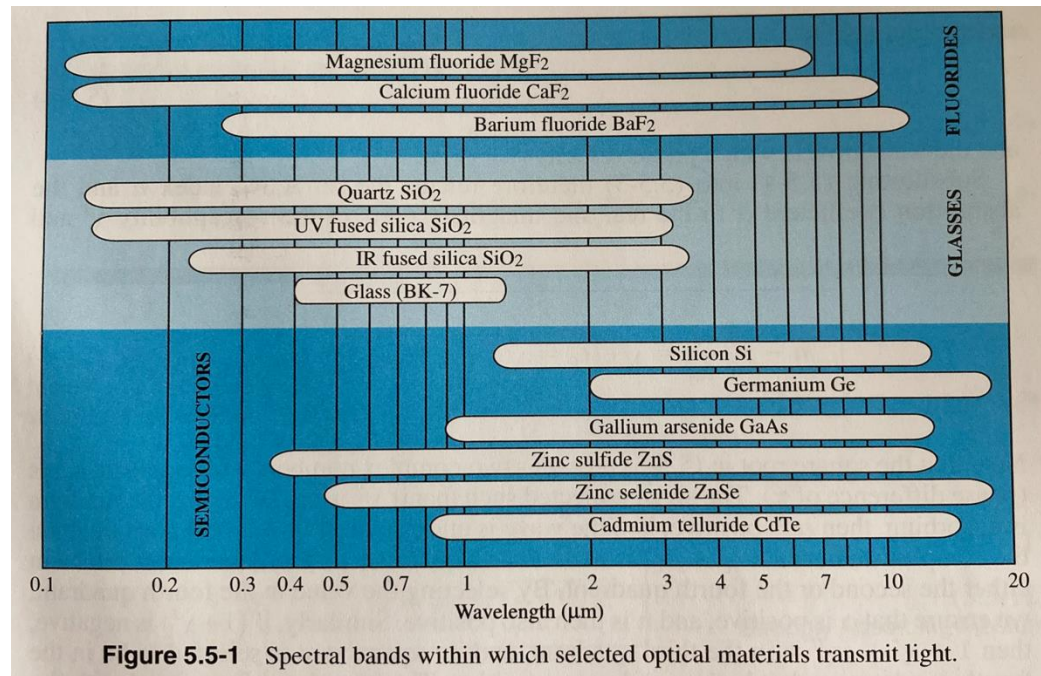
$$c = \frac{1}{\sqrt{\epsilon \mu}}$$

define $n = \frac{c_0}{c} = \sqrt{\frac{\epsilon}{\epsilon_0} \cdot \frac{\mu}{\mu_0}}$ = refractive index

non-magnetic materials

$$n = \sqrt{\frac{\epsilon}{\epsilon_0}} = \sqrt{1 + \chi}$$

Generalized optical constant



Absorption: The imaginary part

Complex refractive index

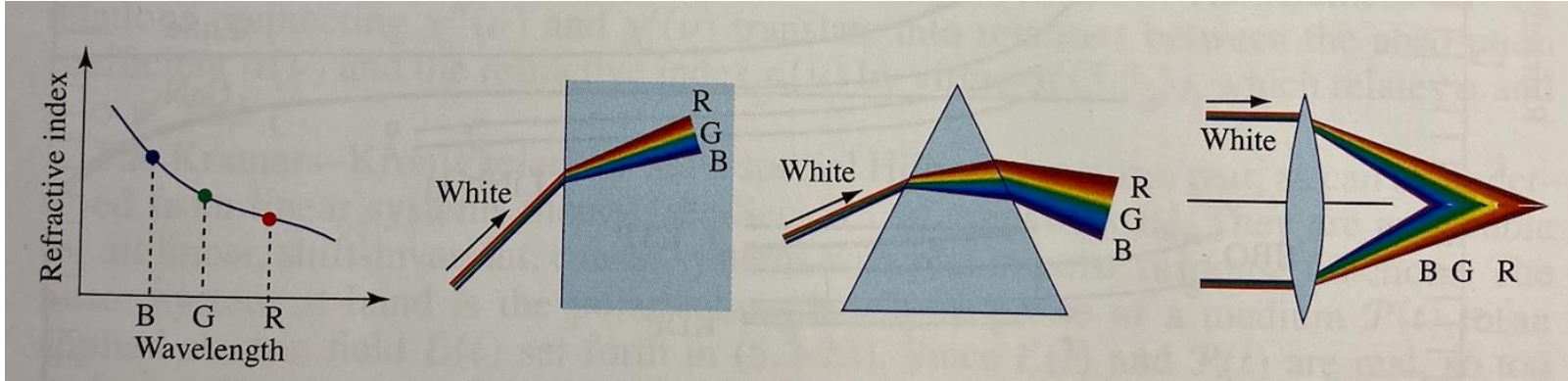
General:

$$n - j\frac{1}{2}\frac{\alpha}{k_o} = \sqrt{\epsilon/\epsilon_o} = \sqrt{1 + \chi' + j\chi''}.$$

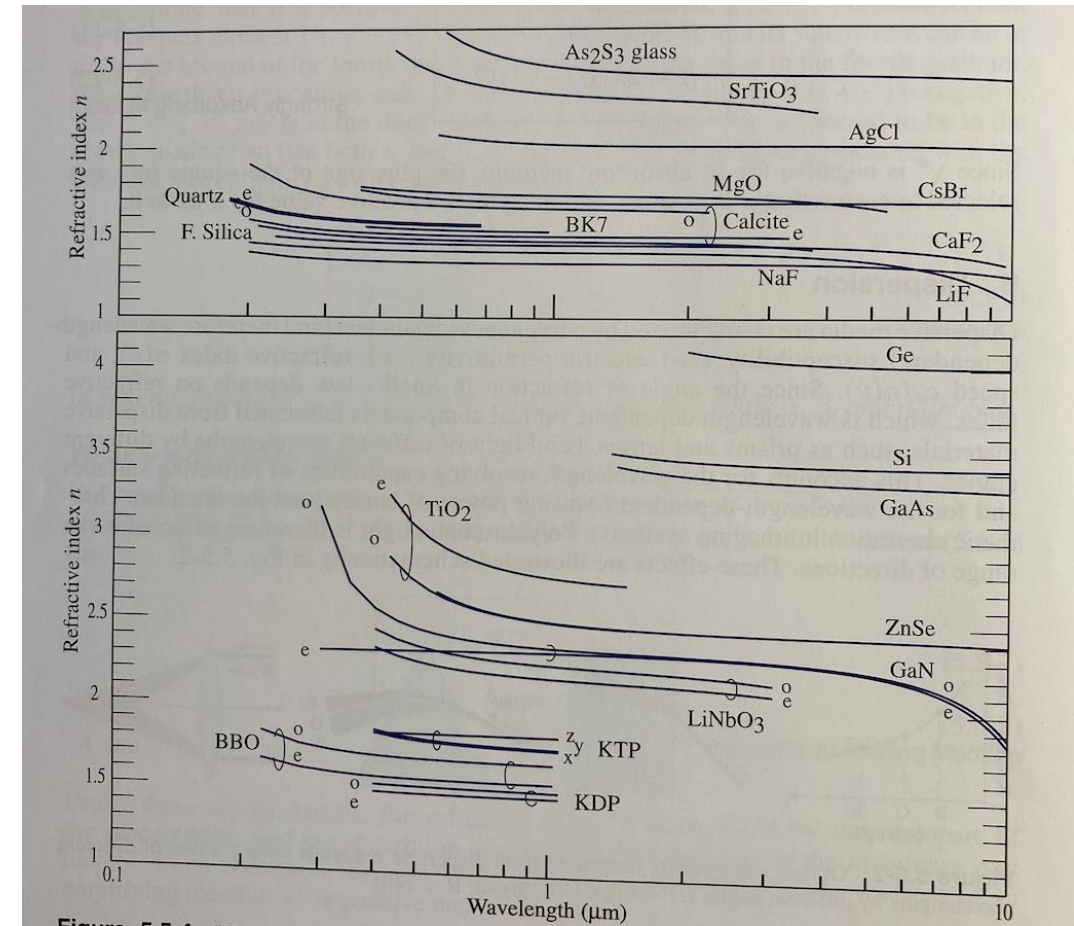
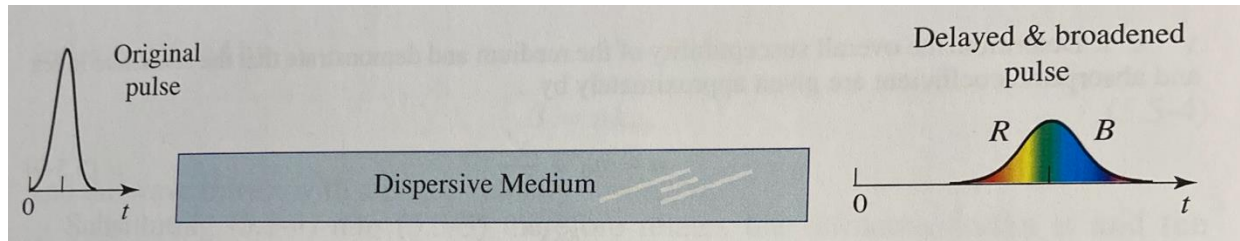
Weakly absorbing media:

$$\begin{aligned} n &\approx \sqrt{1 + \chi'} \\ \alpha &\approx -\frac{k_o}{n}\chi''. \end{aligned}$$

Dispersion: The real part



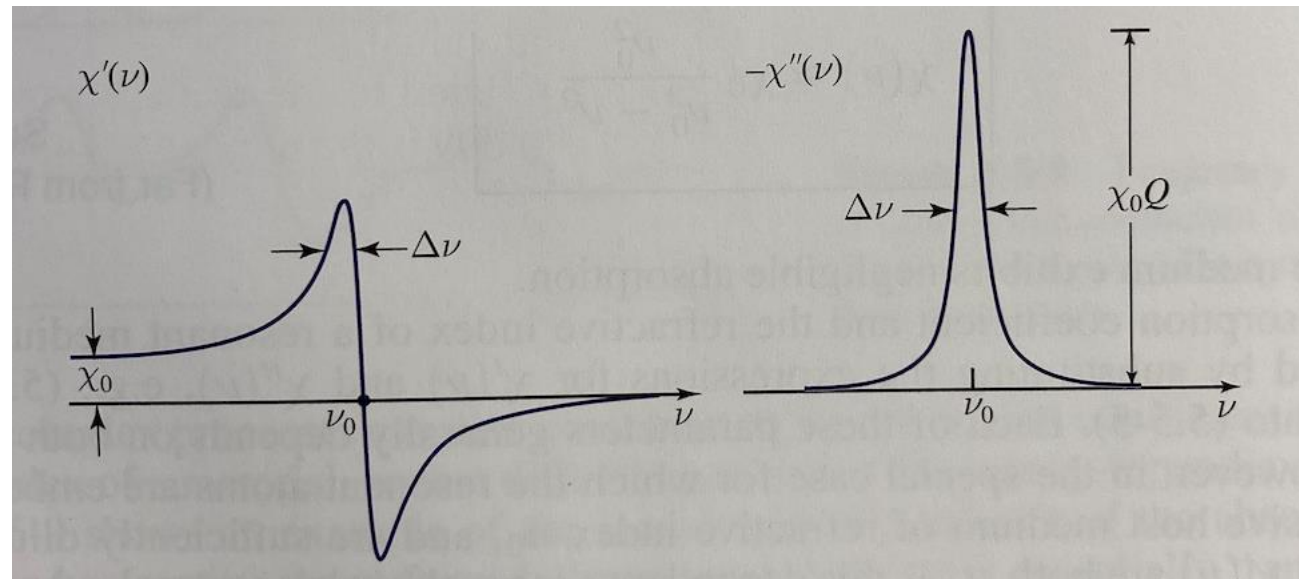
A short pulse in a dispersive medium



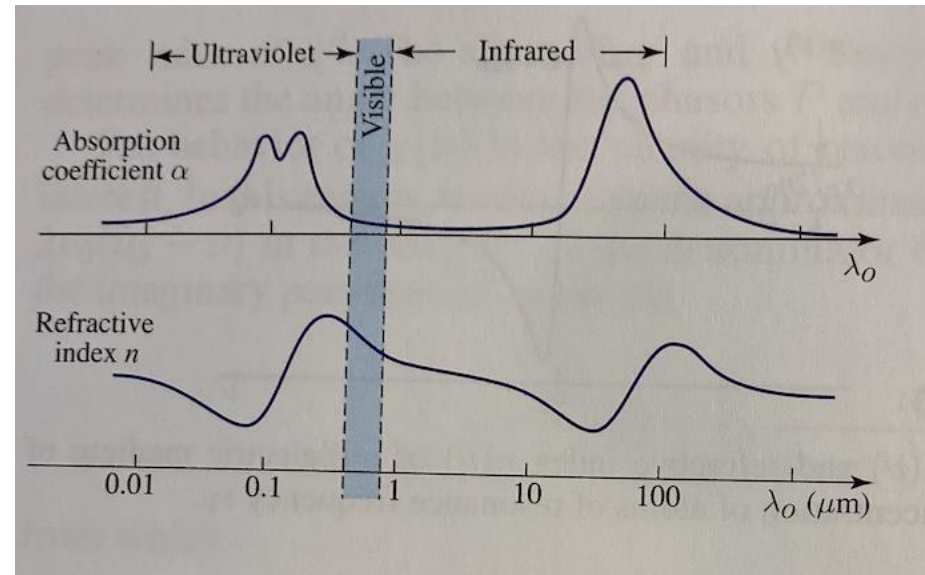
The Kramers-Kronig Relation

$$\chi'(\nu) = \frac{2}{\pi} \int_0^\infty \frac{s\chi''(s)}{s^2 - \nu^2} ds$$
$$\chi''(\nu) = \frac{2}{\pi} \int_0^\infty \frac{\nu\chi'(s)}{\nu^2 - s^2} ds.$$

Resonances and refractive index

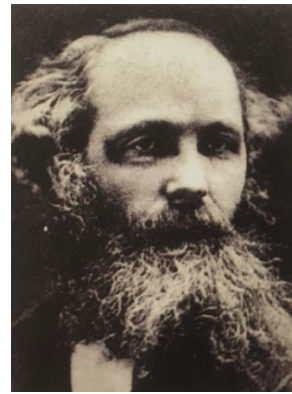


A real optical (transparent) material



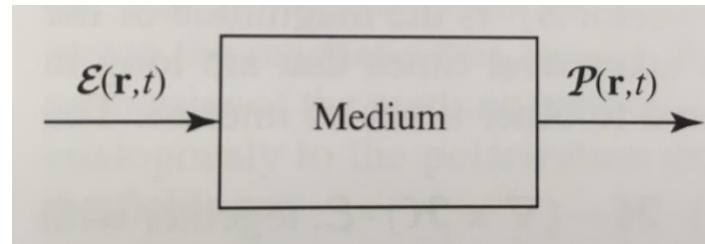
From linear to non-linear

Back to formal description of light as EM wave (from Session 9)



James Maxwell
1831 - 1879

$$\begin{aligned}\nabla \times \mathcal{H} &= \epsilon_o \frac{\partial \mathcal{E}}{\partial t} \\ \nabla \times \mathcal{E} &= -\mu_o \frac{\partial \mathcal{H}}{\partial t} \\ \nabla \cdot \mathcal{E} &= 0 \\ \nabla \cdot \mathcal{H} &= 0,\end{aligned}$$



But stick with linear, nondispersive,
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now:

$$\mathcal{P} = \epsilon_o \chi \mathcal{E},$$

$$\begin{aligned}\nabla \times \mathcal{H} &= \epsilon \frac{\partial \mathcal{E}}{\partial t} \\ \nabla \times \mathcal{E} &= -\mu \frac{\partial \mathcal{H}}{\partial t} \\ \nabla \cdot \mathcal{E} &= 0 \\ \nabla \cdot \mathcal{H} &= 0.\end{aligned}$$

Generalization of susceptibility χ (still linear)

- Inhomogeneous media
- Anisotropic media
- Dispersive media

General:

Interpretation: Dynamic relationship between E and P

- E induces bound electrons in material to oscillate
- Time-dependent Polarization density $P(t)$
- Time-delay between $E(t)$ and $P(t)$

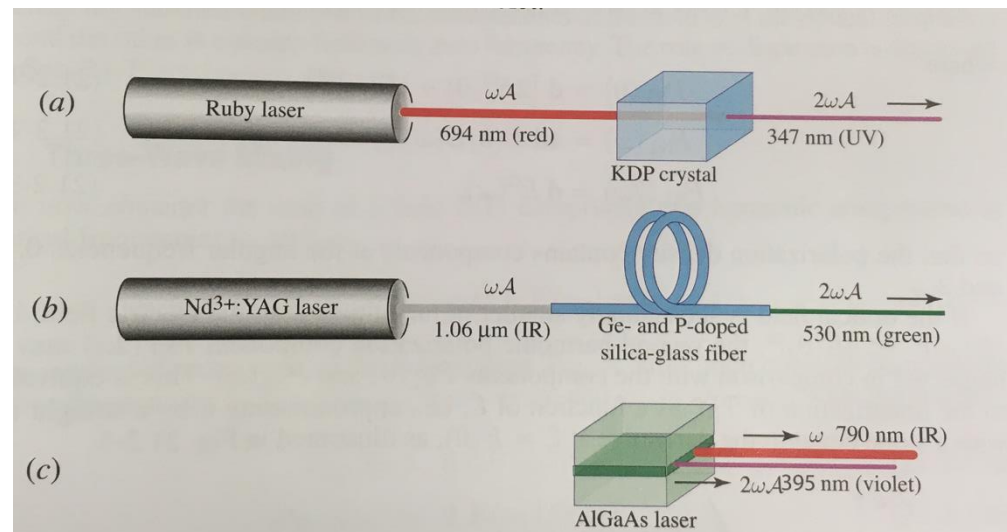
Non-linear optical media

Handwaving:

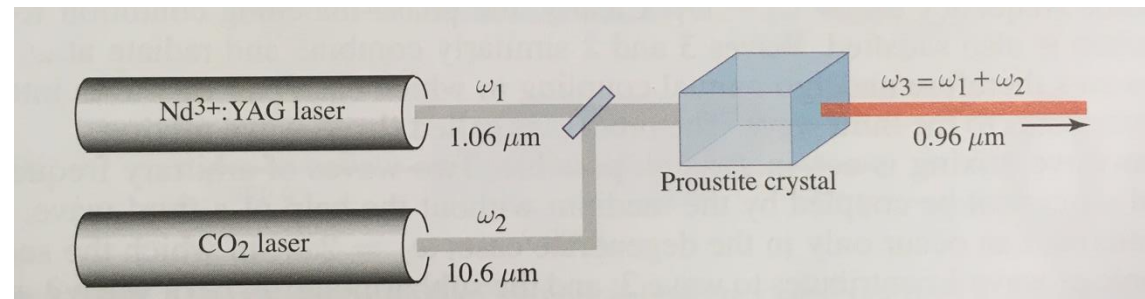
- Linear: restoring force of light induced fields linear (“Hookes law applies”)
- Non-linear: Light induced fields comparable to inter-atomic fields in crystal (“no more linear forces”)
- (Note: fields still weak compared to intra-atomic fields – that is a later topic)

Expanded description of relationship between Polarization density P and Electric field E

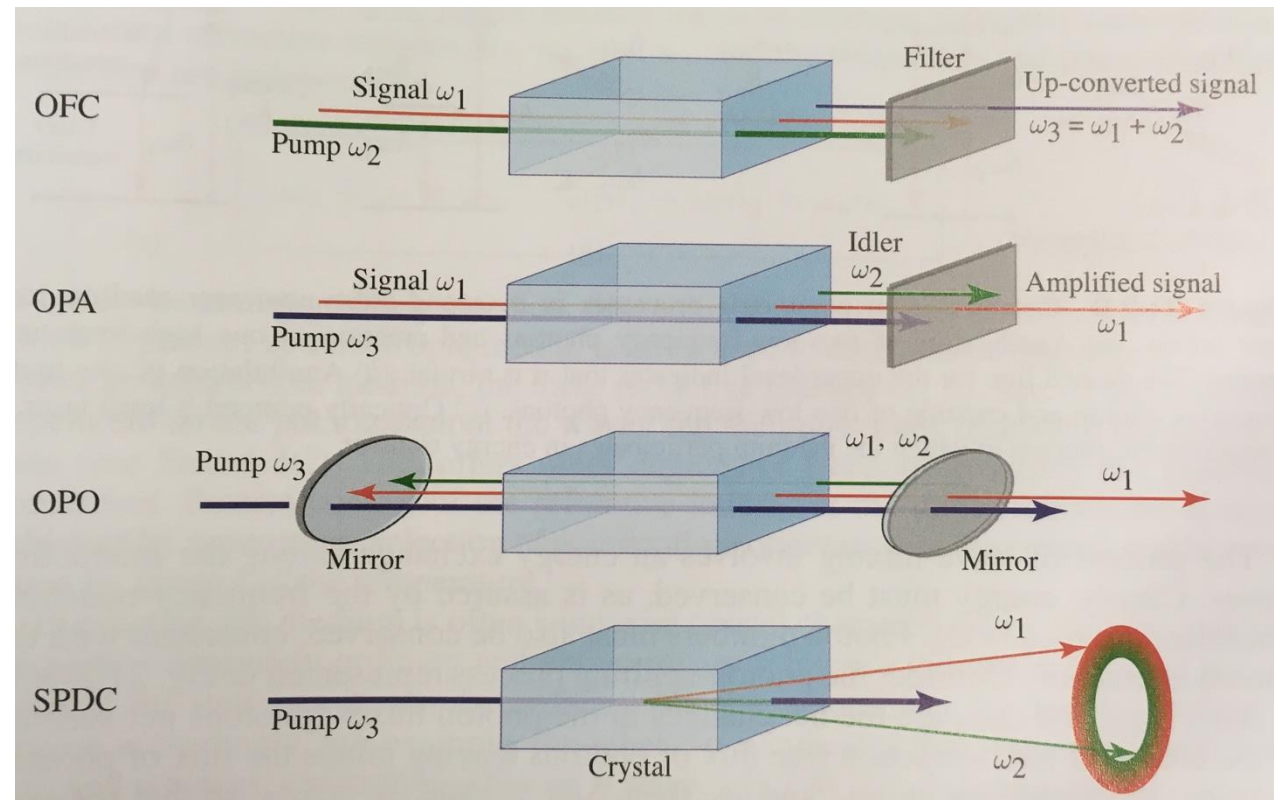
Second order non-linear optics example: Second harmonic generation



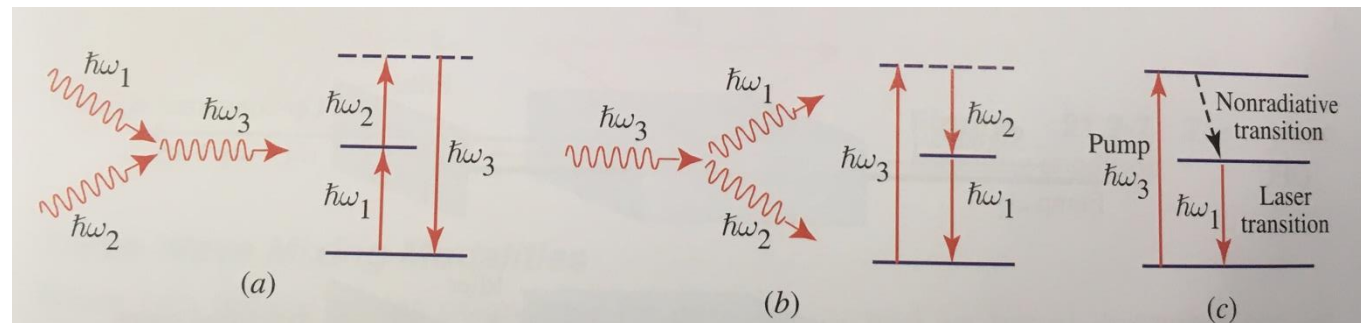
Second order non-linear optics example: Sum frequency generation



Second order non-linear optics example: Optical parametric devices

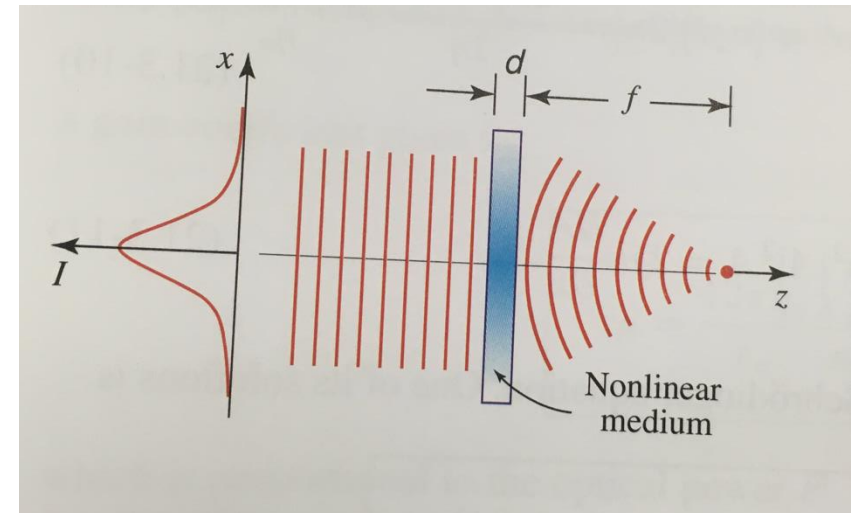


Second order non-linear optics example: Description as photon interaction process



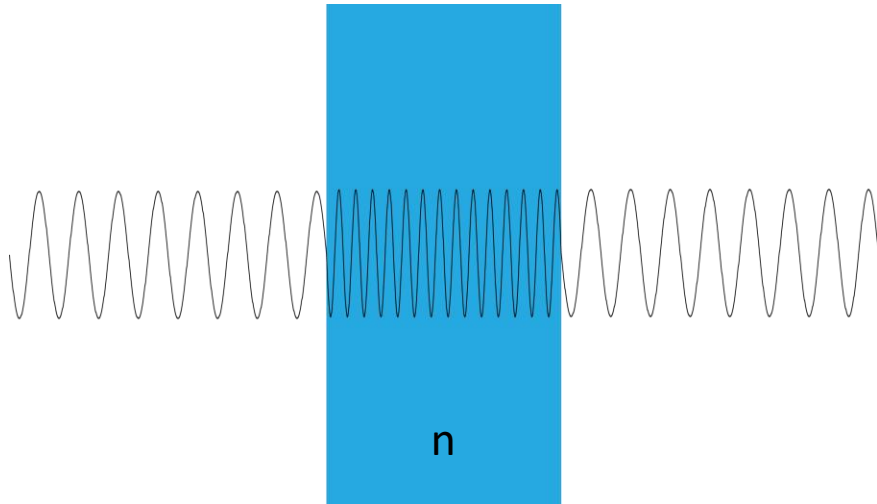
Third order non-linear optics example: Third-harmonic generation

Third order non-linear optics example: Optical Kerr Effect and Self-Focusing

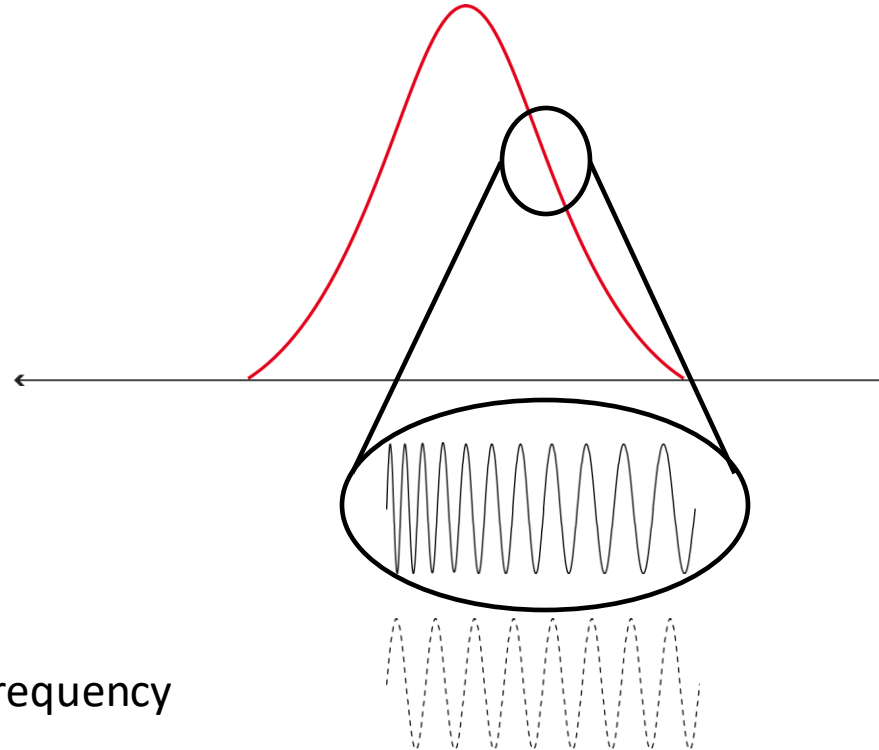


Third order non-linear optics example: Self-phase modulation

normal refractive index



time-dependent refractive index



$$\omega(t) = - \frac{dj(t)}{dt} ; \text{instantaneous frequency}$$

- Conceptually – consider a plane monochromatic wave
- Time-dependent index leads to time-dependent frequency
- New frequency components are generated

The end.