

Structural Analysis

Part III - X-ray tools

Session 3

X-ray scattering and diffraction

Intermezzo: Fourier Transform

Remember

Alternative approach: Fourier Transform (Infrared) Spectrometer

- Not a dispersive measurement but based on interferometry
- Use all wavelength at the same time
- Manipulation in x - cm (real space) to get information in wave numbers $x^{-1} - \text{cm}^{-1}$ (frequencies)
- Measure interferogram, Fourier transform into spectrum

Relationships

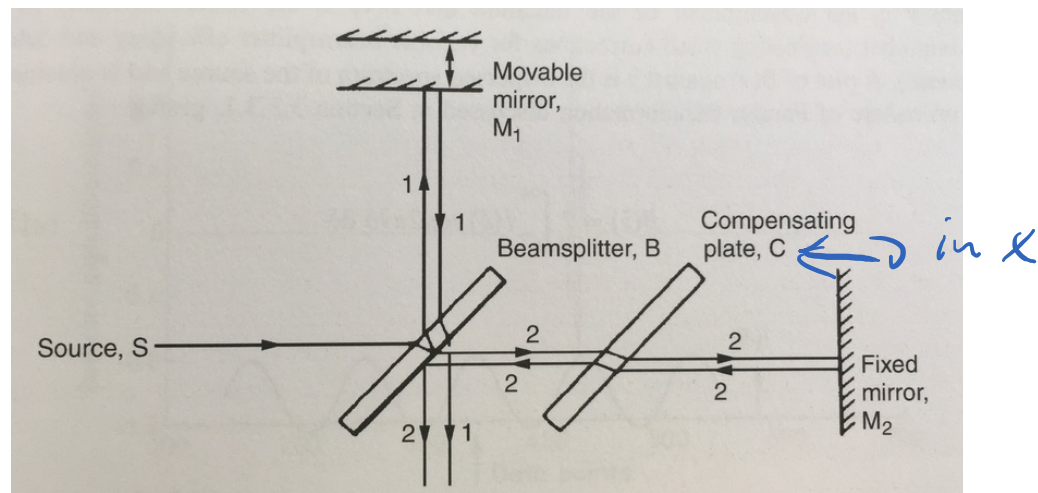
$$\sim \text{cm} \sim \frac{1}{\text{cm}}$$

real $\sim$ inverse
space



Measuring an interferogram

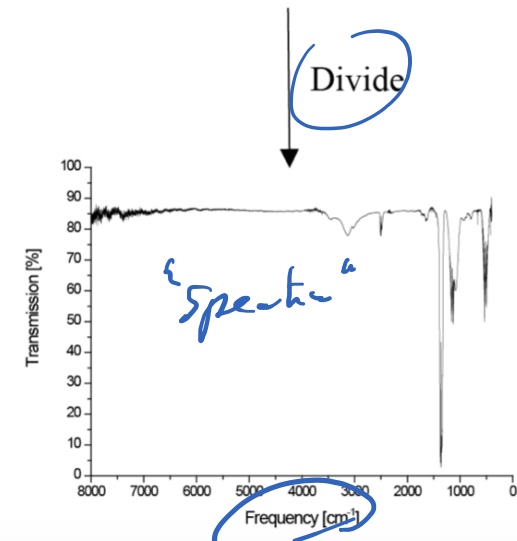
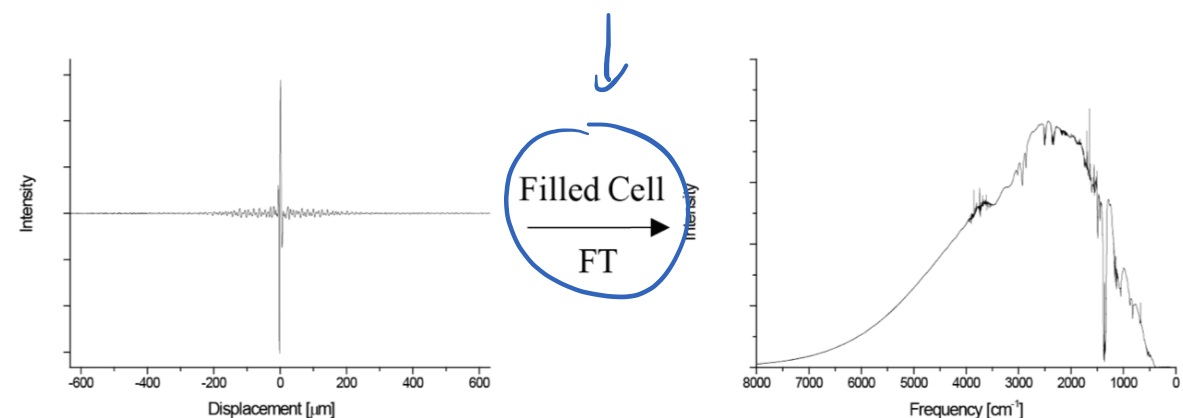
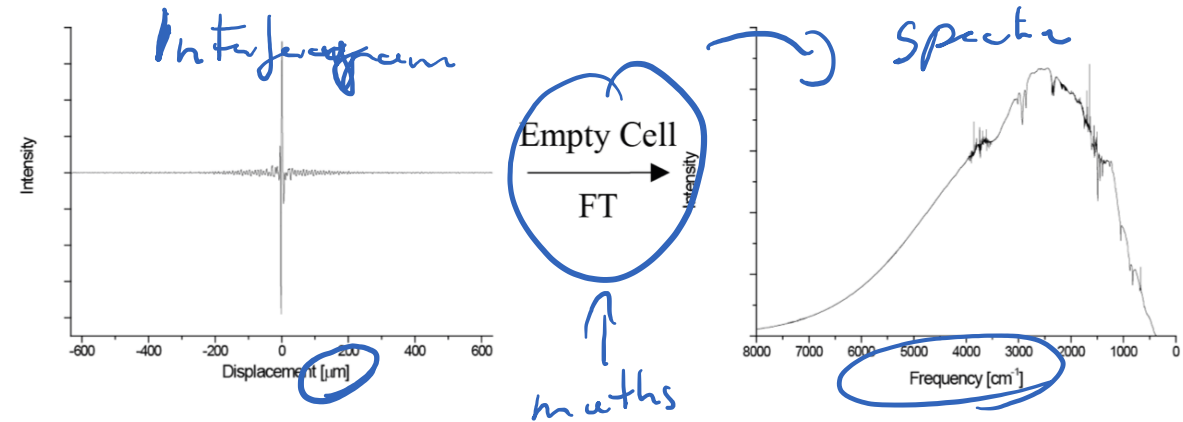
Fourier transform $\rightarrow$ relationship between frequency & real space



Sample

Detector

Computer $\rightarrow$ FT $\rightarrow$ Spectra in $\frac{1}{x}$ [frequency] space
 $\downarrow$
 Energy



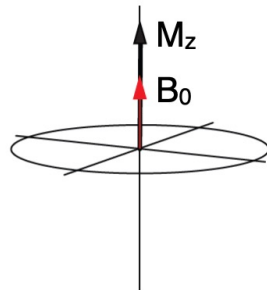
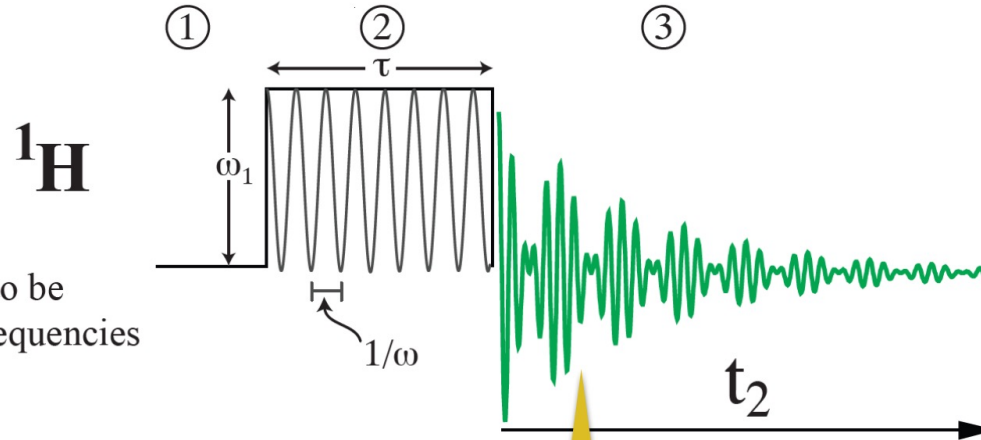
Prof. Emsley part of this class, 5 weeks ago!

Pulsed FTNMR Spectroscopy

$$\omega_1 = -\gamma B_1$$

$$\omega_1 \tau = \pi/2$$

ω = carrier frequency,
chosen by the operator to be
near to the resonance frequencies

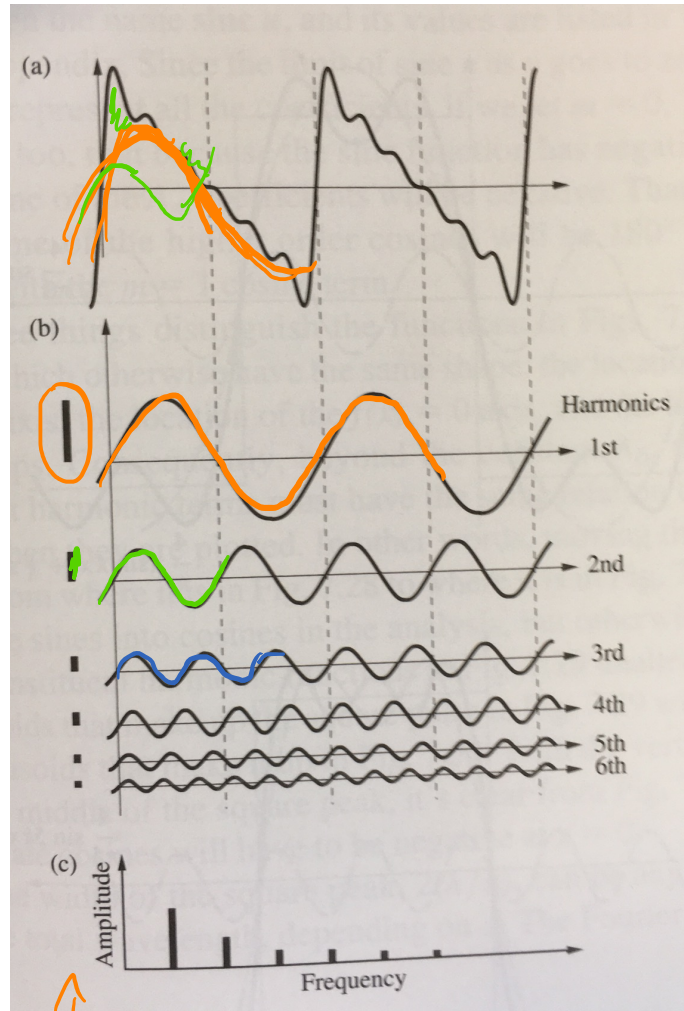


1. Equilibrium. The net magnetization is aligned along the direction of the main field (z-axis).

2. A field is applied in the transverse plane. The magnetization of the ensemble precesses around the field.

3. The field is removed leaving a net transverse component of the ensemble magnetization. This *coherence* then starts to precess around the main field.

Principle idea



↑
Amplitude

Fourier Series

Fourier series $\rightarrow$
infinite sum of harmonic functions that
can describe any / another function

Use $\sum$ of harmonic fct
to approximate any function

use different amplitudes for each harmonic fct
($\rightarrow$ frequencies)

Mathematical description

- The principle idea of the Fourier analysis / transformation is that any function can be represented by an (infinite) series of harmonic functions.
- The Fourier transform decomposes a function into its constituent frequencies.

Thus we can write

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} F(k) e^{-ikx} dk$$

provided that

$$F(k) = \int_{-\infty}^{+\infty} f(x) e^{ikx} dx$$

Fourier transform

frequencies

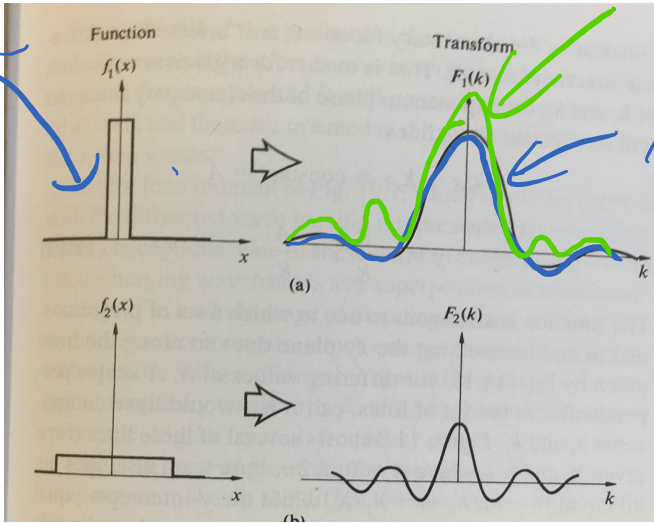
harmonic
fct

Fourier transform is the
"rule how amplitudes & frequencies
come together"

Note Computers do Fourier Transforms really, well
 $\Rightarrow$ FFT $\leftrightarrow$ GPU

Examples $\rightarrow$ selected FT

Square functions



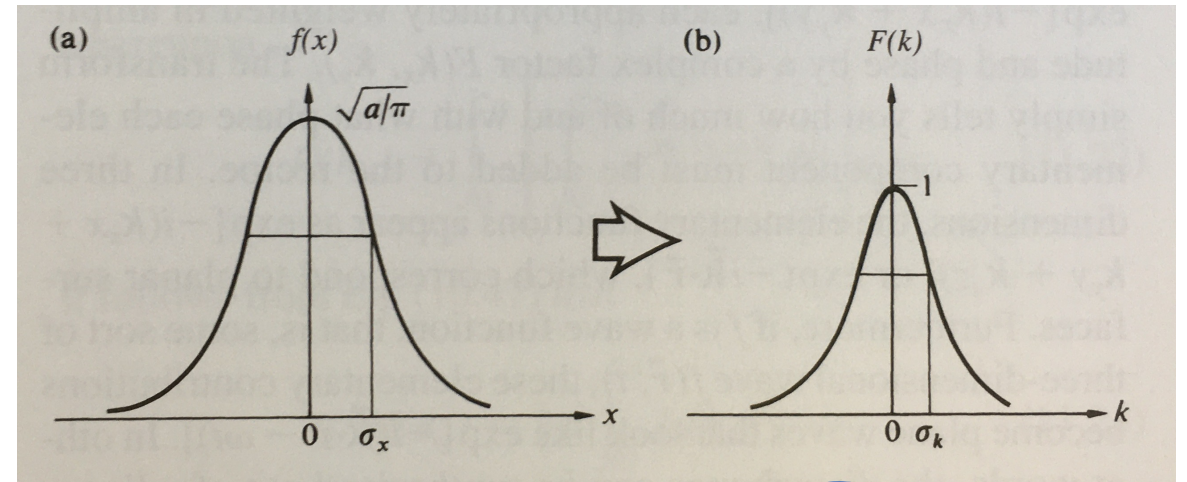
could be transmission how slit

Intensity $|A|^2$

looks like interference pattern

↓
Remember in experiments we do not measure waves but but intensities
 $I = |A|^2$

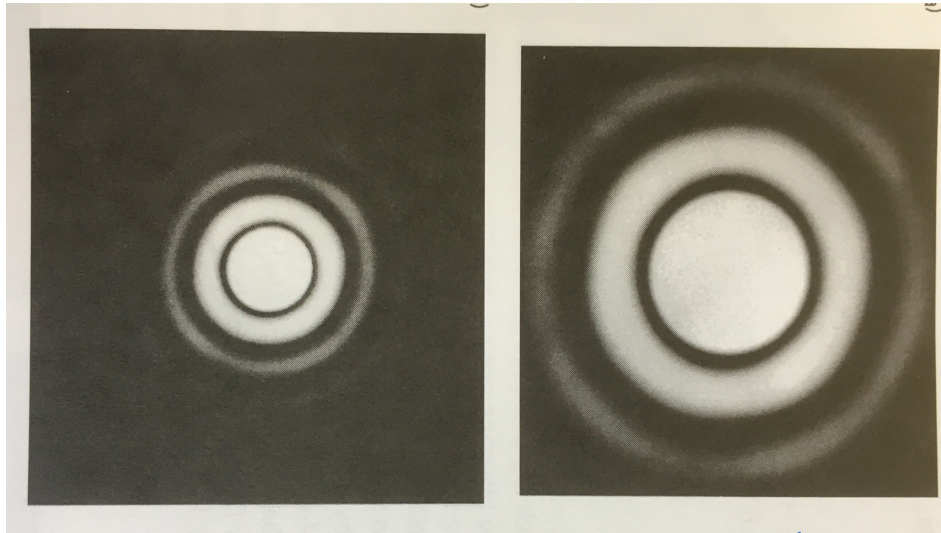
Gauss function



FT "Gaussian" $\rightarrow$ "Gaussian"

Fourier transform of round aperture and airy pattern

Experiment: Diffraction pattern of circular aperture, Airy pattern

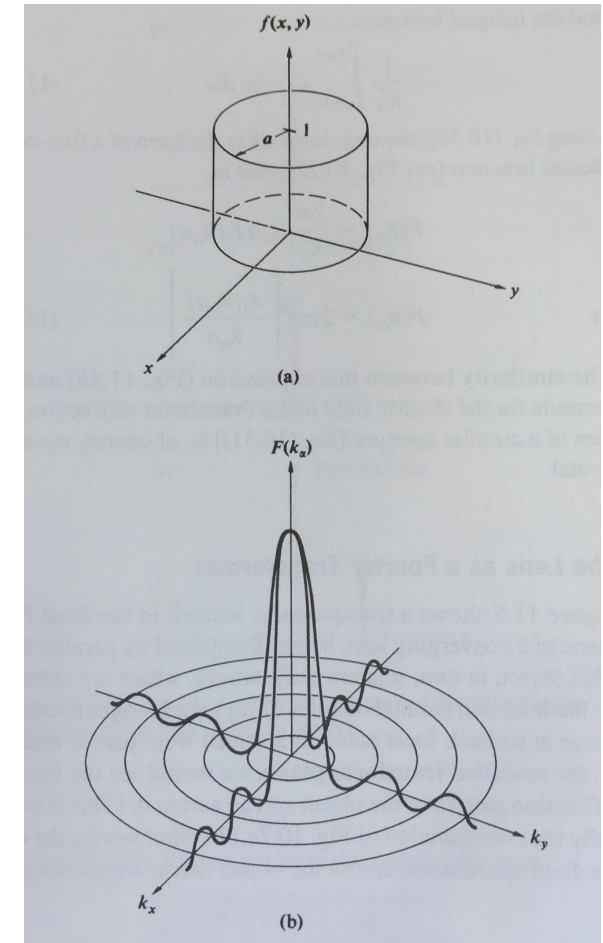


↳ Bessel J_1

General

The diffraction pattern of an object is described by its Fourier transform

Theory: Fourier transform of cylinder or “top-hat” function



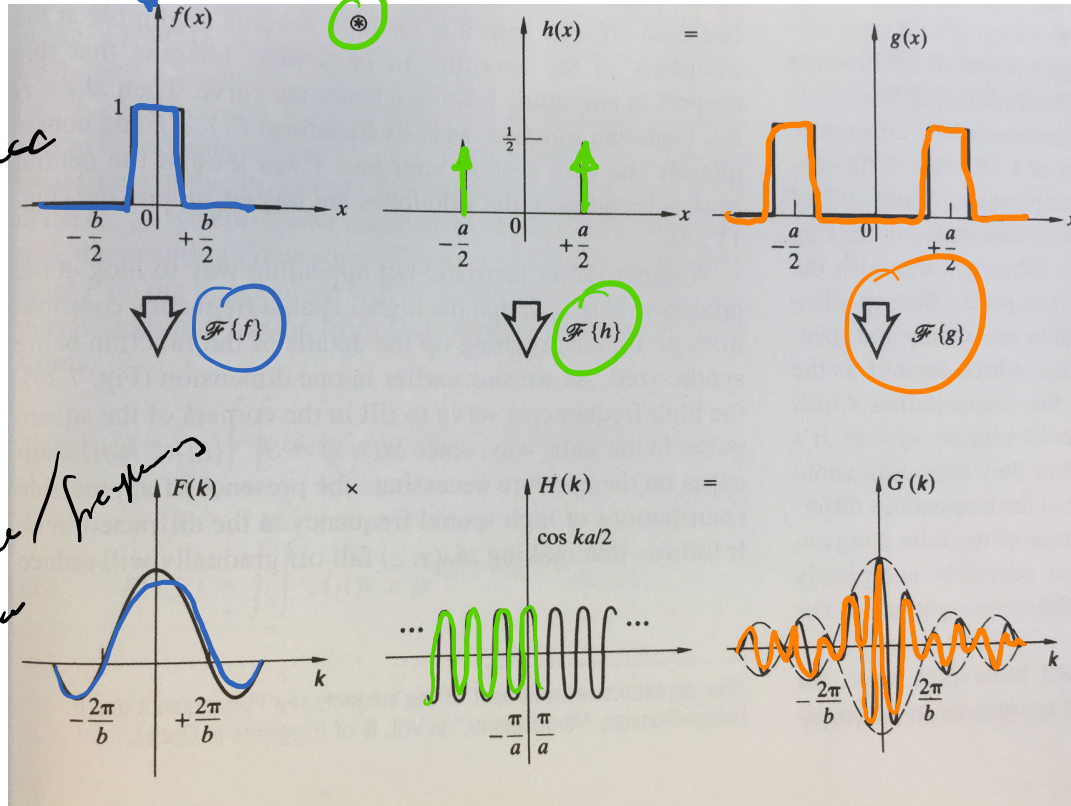
Diffraction as Fourier transform:



single slit

δ -fcn

Real space



Inverse space

sinc fcn
slit width

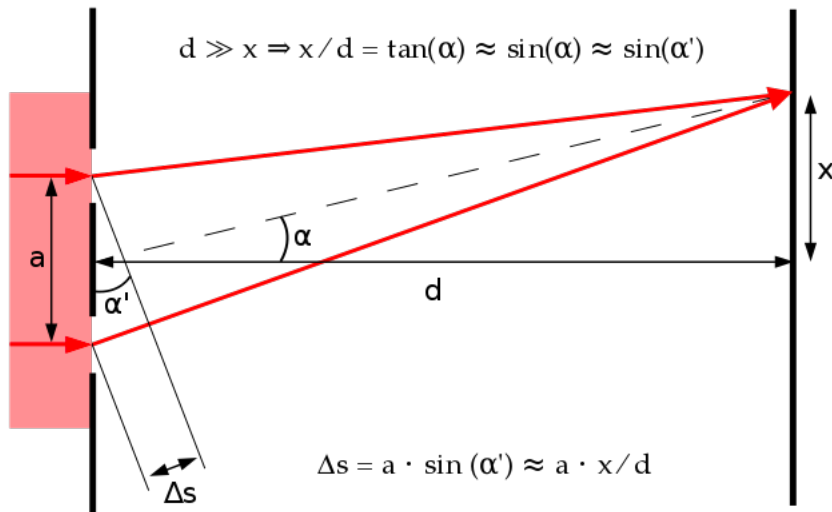
Convolution theorem
FT of a convoluted fcn $\rightarrow$ point product of their FT
From maths

$$FT(f \otimes g) = FT(f) \cdot FT(g)$$

Reconstructed the
double slit experiment

Note measure $I = |A|^2$
double slit

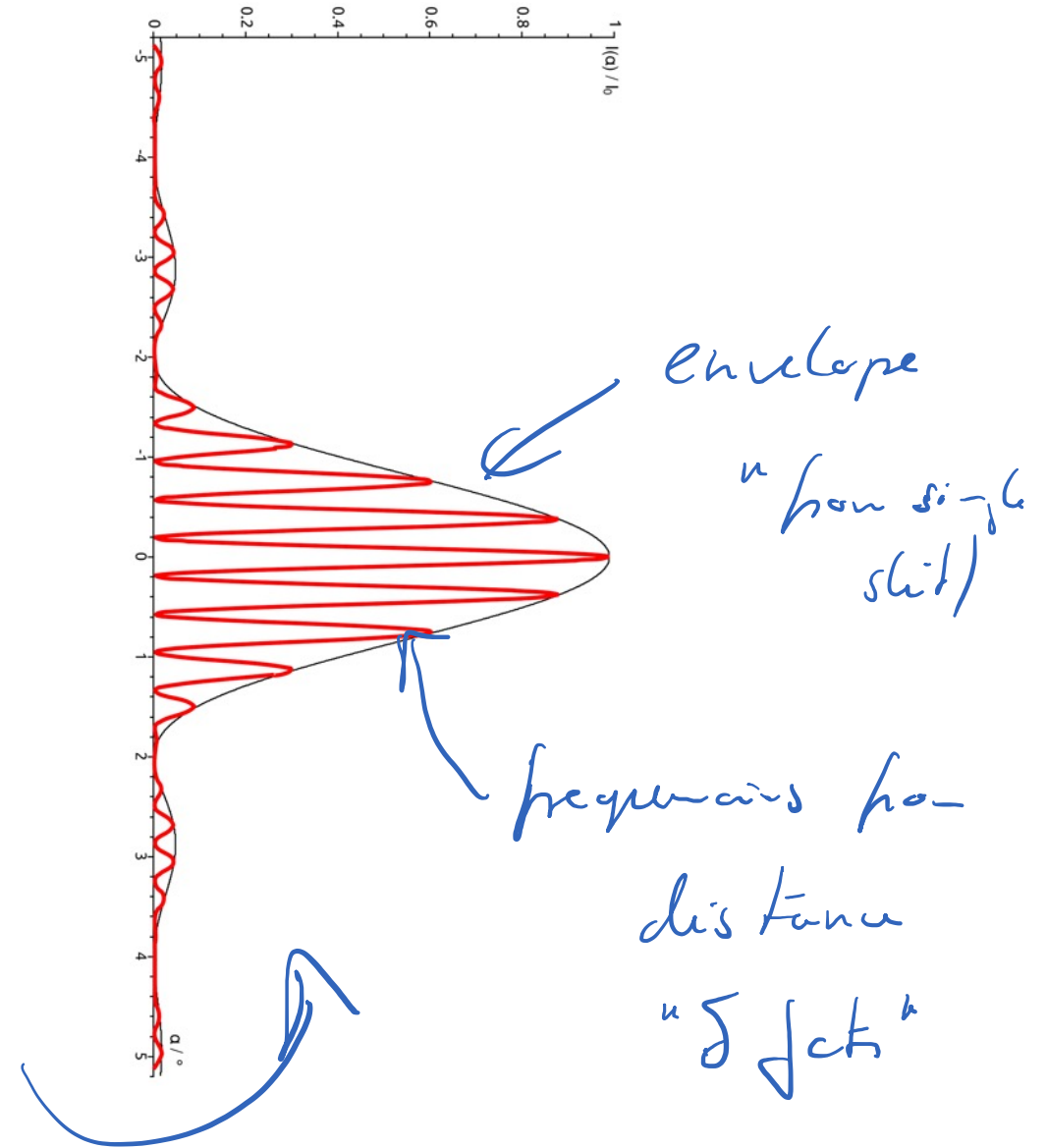
The double slit



↑
double slit

↳ Convolution theorem

$$F(f \otimes g) = F(f) \cdot F(g)$$



Fourier relationships – some examples

Relationship between time & frequency domain
real & imag

$$\sin(x) \quad \text{~~~~~} \rightarrow \quad \delta(x) \quad \begin{array}{c} \uparrow \\ \hline x \end{array}$$

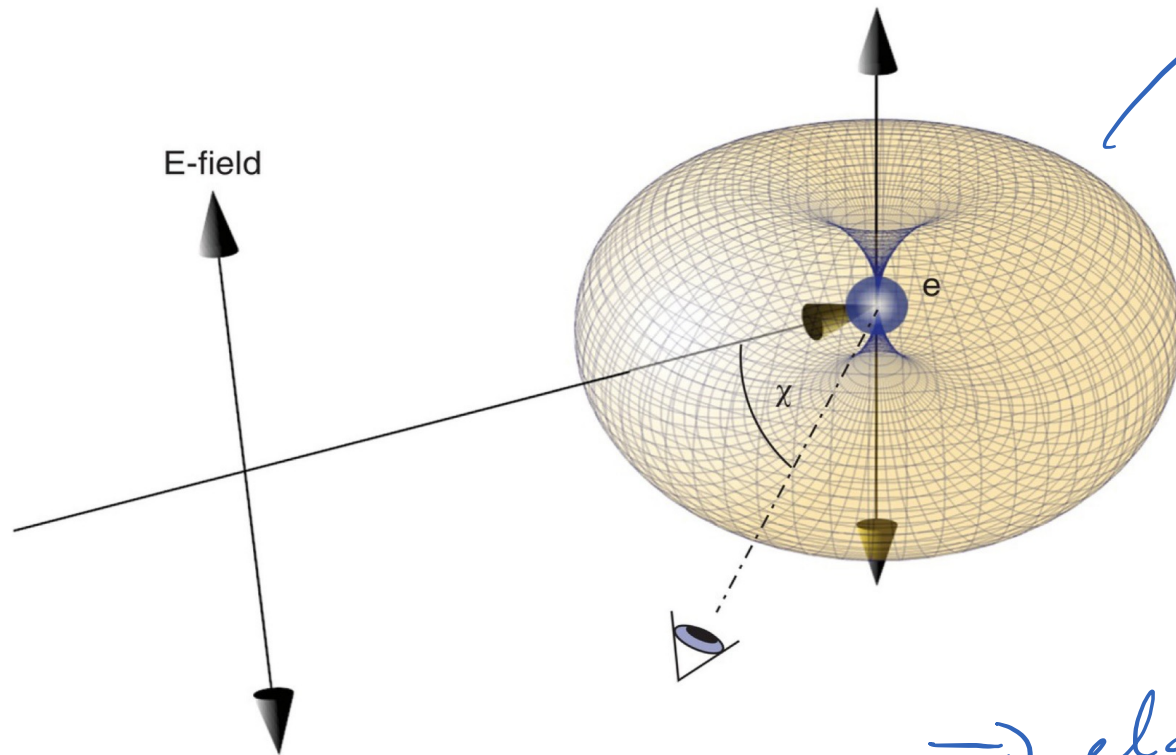
Note

• diffraction pattern $\Rightarrow$ FT of object

• lens is a FT $\rightarrow$ input in focus is $FT(\text{object})$
($\rightarrow$ Opt Method in Chem)

From Session 1

Thomson scattering from a single electron



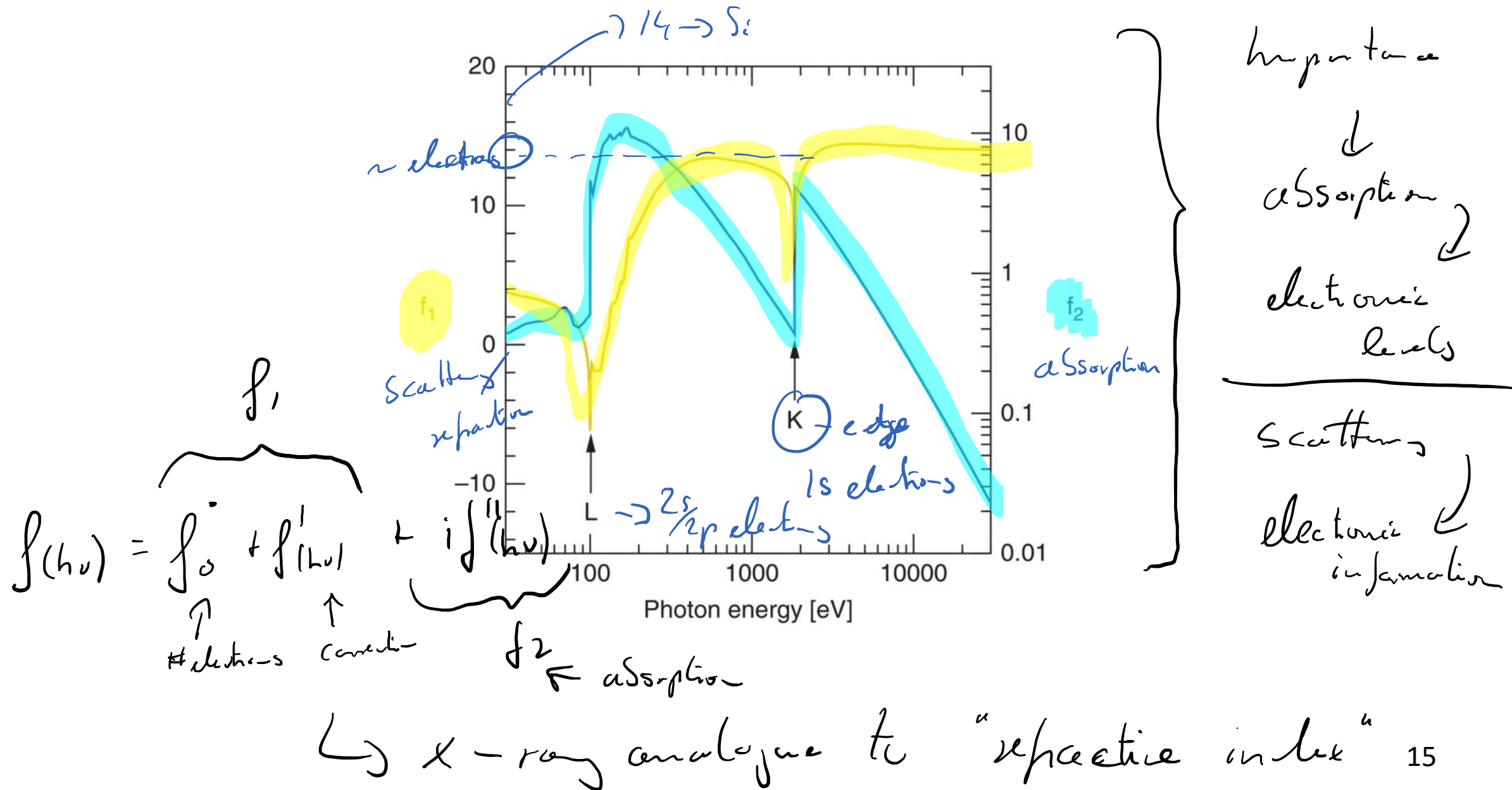
learned x-ray scattering
of electrons
↳ Thomson scattering
↳ classical analog is dipole radiation

⇒ an incoming ^{plane} x-ray wave is converted to spherical wave

⇒ elastic ⇒ no energy change or frequency change
directional change

⇒ can apply principles & ideas from optical regime

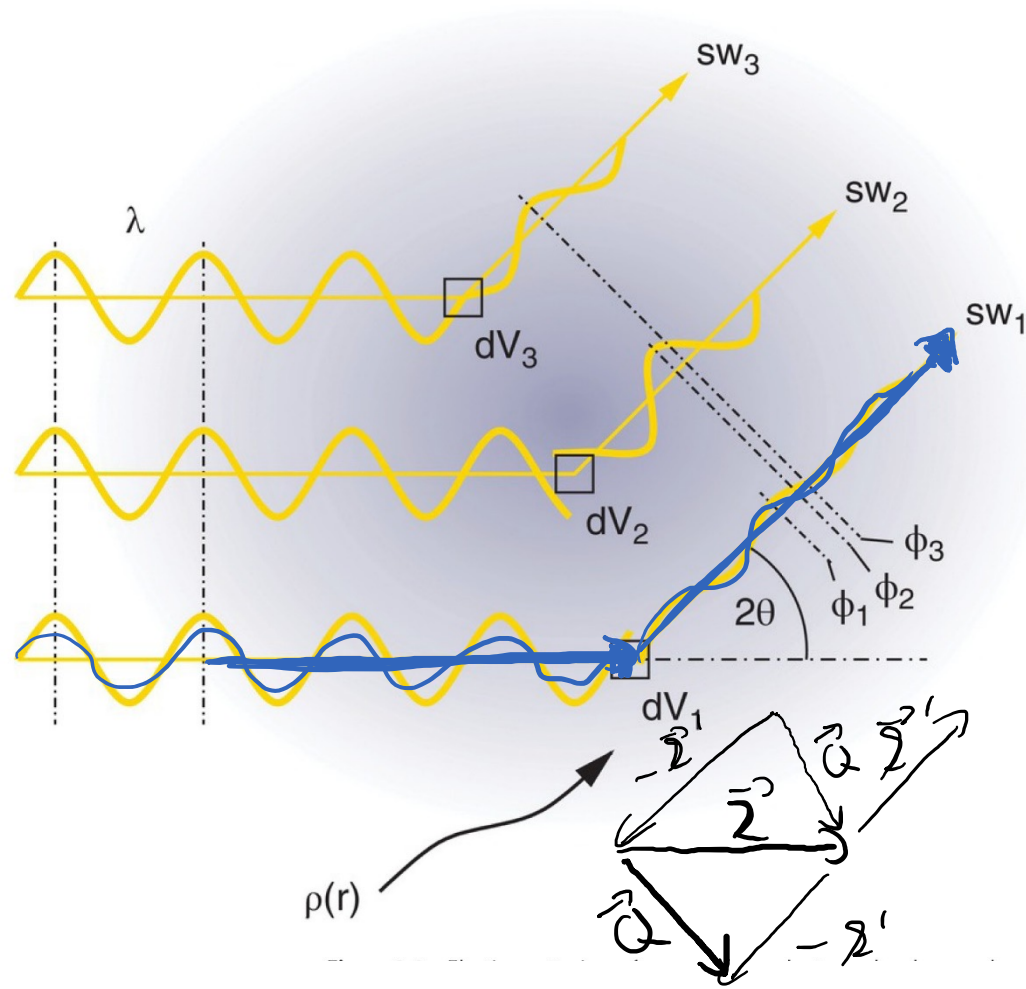
Atomic scattering factors and refractive index



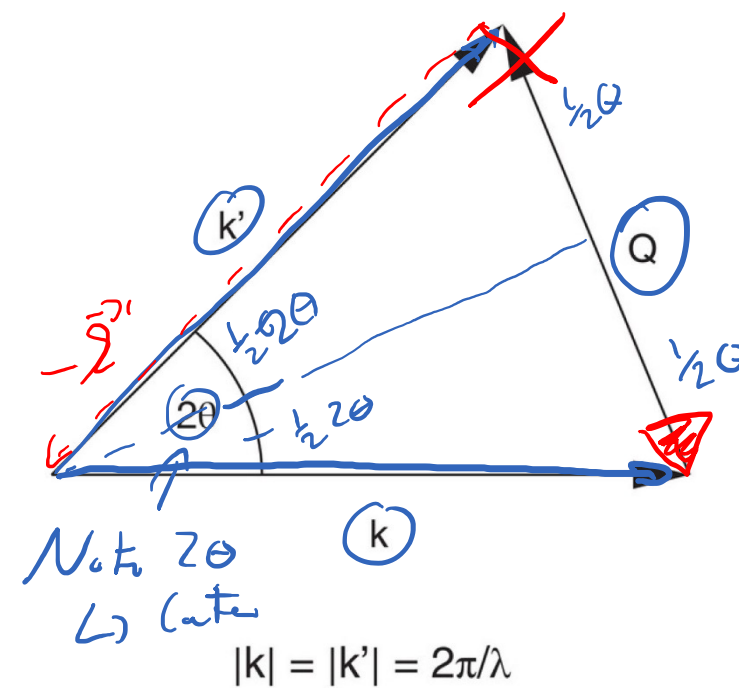
Definition of the scattering vector $\vec{Q}$ $\rightarrow$ Vectorial diffn - between incident & elastically scattered wave vector $\vec{k}, \vec{k}'$

$$\vec{Q} = \vec{k}' - \vec{k}$$

(a)



(b)

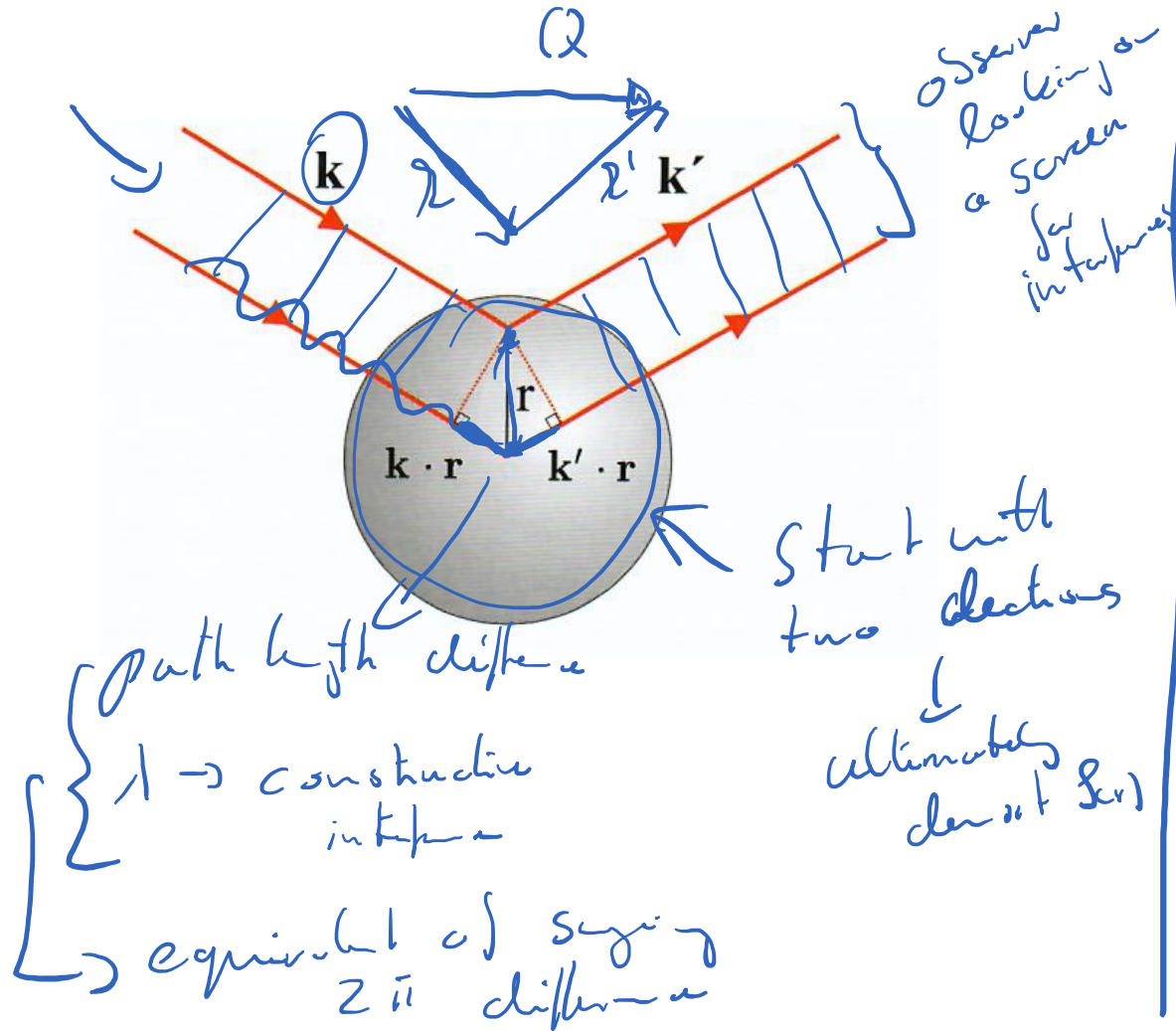


$$\frac{1}{2} Q/k = \sin \theta$$

$$\sin \theta = \frac{Q}{2k}$$

$$Q = 2k \sin \theta$$

Scattering from an electron cloud



→ cloud density $S(r)$

→ constructive interference for
phase difference 2π

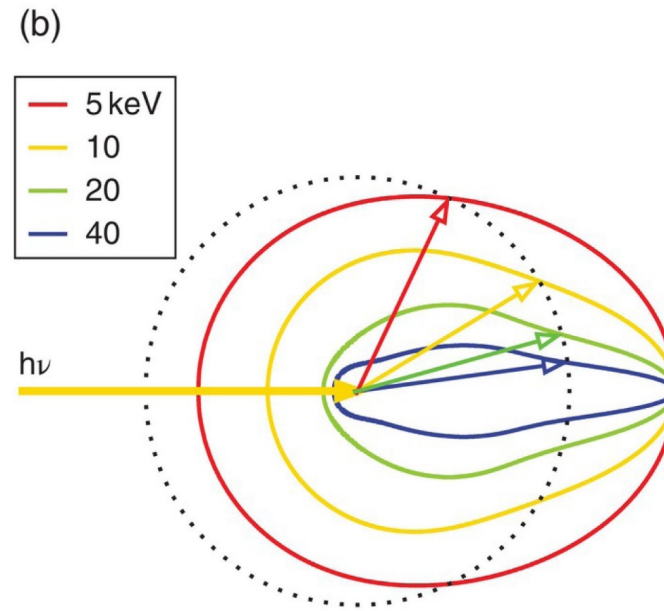
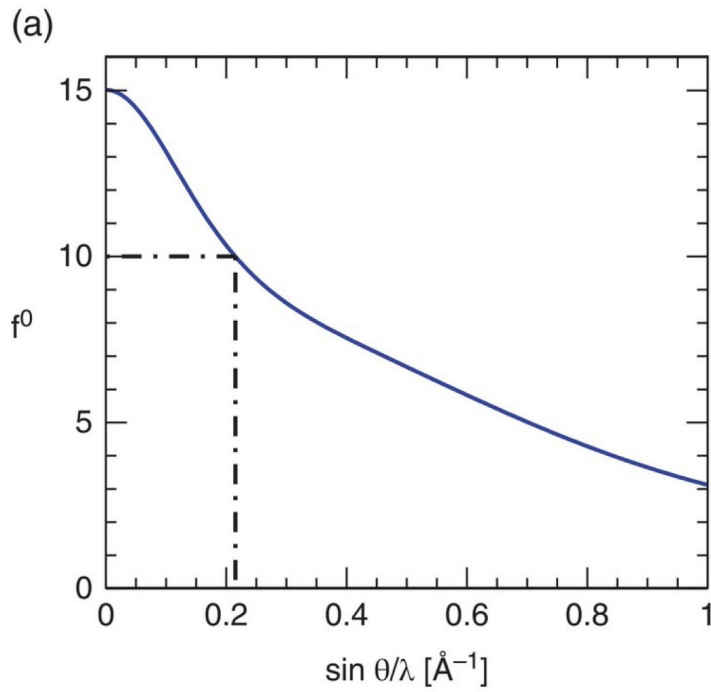
→ $\vec{2} \cdot \vec{r}$ ~ "phase"
(look at projections of $\vec{2}, \vec{r}$)

⇒ At observer the difference
$$\Delta\phi = \vec{2} \cdot \vec{r} - \vec{2}' \cdot \vec{r} = (\vec{2} - \vec{2}') \cdot \vec{r} = \vec{Q} \cdot \vec{r}$$

⇒ elastic $|\vec{2}| = |\vec{2}'|$
$$Q = 2\lambda \sin\theta = \frac{4\pi}{\lambda} \sin\theta$$

⇒ Let's look at volume dr at r
with $e^{i\vec{Q} \cdot \vec{r}}$
- $r_0 S(r) \frac{dr}{\pi \text{ volume element}}$ with $e^{i\vec{Q} \cdot \vec{r}}$
phase factor $e^{i\vec{Q} \cdot \vec{r}}$

Now full cross section / atomic scattering factors



$$- r_0 \int \rho(r) dr \text{ with } e^{iQr}$$

$$- r_0 \int \rho(r) e^{iQr} dr$$

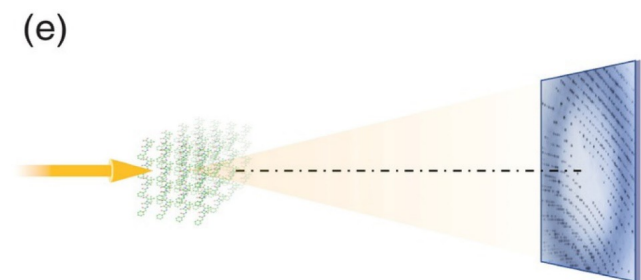
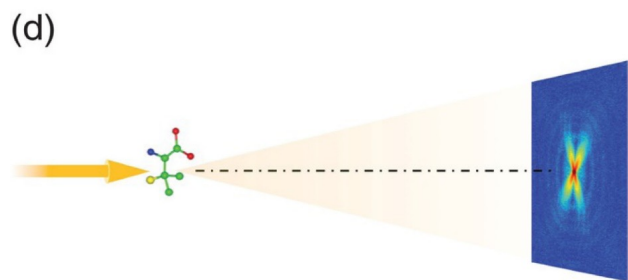
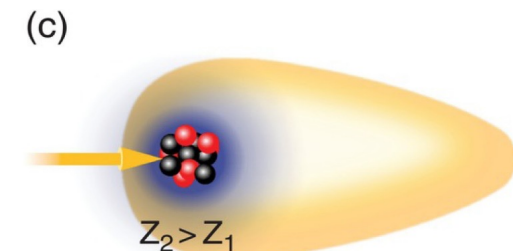
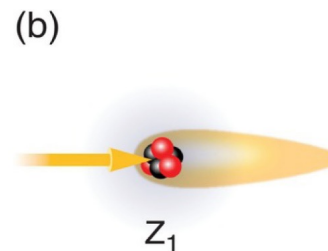
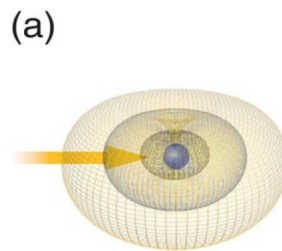
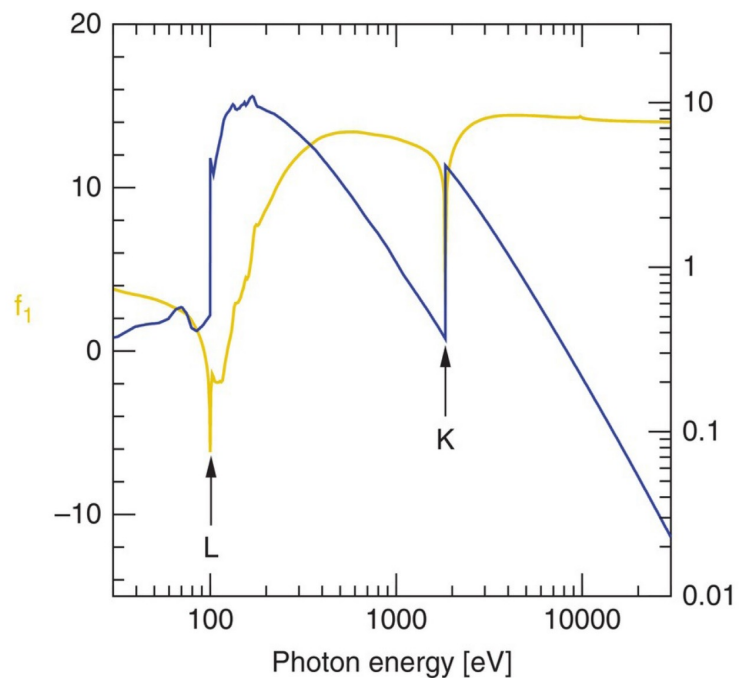
For $Q \rightarrow 0 \Rightarrow Z$
 FT of electron distribution
 "form factor"

$f^0 \rightarrow$ is form factor of atoms
 $\rightarrow Z$ in forward direction
 $\rightarrow$ decreases with increasing scattering angle
 $\Rightarrow$ tabulated

If Q increases, volume elements get out of phase

$\Rightarrow f^0$ is a fct of Q

Atomic scattering factors again



now let $Q \rightarrow 0$
 $Z f - Q = 0$

$$f(Q, h\nu) = \underbrace{f(Q) + f'(h\nu)}_{\text{scattering}} + \underbrace{if''(h\nu)}_{\text{absorption}}$$

From Session 2

Bragg scattering

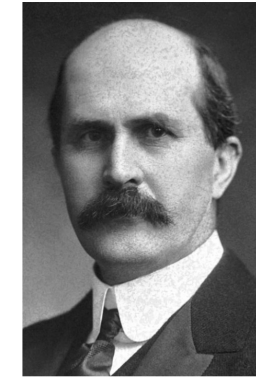


Photo from the Nobel Foundation archive.
Sir William Henry Bragg
Prize share: 1/2

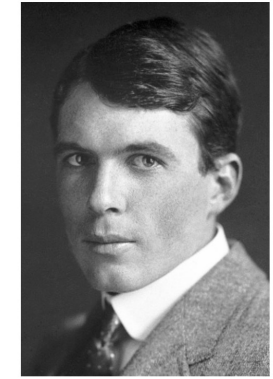
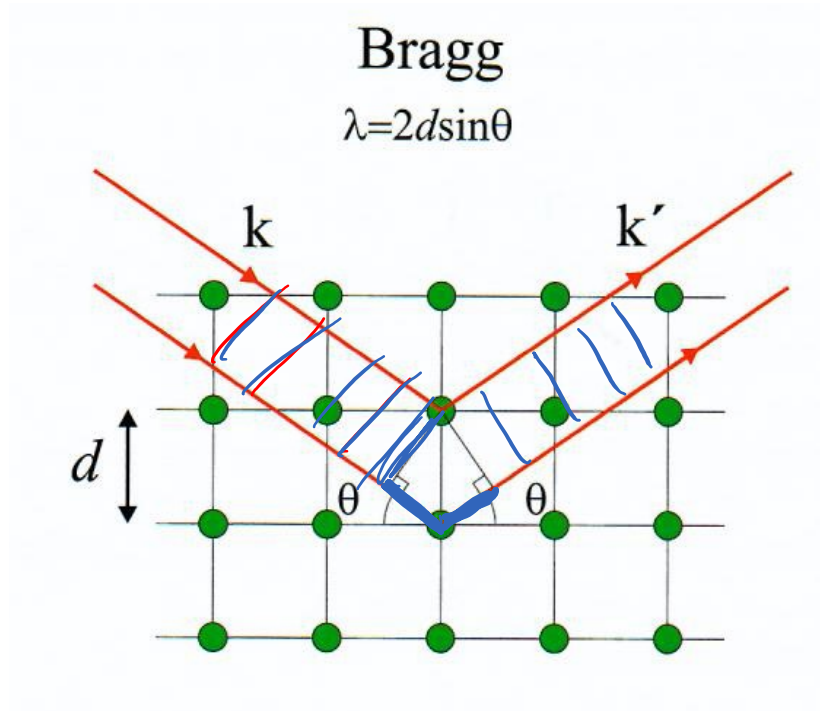
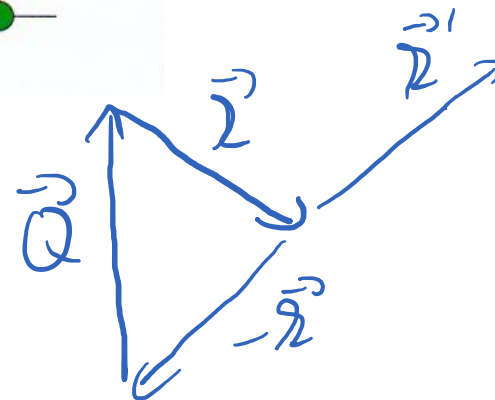
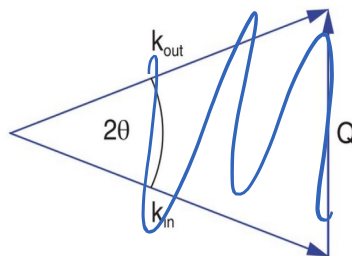


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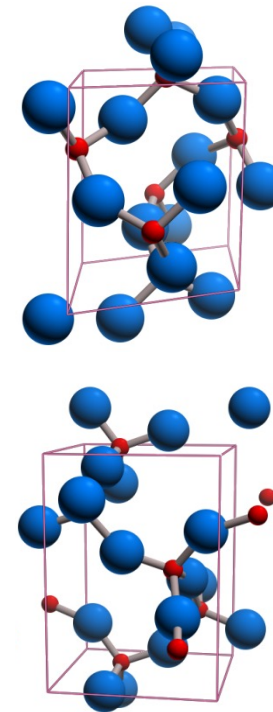
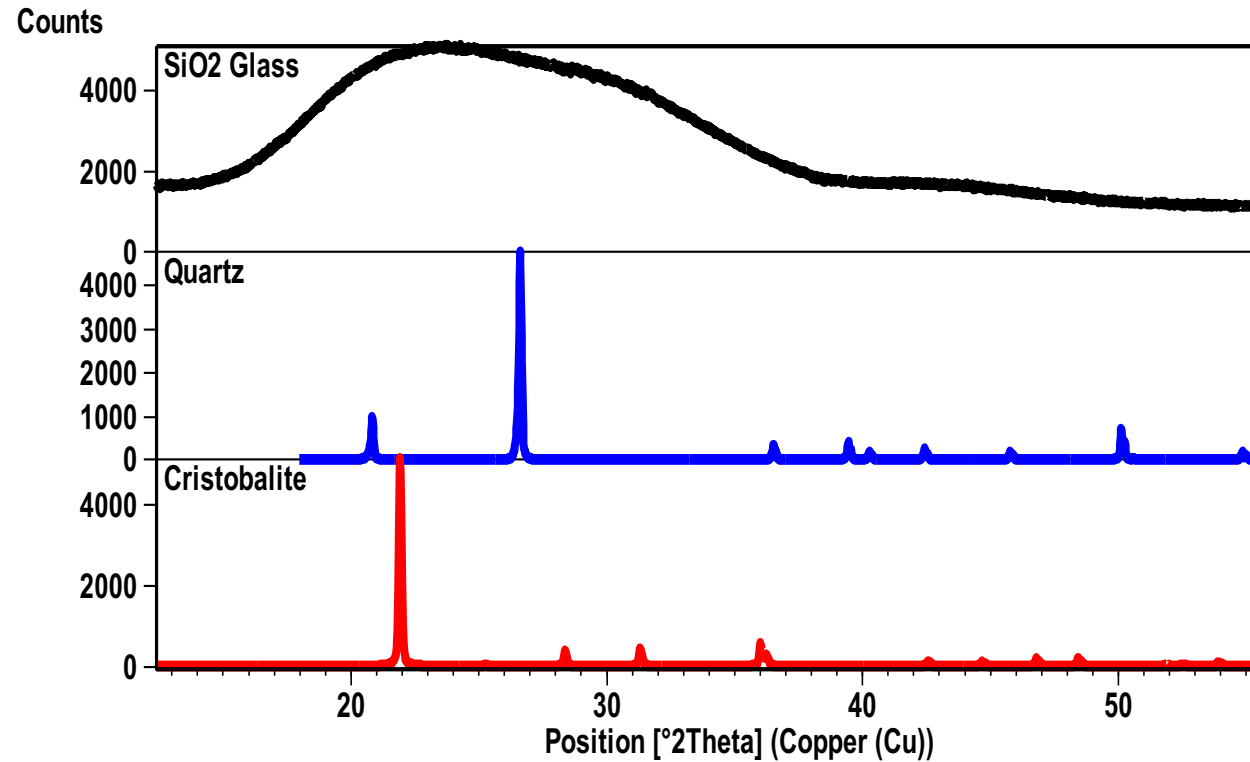


$$\rightarrow \sin \theta = \frac{n\lambda}{2d} \Rightarrow 2d \sin \theta = n\lambda$$

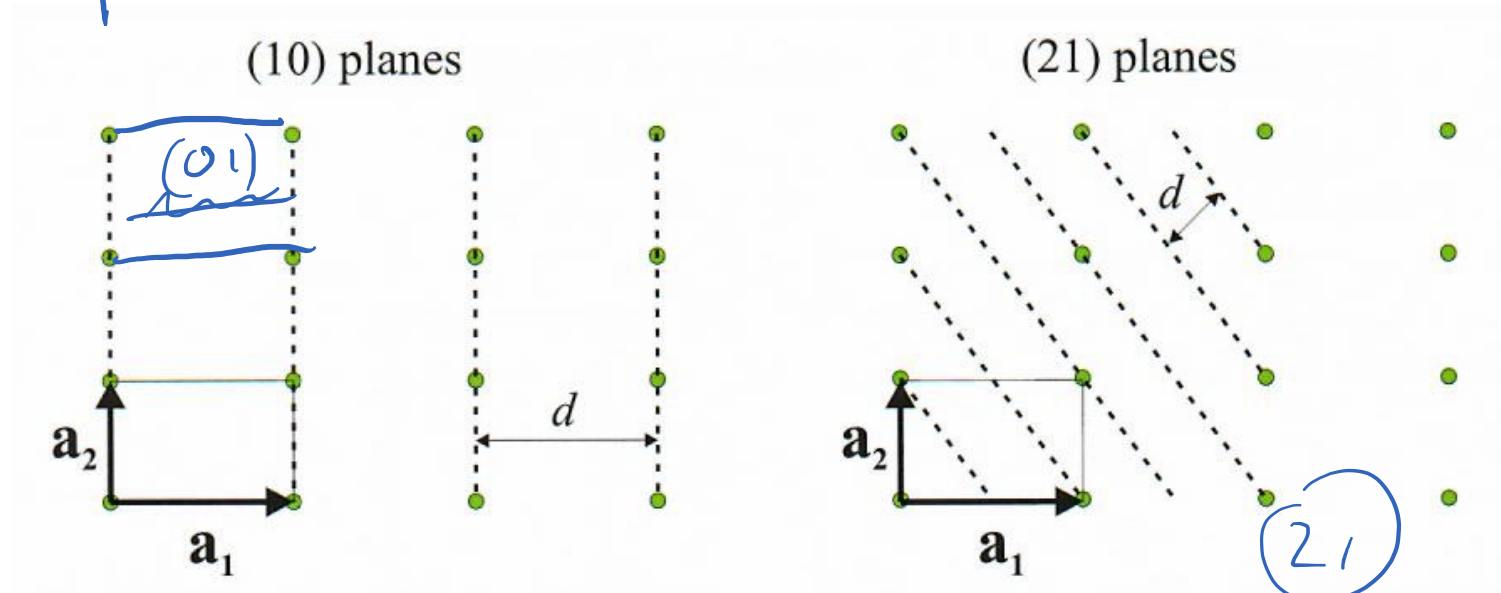
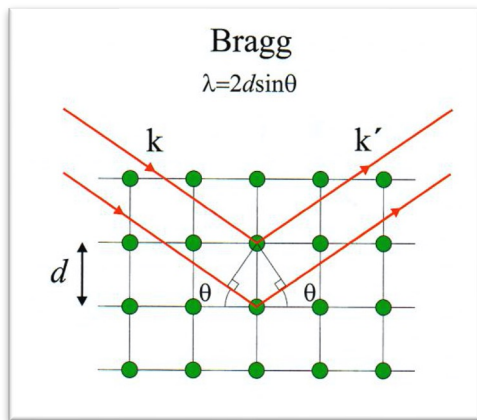


Note $\vec{Q}$ is normal to the lattice planes

Characteristic signals from different – chemically identical – samples

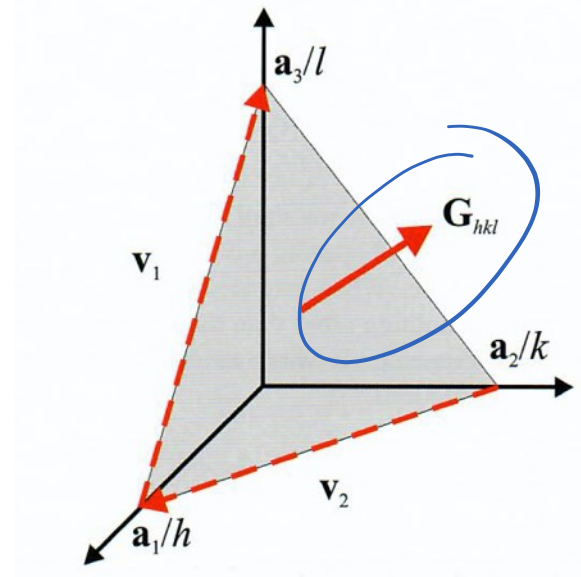
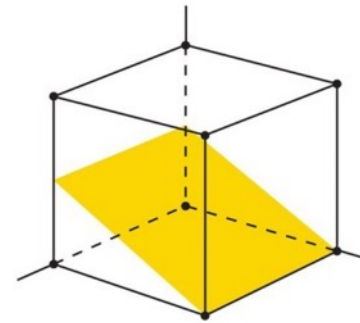
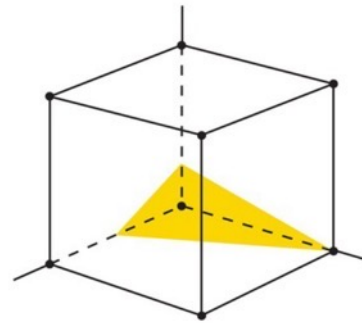
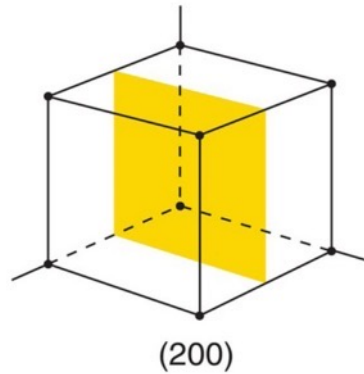
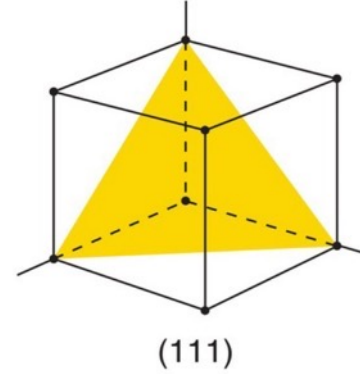
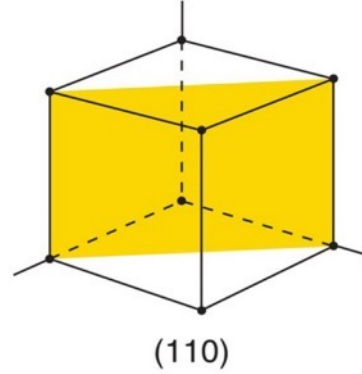
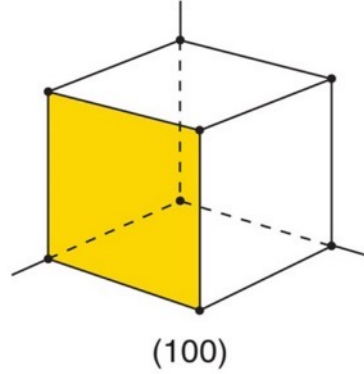


Lattice planes



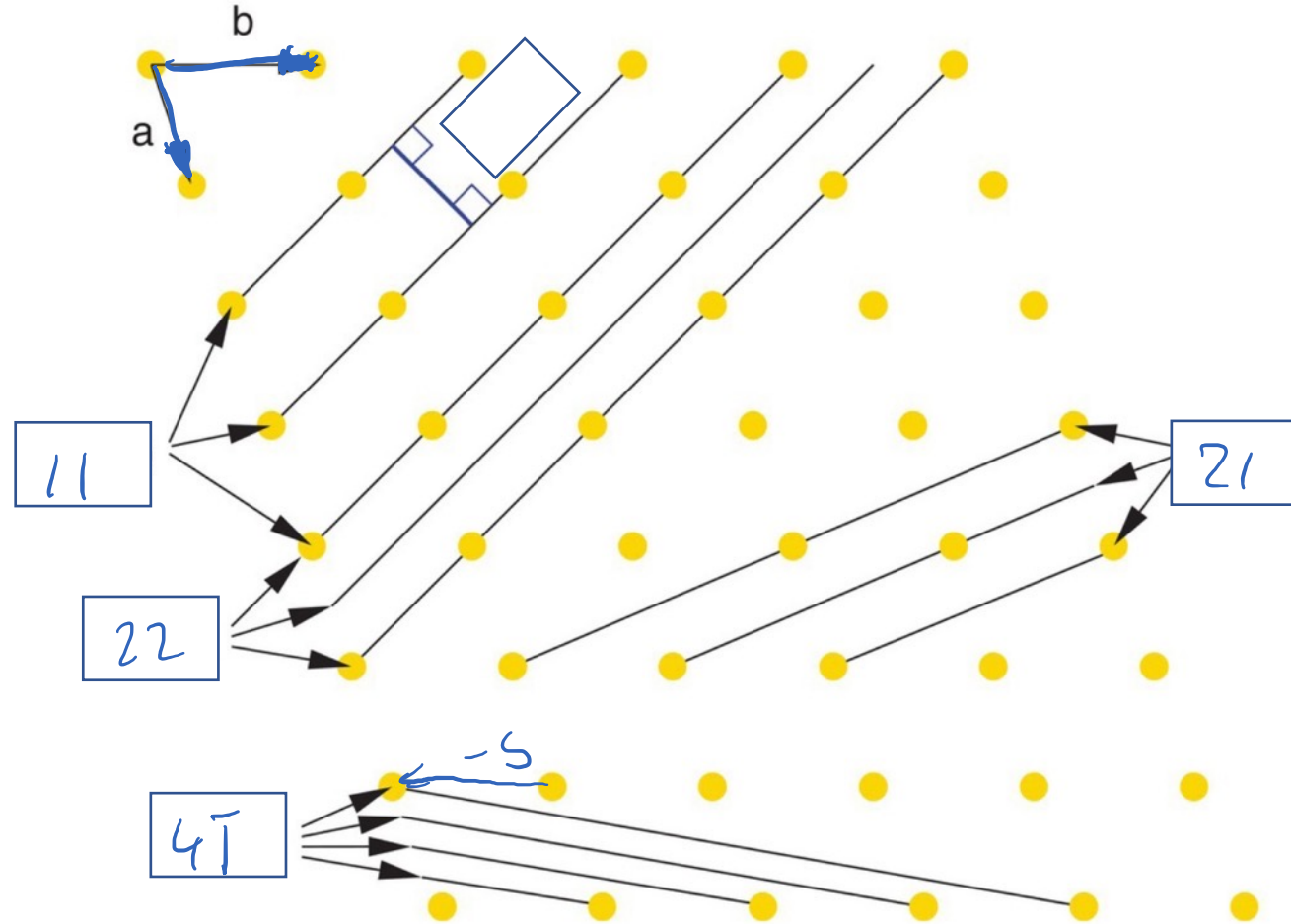
- Find the intercepts of the plane with the respective crystal axis
- Take the reciprocal of these numbers, reduce to smallest integer

Miller indices



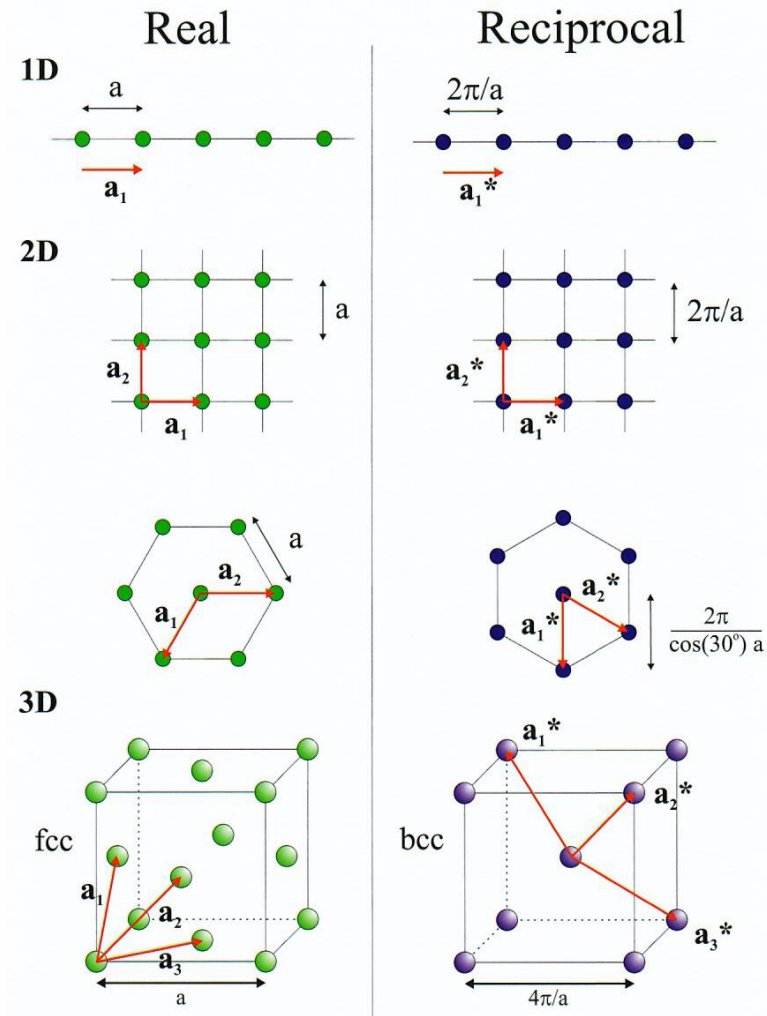
- Find the intercepts of the plane with the respective crystal axis
- Take the reciprocal of these numbers, reduce to smallest integer

Exercise: Identify lattice planes



More on diffraction of single crystals

For each lattice a reciprocal lattice can be defined



→ For each lattice there exists a so-called reciprocal lattice

→ Reciprocal lattice is the Fourier transform of the direct lattice

→ Reciprocal lattice is defined by

$$\vec{G} = h\vec{a}_1^* + k\vec{a}_2^* + l\vec{a}_3^*$$

perpendicular to planes
h k l

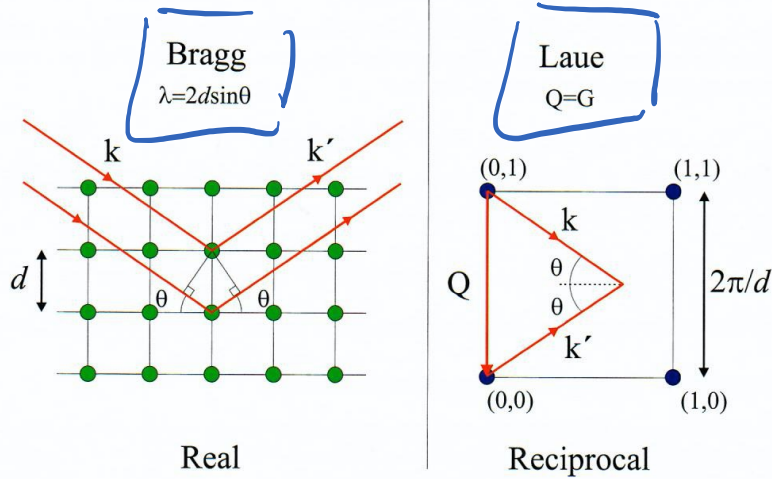
$$|\vec{G}_{hkl}| = \frac{2\pi}{d_{hkl}} \rightarrow \text{lattice planes}$$

Bragg and Laue conditions

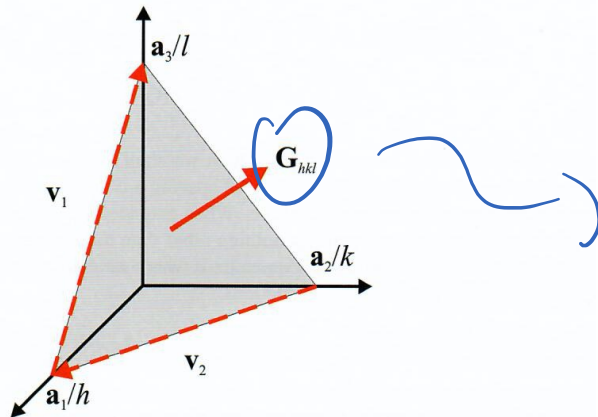
Real

Reciprocal

(a) Equivalence of Bragg and Laue



(b) Miller indices and reciprocal lattice vectors



In reciprocal space

One can show that
Constructive interference
is observed for

$$\vec{Q} = \vec{G} \quad \text{Laue condition}$$

$\vec{G} \rightarrow$ perpendicular to the lattice planes
with Miller indices hkl

$$|\vec{G}| = \frac{2\pi}{d_{hkl}} \Rightarrow d = \text{lattice spacing of } hkl \text{ planes}$$



Photo from the Nobel Foundation archive.
Max von Laue
Prize share: 1/1

The Ewald Sphere

→ elegant geometric construction of the relationships

→ wave vector $\vec{k}$

→ diffraction angle

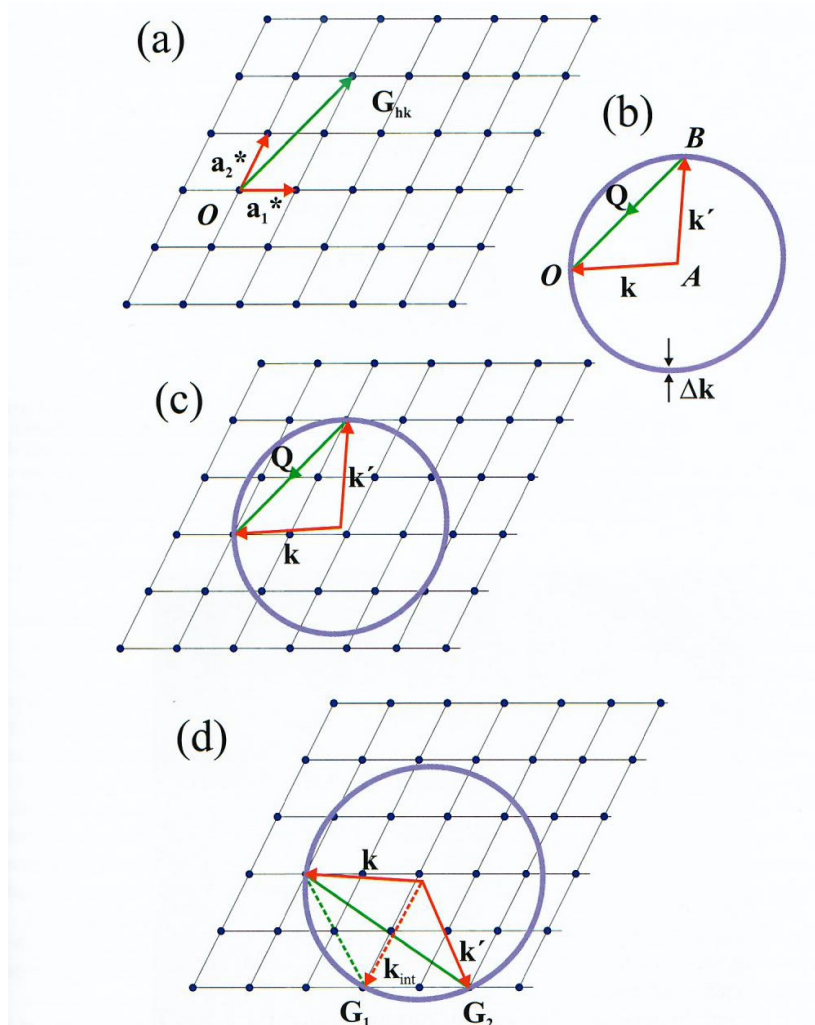
→ in the reciprocal lattice

• Laue condition $\vec{Q} = \vec{G}_{hkl}$

• Create sphere (circle) with radius $|\vec{k}|$

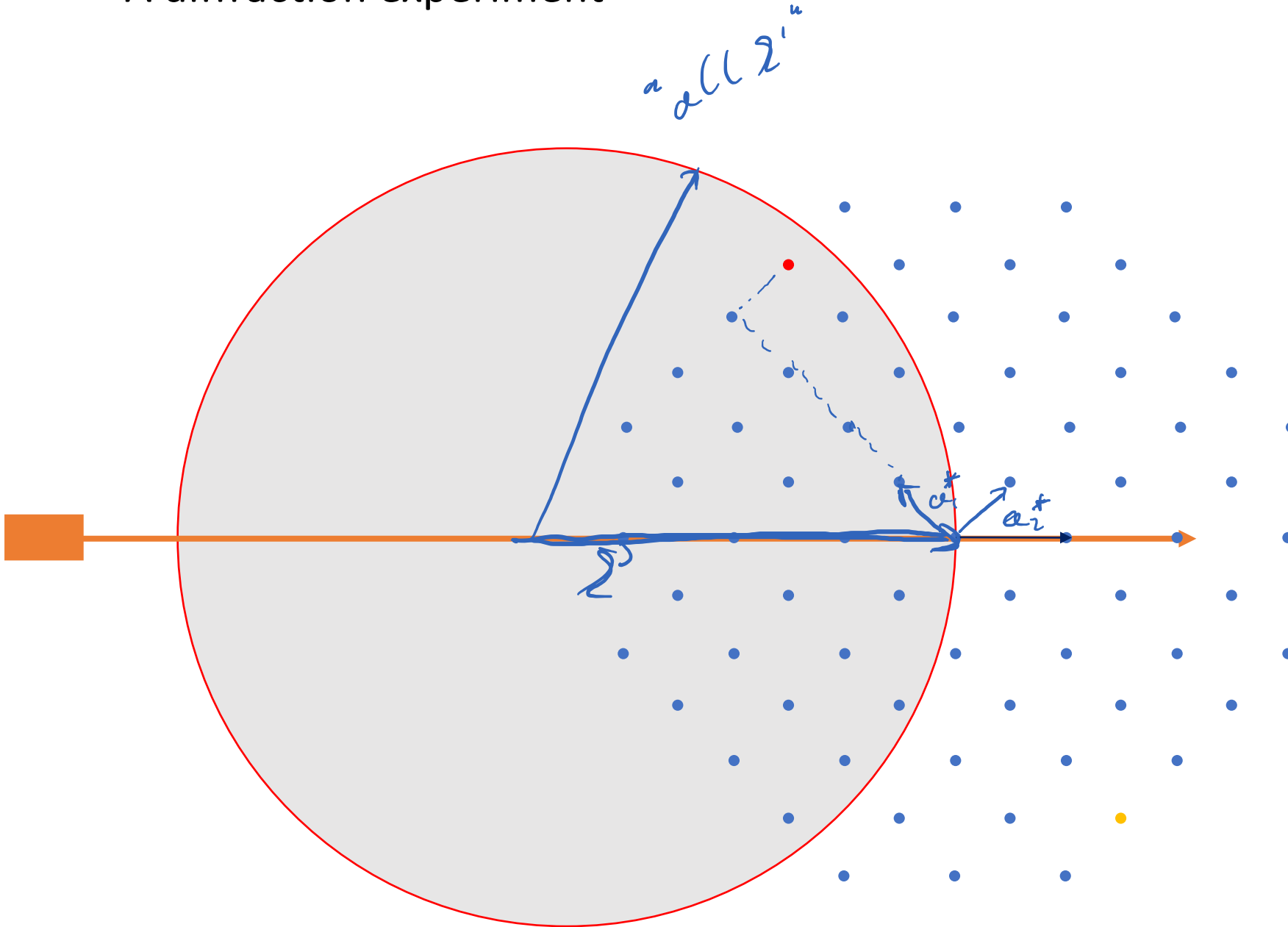
• Move sphere on "origin" O

• Monochromatic light can be scattered constructively ("signal") whenever sphere overlaps with reciprocal lattice point as here $\vec{Q} = \vec{G}$ ✓



A Laue diffraction experiment

A diffraction experiment

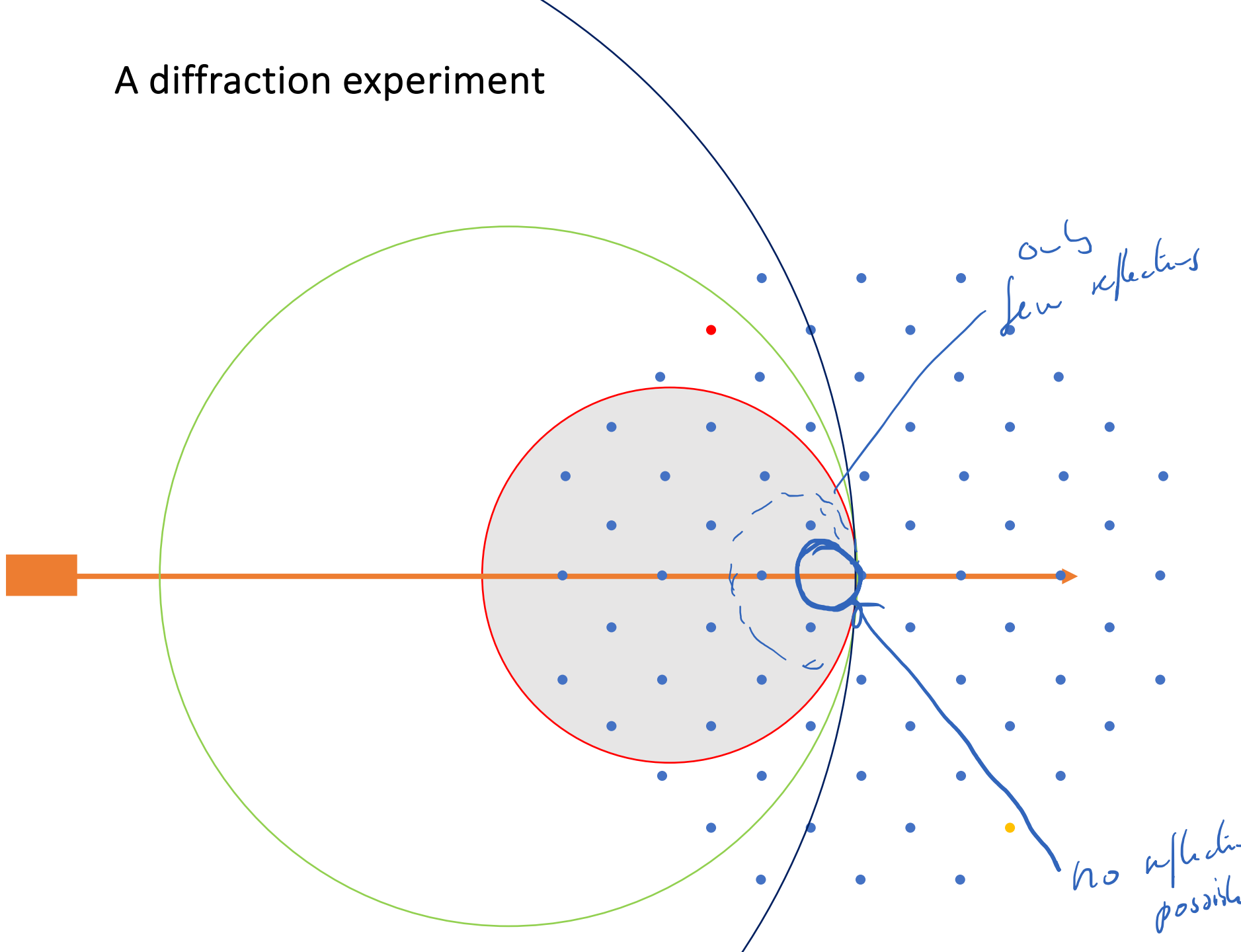


reciprocal lattice
fixed with respect
to real lattice

↓
rotate crystal
until we see
reflections

↓
Reflection angles
from respective
points

A diffraction experiment

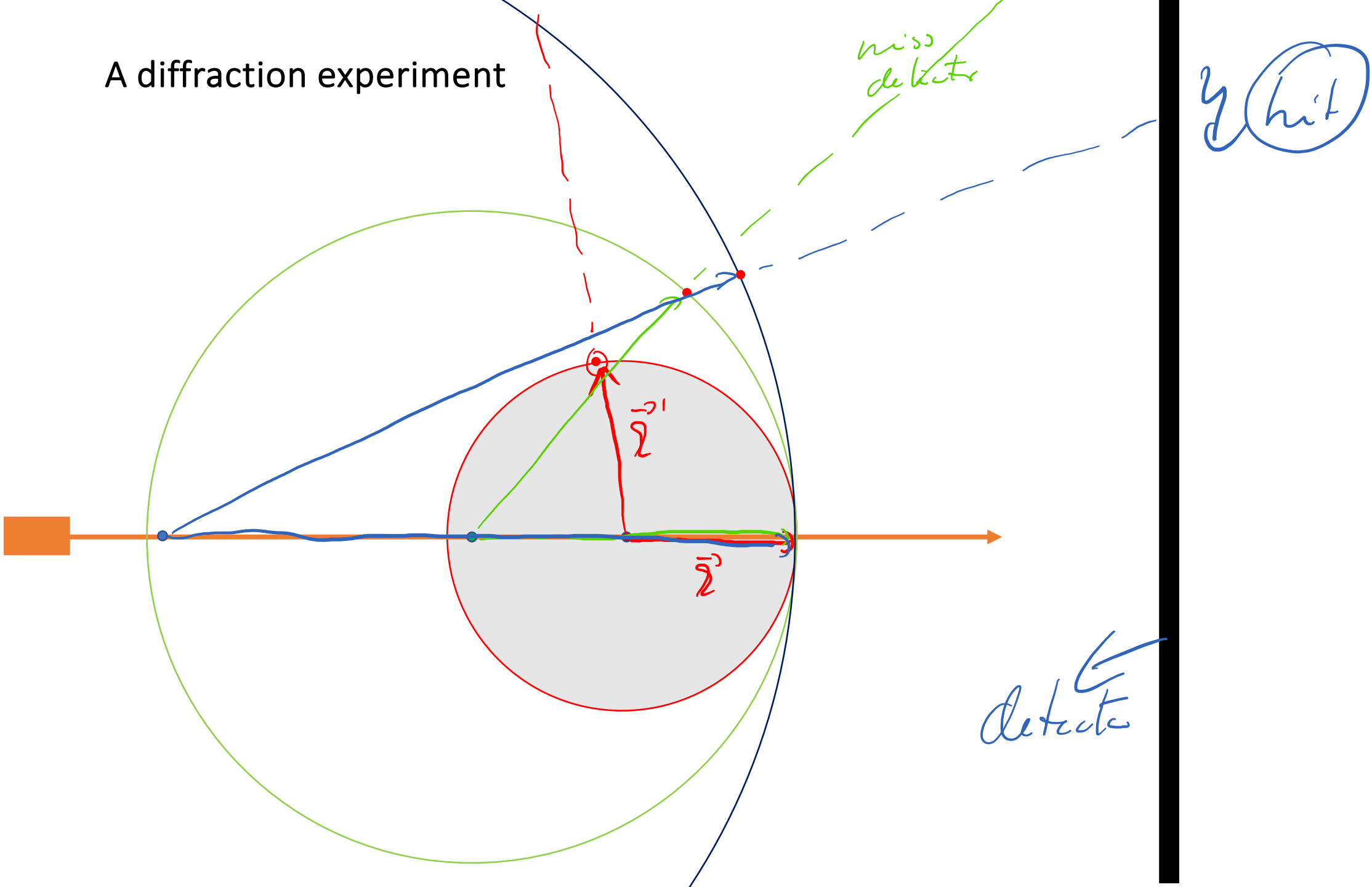


See dependence
of reflections as
fct of $|\vec{Q}|$
i.e. energy

Radius of sphere
 $\sim |\vec{Q}|$
 $\sim h\nu$

More accessible
reflections for
higher energies
 $\downarrow$
higher resolution

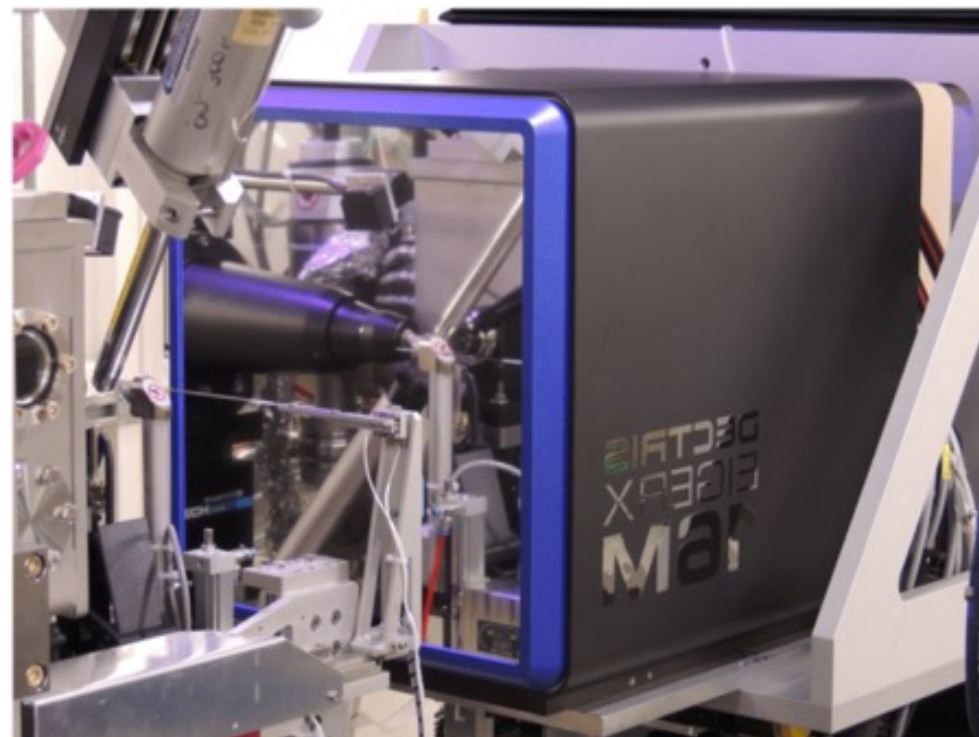
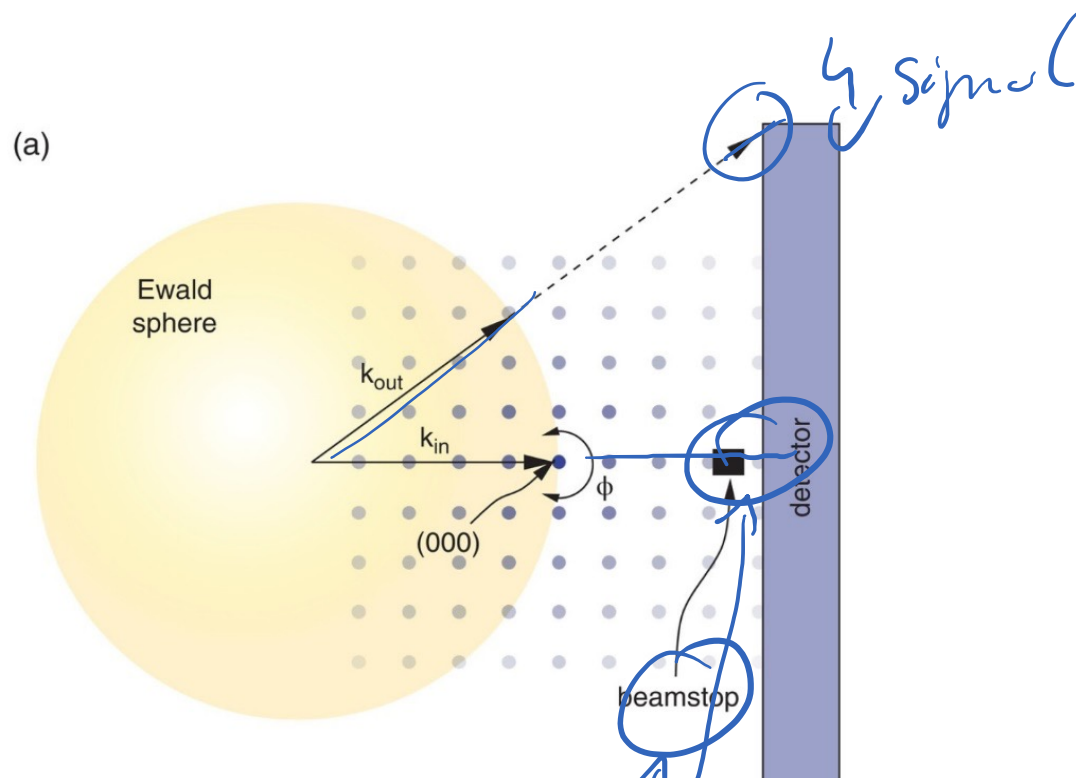
A diffraction experiment



A diffraction experiment

In CoS

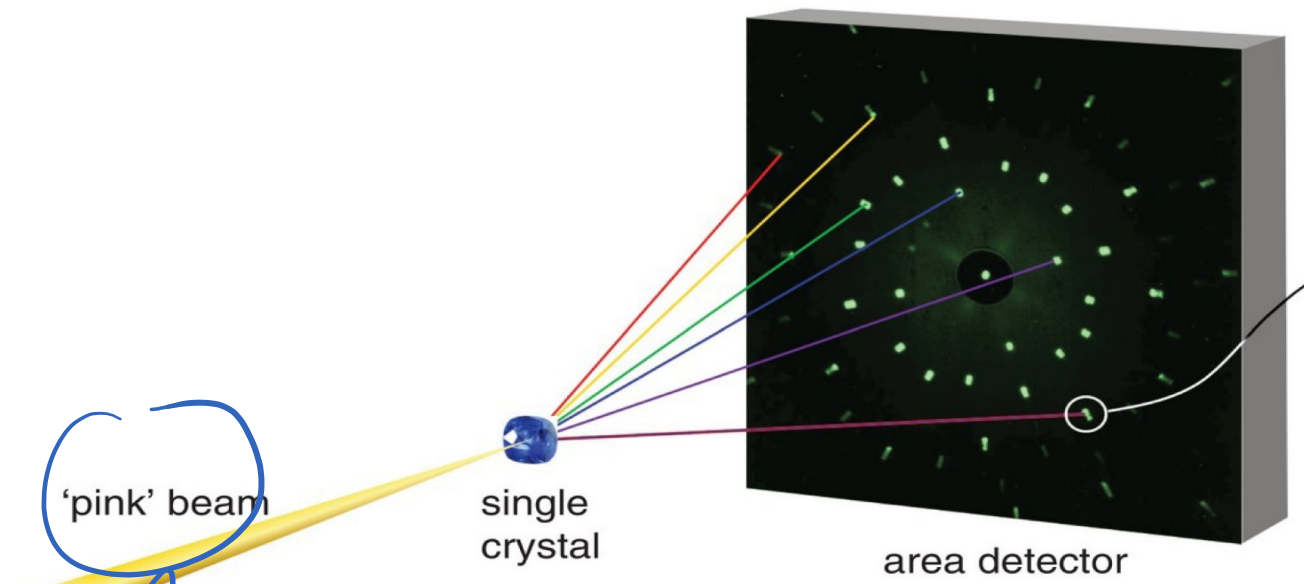
make source & small detector and
 $\rightarrow$ quite inefficient



important so that
 you don't kill
 detector

At S.R. facilities $\rightarrow$ rotate sample
 use an detector
 $\Rightarrow$ efficient

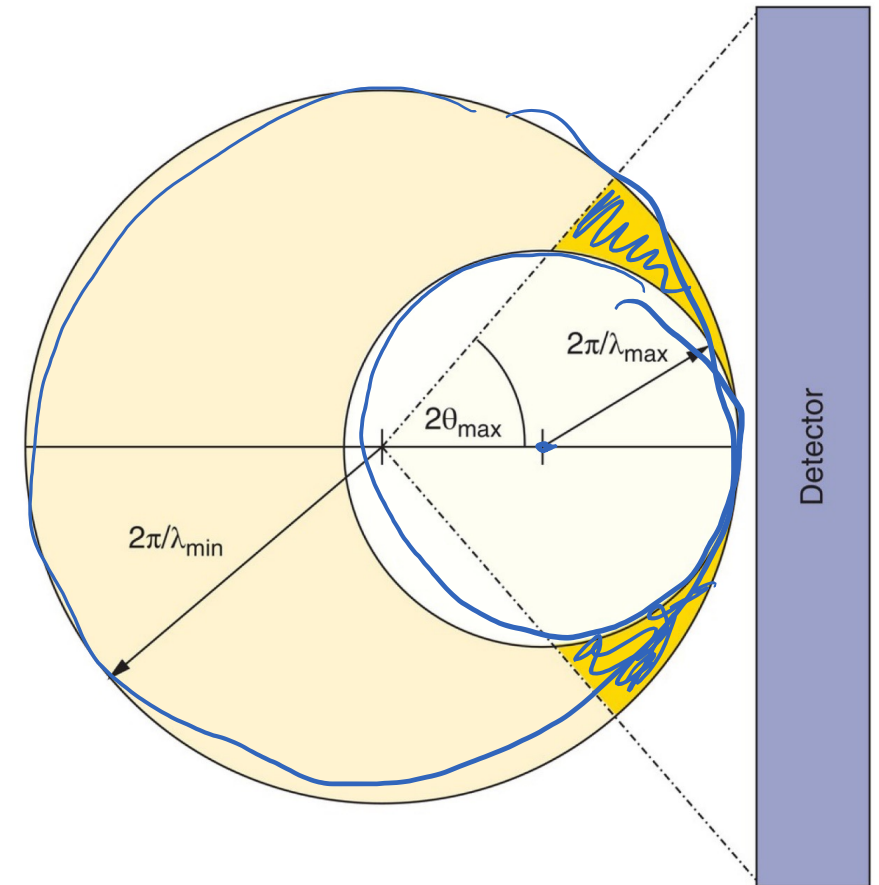
Laue method



"bandwidth"

pink beam

↓
all reflections within
bandwidth are
visible



The end